

Mathematical Foundation of Kinetic Monte Carlo

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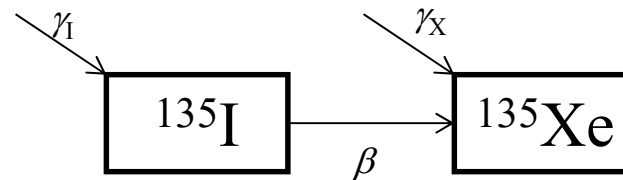
Contents

1. Application of KMC for Solving the Depletion Equation
2. Mathematical Foundation of KMC



Problem: Depletion Eq. for Xe-135

- A simple decay chain of I-135 & Xe-135 can be shown as below.



- And the corresponding depletion equations can be written by

$$\frac{dN_I}{dt} = \gamma_I \Sigma_f \phi - \lambda_I N_I, \quad \text{----- (A.1)}$$

$$\frac{dN_X}{dt} = \gamma_X \Sigma_f \phi + \lambda_I N_I - (\lambda_X + \sigma_X \phi) N_X \quad \text{----- (A.2)}$$

- And from initial conditions of the clean state at $t=0$, the two nuclide densities can be obtained by

$$N_I(t) = \frac{\gamma_I \Sigma_f \phi}{\lambda_I} (1 - e^{-\lambda_I t}) \quad \text{----- (A.3)}$$

$$N_X(t) = \frac{\gamma_X \Sigma_f \phi + \gamma_I \Sigma_f \phi}{\lambda_X + \sigma_X \phi} (1 - e^{-(\lambda_X + \sigma_X \phi)t}) - \frac{\gamma_I \Sigma_f \phi}{\lambda_X + \sigma_X \phi - \lambda_I} (e^{-\lambda_I t} - e^{-(\lambda_X + \sigma_X \phi)t}) \quad \text{----- (A.4)}$$

Densities of I-135 & Xe-135

- For simplicity, Eqs. (A.1) & (A.2) can be rewritten by

$$\frac{dN_I}{dt} = k_I^C - \kappa_I^D N_I, \quad \text{..... (A.5)}$$

$$\frac{dN_X}{dt} = k_X^C + \kappa_I^D N_I - \kappa_X^D N_X \quad \text{..... (A.6)}$$

- Solve Eqs. (A.5) & (A.6) under the following conditions[*] by the kinetic Monte Carlo method.

$$k_I^C = \gamma_I \Sigma_f \phi = (0.056) \times (3 \times 10^{-3} \text{ cm}^{-1}) \times (10^4 \text{ cm}^{-2} \text{ sec}^{-1}) = 1.68 / \text{cc} \cdot \text{sec}$$

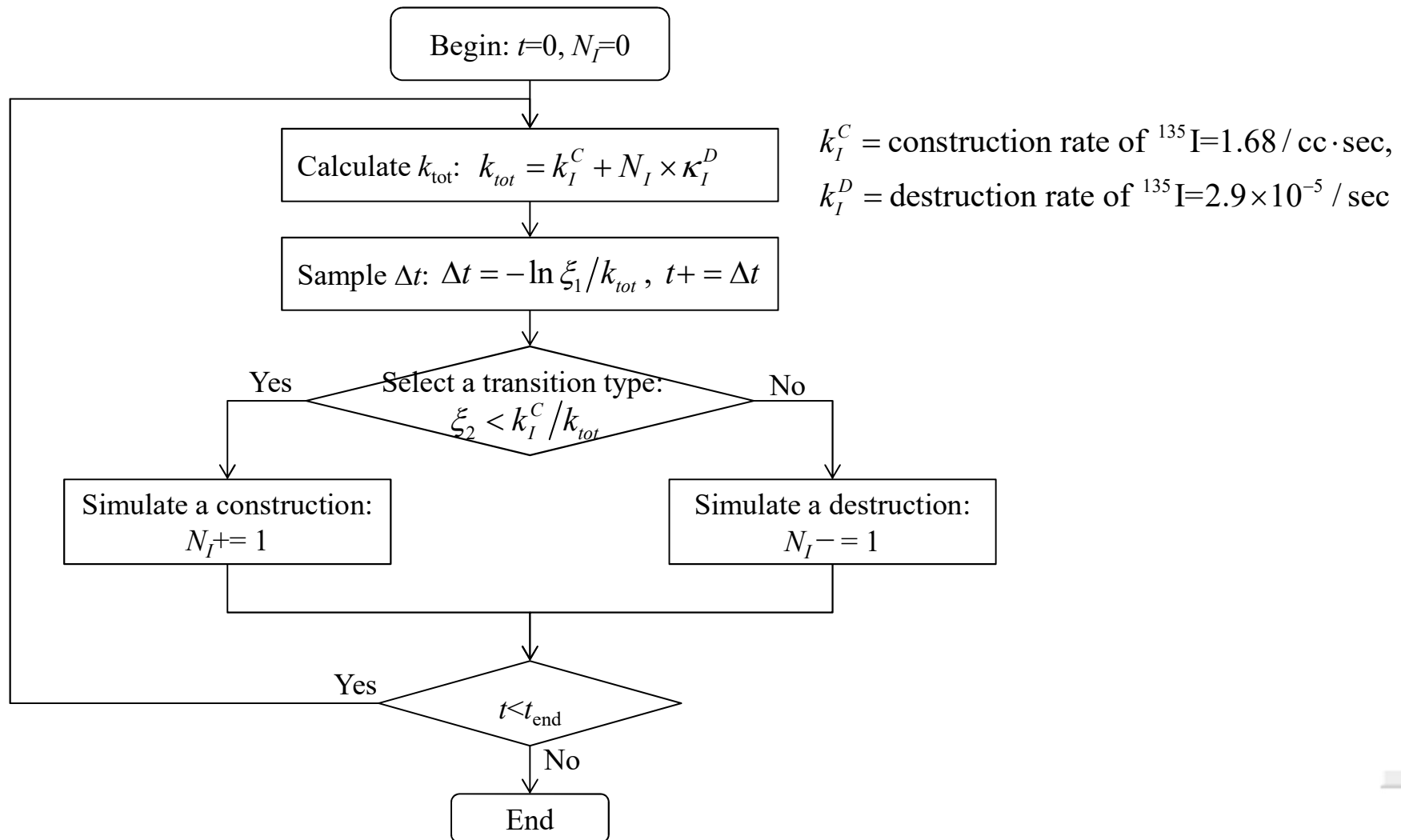
$$\kappa_I^D = \lambda_I = 2.9 \times 10^{-5} / \text{sec}$$

$$k_X^C = \gamma_X \Sigma_f \phi = (0.003) \times (3 \times 10^{-3} \text{ cm}^{-1}) \times (10^4 \text{ cm}^{-2} \text{ sec}^{-1}) = 9 \times 10^{-2} / \text{cc} \cdot \text{sec}$$

$$\begin{aligned} \kappa_X^D &= \lambda_X + \sigma_X \phi = 2.1 \times 10^{-5} \text{ sec}^{-1} + (3.5 \times 10^{-18} \text{ cm}^2) \times (10^4 \text{ cm}^{-2} \text{ sec}^{-1}) \\ &= 2.1000000035 \times 10^{-5} / \text{sec} \end{aligned}$$

[*] D. L. Hetrick, *Dynamics of Nuclear Reactors*, American Nuclear Society, Inc., La Grange Park, IL (1993).

CASE #1: Density of I-135 – Algorithm



CASE #1: Density of I-135 – Implementation

```
do {
  //Calculate the total conversion rate
  totalLambda = genRateI + lambdaI*amountI;

  //Sample delta_time
  deltaT = -log(RNG.GetRN())/totalLambda;

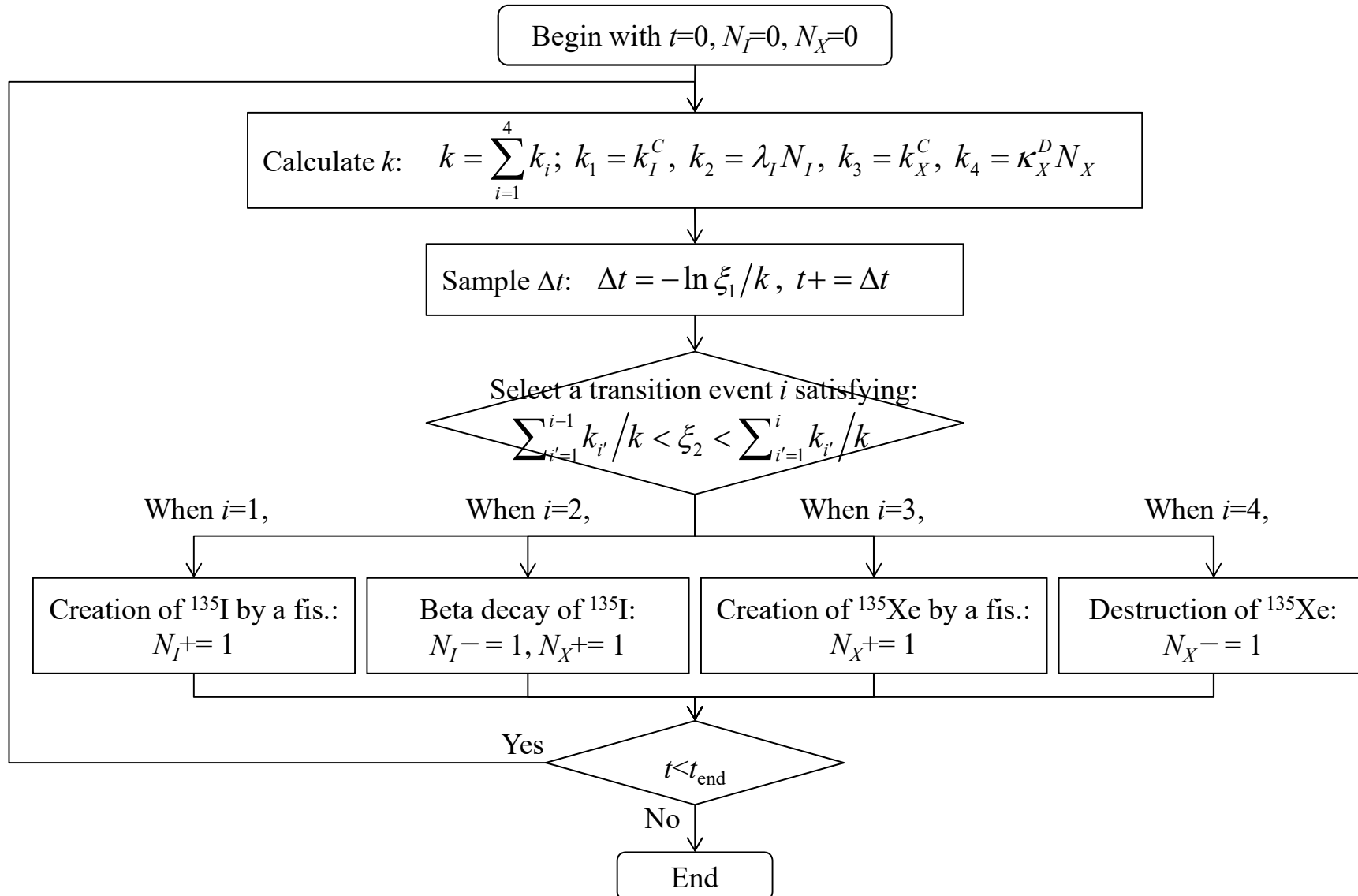
  //Select a conversion type
  if(RNG.GetRN() < genRateI/totalLambda) {
    //Simulate a generation of I-135
    amountI += 1.;
  }
  else {
    //Simulate a destruction of I-135
    amountI -= 1.;
  }

  //Update time
  time += deltaT;

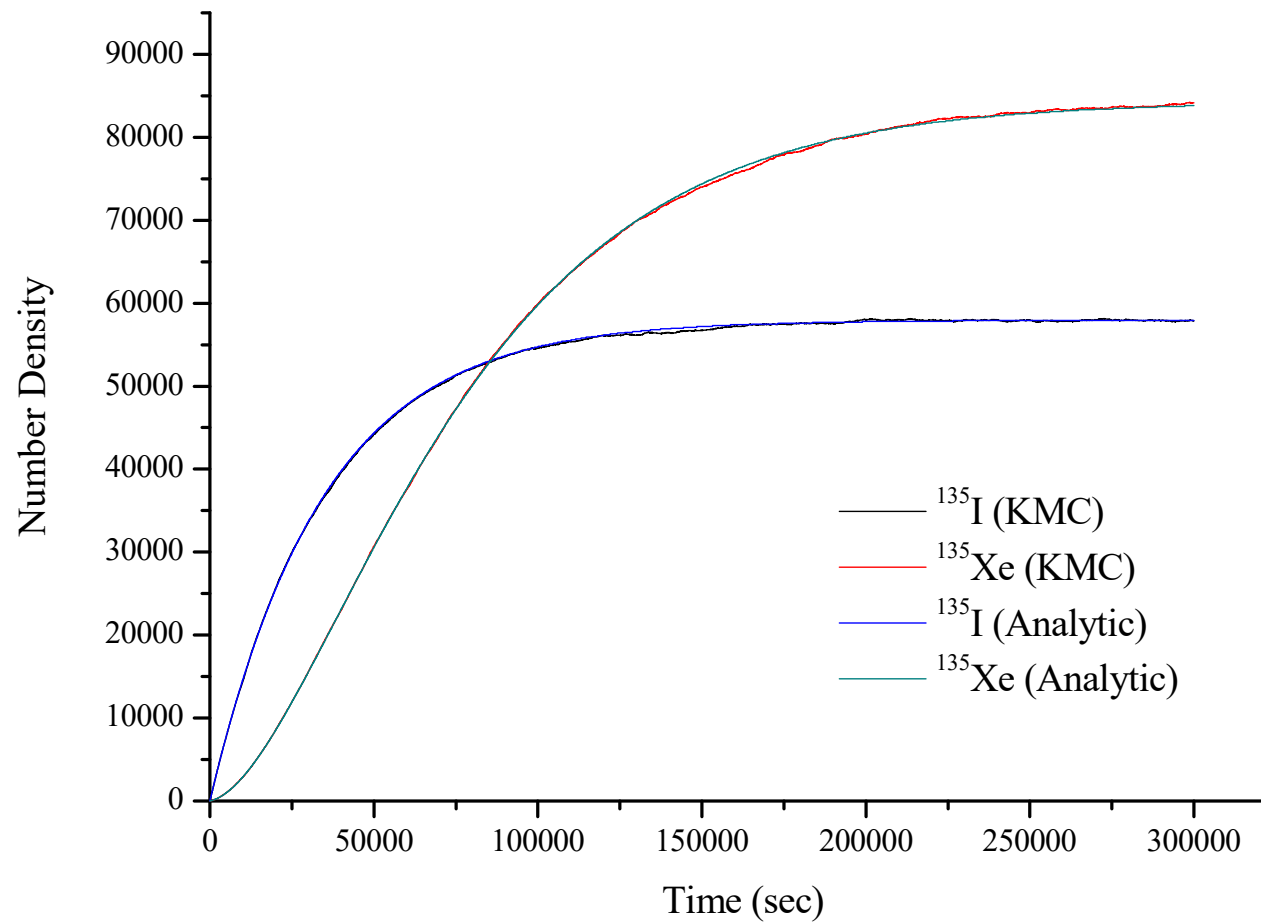
  if(time >= chkTime) {
    out << setw(10) << time << setw(15) << amountI << endl;
    chkTime += binTime;
  }

  //Check the end of simulation time
  if(time > 1.e-3) bEnd = true;
}while(!bEnd);
```

Algorithm for Xe-135



Numerical Results



Mathematical Foundation of Kinetic Monte Carlo Method



Kinetic Monte Carlo vs. Poisson Process

- K. A. Fichthorn and W. H. Weinberg [1] presented the theoretical basis for a dynamical Monte Carlo method in terms of the theory of Poisson process.
- They showed that if
 1. A “dynamical hierarchy” of transition probabilities is created which also satisfy the detailed-balance criterion;
 2. Time increments upon successful events are calculated appropriately;
 3. The effective independence of various events comprising the system can be achieved,then Monte Carlo methods may be utilized to simulate the Poisson process.

[1] K. A. Fichthorn and W.H Weinberg, “Theoretical Foundations of Dynamical Monte Carlo Simulations,” *J. Chem. Phys.*, **95**(2), 1090 (1991).

Kinetic Monte Carlo vs. Poisson Process (Contd.)

- They stated that under a dynamical interpretation, the Monte Carlo method provides a numerical solution to the Master equation:

$$\frac{\partial P(\mathbf{X}, t)}{\partial t} = \sum_{\mathbf{X}'} K(\mathbf{X}' \rightarrow \mathbf{X}) P(\mathbf{X}', t) - \sum_{\mathbf{X}'} K(\mathbf{X} \rightarrow \mathbf{X}') P(\mathbf{X}, t) \quad \text{..... (1)}$$

where \mathbf{X} and \mathbf{X}' are successive states of the system.

$P(\mathbf{X}, t)$ is the probability that the system is in state \mathbf{X} at time t .

$K(\mathbf{X}' \rightarrow \mathbf{X})$ is the probability per unit time that the system will undergo a transition from state \mathbf{X}' to state \mathbf{X} .

- The kinetic Monte Carlo method is explained as
 - The solution of the Master equation is achieved computationally by choosing randomly among various possible transitions with appropriate probabilities.
 - Upon each successful transition, time is typically incremented in integral units of Monte Carlo steps.

Mathematical Foundation of KMC

- We are going to derive a mathematical formulation corresponding to the kinetic Monte Carlo method from the general kinetic equation as

$$\frac{\partial P(\mathbf{X}, t)}{\partial t} = \sum_{\mathbf{X}'} k_{\mathbf{X}' \rightarrow \mathbf{X}} P(\mathbf{X}', t) - \sum_{\mathbf{X}'} k_{\mathbf{X} \rightarrow \mathbf{X}'} P(\mathbf{X}, t), \quad \text{..... (2)}$$

and the initial condition at $t=0$ given by

$$Q(\mathbf{X}) \equiv P(\mathbf{X}, 0) \quad \text{..... (3)}$$

- For simplicity, we rewrite Eq. (2) as

$$\frac{\partial P(\mathbf{X}, t)}{\partial t} + k_{\mathbf{X}} P(\mathbf{X}, t) = S(\mathbf{X}, t); \quad \text{..... (4)}$$

$$k_{\mathbf{X}} = \sum_{\mathbf{X}'} k_{\mathbf{X} \rightarrow \mathbf{X}'}, \quad \text{..... (5)}$$

$$S(\mathbf{X}, t) = \sum_{\mathbf{X}'} k_{\mathbf{X}' \rightarrow \mathbf{X}} P(\mathbf{X}', t). \quad \text{..... (6)}$$

Derivation of KMC Algorithm

- Equation (4) with the initial condition of Eq. (3) is seen to be a first-order linear partial differential equation which has a unique solution. By introducing an integrating factor, Eq. (4) becomes

$$\frac{\partial}{\partial t} \left[e^{k_{\mathbf{x}}t} \cdot P(\mathbf{X}, t) \right] = e^{k_{\mathbf{x}}t} \cdot S(\mathbf{X}, t) \quad \dots\dots\dots (7)$$

- Then an integration of Eq. (7) from $t=0$ yields

$$P(\mathbf{X}, t)e^{k_{\mathbf{x}}t} - Q(\mathbf{X}) = \int_0^t e^{k_{\mathbf{x}}t'} \cdot S(\mathbf{X}, t') dt'$$

$$\Rightarrow P(\mathbf{X}, t) = Q(\mathbf{X})e^{-k_{\mathbf{x}}t} + \int_0^t e^{-k_{\mathbf{x}}(t-t')} \cdot S(\mathbf{X}, t') dt' \quad \dots\dots\dots (8)$$

- Eq. (8) implies that the probability of state \mathbf{X} is made up of the state appeared in the previous time multiplied by the attenuation factor of $e^{-k_{\mathbf{x}}(t-t')}$.

Transition Probability Equation

- The transition probability can be defined by

$$\Psi(\mathbf{X}, t) \equiv k_{\mathbf{X}} P(\mathbf{X}, t) \quad \text{----- (9)}$$

- Then by multiplying $k_{\mathbf{X}}$ on the both sides of Eq. (8), it can be expressed as

$$\Psi(\mathbf{X}, t) = k_{\mathbf{X}} e^{-k_{\mathbf{X}} t} Q(\mathbf{X}) + \int_0^t k_{\mathbf{X}} e^{-k_{\mathbf{X}}(t-t')} \cdot S(\mathbf{X}, t') dt' \quad \text{----- (10)}$$

- By introducing the time-flight kernel, T defined by

$$T(t' \rightarrow t | \mathbf{X}) = k_{\mathbf{X}} e^{-k_{\mathbf{X}}(t-t')} \quad \text{----- (11)}$$

Eq. (10) becomes

$$\Psi(\mathbf{X}, t) = \hat{Q}(\mathbf{X}, t) + \int_0^t T(t' \rightarrow t | \mathbf{X}) \cdot S(\mathbf{X}, t') dt'; \quad \text{----- (12)}$$

$$\hat{Q}(\mathbf{X}, t) = T(0 \rightarrow t | \mathbf{X}) Q(\mathbf{X}) \quad \text{----- (13)}$$

$\hat{Q}(\mathbf{X}, t)$ is named the first transition source.

Transition Density Equation (Contd.)

- And we define another kernel named the event kernel as

$$C(\mathbf{X}' \rightarrow \mathbf{X}) = \frac{k_{\mathbf{X}' \rightarrow \mathbf{X}}}{k_{\mathbf{X}'}} \quad \text{----- (14)}$$

- Using the event kernel C and the transition probability Ψ , $S(\mathbf{X}, t)$ of Eq. (6) can be expressed as

$$\begin{aligned} S(\mathbf{X}, t) &= \sum_{\mathbf{X}'} k_{\mathbf{X}'} C(\mathbf{X}' \rightarrow \mathbf{X}) P(\mathbf{X}', t) \\ &= \sum_{\mathbf{X}'} C(\mathbf{X}' \rightarrow \mathbf{X}) \Psi(\mathbf{X}', t) \quad \text{----- (15)} \end{aligned}$$

- The introduction of Eq. (15) into Eq. (12) yields

$$\Psi(\mathbf{X}, t) = \hat{Q}(\mathbf{X}, t) + \int_0^t \sum_{\mathbf{X}'} K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) \Psi(\mathbf{X}', t') dt' \quad \text{----- (16)}$$

$$K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) = T(t' \rightarrow t | \mathbf{X}) \cdot C(\mathbf{X}' \rightarrow \mathbf{X}) \quad \text{----- (17)}$$

Series Solution

$$\Psi(\mathbf{X}, t) = \hat{Q}(\mathbf{X}, t) + \int_0^t \sum_{\mathbf{X}'} K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) \Psi(\mathbf{X}', t') dt' \quad \text{..... (16)}$$

- Let's consider the solution of Eq. (16) obtained by iteration; thus

$$\psi_0(\mathbf{X}, t) = \hat{Q}(\mathbf{X}, t)$$

$$\psi_1(\mathbf{X}, t) = \int_0^t \sum_{\mathbf{X}'} K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) \psi_0(\mathbf{X}', t') dt'$$

$$\vdots \quad \quad \quad \vdots$$

$$\psi_n(\mathbf{X}, t) = \int_0^t \sum_{\mathbf{X}'} K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) \psi_{n-1}(\mathbf{X}', t') dt'$$

Clearly ψ_0 is the first-transition source. ψ_1 means the transition probability from the second-transition state. Similarly, ψ_2 indicates the contribution of the third-

transition state, and so on. If the series $\sum_{j=0}^{\infty} \psi_j(\mathbf{X}, t)$ converges, it represents a solution to Eq. (16).

Neumann Series Solution

$$\Psi(\mathbf{X}, t) = \hat{Q}(\mathbf{X}, t) + \int_0^t \sum_{\mathbf{X}'} K(\mathbf{X}', t' \rightarrow \mathbf{X}, t) \Psi(\mathbf{X}', t') dt' \quad \text{----- (16)}$$

- The solution of Eq. (16) can be expressed by the Neumann series:

$$\Psi(\mathbf{X}, t) = \sum_{j=0}^{\infty} \psi_j(\mathbf{X}, t); \quad \text{----- (18)}$$

$$\psi_j(\mathbf{X}, t) = \int_0^t dt_0 \sum_{\mathbf{X}_0} K_j(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t) \hat{Q}(\mathbf{X}_0, t_0), \quad \text{----- (19)}$$

$$K_0(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t) = \delta(\mathbf{X}_0 - \mathbf{X}) \delta(t_0 - t),$$

$$K_1(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t) = K(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t),$$

$$K_2(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t) = \int_{t_0}^t dt_1 \sum_{\mathbf{X}_1} K(\mathbf{X}_1, t_1 \rightarrow \mathbf{X}, t) K(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}_1, t_1),$$

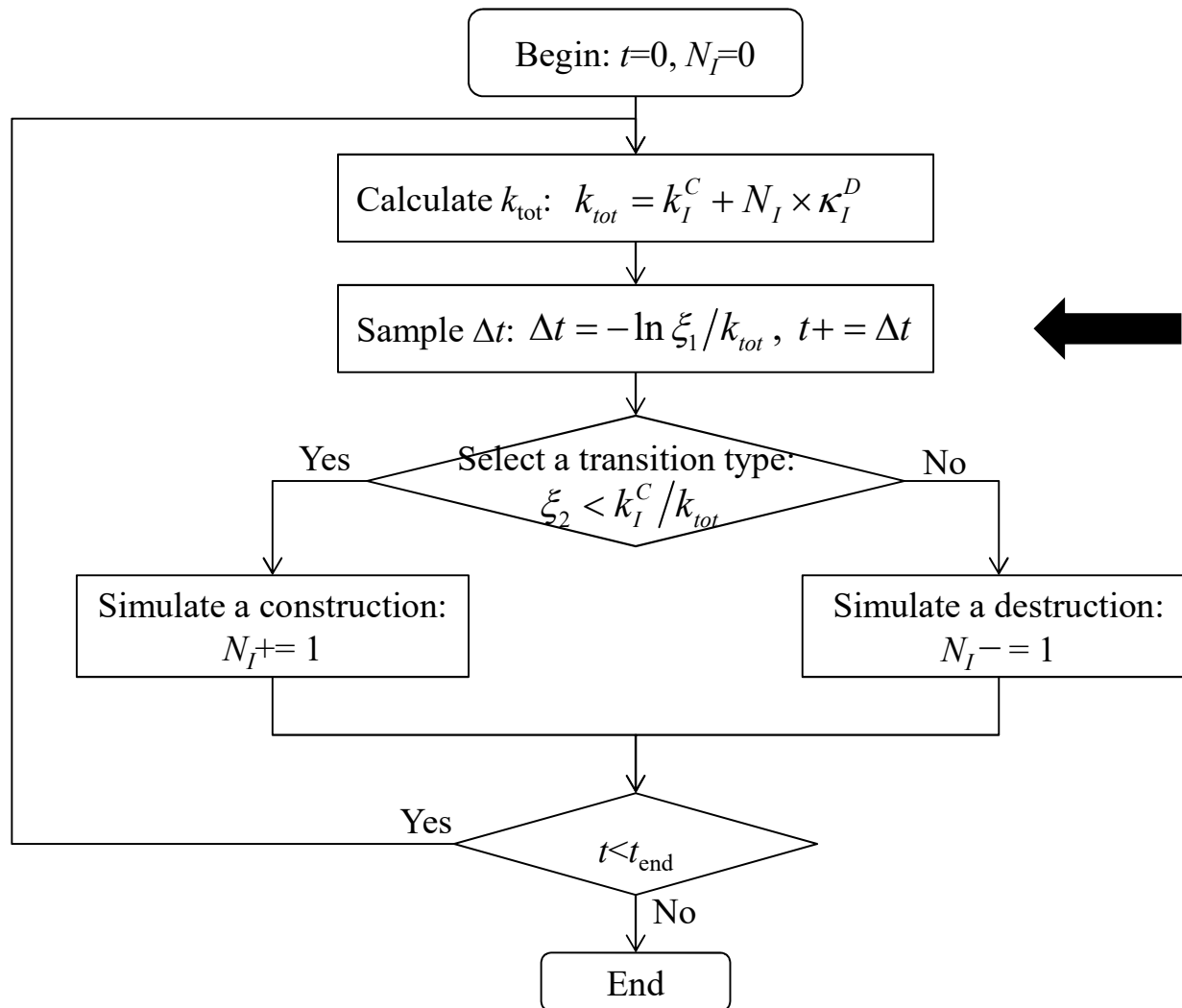
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$$K_j(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}, t) = \int dt_1 \sum_{\mathbf{X}_1} \cdots \int dt_{j-1} \sum_{\mathbf{X}_{j-1}} K(\mathbf{X}_{j-1}, t_{j-1} \rightarrow \mathbf{X}, t) K(\mathbf{X}_{j-2}, t_{j-2} \rightarrow \mathbf{X}_{j-1}, t_{j-1}) \cdots K(\mathbf{X}_0, t_0 \rightarrow \mathbf{X}_1, t_1) \quad \text{----- (20)}$$

From Eq. (18), we can find that the transition probability is the sum of the contributions from transition at (\mathbf{X}, t) first and after a transition or more.

- The Monte Carlo kinetic analysis is based on Eqs. (18) & (19).

Interpretation of KMC by MC Sim. of Series Sol.



$$T(\tau) = k_{tot} e^{-k_{tot}\tau} \Rightarrow \int_0^\tau k_{tot} e^{-k_{tot}\tau'} d\tau' = \xi_1$$

$$\Rightarrow \tau = -\frac{\ln \xi_1}{k_{tot}}$$

$$C(\mathbf{X}' \rightarrow \mathbf{X}) = \frac{k_{\mathbf{X}' \rightarrow \mathbf{X}}}{k_{\mathbf{X}'}}$$