

## Chapter 14

# The Classical Statistical Treatment of an Ideal Gas

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# 14.1 Thermodynamic Properties from the Partition Function

$$S = \frac{U}{T} + Nk(\ln Z - \ln N + 1) \quad \text{for Maxwell-Boltzmann Distribution}$$

$\sum g_j e^{-\varepsilon_j/kT}$        $6 \times 10^{23}$  molecules       $S = k \ln W$

Partition function

Helmholtz function     $F = U - TS = -NkT(\ln Z - \ln N + 1)$

cf.     $dF = -SdT - PdV + \mu dN$

$$\mu = \left( \frac{dF}{dN} \right)_{T,V} \quad \mu = -kT(\ln Z - \ln N)$$

# 14.1 Thermodynamic Properties from the Partition Function

## (1) Internal energy

$$U = \sum N_j \varepsilon_j = \sum g_j \frac{N}{Z} e^{-\varepsilon_j/kT} \varepsilon_j$$

Partition function  $Z = \sum g_j e^{-\varepsilon_j/kT}$

$$\left. \frac{\partial Z}{\partial T} \right)_V = \sum g_j \frac{\varepsilon_j}{kT^2} e^{-\varepsilon_j/kT}$$

$\varepsilon_j$  fixed 

$$\therefore U = \frac{N}{Z} kT^2 \left. \frac{\partial Z}{\partial T} \right)_V = NkT^2 \left. \frac{\partial \ln Z}{\partial T} \right)_V$$

# 14.1 Thermodynamic Properties from the Partition Function

## (2) Gibbs function

$$G = \mu N = -NkT(\ln Z - \ln N)$$

## (3) Enthalpy

$$H = G + TS$$

$$= -NkT(\ln Z - \ln N) + T \left[ \frac{U}{T} + Nk(\ln Z - \ln N + 1) \right]$$

$$= U + NkT$$

$$= NkT \left[ 1 + T \left( \frac{\partial \ln Z}{\partial T} \right)_V \right]$$

# 14.1 Thermodynamic Properties from the Partition Function

## (4) Pressure

$$P = - \left( \frac{dF}{dV} \right)_{T,N} = NkT \left( \frac{\partial \ln Z}{\partial V} \right)_T$$

cf.  $dF = -SdT - PdV + \mu dN$

# 14.2 Partition Function for a Gas

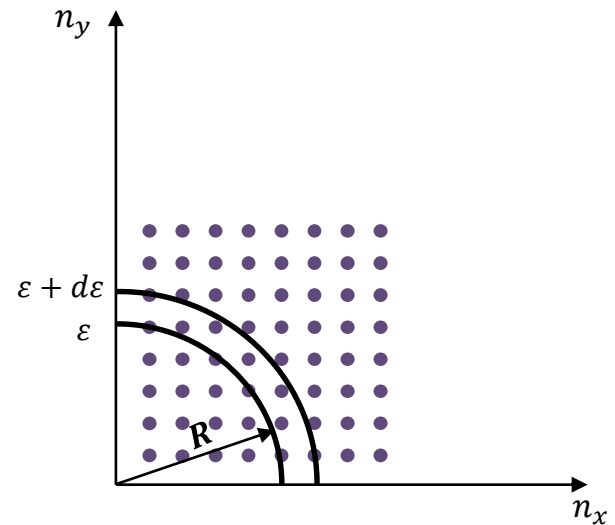
- Partition function

$$\begin{aligned}
 Z &= \sum g_j e^{-\varepsilon_j/kT} \\
 &= \int g(\varepsilon) e^{-\varepsilon/kT} d\varepsilon = \int_0^\infty \frac{4\sqrt{2}\pi V}{h^3} m^{3/2} \varepsilon^{1/2} e^{-\varepsilon/kT} d\varepsilon \\
 &= \frac{4\sqrt{2}\pi V}{h^3} m^{3/2} \int_0^\infty \varepsilon^{1/2} e^{-\varepsilon/kT} d\varepsilon \\
 &= V \left( \frac{2\pi m k T}{h^2} \right)^{3/2} \frac{kT}{2} \sqrt{\pi k T}
 \end{aligned}$$

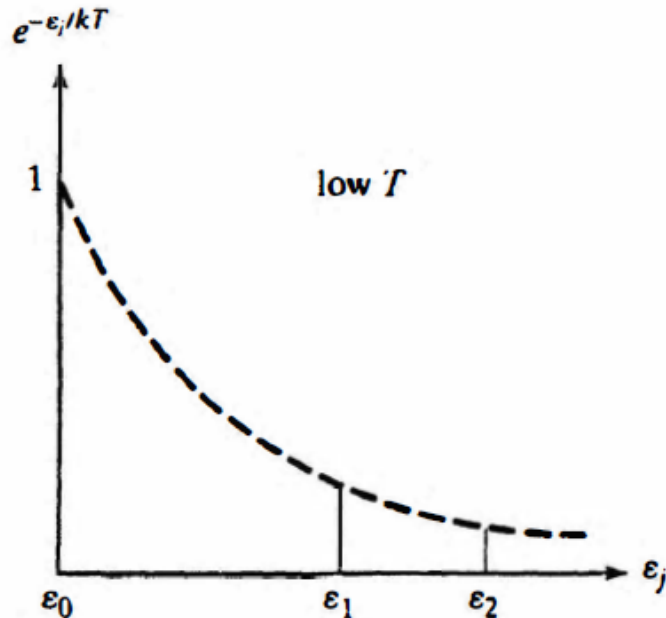
cf.

$g(\varepsilon)d\varepsilon$  : number of quantum state  $\varepsilon < \varepsilon + d\varepsilon$

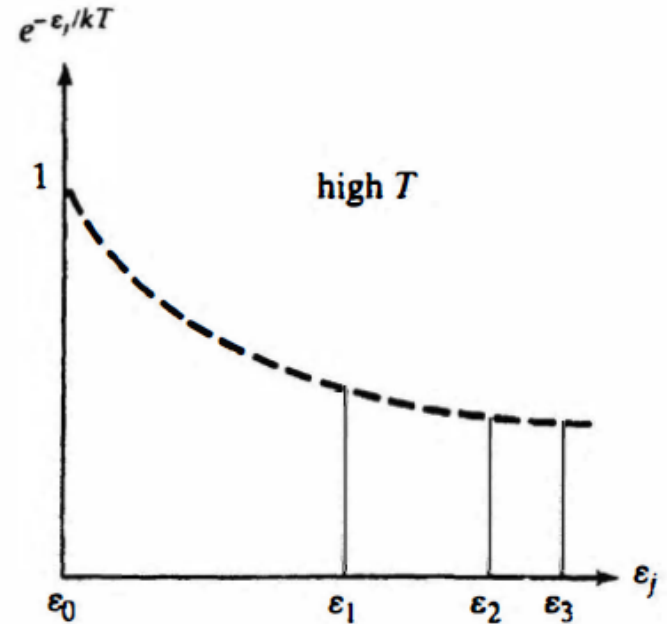
$$g(\varepsilon)d\varepsilon = \frac{\pi}{6} V \left( \frac{2^{9/2} m^{2/3}}{h^3} \right) \frac{3}{2} \varepsilon^{1/2} d\varepsilon = \frac{4\sqrt{2}\pi V}{h^3} m^{3/2} \varepsilon^{1/2} d\varepsilon$$



## 14.2 Partition Function for a Gas



(a)



(b)

Successive terms of the partition function sum for (a) low temperature; and (b) high temperature. The energy level spacings are comparatively large for a small volume.

## 14.3 Properties of a Monatomic Ideal gas

The partition function depends on both volume and temperature

$$\ln Z = \frac{3}{2} \ln T + \ln V + \frac{3}{2} \ln \left( \frac{2\pi mk}{h^2} \right)$$

$$\left( \frac{\partial \ln Z}{\partial V} \right)_T = \frac{1}{V} \quad \left( \frac{\partial \ln Z}{\partial T} \right)_V = \frac{3}{2} \cdot \frac{1}{T}$$

$$P = - \left( \frac{\partial F}{\partial V} \right)_T = NkT \left( \frac{\partial \ln Z}{\partial V} \right)_T = \frac{NkT}{V} \quad \longrightarrow \quad PV = NkT$$

$$U = NkT^2 \left( \frac{\partial \ln Z}{\partial T} \right)_V = NkT^2 \frac{3}{2T} \quad \longrightarrow \quad U = \frac{3}{2} NkT$$



# 14.3 Properties of a Monatomic Ideal gas

The calculation of entropy

$$S = \frac{U}{T} + Nk(\ln Z - \ln N + 1)$$

$$= \frac{U}{T} + Nk \left( \frac{3}{2} \ln T + \ln V + \frac{3}{2} \ln \left( \frac{2\pi mk}{h^2} \right) - \ln N + 1 \right)$$

$$= \frac{5}{2} Nk + Nk \ln \left[ \frac{V(2\pi mkT)^{3/2}}{Nh^3} \right] \quad \text{Sackur-Tetrode Equation}$$

## 14.3 Properties of a Monatomic Ideal gas

Statistical approach

$$s = S/n, \quad Nk = n\bar{R} \quad c_v = \frac{3}{2}R$$

$$s = c_v \ln T + R \ln v + s_0$$

Classical approach

$$Tds = du + Pdv$$

$$ds = c_v \frac{dT}{T} + R \frac{dv}{v}$$

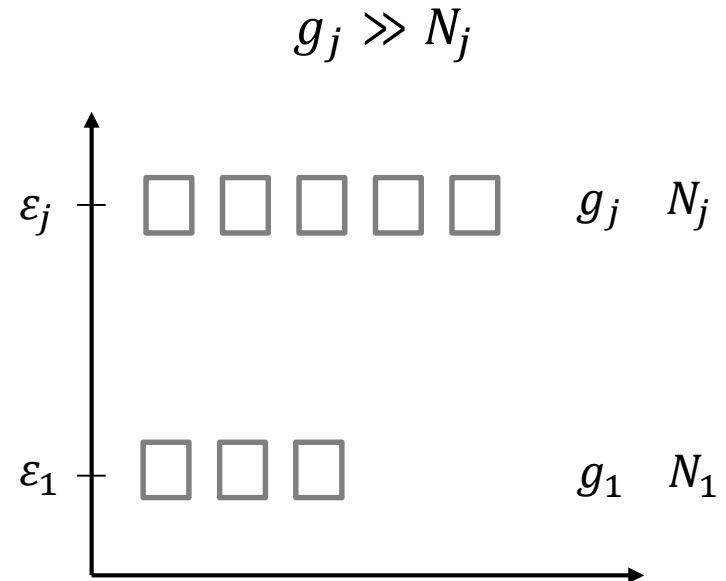
# 14.4 Applicability of MB Distribution

Maxwell-Boltzmann distribution: valid under dilute gas

$$Z = V n_Q = V \left( \frac{2\pi m k T}{h^2} \right)^{3/2}$$

Maxwell-Boltzmann distribution

$$\frac{N_j}{g_j} = \frac{N}{Z} e^{-\varepsilon_j/kT} = \left( \frac{N}{V} \right) \frac{1}{n_Q} e^{-\varepsilon_j/kT}$$



## 14.4 Applicability of MB Distribution

For Helium,

$$m = 6.65 \times 10^{-27} \text{ kg}$$

$$T = 273 \text{ K}$$

$$k = 1.381 \times 10^{-23} \text{ J/K}$$

$$h = 6.626 \times 10^{-34} \text{ J}\cdot\text{s}$$

$$\longrightarrow n_Q = 7 \times 10^{30} \text{ m}^{-3}$$

Number of particles  
per unit volume

$$\frac{N}{V} = \frac{6 \times 10^{26}}{22.4 \text{ m}^3} \approx 3 \times 10^{25} \text{ m}^{-3}$$

$$\frac{N_j}{g_j} = \frac{N}{Z} e^{-\varepsilon_j/kT} = \left(\frac{N}{V}\right) \frac{1}{n_Q} e^{-\varepsilon_j/kT} \sim 4 \times 10^{-6}$$

# 14.5 Distribution of Molecular Speed

$$f(\varepsilon) = \frac{N(\varepsilon)}{g(\varepsilon)} = \frac{N}{Z} e^{-\varepsilon/kT} \quad \text{Maxwell-Boltzmann distribution}$$

$N(\varepsilon)d\varepsilon$  : # of particles  $\varepsilon < \varepsilon + d\varepsilon$

For an ideal gas,

$$Z = V \left( \frac{2\pi mkT}{h^2} \right)^{3/2}$$

The density of states is given by

$$g(\varepsilon) = \frac{4\sqrt{2}\pi V}{h^3} m^{3/2} \varepsilon^{1/2}$$

# 14.5 Distribution of Molecular Speed

$$\begin{aligned} N(\varepsilon)d\varepsilon &= \frac{Ne^{-\varepsilon/kT}}{V \left(\frac{2\pi mkT}{h^2}\right)^{3/2}} \cdot \frac{4\sqrt{2}\pi V}{h^3} m^{3/2} \varepsilon^{1/2} d\varepsilon \\ &= \frac{2\pi N}{(NkT)^{3/2}} \varepsilon^{1/2} e^{-\varepsilon/kT} d\varepsilon \end{aligned}$$

$\varepsilon$  is the single particle kinetic energy. Thus

$$\varepsilon = \frac{1}{2}mv^2, \quad d\varepsilon = mv dv$$

$$\varepsilon^{1/2} d\varepsilon = \frac{(mv^2)^{1/2}}{\sqrt{2}} mv dv = \frac{1}{\sqrt{2}} m^{3/2} v^2 dv$$

# 14.5 Distribution of Molecular Speed

$N(v)dv$  : # of particles  $v < v < v + dv$

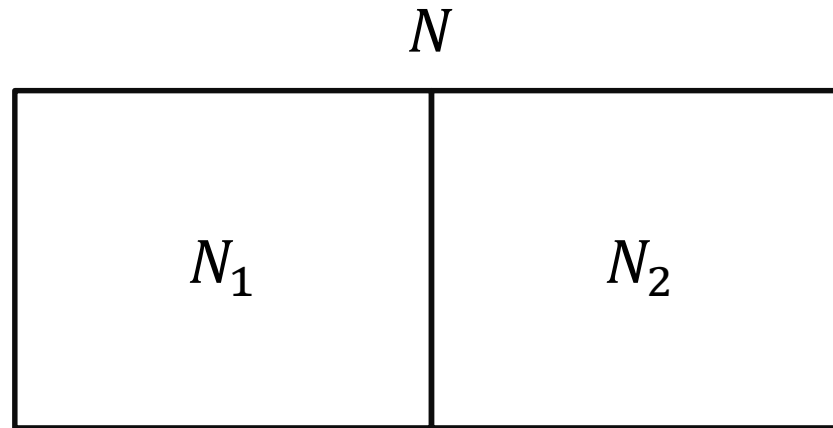
$$\begin{aligned} N(v)dv &= \frac{2\pi N}{(2\pi kT)^{3/2}} \cdot \frac{1}{\sqrt{2}} m^{3/2} v^2 dv \cdot e^{-mv^2/2kT} \\ &= 4\pi N \left( \frac{m}{2\pi kT} \right)^{3/2} v^2 e^{-mv^2/2kT} dv \end{aligned}$$

cf. from kinetic theory,

$$dN_v = \underbrace{(N\alpha^3 e^{-\beta^2 v^2})}_{\rho} \times \underbrace{(4\pi v^2 dv)}_V = 4\pi N\alpha^3 v^2 e^{-\beta^2 v^2} dv$$

$$dN_v = 4\pi N \left( \frac{m}{2\pi kT} \right)^{3/2} v^2 e^{-mv^2/2kT} dv$$

# 14.7 Entropy Change of Mixing



$$\Delta S = -n\bar{R}(x_1 \ln x_1 + x_2 \ln x_2)$$

- **Effect of mixing**

$$\Delta W = \frac{N!}{N_1! N_2!} = \frac{N!}{(x_1 N)! (x_2 N)!}$$



## 14.7 Entropy Change of Mixing

$$\Delta S = k \ln W = k[\ln N! - \ln(x_1 N)! - \ln(x_2 N)!]$$

$$\begin{aligned}\frac{\Delta S}{k} &= N \ln N - N - \{x_1 N \ln(x_1 N) - x_1 N\} - \{x_2 N \ln(x_2 N) - x_2 N\} \\ &= -N(x_1 \ln x_1 + x_2 \ln x_2)\end{aligned}$$