

Fundamentals





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Objectives

- Present concept of fluidity
- Introduce fundamental properties of fluid





Problems of water supply

flood prevention navigation water power

irrigation

 \rightarrow need to know fluid phenomena





• d'Alembert (1744)

"The theory of fluids must necessarily be based upon experiment"

- d'Alembert paradox { theory ideal, inviscid fluid practice real fluid (viscous)
- Two schools theoretical group → hydrodynamics practical group → hydraulics
- Navier and Stokes
 - \rightarrow general equations for viscous fluid \rightarrow equation of motion





[Re] Navier-Stokes equation

Claude-Louis Navier (1785-1836, French engineer) and George Gabriel Stokes (1819-1903, UK mathematician & physicist)

- one continuity equation + three momentum equations
- model the weather, ocean currents, water flow in a pipe, the air's flow around a wing, and motion of stars inside a galaxy
- design of aircraft and cars, the study of blood flow, the design of power stations, the analysis of pollution,
- exact solution one of the seven most important open problems in mathematics





1.2 Historical Perspective

- New problems in modern times
 - Dispersion of man's wastes in lakes, rivers, and oceans

→ Environmental Fluid Mechanics (Hydraulics)

• state: solid

	[liquid	increasing spacing	increasing inter
fluid	gaseous	and latitude of	molecular cohesive
	└ plasma 🛛 ↓	particle motion	force

• fluid – continuum \rightarrow no voids or holes











otropo	strain			
511855	solid	fluid		
tension		unable to support tension (surface tension)		
compression	elastic deformation	elastic deformation (compressible fluid)		
shear (tangential forces)	→ permanent distortion	permanent distortion or flo w (change shape) to infini tesimal shear stress		











stress	real fluid (vis	ideal fluid (non-viscous fluid)	
	in motion	at rest	at rest and in motion
compression (pressure)	Ο	Ο	Ο
shear	0	×	×





incompressible fluid	compressible fluid
① Compressibility is of small important.	① Compressibility is predominant.
② Liquids and gases may be treated similarly.	② Behavior of liquids and gases is quite dissimilar.
③ Fluid problems may be solved with the principles of mechanics.	③ Thermodynamics and heat transfer concepts must be used as well as principles of mechanics.





- * Fluid does not resist any small shearing stress \rightarrow "Flow occurs"
- Properties of pressure (compression)
- ① Pressure must be transmitted to solid boundaries normal to those boundaries.
- ② At a point, pressure has the same magnitude in all directions.
- \rightarrow Pressure is a <u>scalar quantity</u>.











[Pf]

Apply Newton's law for static equilibrium

$$\sum F_x = p_1 dz - p_3 ds \sin \theta = 0 \tag{a}$$

 $\sum \vec{F} = 0$

$$\sum F_z = p_2 dx - \rho g dx dz / 2 - p_3 ds \cos \theta = 0$$
 (b)

Substitute following relations into Eq. (a) & (b)

$$dx = ds \cos \theta$$

$$dz = ds \sin \theta$$

$$\therefore \quad (a): \ p_1 ds \sin \theta - p_3 ds \sin \theta = 0 \rightarrow p_1 = p_3$$





(b):
$$p_2 ds \cos \theta - \rho g \frac{dz}{2} ds \cos \theta - p_3 ds \cos \theta = 0$$

 $\therefore \quad p_2 = p_3 + \frac{1}{2} \rho g dz$

As
$$dz \rightarrow 0$$
 then $p_2 \approx p_3$

$$\therefore$$
 $p_1 = p_2 = p_3$ at a point $(dx = dz = 0)$





- SI units SI system metric system
- Frequency (f): hertz (HZ = s⁻¹)
- Force, F
- → introduce Newton's 2nd law of motion

F = ma

Force = mass × acceleration

$$a = v / t = L / t^{2} \qquad \begin{bmatrix} Lt^{-2} \end{bmatrix} \qquad (m/s^{2})$$
$$v = L / t \qquad \begin{bmatrix} Lt^{-1} \end{bmatrix} \qquad (m/s)$$

 $\therefore F \rightarrow 1 \text{kg} \cdot \text{m/s}^2 = 1 \text{N}(Newton)$





Dimension	SI unit	English system (FSS)
Length (L)	metre (m)	feet (ft)
Mass (M)	kilogram (kg)	slug (-)
Time (t)	second (s)	second (s)
Temp. (T)	kelvin (K)	degree Rankine (°R)





• Energy, E (work)

$$E = FL \rightarrow \text{kg} \cdot \text{m}^2/\text{s}^2 = J(Joule)$$

• Power, P

$$P = E / t \rightarrow J / s = \mathrm{kg} \cdot \mathrm{m}^2 / \mathrm{s}^3$$

• Pressure, p; Stress, σ, τ

$$p = F / A \rightarrow N/m^2 = Pa (pascal) = kg/m \cdot s^2$$

• Temperature, T: degree Celsius (°C)





- Density, ρ
- = mass per unit volume
- ~ depends on the number of molecules per unit of volume
- ~ decreases with increasing temperature

$$\rho = \frac{M}{V} \rightarrow \text{kg/m}^3$$



- Specific weight (weight density), γ
- = weight (force) per unit volume

$$\gamma = \frac{W}{V} \rightarrow N/m^3 = kg/m^2 \cdot s^2$$





[Re]

 $W = Mg \quad \text{(Newton's 2^{nd} law of motion)}$ g = acceleration due to gravity $\therefore \gamma = \rho g \qquad (1.1)$

- Specific volume=volume per unit mass= $1/\rho$
- Specific gravity, s.g., ~ r. d. (relative density)
 - = ratio of density of a substance to the density of water at a specified

temperature and pressure

$$s.g. = \frac{\rho_f}{\rho_w} = \frac{\gamma_f}{\gamma_w}$$





[Re] s.g. of sea water = 1.03

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s.g. of soil = 2.65
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- s.g. of mercury = 13.6
- Advantage of SI system and English FSS system

① It distinguishes between force (F) and mass (M).

② It has no ambiguous definitions.





	SI	English system
ρ	1,000 kg/m ³	1.94 slugs/ft ³
γ	9,806 N/m ³	62.4 lb/ft ³
g	9.81 m/s²	32.2 ft/s ²





- Greek Alphabet
 - lpha Alpha angle
 - β Beta [beitə]
 - γ, Γ Gamma specific weight, circulation
 - δ, Δ Delta thickness of boundary layer

angle

eddy viscosity, height of surface roughness

- *E* Epsilon
 - ζ Zeta
 - η Eta
- $heta, \Theta$ Theta
- *l* lota [aioutə]
- K Kappa [kæpə]



- λ, Λ Lambda
- *μ* Mu [mju:]
- ν Nu
- ξ Xi [gzai, ksai]
- *o* Omicron
- π Pi [pai]
- ho Rho
- σ, Σ Sigma

au Tau

υ,Υ Upsilon



dynamic viscosity kinematic viscosity

mass density

vorticity

Sigma Xi, Scientific Research Society, 1886 honor society for scientists & engineers

shear



- φ, Φ Phi [fai]Phi Beta Kappa χ Chi [kai] ψ, Ψ Psi [psai, sai]stream function
- ω, Ω Omega angular velocity

Prefixes

Е	exa	10 ¹⁸
Ρ	peta	10 ¹⁵
Т	tera	10 ¹²
G	giga	10 ⁹
М	mega	106





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1.4 Units and Density

	М	mega	10 ⁶
	k	kilo	10 ³
	h	hecto	10 ²
	da	deca	10 ¹
	d	deci	10 ⁻¹
	С	centi	10 ⁻²
	m	milli	10 ⁻³
	μ	micro	10 ⁻⁶
	n	nano	10 ⁻⁹
	р	pico	10 ⁻¹²
	f	femto	10 ⁻¹⁵
	а	atto	10 ⁻¹⁸
S // C			



1.5 Compressibility, Elasticity

• Elastic behavior to compression

- Compressibility ≡ change in volume due to change in pressure solid - modulus of elasticity, E (N/m²)
- fluid bulk modulus











• Stress-strain curve (E[↑], difficult to compress)



= modulus of compressibility (m²/N)





[Re] large E/small C \rightarrow less compressible

- incompressible fluid (inelastic): $E = \infty, C \ll 1$
- ightarrow constant density ho =const.

~ water

- compressible fluid
- \rightarrow changes in density \rightarrow variable density
- ~ gas





Pressure	Temperature, °C				
10 ⁶ N/m ²	0°	20°	50°	100°	150°
0.1	1950	2130	2210	2050	
10.0	2000	2200	2280	2130	1650
30.0	2110	2320	2410	2250	1800
100.0	2530	2730	<u>2840</u>	2700	2330





1.5 Compressibility, Elasticity

- *E* increases as pressure increases.
- *E* is maximum at about 50 °C.
- \rightarrow The water has minimum compressibility at about 50 °C.
- For the case of a fixed mass of liquid at constant temperature

$$E = -V_1 \frac{dp}{dV}$$
$$\frac{\Delta V}{V_1} \approx -\frac{\Delta p}{E}$$

$$\frac{V_2 - V_1}{V_1} \approx -\frac{p_2 - p_1}{E}$$





Compressibility		Modulus of Elasticity, E (kPa)		
steel	1/80 of water	water	2,170,500	
mercury	1/12.5 of water	sea water	2,300,000	
nitric acid	6 of water	mercury	26,201,000	





[Ex] For water; $E = 2,200 \times 10^6 \text{ Pa}$ @ 20°C

$$p_{2} = p_{1} + 7 \times 10^{6} \text{ Pa}$$

$$\Delta p = 7 \times 10^{6} \text{ Pa} / 101.3 \times 10^{3} \text{ Pa} \approx 707 \text{ P} \text{H}$$

$$\therefore \frac{V_{2} - V_{1}}{V_{1}} \approx -\frac{p_{2} - p_{1}}{E} = -0.0032$$

$$\therefore V_{2} = (1 - 0.0032) V_{1}$$

$$\Delta V \approx 0.3\% \text{ decrease}$$

 \rightarrow water is incompressible





[Re] From Wikipedia

Viscosity is a measure of the <u>resistance of a fluid</u> which is being deformed by either shear stress or tensile stress.

Viscosity ~ "thickness" or "internal friction"

- water ~ "thin", having a lower viscosity
- honey ~ "thick", having a higher viscosity

The less viscous the fluid is, the greater its ease of movement (fluidity). Viscosity describes a fluid's <u>internal resistance</u> to flow and may be thought of as a measure of fluid friction.







1.6 Viscosity

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For example, high-viscosity felsic magma will create a tall, steep stratovolcano, because it cannot flow far before it cools, while lowviscosity mafic lava will create a wide, shallow-sloped shield volcano. All real fluids (except superfluids) have some resistance to stress and therefore are **viscous**, but a fluid which has no resistance to shear stress is known as an **ideal fluid** or **inviscid fluid**.

[Re] super fluid – a fluid having frictionless flow, and other unusual properties




• Two types of fluid motion (real fluid)

1) laminar flow:

- viscosity plays a dominant role
- fluid elements or particles slide over each other in layers (laminar)
- molecular diffusion

[Ex] flow in a very small tube, a very thin flow over the pavement, flow in the laminar flow table





2) turbulent flow:

- random or chaotic motion, eddies of various sizes are seen
- common in nature (streams, rivers, pipes)
- large scale mixing between the layers

[Ex] flows in the water supply pipe, flows in the storm sewer pipe, flows in the and canals and streams

• Reynolds number

 $\operatorname{Re} = \frac{Vd}{v}$ Diameter of pipe, Depth of stream

where V = flow velocity; d = characteristic length; n = kinematic viscosity

















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24. Circular cylinder at R=1.54. At this Reynolds number the streamline pattern has clearly lost the fore-and-aft symmetry of figure 6. However, the flow has not yet separated at the rear. That begins at about R=5,

though the value is not known accurately. Streamlines are made visible by aluminum powder in water. Photograph by Sadatoshi Taneda



25. Sphere at R=9.8. Here too, with wall effects negligible, the streamline pattern is distinctly asymmetric, in contrast to the creeping flow of figure 8. The fluid is evidently moving very lowly at the rear, making it difficult to estimate the onset of separation. The flow is presumably attached here, because separation is believed to begin above R=20. Streamlines are shown by magnesium cuttings fluminated in water. Photograph by Maddeline Contanceus and Michele Fixpard

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Separation, eddy formation

40. Circular cylinder at R=9.6. Here, in contrast to figure 24, the flow has clearly <u>separated to form a pair of recirculating eddler</u>. The cylinder is moving through a tank of water containing aluminum powder, and is illuminated

by a sheet of light below the free surface. Extrapolation of such experiments to unbounded flow suggests separation at R=4 or 5, whereas most numerical computations give R=5 to 7. Photograph by Sadatoshi Taneda

Growth of eddy



4L Circular cylinder at R=13.1. The standing eddies become elongated in the flow direction as the speed increases. Their length is found to increase linearly with Reynolds number until the flow becomes unstable above R=40. Tanzda 1956a

42. Circular cylinder at R=26. The downstream distance to the cores of the eddies also increases linearly with Reynolds number. However, the lateral distance between the cores appears to grow more nearly as the square root. Photograph by Sadatoshi Taneda







43. Circular cylinder at R=24.3. A different view of the flow is obtained by moving a cylinder through oil. Tiny magnesium cuttings are illuminated by a sheet of light from an arc projector. The two dark wedges below the circle are an optical effect. The lengths of the particle trajectories have been measured to find the velocity field to within two per cent. Contanceau & Bouard 1977



44. Circular cylinder at R=30.2. The flow is here still completely steady with the recirculating wake more than one diameter long. The walls of the tank, 8 diameters away, have little effect at these speeds. Photograph by Madeleine Contanceau and Roger Bouard

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47. Circular cylinder at R=2000. At this Reynolds number one may properly speak of a boundary layer. It is laminar over the front, separates, and breaks up into a turbulent wake. The separation points, moving forward as the Reynolds number is increased, have now attained their upstream limit, ahead of maximum thickness. Visualization is by air bubbles in water. ONERA photograph, Werlé & Gallon 1972



48. Circular cylinder at R=10,000. At five times the speed of the photograph at the top of the page, the flow pattern is scarcely changed. The drag coefficient consequently remains almost constant in the range of Reynolds number spanned by these two photographs. It drops later when, as in figure 57, the boundary layer becomes turbulent at separation. *Photograph by Thomas Corke and Hassan Nagib*





- laminar flow
- strain = relative displacement

$$=\frac{d_2-d_1}{dy}=\frac{dvdt}{dy}=\frac{dv}{dy}dt$$

[Re]
$$d_2 = v_2 dt; d_1 = v_1 dt$$

 $d_2 - d_1 = (v_2 - v_1) dt$

[Cf] solid mechanics



total angular displacement











• Experiment has shown that, in many fluids, shearing (frictional) stress per unit of contact area, τ is proportional to the <u>time rate of relative</u> <u>strain</u>.

$$\therefore \tau \propto \frac{dv}{dy} dt / dt = \frac{dv}{dy} \quad \text{(velocity gradient)}$$
$$\tau = \mu \frac{dv}{dy} \rightarrow \text{Newton's equation of viscosity} \quad (1.2)$$

where μ = coefficient of viscosity = dynamic (absolute) viscosity





- viscosity = measure of fluid's resistance to shear or angular deformation
- = internal resistance of a fluid to motion (fluidity)
- [Re] Friction forces result from
- <u>cohesion</u> for liquid
- <u>momentum interchange</u> between molecules for gas

[Re] angular deformation due to tangential stress











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1.6 Viscosity

- rate of angular deformation
- (i) displacement of AB relative to CD

$$= \left(u + \frac{du}{dy} \Delta y \right) \Delta t - u \Delta t = \frac{du}{dy} \Delta y \Delta t$$

(ii) angular displacement of AC

$$=\frac{du}{dy}\Delta y\Delta t / \Delta y = \frac{du}{dy}\Delta t$$

(iii) time rate of <u>angular deformation</u> $= \frac{du}{dy} \Delta t / \Delta t = \frac{du}{dy}$ $\tau = \mu \frac{dv}{dy}$





• dynamic viscosity, μ

$$\tau = F / A$$

$$[\tau] = \left[MLT^{-2} / L^{2} \right] \left[ML^{-1}T^{-2} \right] = kg/(m \cdot s^{2}) = Pa$$

$$\left[\frac{dv}{dy} \right] = \left[\frac{LT^{-1}}{L} \right] = \left[T^{-1} \right]$$

$$\therefore \left[\mu \right] = \left[\tau / \frac{dv}{dy} \right] = \left[\frac{ML^{-1}T^{-2}}{T^{-1}} \right] = \left[ML^{-1}T^{-1} \right] = kg/m \cdot s = N \cdot s / m^{2} = Pa \cdot s$$

$$\Rightarrow 1 poises(Poiseuille) = 10^{-1} Pa \cdot s$$





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• kinematic viscosity, ν

$$v = \frac{\mu}{\rho}$$

$$\left[v\right] = \left[\frac{ML^{-1}T^{-1}}{ML^{-3}}\right] = \left[L^{2}T^{-1}\right] = m^{2}/s$$
(1.3)

$$1 \text{ m}^2/\text{s} = 10^4 \text{ stokes} = 10^6 \text{ centistokes}$$

- Remarks on Eq. (1.2)
- (1) τ , μ are independent of pressure. [Cf] friction between two moving solids
- ② Shear stress τ (even smallest τ) will cause flow (velocity gradient).





③ Shearing stress in viscous fluids <u>at rest</u> will be zero.

$$\frac{dv}{dy} = 0 \rightarrow \tau = 0 \quad \text{regardless of } \mu$$

④ At solid boundary, dv/dy ≠∞ (→ τ ≠∞ (no infinite shear))
→ Infinite shearing stress between fluid and solid is not possible.
⑤ Eq. 1.2 is limited to laminar (non-turbulent) fluid motion in which

viscous action is predominant.

[Cf] turbulent flow





- ⁽⁶⁾ Velocity at a solid boundary is zero.
 - \rightarrow No slip condition (continuum assumption)
- Newtonian and non-Newtonian fluids
- i) Newtonian fluid ~ water
- ii) Non-Newtonian fluid ~ plastic, blood, suspensions, paints, polymer solutions → rheology











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Non-Newtonian fluid

1)
$$\tau - \tau_1 = \mu \frac{dv}{dy}$$
 plastic, τ_1 = threshold

2)
$$\tau = K\left(\frac{dv}{dy}\right)$$
 $n > 1$ Shear-thickening fluid $n < 1$ Shear-thinning fluid

 Couette flow: <u>laminar flow</u> in which the shear stress is constant thin fluid film between two large flat plates thin fluid film between the surfaces of coaxial cylinders

$$\frac{dv}{dy} = \frac{V}{h} \sim \text{linear velocity gradient}$$
$$\tau = \mu \frac{V}{h} \sim \text{constant}$$









Turbulent flow

$$\tau = (\mu + \varepsilon) \frac{dv}{dy}$$

\mathcal{E} = eddy viscosity = viscosity due to turbulent factor





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	gas	liquid
main cause of viscosity	exchange of molecule's momentum → interchange of molecules between the fluid layers of different velocities	intermolecular cohesion
effect of temperature variation	temp↑ → molecular activity↑ → viscosity↑ → shearing stress↑	temp↑ → cohesion↓ → viscosity↓ → shear stress↓





[Re] Exchange of momentum

fast-speed layer (FSL)

Two layers tend to stick together as if there is so me viscosity be tween two.



molecules from FSL speed up molecules in LSL

molecules from LSL slow down molecules in FSL $^{
m J}$

low-speed layer (LSL)





1) exchange of momentum : exchange momentum in either direction from high to low or from low to high momentum due to random motion of molecules

2) transport of momentum : transport of momentum from layers of high mome

(high velocity, mv) to layers of low momentum





surface tension

- occur when the liquid surfaces are in contact with another fluid (air) or solid

- f_n (relative sizes of intermolecular cohesive and adhesive forces to another body)

- as temp $\uparrow \rightarrow$ cohesion $\downarrow \rightarrow \sigma \downarrow \implies$ Table A2.4b, p. 694
- some important engineering problems related to surface tension
 - capillary rise of liquids in narrow spaces
 - mechanics of bubble formation
 - formation of liquid drops
 - small models of larger prototype \rightarrow dam, river model





- surface tension, σ (F / L, N/m)
 - force per unit length
 - force attracting molecules away from liquid

Consider static equilibrium

 $\sum F = 0$ (forces normal to the element a, b, c, d)

$$(p_i - p_0)dxdy = 2\sigma dy \sin \alpha + 2\sigma dx \sin \beta$$

where p_i = pressure inside the curvature; p_o = pressure inside the curvature $\sin \alpha = \frac{dx}{2R_1}, \quad \sin \beta = \frac{dy}{2R_2} \quad [dx = 2(R_1 \sin \alpha)]$ $\therefore p_i - p_0 = \sigma \left(\frac{1}{R_1} + \frac{1}{R_2}\right)$ (1.4)











- Cylindrical capillary tube
 - due to both cohesion and adhesion
 cohesion < adhesion → rise (water)
 cohesion > adhesion → depression (mercury)











For a small tube, given conditions are as follows

$$R_1 = R_2 = R$$
 (liquid surface \approx section of sphere) \leftarrow Ch. 2
 $p_0 = -\gamma h$ (hydrostatic pressure)
 $p_i = 0$ (atmospheric)

Substitute above conditions into Eq. 1.15: $p_i - p_0 = \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right)$ (1.4)

$$\therefore \gamma h = \sigma \frac{2}{R}$$

By the way,
$$r = R \cos \theta$$

 $\therefore \gamma h = \sigma \frac{2}{r/\cos \theta} = \frac{2\sigma \cos \theta}{r}$
 $h = \frac{2\sigma \cos \theta}{\gamma r}$
SEHLAB

(1.5)
```
in which h = \text{capillary rise} \rightarrow r \uparrow \rightarrow h \downarrow
```

 θ = angle of contact

r = radius of tube \leq 2.5 mm for spherical form

[Ex] water and mercury \rightarrow Fig. 1.11

If r > 12 mm, *h* is negligible for water.





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- Pressure measurement using tubes in hydraulic experiments
- → Ch.2 manometer
- ~ capillarity problems can be avoided entirely by providing tubes large enough to render the capillarity correction negligible.
- Fomation of curved surface, droplet
 - At free liquid surface contacting the air, cohesive forces at the outer layer are not balanced by a layer above.
 - \rightarrow The surface molecules are pulled tightly to the lower layer.
 - \rightarrow Free surface is curved.
- [Ex] Surface tension force supports small loads (water strider).













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[IP 1.10] For a droplet of water (20 °C), find diameter of droplet Given: $p_i - p_0 = 1.0$ kPa

At 20°C, $\sigma = 0.0728$ N/m \leftarrow App. 2

[Sol]

$$p_i - p_0 = \sigma \left(\frac{1}{R_1} + \frac{1}{R_2} \right) = \frac{2\sigma}{R}$$
 (1.4)
 $\therefore 1 \times 10^3 \text{ N/m}^2 = 2(0.0728) \cdot \frac{1}{R}$

 $\therefore R = 0.000146 \text{m} = 0.146 \text{mm} \rightarrow d = 0.292 \text{mm}$





[IP 1.11] Find height of capillary rise in a clean glass tube of 1 mm diameter if the water temperature is 10°C or 90°C.

[Sol]

From App. 2 Table A 2.4b;

- @ 10°C σ =0.0742 N/m, γ = 9.804 kN/m³
- @ 90°C σ =0.0608 N/m, γ = 9.466 kN/m³





Use Eq. 1.16

For water,
$$\theta = 0^{\circ}$$

$$h = \frac{2\sigma\cos\theta}{\gamma r}$$
 (1.5)

$$\therefore h_{10} = \frac{2(0.0742)(1)}{9804(0.0005)} = 0.030 \text{m} = 30 \text{mm}$$

$$h_{90} = \frac{2(0.0608)(1)}{9466(0.0005)} = 0.026 \text{m} = 26 \text{mm}$$





- vapor pressure = partial pressure exerted by ejected molecules of liquid
 - \rightarrow Table A2.1 and A2.4b
- liquids ~ tend to vaporize or evaporate due to molecular thermal vibrations (molecular activity)
 - \rightarrow change from liquid to gaseous phase
- temperature $\uparrow \rightarrow$ molecular activity $\uparrow \rightarrow$ vaporization $\uparrow \rightarrow$ vapor pressure \uparrow





• volatile liquids:

~ easy to vaporize → high vapor pressure gasoline: $p_v = 55.2$ kPa at 20 °C water: $p_v = 2.34$ kPa at 20 °C mercury: $p_v = 0.00017$ kPa at 15.6 °C

- mercury : low vapor pressure and high density = difficult to vaporize
 - \rightarrow suitable for pressure-measuring devices





• Cavitation: App. 7 (p. 672)

In a flow fluid wherever the local pressure falls to the vapor pressure of the liquid, local vaporization occurs.

- \rightarrow <u>Cavities</u> are formed in the low pressure regions.
- High velocity region
- \rightarrow The cavity contains a swirling mass of droplets and vapor.
- \rightarrow Cavities are swept downstream into a region of high pressure.
- \rightarrow Then, cavities are collapses suddenly.





 \rightarrow surrounding liquid rush into the void together

 \rightarrow it causes <u>erosion (pitting)</u> of solid boundary surfaces in machines, and vibration

 \rightarrow boundary wall receives a blow as from a tiny hammer











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- Prevention of cavitation
- ~ cavitation is of great importance in the design of high-speed hydraulic machinery such as
- turbines, pumps, in the overflow and underflow structures of high dams, and in high-
- speed motion of underwater bodies (submarines, hydrofoils).
 - \rightarrow design improved forms of boundary surfaces
 - \rightarrow predict and control the exact nature of cavitation \rightarrow set limits





Body

cavitation

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1.8 Vapor Pressure



(a)





- Boiling:
 - = rapid rate of vaporization caused by an increase in temperature
 - = formation of vapor bubbles throughout the fluid mass
 - ~ occur (whatever the temperature) when the external absolute pressure

imposed on the

liquid is equal to or less than the vapor pressure of the liquid

~ boiling point = f (imposed pressure, temp.)

 $p_{atm} \leq p_v \rightarrow \text{boiling occurs}$





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1.8 Vapor Pressure

		Ta	able A 4b	Table A 2.5b	
altitude (El. m)	Temp. (°C)	p _ν (kPa), absolute	p _{atm} (kPa), absolute	boiling point (°C)	remark
m.s.l.	100	101.3	101.3	100	
12,000	60	19.9	19.4	60	undercooked





• Evaporation: When the space surrounding the liquid is too large, the liquid continues to P_v

vaporize until the liquid is gone and only vapor remains at a pressure less than or equal.

[IP 1.12] For a vertical cylinder of diameter 300 mm, find min. force that will cause the water boil.





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1.8 Vapor Pressure







[Sol] From Table A2.4b; P_{ν} =31.16 kPa at 70 °C

For water to boil; $p' \leq p_v = 31.16$,

$$\therefore p' = 100 - \frac{F}{A} = 31.16$$

:.
$$F = (100 - 31.16) \frac{\pi (0.3)^2}{4} = 4.87 \text{ kN}$$





Homework Assignment # 1

Due: 1 week from today

- Prob. 1.2
- Prob. 1.10
- Prob. 1.27
- Prob. 1.46
- Prob. 1.49
- Prob. 1.58
- Prob. 1.69
- Prob. 1.72

Prob. 1.82



