

# Quantum Theory

**Wave-particle duality**

**Schrödinger equation**

**Born interpretation: normalization, quantization**

Reading: Atkins, Ch. 8 (p. 249)

# Wave-particle duality

	<u>C.M.</u>	<u>Exp.</u>
electromagnetic radiation	“wave”	→ also “particle” characteristics
electron	“particle”	→ also “wave” characteristics

## (1) Particle characteristics of electromagnetic radiation

- spectrum: electromagnetic radiation of frequency  $\nu$  possesses the energies of  $0, h\nu, 2h\nu \rightarrow 0, 1, 2$  particles, each particle having  $h\nu$  energy: “**Photon**”

e.g., yellow light (560 nm) of 100 W lamp in 0.1 s (efficiency 100 %)

→ number of photons:  $N = E/h\nu = Pt/h(c/\lambda) = 2.8 \times 10^{20}$

(40 min  $\rightarrow$  1 mol photons)

- photoelectric effect: light energy =  $nh\nu \rightarrow$  particle-like collisions of light

## (2) Wave characteristics of particles

- electron diffraction (Davisson & Germer Exp. (1925)): a characteristic property of waves

→ waves interfere constructively and destructively in different directions:  
wave-like property of electron, molecular hydrogen

# de Broglie relation

- particle  $\rightarrow$  wave-like  
wave  $\rightarrow$  particle-like  $\Rightarrow$  “wave-particle duality”
- 1924, de Broglie (France) suggested that any particle travelling with a linear momentum  $p$  ( $=mv$ ) should have a wavelength of

$$\lambda = h/p: \text{ de Broglie relation}$$

- particle with high momentum  $\rightarrow$  short wavelength  
Macroscopic body: high momentum  $\rightarrow$  wavelength are undetectably small:  
wave-like properties can not be observed

e.g., golf ball, 45 g, velocity 30 m/s  $\rightarrow \lambda = 4.9 \times 10^{-22}$  pm (no wave-like)

- wave-particle duality & quantized energy  $\Rightarrow$  **new mechanics** needed  
(cf. Classical mechanics treated particles and waves as entirely separate entities)

- **wave in new mechanics replaces classical concept of trajectory:** rather than travelling along a definite path, a particle is distributed through space like a wave  
⇒ “wavefunction” ( $\Psi$ ,  $\psi$ )

# Schrödinger equation

- Schrödinger equation (1926): Austrian physicist
- He proposed an equation for finding the wavefunction of any system
- time-independent Schrödinger equation  
particle mass  $m$  moving in 1-dimensional with energy  $E$

$V(x)$ : potential energy of the particle at point  $x$

$E$ : total energy

$\hbar$  (h-cross or h-bar) =  $h/2\pi = 1.05457 \times 10^{-34}$  Js

3-D

In general case, Schrödinger equation

$$\mathbf{H}\Psi = \mathbf{E}\Psi$$

H: hamiltonian operator

Time-dependent Schrödinger equation



cf) Derive Schrödinger equation

# Born interpretation of the wavefunction

- “wavefunction” contains all the dynamic information about the system
- Max Born: interpretation of the wavefunction in terms of the location of the particle

cf: the wave theory of light: square of amplitude of electromagnetic wave = intensity: probability of finding a photon in the region

## 1-D

- if particle has  $\Psi$  at  $x$ , the probability of finding the particle between  $x$  and  $x + dx$  is proportional to  $|\Psi|^2 dx$

$|\Psi|^2 dx = \Psi^* \Psi$  if  $\Psi$  is complex:  $|\Psi|^2$  “probability density”

$\Psi$ : probability amplitude

## 3-D

$\Psi$  at  $r \rightarrow$  probability of finding the particle in  $d\tau = dx dy dz \Rightarrow |\Psi|^2 d\tau$

## (a) Normalization

Schrodinger equation  $\rightarrow N\Psi$ : all probability of the particle must be 1  
 $\rightarrow$  possible to find “normalization constant”  $N$

probability:  $(N\Psi^*)(N\Psi)dx$   
 $\Rightarrow N^2\int\Psi^*\Psi dx = 1 \Rightarrow N = 1/[\int\Psi^*\Psi dx]^{1/2}$

where the integral is over all the space (from  $-\infty$  to  $+\infty$ )

We can find  $N$  and ‘normalize’ the wavefunction  
 $\rightarrow$  normalized wavefunction:  $\int\Psi^*\Psi dx = 1$  or  $\int\Psi^*\Psi d\tau = 1$

$$d\tau = dx dy dz$$

in spherical polar coordinates,  $r, \theta, \phi$

$$x = r \sin \theta \cos \phi, \quad y = r \sin \theta \sin \phi, \quad z = r \cos \theta$$

$$d\tau = r^2 \sin \theta dr d\theta d\phi, \quad r: 0 \rightarrow \infty, \quad \theta: 0 \rightarrow \pi, \quad \phi: 0 \rightarrow 2\pi$$

## (b) Quantization

$\int \Psi^* \Psi d\tau = 1 \Rightarrow$  severe restrictions on the acceptability of wavefunctions

(i)  $\Psi$  must not be infinite anywhere

if it were  $\Rightarrow N \int \Psi^* \Psi = \infty = 1 \Rightarrow N \infty = 1 \Rightarrow N = 0$  (x)

cf: acceptable: infinite  $\Psi$  over infinitesimal since  $\int \Psi^* \Psi$  is finite  
(infinitely high x infinitely narrow = finite area) e.g., a particle at a single, precise point

(ii)  $|\Psi|^2 = \Psi^*\Psi$ : probability of finding the particle  $\Rightarrow$  wavefunction ( $\Psi$ ) must be single-valued

(iii)  $\Psi$ : 2<sup>nd</sup>-order differential equation  $\Rightarrow$  2<sup>nd</sup> derivative should exist:  $\Psi$  should be continuous

1<sup>st</sup> derivative (slope) should also be continuous

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- $\therefore \Psi$  must be continuous, have a continuous slope, be single-valued, and be finite everywhere, cannot be zero everywhere (particle must be somewhere)
- $\Rightarrow$  the energy of a particle is quantized (acceptable solutions of the Schrödinger equation for these severe restrictions at only certain energies)