

Introduction to Nuclear Fusion

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Tokamak stability

Tokamak Stability

- Considering plasma states which are not in perfect thermodynamic equilibrium (no exact Maxwellian distribution, e.g. non-uniform density), even though they represent equilibrium states in the sense that the force balance is equal to 0 and a stationary solution exists, means their entropy is not at the maximum possible and hence free energy appears available which can excite perturbations to grow:
unstable equilibrium state
- The gradients of plasma current magnitude and pressure are the destabilising forces in connection with the bad magnetic field curvature: The ratio of these two free energies turns out to be β_p

Tokamak Stability

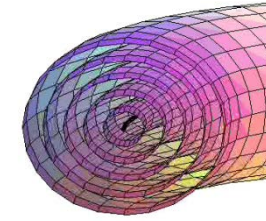
- **Ideal MHD instabilities**

- current driven (kink) instabilities
 - internal modes
 - external modes
- pressure driven instabilities
 - interchange modes
 - ballooning modes
- current+pressure driven: Edge Localised Modes (ELMs)
- vertical instability

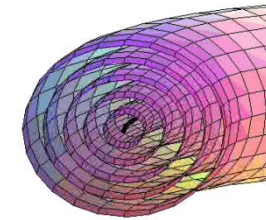
- **Resistive MHD instabilities**

- current driven instabilities
 - tearing modes
 - neoclassical tearing modes (NTMs)
- nonlinear modes
 - sawtooth
 - disruption

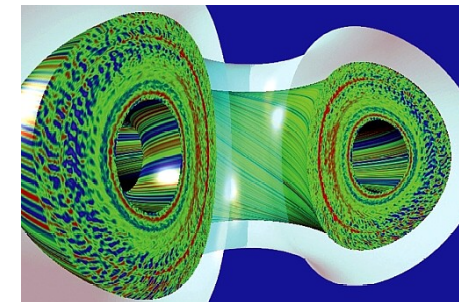
- **Microinstabilities - Turbulence**



Flux conservation
Topology unchanged



Reconnection of field lines
Topology changed



Ideal MHD instabilities in a Tokamak

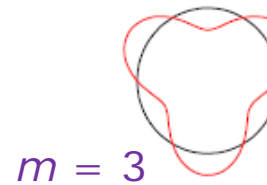
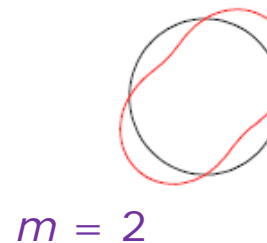
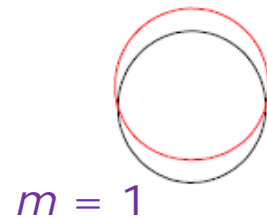
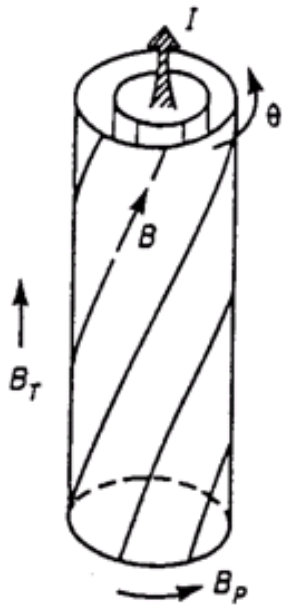
Ideal MHD Instabilities

- **The most Virulent Instabilities**
 - fast growth (microseconds)
 - the possible extension over the entire plasma

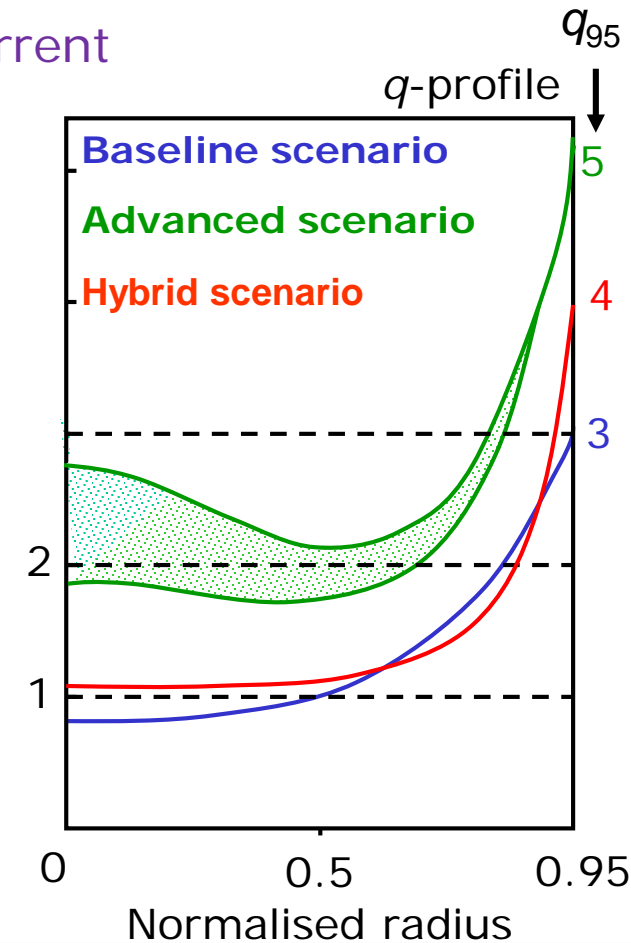
Ideal MHD Instabilities

- Kink modes

- Causing a contortion of the helical plasma column
- Driven by the radial gradient of the toroidal current
- External kind modes:
 - Fastest and most dangerous
 - Arising mainly when $q_a < 2$



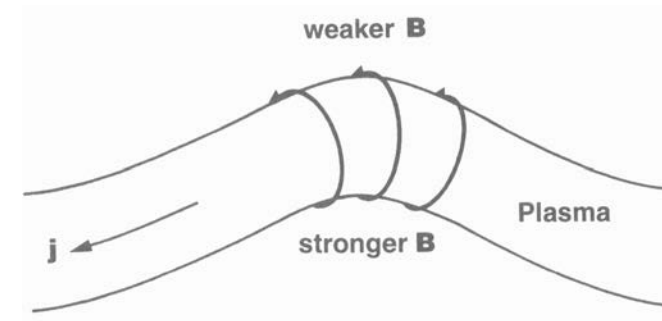
m : poloidal mode number



Ideal MHD Instabilities

- Kink modes

- Stabilising effect by the conducting wall and strong toroidal magnetic field



$$q_a = \frac{aB_\phi}{R_0 B_\theta} = \frac{aB_\phi}{R_0 \mu_0 I_p / 2\pi a} \propto \frac{B_\phi}{I_p}$$

↑ stabilising
 Determining plasma current limit
 set by kink instabilities → **safety** factor
 ↓ destabilising

$q_a > 1$ Kruskal-Shafranov criterion:
 stability condition for external kink mode for the worst case

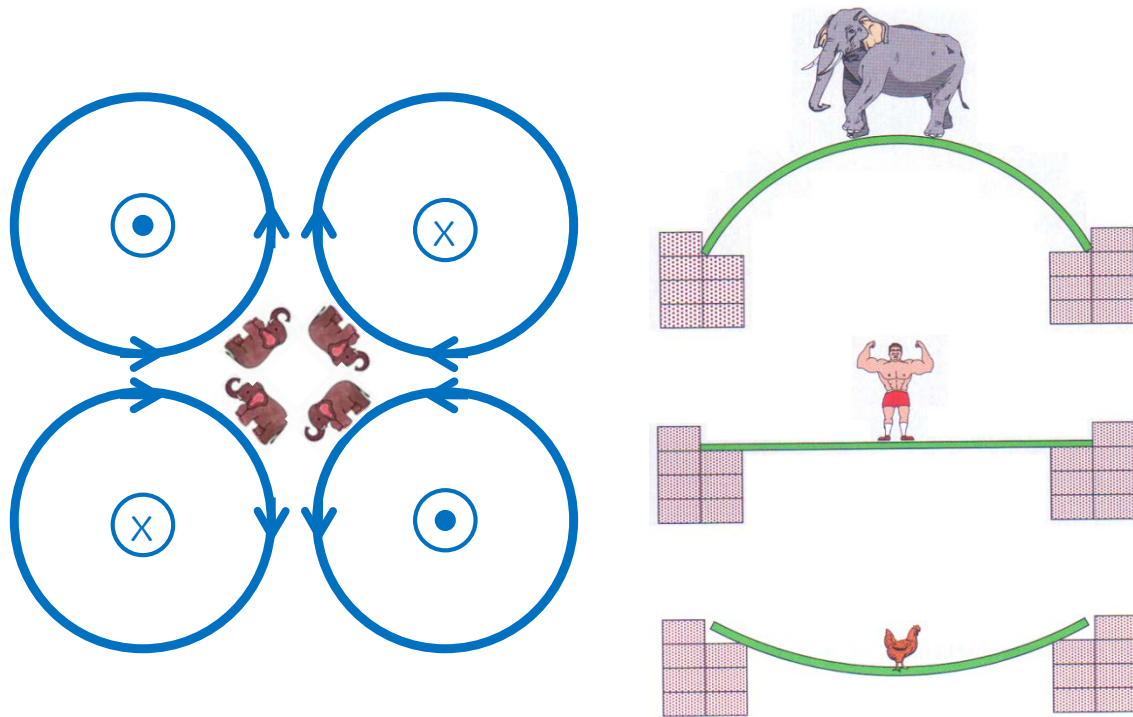
Imposing an important constraint on tokamak operation:
 toroidal current upper limit: Kruskal-Shafranov current ($I < I_{KS}$)

$$q_a = \frac{aB_\phi}{R_0 B_p} = \frac{2\pi a^2 B_\phi}{\mu_0 R_0 I_{KS}} = 1 \quad I_{KS} \equiv 2\pi a^2 B_\phi / \mu_0 R_0 = 5a^2 B_\phi / R_0 \text{ [MA]}$$

Ideal MHD Instabilities

- Interchange modes

- A toroidally confined plasma sees 'bad' convex curvature of the helical magnetic field lines on the outboard side of the torus.

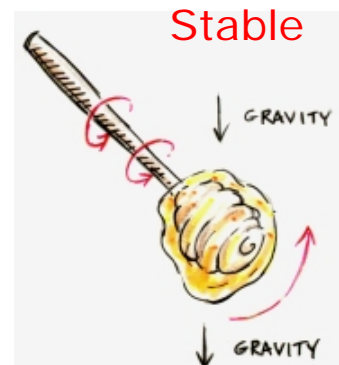
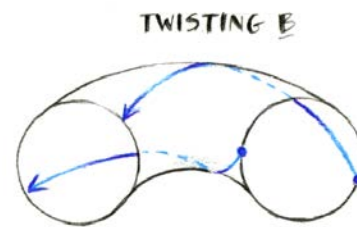
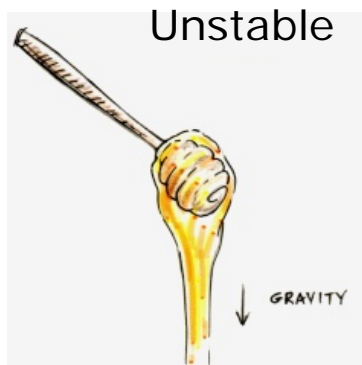
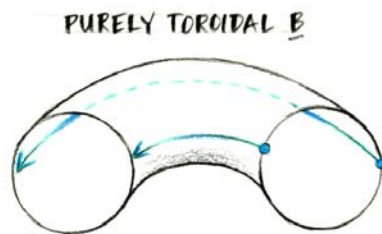


F. F. Chen, "An Indispensable Truth", Springer (2011)

Ideal MHD Instabilities

- Interchange modes

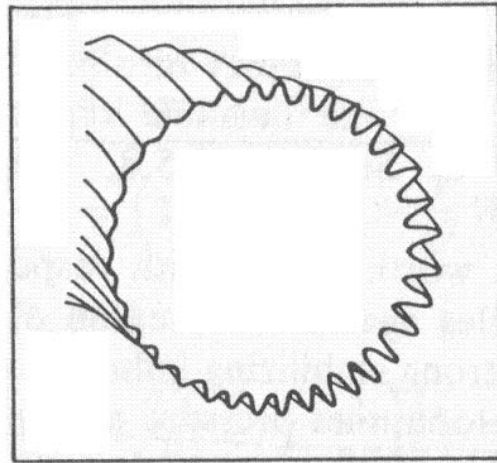
- A toroidally confined plasma sees 'bad' convex curvature of the helical magnetic field lines on the outboard side of the torus.
- The average curvature of **B**-field lines over a full poloidal rotation is 'good' for windings with a rotational transform $I \leq 2\pi$, i.e., $q \geq 1$.
- Interchange perturbations do not grow in normal tokamaks if $q \geq 1$.



Ideal MHD Instabilities

- **Ballooning modes**

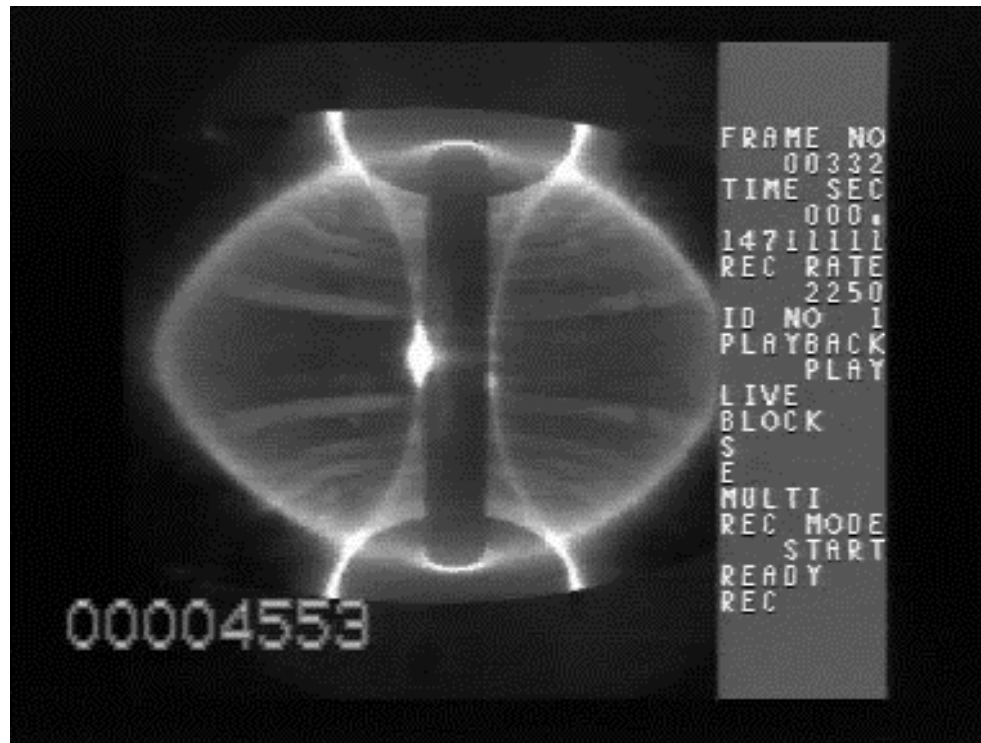
- locally grow in the outboard bad curvature region: ballooning modes
- A high local pressure gradient is responsible for driving the ballooning instability.
- Can be suppressed almost everywhere in the plasma by establishing appropriate pressure profiles and appropriate magnetic field line windings.



Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

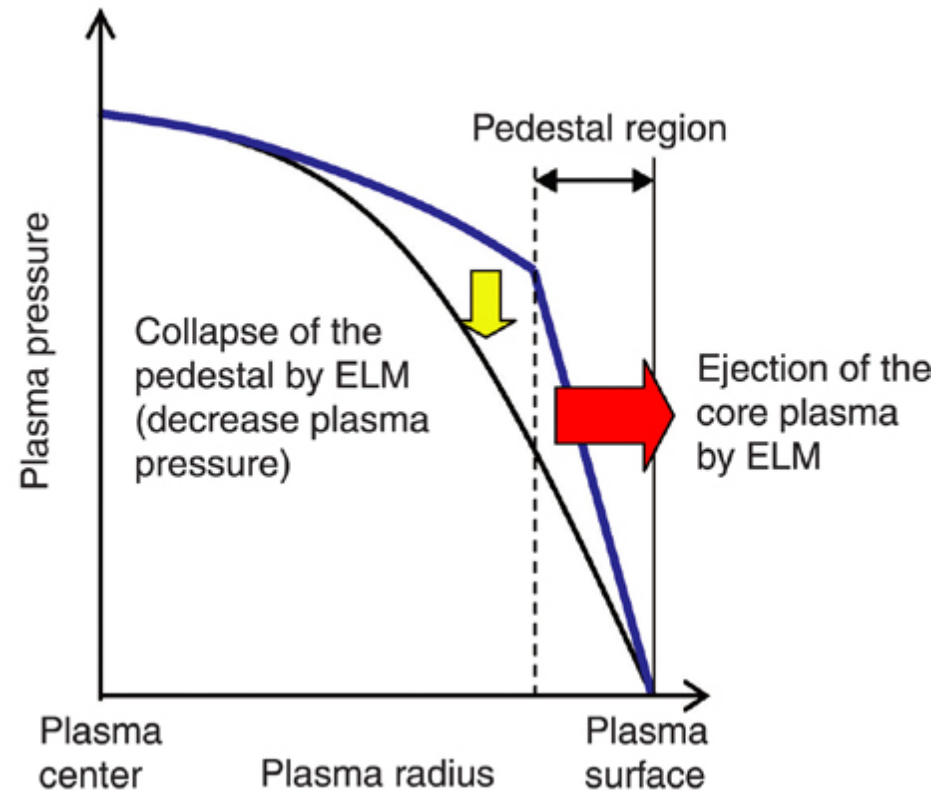
- current driven (peeling mode) and pressure driven (ballooning mode) combined instability



Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

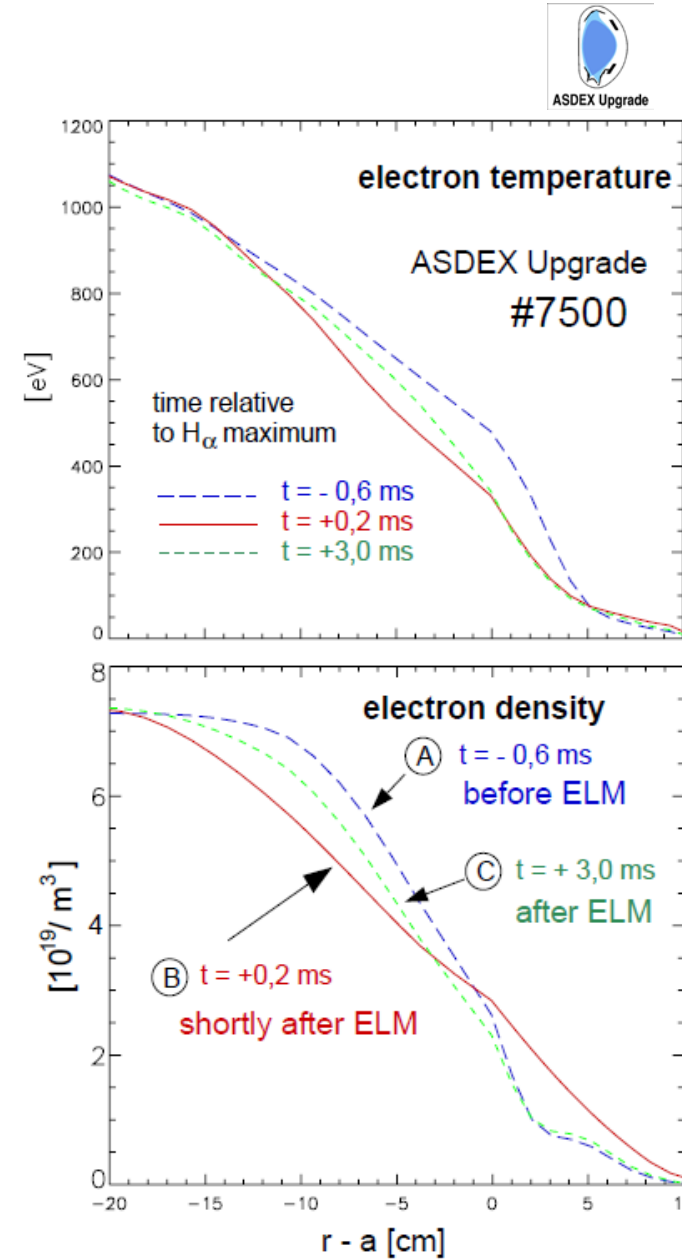
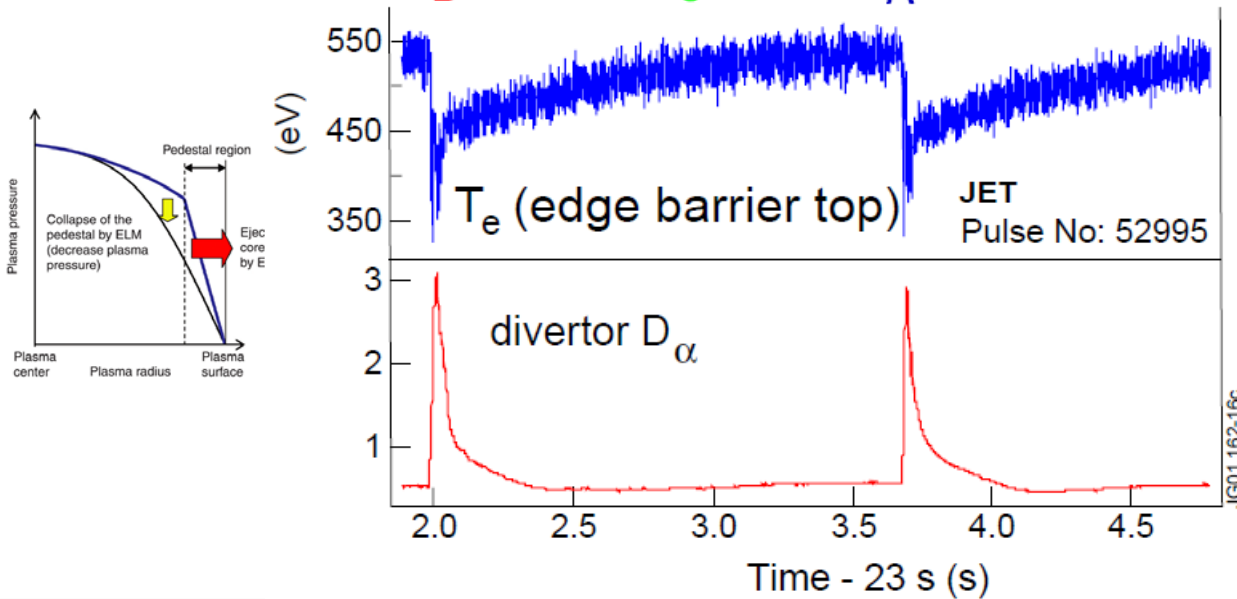
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Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

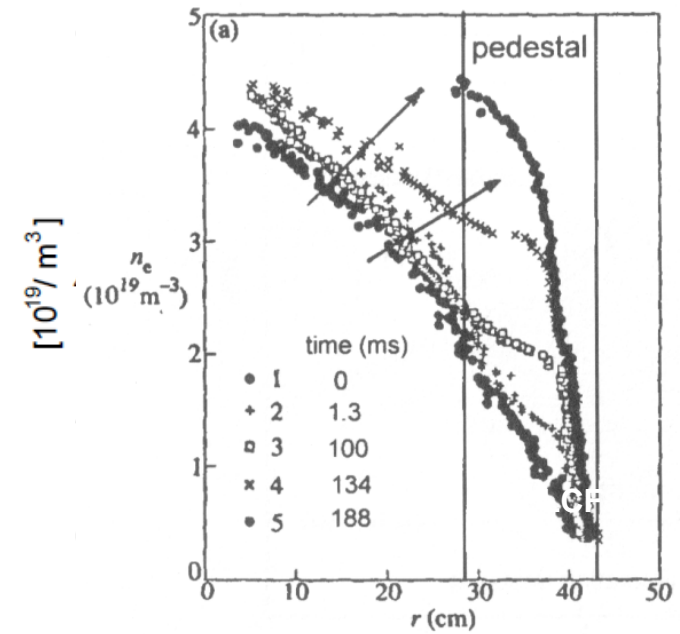
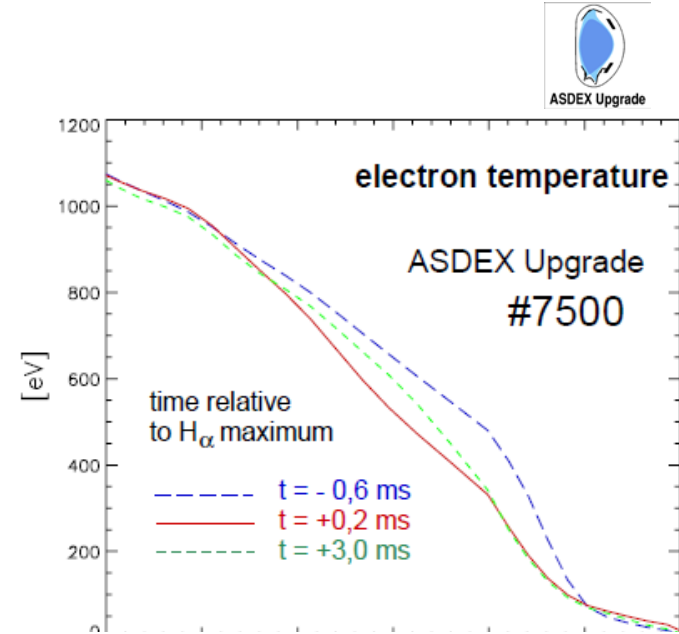
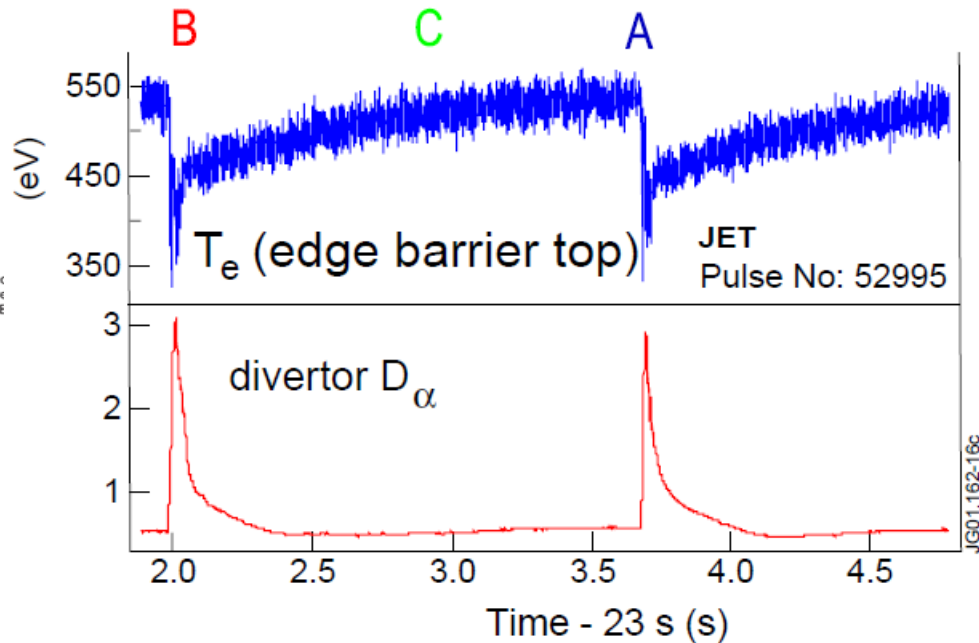
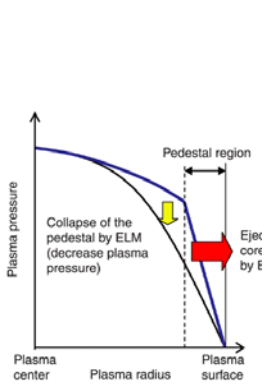
- A. Critical ∇p in H-mode barrier region reached
→ short unstable phase (ELM event)
- B. Energy and particle loss reduces gradients.
- C. Gradients build up during reheat/refuelling phase.



Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

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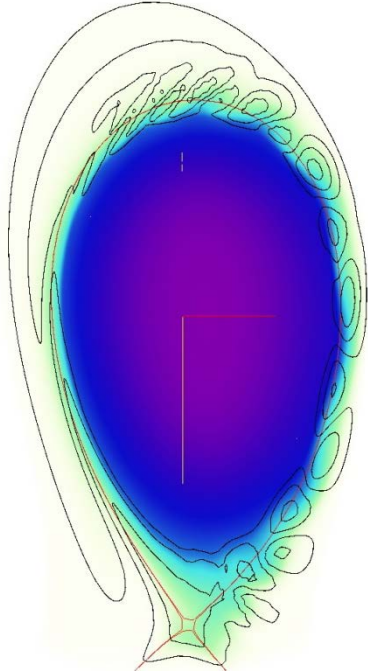


Ideal MHD Instabilities

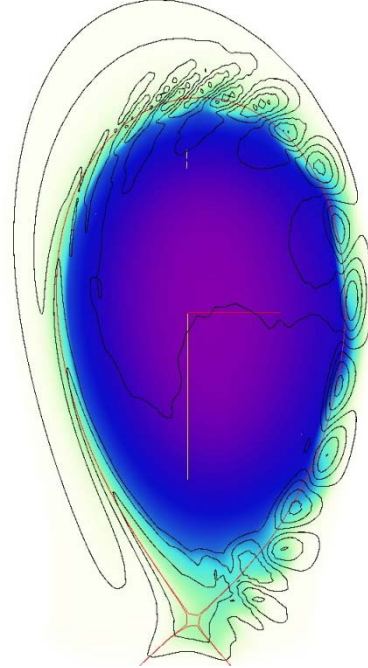
- Edge Localised Modes (ELMs)

- Non-linear MHD simulations with JOEUK

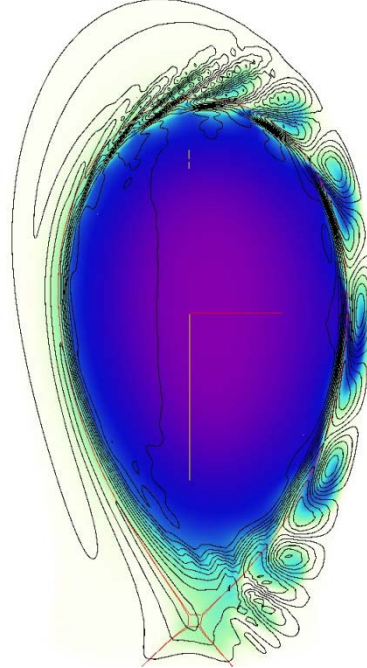
$t = 2650 \tau_A^{-1}$



$t = 2700 \tau_A^{-1}$



$t = 2890 \tau_A^{-1}$



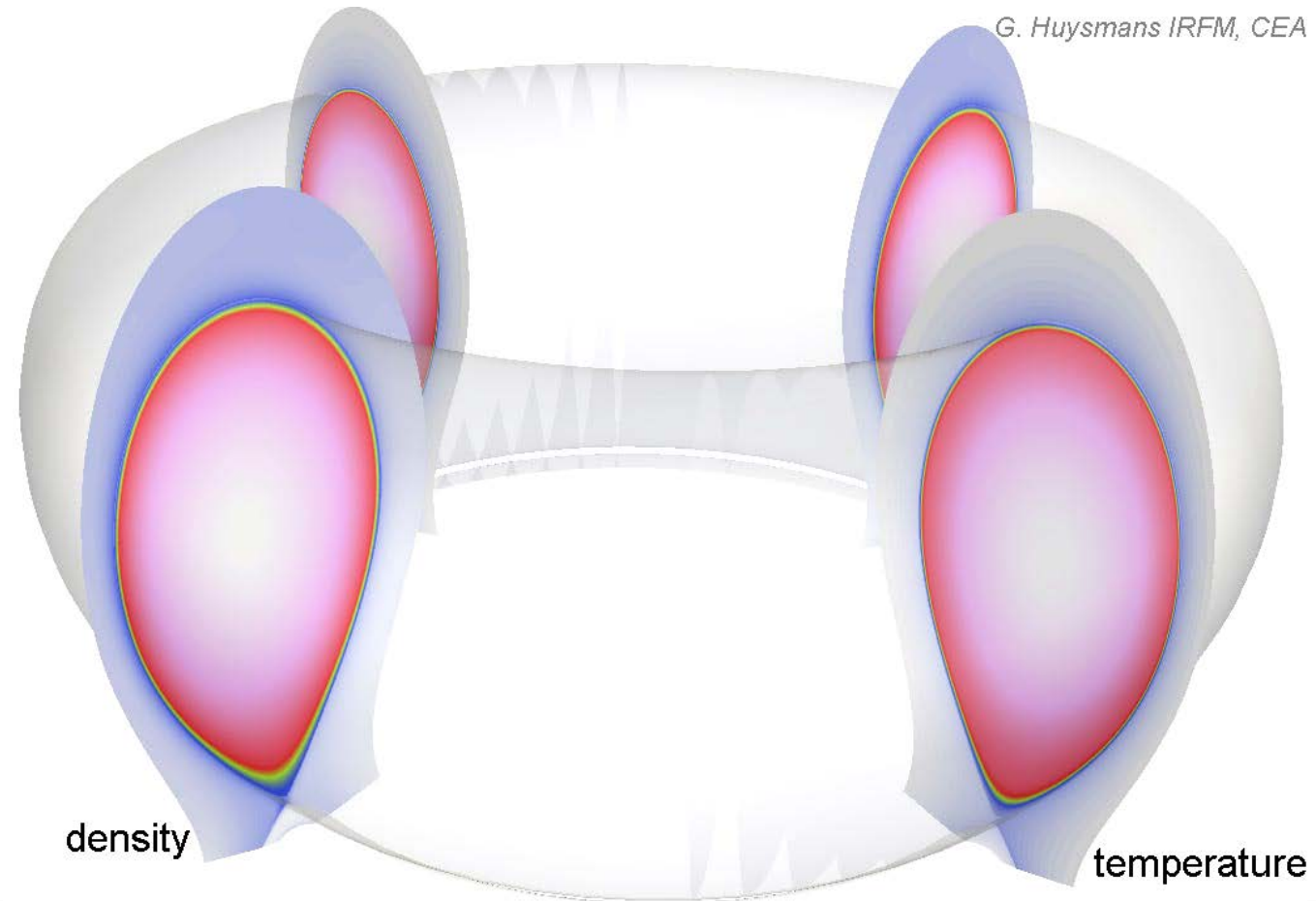
Evolution of ballooning mode



Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

- Non-linear MHD simulations with JOEKE

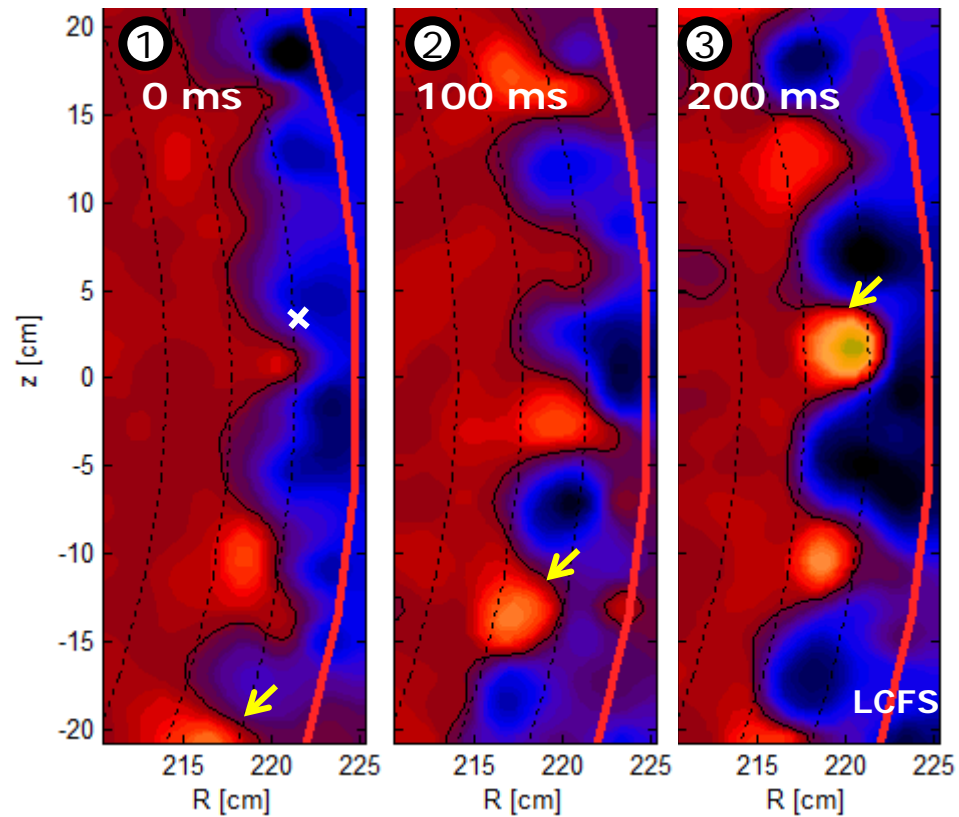


Ideal MHD Instabilities

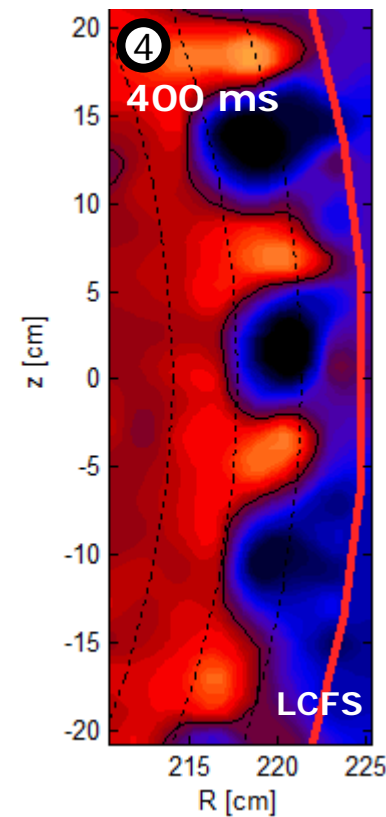
- Edge Localised Modes (ELMs)

- Standard ELM dynamics in the KSTAR visualized by ECEI

(1) Initial Growth



(2) Saturation

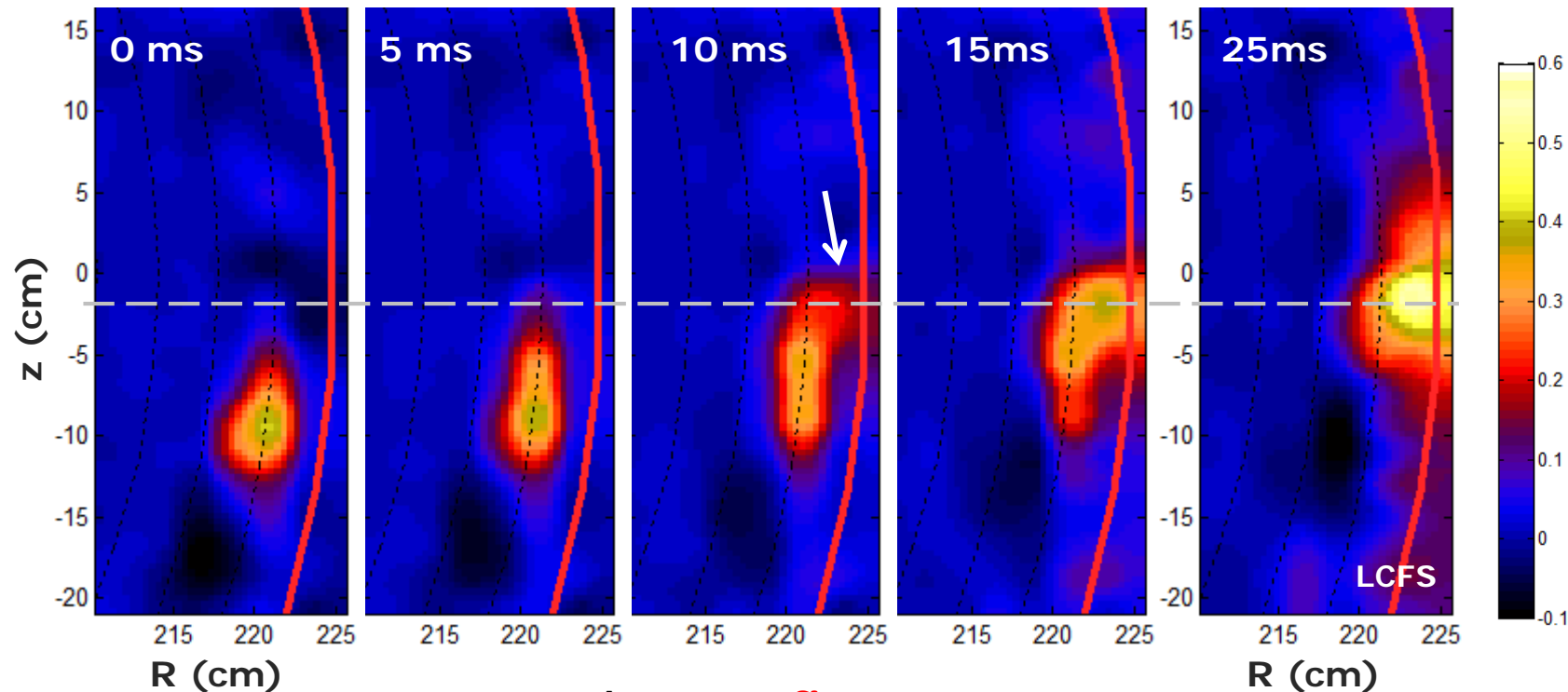


Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

- Standard ELM dynamics in the KSTAR visualized by ECEI

(3) ELM crash



Filaments elongate poloidally



A narrow **finger**-like structure develops



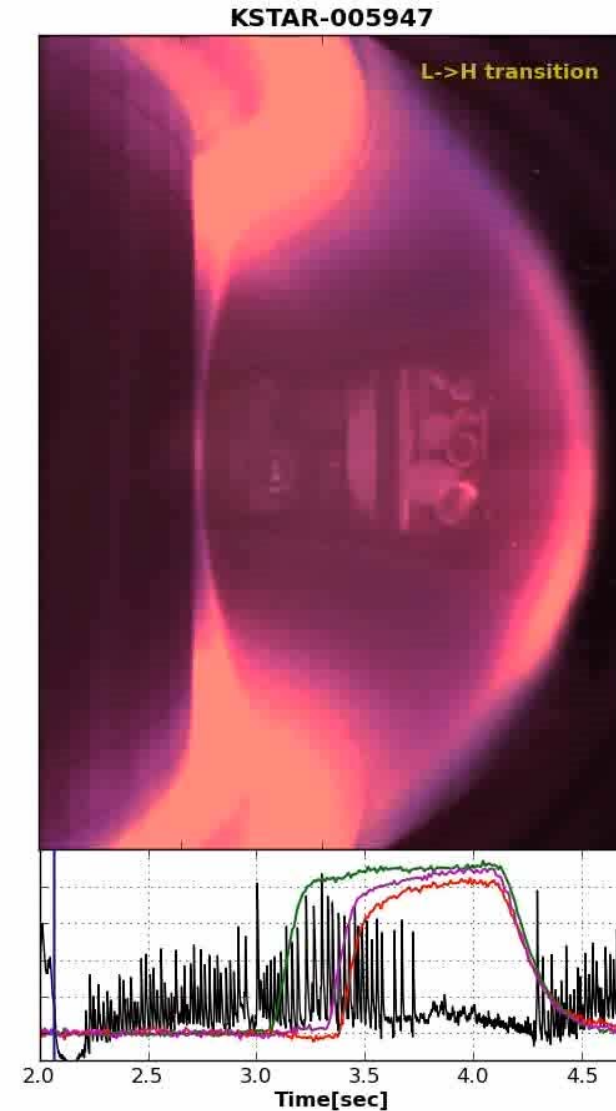
Particles/heat transport through the finger

Ideal MHD Instabilities

- Edge Localised Modes (ELMs)

- Mitigation by 3D magnetic perturbation

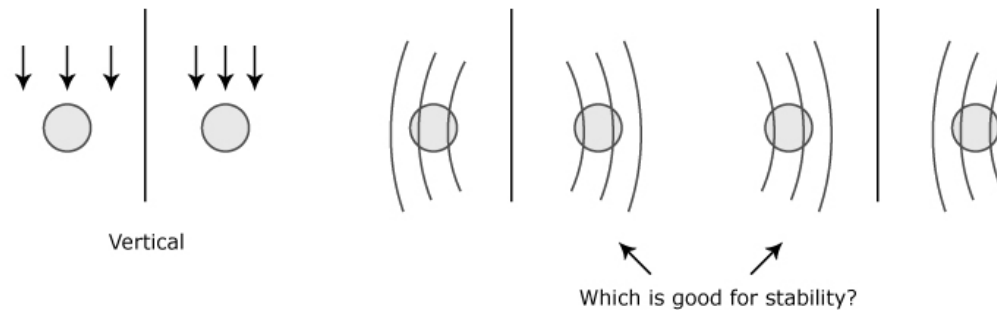
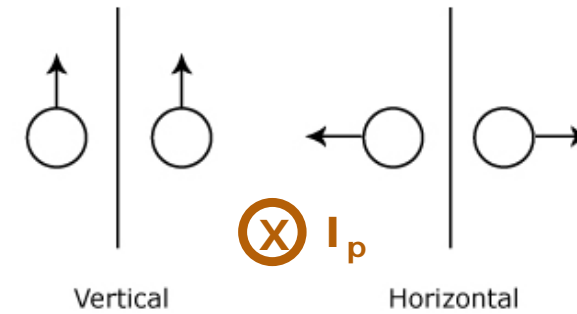
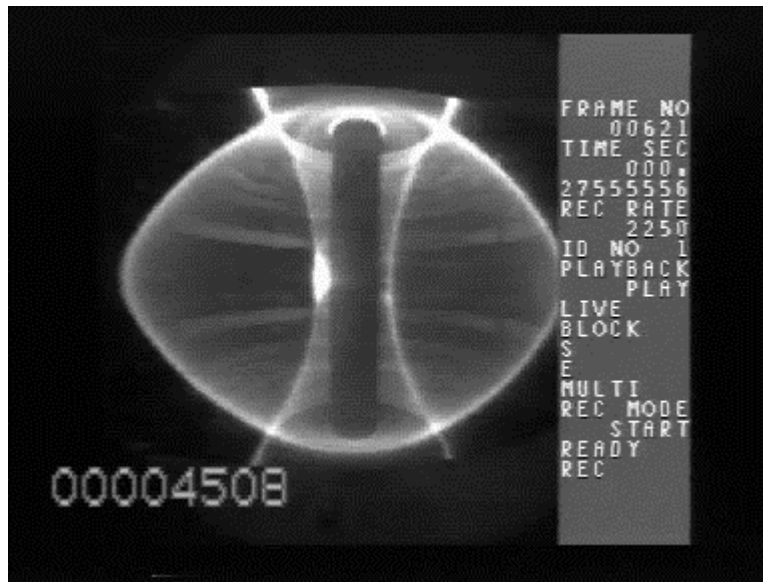
- COMPASS-D ($n = 1$): triggered (2001)
 - DIII-D ($n = 3$): suppressed (2004)
 - JET ($n = 1$ or 2): mitigated (2007)
 - NSTX ($n = 3$): triggered (2010)
 - MAST ($n = 3$): mitigated (2011)
 - ASDEX Upgrade ($n = 2$): mitigated/suppressed (2011)
 - KSTAR ($n = 1$): ELMs suppressed (2011)



Ideal MHD Instabilities

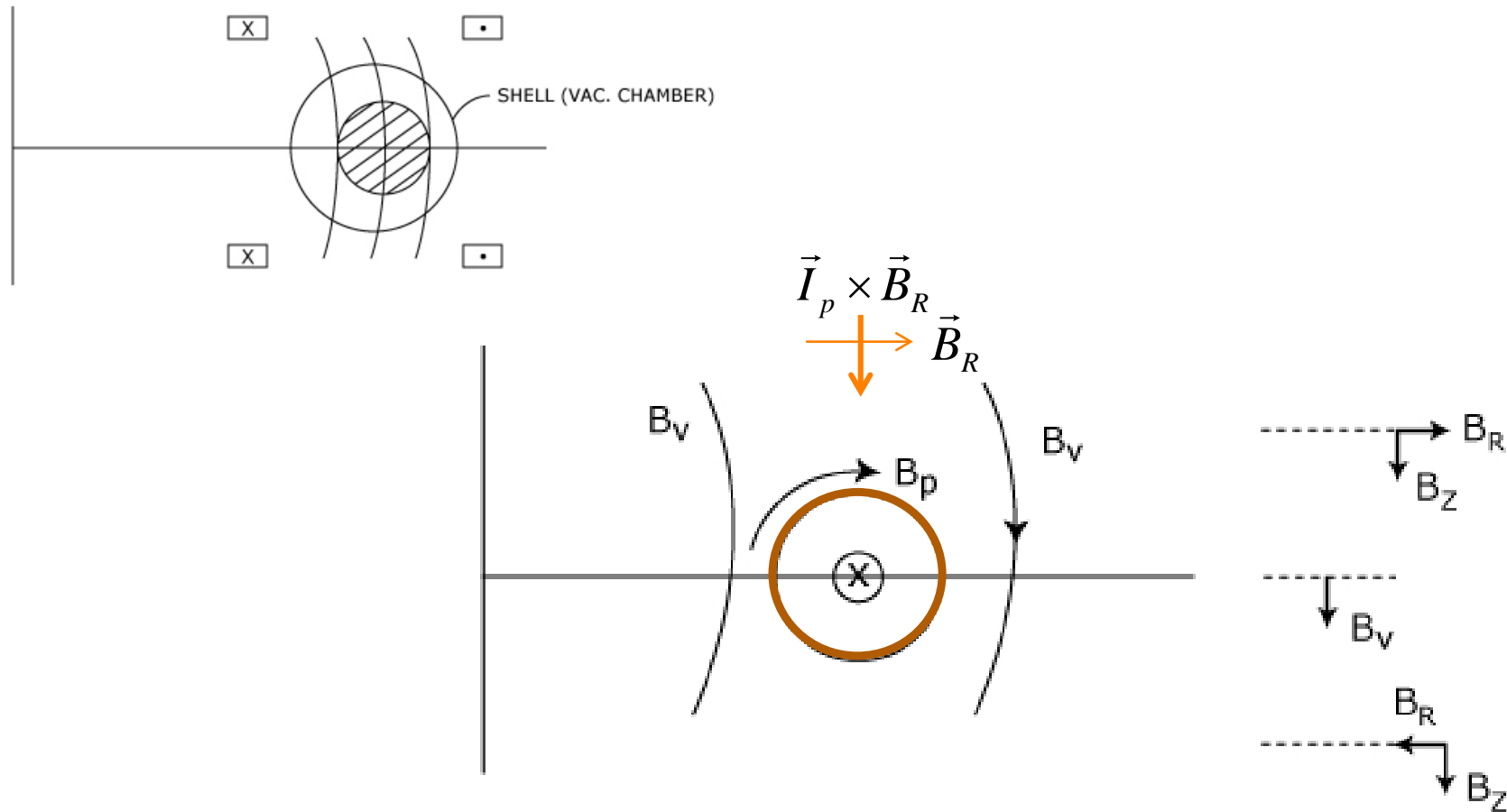
- Vertical Instability

- Macroscopic vertical motion of the plasma towards the wall



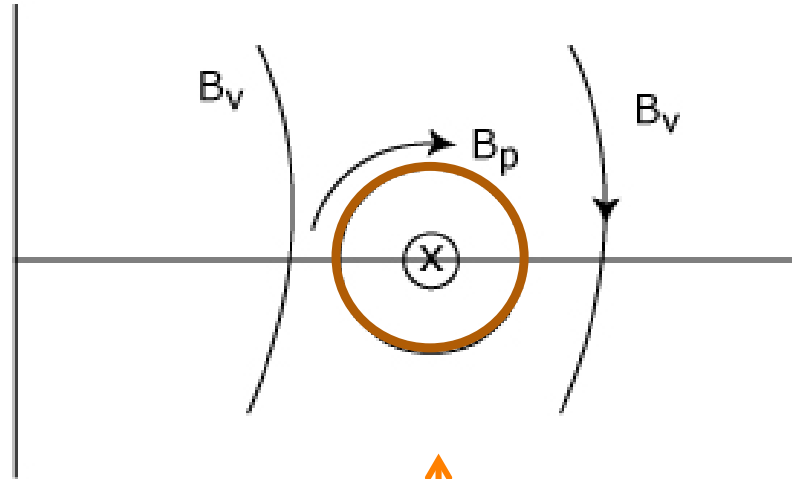
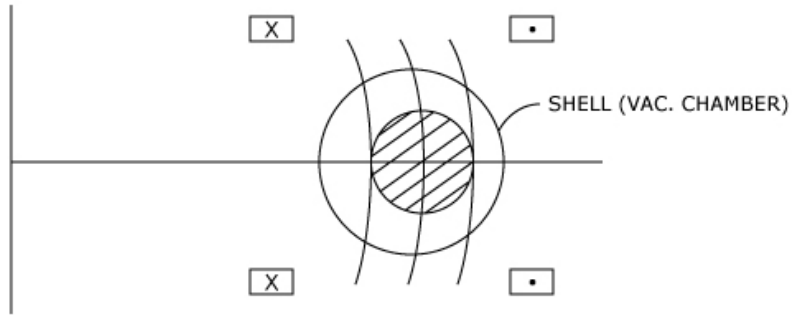
Ideal MHD Instabilities

- Vertical Instability



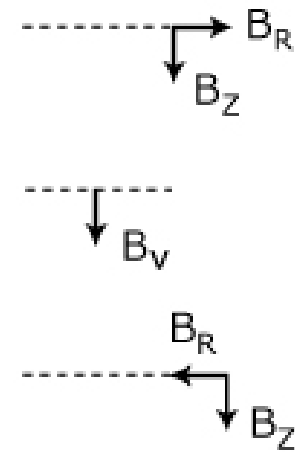
Ideal MHD Instabilities

- Vertical Instability



$$\vec{I}_p \times \vec{B}_R$$

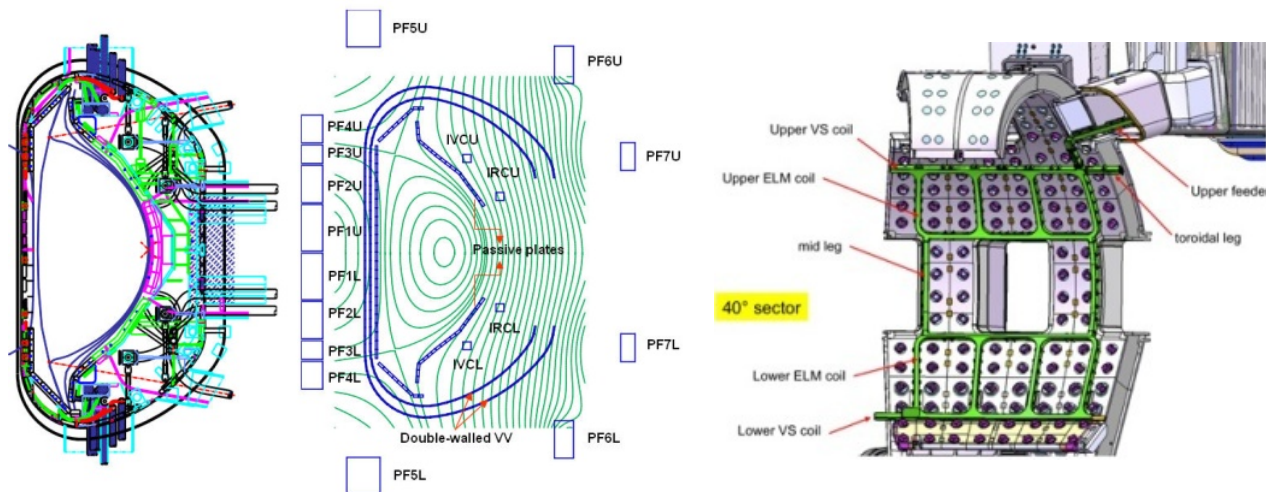
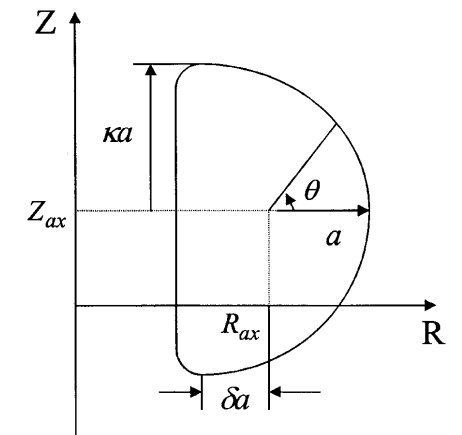
A vector diagram showing the cross product $\vec{I}_p \times \vec{B}_R$. The poloidal current \vec{I}_p is represented by a horizontal arrow pointing to the left, and the radial magnetic field \vec{B}_R is represented by a vertical arrow pointing upwards. The resulting vector is a vertical arrow pointing downwards.



Ideal MHD Instabilities

• Vertical Instability

- For a circular cross sections a moderate shaping of the vertical field should provide stability.
- For noncircular tokamaks, vertical instabilities produce important limitations on the maximum achievable elongations.
- Even moderate elongations require a conducting wall or a feedback system for vertical stability.



KSTAR & ITER