





Prepared by Jin Hwan Hwang





1. Form Drag and Skin Friction

- Bed forms can have a profound influence on flow resista nce, and thus on sediment transport in an alluvial chann el.
- Consider, the case of normal flow in a wide rectangular c hannel. In the presence of bed forms must be

$$\overline{\tau}_b = \rho g H S$$

- Where overbar tau is an effective boundary shear stress, where the overbar denotes averaging over the bed forms , and can be defined as the streamwise drag force per u nit area.
- H represents the depth averaged over the bed forms.





The effective boundary shear stress consists of two majo r sources of "skin friction (local shear stress)", and "form drag (pressure)".

$$\overline{\tau}_{b} \equiv \tau_{bs} + \tau_{bf}$$

 au_{bs} : Skin friction

$$au_{bf}: Form drag$$

- The form drag results from the net pressure distribution o ver an entire bed form.
- Drag is not directly on the particle itself (since normal to surface).
- Drag appears when flow separates in the lee of the crest , and this one is substantial (so, skin friction dominant)





1. Form Drag and Skin Friction

- The part of the effective shear stress that governs sedim ent transport is thus seen to be the skin friction.
- Any equation of the bed-load is to be applied, it is neces sary to replace the Shields stress by the Shields stress a ssociated with skin friction only:

$$\tau_s^* = \frac{\tau_{bs}}{\rho g R D}$$





2. Shear Stress Partitions

- Einstein Partition
 - Einstein was the first to recognize the necessity of distingui shing between skin friction and form drag.

 $\overline{\tau}_b = \rho C_f U^2$

- Where C_f represents a resistance coefficient that includes both skin friction and form drag.

$$\tau_{bs} = \rho C_{fs} U^2$$

- Where C_{fs} is the frictional resistance coefficient that would result if bed forms were absent. In rough turbulent flow

$$C_{fs} = \left[\frac{1}{\kappa} \ln\left(11\frac{H_s}{k_s}\right)\right]^{-1}$$

- H_s denotes the depth that would result in the absence of b edforms





 The purpose of the calculation is to get H_s. In normal flo w (steady, uniform)

$$\overline{\tau}_{b} = \rho C_{f} U^{2} = \rho g H S$$
$$\tau_{bs} = \rho C_{fs} U^{2} = \rho g H_{s} S$$

Now,

$$H_{s} = \frac{U^{2}}{gS} \left[\frac{1}{\kappa} \ln \left(11 \frac{H_{s}}{k_{s}} \right) \right]^{-2}$$

If form drag is

$$\tau_{bf} = \rho C_{ff} U^2 = \rho g H_f S$$

$$\tau_{bf} = \overline{\tau}_b - \tau_{bs}$$



2. Shear Stress Partitions

Then the coefficients are

$$C_f = C_{fs} + C_{ff}$$
$$H = H_s + H_f$$





Example Calculations

- Consider a sand-bed stream at a given cross section with a slope of 0.0004, a mean depth of 2.9m, a median bed sediment size of 0.35mm and a discharge per unit width $q = U \times H = 4.4 m^2 / s$. Assume that the flow is under near -normal condition. Compute values of τ_{bs} , τ_{bf} , C_{fs} , C_{ff} , H_s and H_f
- Sol) The mean flow velocity is given by U=4.4/2.9=1.52m/s An appropriate esimate of k_s for sand-bed steam is k_s=2.5D₅₀. First, H_s=1.047, then H_f. From the each depth, you may can each shear stress. After this, you can get friction coefficietns of each.





2. Shear Stress Partitions

Finally the form-induced Shield stress is

$$\tau_f^* = \frac{\tau_{bf}}{\rho g R D}$$

- The previous case (Example) only 30% of total Shield str ess is skin Shield stress (which contributes to the transp ort of sediment).
- Einstein partition needs the *velocity information*.



 Nelson and Smith (1989) consider flow over a dune; the f low is taken to separate in the lee of the dune.

$$D_{ffs} = \frac{1}{2} \rho C_D B \Delta U_r^2$$

 D_{ffs} denotes that potion of the streamwise drag force D_{fs} that is due to form drag, *B* is the width, U_r is a reference velocity. C_D is site specific but in the Columbia River 0.21.

$$\tau_{bf} = \frac{1}{2} \rho C_D \frac{\Delta}{\lambda} U_r^2 = \frac{D_{ffs}}{B\lambda}$$



- The reference velocity U_r is defined to be the mean veloc ity that would prevail between z=k_s and z=∆ if the bed for ms were not there.
- From the logarithmic profile

$$\frac{U_r}{\sqrt{\tau_{bs} / \rho}} = \frac{1}{\kappa} \left[\ln \left(30 \frac{\Delta}{k_s} \right) - 1 \right]$$

• It is assumed that a rough logarithmic law with roughnes s k_s prevail from $z=k_s$ and $z=\Delta$ and from $z=\Delta$ to z=H, k_c r epresents an another roughness, presenting a composit e roughness length, including the effects of both skin or grain friction and form drag.



Now there are two velocity distributions exist

$$\frac{\overline{u}(z)}{\sqrt{\tau_{bs}/\rho}} = \frac{1}{\kappa} \ln\left(30\frac{z}{k_{s}}\right), \qquad k_{s} < z < \Delta$$
$$\frac{\overline{u}(z)}{\sqrt{(\tau_{bs} + \tau_{bf})/\rho}} = \frac{1}{\kappa} \ln\left(30\frac{z}{k_{c}}\right), \quad \Delta < z < H$$

• Nelson and Smith match the above two laws at the level $z=\Delta$ (the top of the dune).

$$\frac{\tau_{bs} + \tau_{bf}}{\tau_{bs}} = \left[\frac{\ln(30\Delta / k_s)}{\ln(30\Delta / k_c)}\right]^2$$



- The partition requires a prior knowledge of total boundar y shear stress as well as roughness height and dune hei ght, and wave length
- From the partition

$$\tau_{bf} = \overline{\tau}_{b} - \tau_{bs} = \frac{1}{2} C_{D} \frac{\Delta}{\lambda \kappa^{2}} \left[\ln \left(30 \frac{\Delta}{k_{s}} \right) - 1 \right]^{2} \tau_{bs}$$

 This equation will give skin shear stress and so from fricti on. With this information using the previous equation of

$$\frac{\tau_{bs} + \tau_{bf}}{\tau_{bs}} = \left[\frac{\ln(30\Delta / k_s)}{\ln(30\Delta / k_c)}\right]^2$$

You can get the composite roughness.



2. Example

• Chosen to rather similar to the previous one, let $H = 2.9m; S = 0.0004; k_s = 2.5D_{50}; D_{50} = 0.35mm,$ $\Delta = 0.4m; \lambda = 15m$

• Then
$$\overline{\tau}_{b} = \rho g H S$$
, from
 $\overline{\tau}_{b} - \tau_{bs} = \frac{1}{2} C_{D} \frac{\Delta}{\lambda \kappa^{2}} \left[\ln \left(30 \frac{\Delta}{k_{s}} \right) - 1 \right]^{2} \tau_{bs}$

• $\tau_{bf} = \overline{\tau}_b - \tau_{bs}$, now you can get k_c from, $\frac{\tau_{bs} + \tau_{bf}}{\tau_{bs}} = \left[\frac{\ln(30\Delta / k_s)}{\ln(30\Delta / k_c)}\right]^2$

• And you can get, $C_f = C_{fs} + C_{ff}$ $H = H_s + H_f$



2. Example

 In computing friction coefficients, the flowing relationship was used for the depth-averaged velocity

$$\frac{U}{\sqrt{\left(\tau_{bs}+\tau_{bf}\right)/\rho}} = \frac{1}{\kappa} \ln\left(11\frac{H}{k_c}\right)$$

- Now you can calculate the depth-average flow velocity.
- In another way, if you know the mean depth averaged ve locity, you can get friction coefficient for form drag as

$$k_{c} = \frac{11H}{e^{\kappa U/u_{*}}} \quad \left(u_{*} = \sqrt{\left(\tau_{bs} + \tau_{bf}\right)/\rho}\right)$$

 The Nelson-Smith does not require the assumption o f quasi-normal flow (Einstein method)



3. Empirical Formulas for Stage-Discharge Relations

- To use the partition method, it is necessary to know in ad vance the total effective boundary shear stress.
- Therefore, partition itself cannot say anything.
- The Einstein-Barbarossa Method
 - For the case of dune resistance in sand-bed streams.
 - Empirical relation of the following form:

$$C_{ff} = f(\tau_{s35}^{*})$$
$$\tau_{s35}^{*} = \frac{\tau_{bs}}{\rho g R D_{35}}$$



- 3. Empirical Techniques for Stage-Discharge Relations
 - The friction coefficient for the bed forms declines for incr easing shear.
 - Increased intensity of flow causes a decrease in form dra g (transition from dunes to flat bed).





- 3. Empirical Techniques for Stage-Discharge Relations
 - We need to find a relation between H and water discharg e Q.
 - It is assumed that the river slope S and the sizes D_{50} and D_{35} are known. The river is taken to be sufficiently wide s o that the hydraulic radius $R_h=H$: (Hydraulic radius can b e sued for the depth).
 - The channel width is

B = B(H)



3. Empirical Techniques for Stage-Discharge Relations

$$H_{s} \rightarrow C_{fs} \qquad \left(C_{fs} = \left[\frac{1}{\kappa}\ln\left(11\frac{H_{s}}{k_{s}}\right)\right]^{-2}\right)$$

$$C_{fs}, H_{s} \rightarrow U \left(U^{2} = \frac{gH_{s}S}{C_{fs}}\right)$$

$$H_{s} \rightarrow \tau_{bs} \rightarrow \tau_{s35}^{*} \qquad \left(\tau_{bs} = \rho gH_{s}S, \ \tau_{s35}^{*} = \frac{\tau_{bs}}{\rho gRD_{35}}\right)$$

$$\tau_{s35}^{*} \rightarrow C_{ff} \qquad \text{Use diagram}$$

$$C_{ff}, U \rightarrow H_{f} \qquad \left(\tau_{bf} = \rho C_{ff}U^{2} = \rho gH_{f}S\right)$$

$$H = H_{s} + H_{f}$$

$$Q = UHB$$

 So result may be plotted in terms o f H versus Q for t he desired depthdischarge relatio n.



- 3. Empirical Techniques for Stage-Discharge Relations
 - With the previous value, we may have the bed-load tran sport rate.

$$\tau_s^* = \frac{\tau_{bs}}{\rho g R D_{50}}$$

• The volumetric bed load transport rate Q_b is

$$Q_b = q_b B$$

 This is for the depth-discharge prediction in the sand-be d streams.



- 3. Stage-Discharge Relations (Engelund-Hansen Method)
- Sand-bed streams.
- More accurate than the method of Einstein and Barbaro ssa.
- The method assumes quasi-uniform material; it is neces sary to know only a single grain size *D*. Roughness heig ht k_s is $k_s = 2.5D_{50}$
- The method uses the Einstein partition. Skin friction is c omputed

$$C_{fs} = \left[\frac{1}{\kappa} \ln\left(11\frac{H_s}{k_s}\right)\right]^{-2}$$
$$\tau^* = \frac{\overline{\tau}_b}{\rho Rg D}; \qquad \tau^*_s = \frac{\tau_{bs}}{\rho Rg D}$$



- 3. Stage-Discharge Relations (Engelund-Hansen Method)
- The previous equation have two branches.
- The lower-regime branch is given by

 $\tau_{s}^{*} = 0.064 + 0.4(\tau^{*})^{2}$

Upper branch satisfies the relation

 $au_s^* = au^*$

over a range; this implies on upper-regime plane bed. F or higher values of Shields stress again exceeds τ_s^* implying form drag due to the development of antidunes.

 The procedure rather closely parallels that of the Einstei n Barbarossa method. It is assumed that values of S an d D are know, as well as cross-sectional geometry. Seoul National University

- 3. Stage-Discharge Relations (Engelund-Hansen Method)
- Values of H_s are selected, ranging from a low value to near b ank-full. The calculation then proceeds as follows

$$H_{s} \rightarrow C_{fs} \rightarrow U \qquad \left(C_{fs} = \left[\frac{1}{\kappa} \ln \left(11 \frac{H_{s}}{k_{s}} \right) \right]^{-2}, U^{2} = \frac{gH_{s}S}{C_{fs}} \right)$$
$$H_{s} \rightarrow \tau_{bs} \rightarrow \tau_{s}^{*} \qquad \left(\tau_{bs} = \rho gH_{s}S, \ \tau_{s}^{*} = \frac{\tau_{bs}}{\rho gRD} \right)$$
$$\tau_{s}^{*} \rightarrow \tau^{*} \qquad \left(\tau_{s}^{*} = f\left(\tau^{*}\right), \tau_{s}^{*} = 0.006 + 0.4\left(\tau^{*}\right)^{2} \right)$$
$$\tau^{*} \rightarrow \overline{\tau}_{b} \rightarrow H \qquad \left(\tau_{s}^{*} = \frac{\tau_{bs}}{\rho RgD}, \overline{\tau}_{b} = \rho gHS \right)$$
$$Q = UHB$$

 The value of τ_s^{*} may then be used to calculate bedload trans port rates, in a fashion completely analogous to the procedur e outlined for the Einstein-Barbarossa method.