

## **Potential Sweep Methods (Ch. 6)**

**Nernstian (reversible) systems**

**Totally irreversible systems**

**Quasireversible systems**

**Cyclic voltammetry**

**Multicomponent systems & multistep charge transfers**

## **Introduction**

Linear sweep voltammetry (LSV)

Cyclic voltammetry (CV)

## Nernstian (reversible) systems

### Solution of the boundary value problem

O + ne = R (semi-infinite linear diffusion, initially O present)

$$E(t) = E_i - vt$$

Sweep rate (or scan rate):  $v$  (V/s)

Rapid e-transfer rate at the electrode surface

$$C_O(0, t)/C_R(0, t) = f(t) = \exp[nF (E_i - vt - E^{0'})/RT]$$

$$i = nFAC_O^*(\pi D_O \sigma)^{1/2} \chi(\sigma t)$$

$$\sigma = (nF/RT)v$$

## Peak current and potential

Peak current:  $\pi^{1/2}\chi(\sigma t) = 0.4463$

$$i_p = 0.4463(F^3/RT)^{1/2}n^{3/2}AD_O^{1/2}C_O^*v^{1/2}$$

At 25°C, for A in cm<sup>2</sup>, D<sub>O</sub> in cm<sup>2</sup>/s, C<sub>O</sub><sup>\*</sup> in mol/cm<sup>3</sup>, v in V/s → i<sub>p</sub> in amperes

$$i_p = (2.69 \times 10^5)n^{3/2}AD_O^{1/2}C_O^*v^{1/2}$$

Peak potential, E<sub>p</sub>

$$E_p = E_{1/2} - 1.109(RT/nF) = E_{1/2} - 28.5/n \text{ mV at } 25^\circ\text{C}$$

Half-peak potential, E<sub>p/2</sub>

$$E_{p/2} = E_{1/2} + 1.09(RT/nF) = E_{1/2} + 28.0/n \text{ mV at } 25^\circ\text{C}$$

E<sub>1/2</sub> is located between E<sub>p</sub> and E<sub>p/2</sub>

$$|E_p - E_{p/2}| = 2.20(RT/nF) = 56.5/n \text{ mV at } ^\circ\text{C}$$

For reversible wave, E<sub>p</sub> is independent of scan rate, i<sub>p</sub> is proportional to v<sup>1/2</sup>

## Spherical electrodes and UMEs

Spherical electrode (e.g., a hanging mercury drop)

$$i = i(\text{plane}) + nFAD_0C_0 * \phi(\sigma t)/r_0$$

$\phi(\sigma t)$ : tabulated function (Table 6.2.1)

For large  $v$  in conventional-sized electrode  $\rightarrow i(\text{plane}) \gg 2^{\text{nd}}$  term

Same for hemispherical & UME at fast scan rate

For UME at very small  $v$ :  $r_0$  is small  $\rightarrow i(\text{plane}) \ll 2^{\text{nd}}$  term

$\rightarrow$  voltammogram is a steady-state response independent of  $v$

$\rightarrow v \ll \text{RTD}/nFr_0^2$

$r_0 = 5 \mu\text{m}$ ,  $D = 10^{-5} \text{ cm}^2/\text{s}$ ,  $T = 298 \text{ K} \rightarrow$  steady-state voltammogram at  $v < 1 \text{ V/s}$

$r_0 = 0.5 \mu\text{m} \rightarrow$  steady-state behavior up to  $10 \text{ V/s}$

Transition from typical peak-shaped voltammograms at fast  $v$  to steady-state voltammograms at small  $v$

cf. For potential sweep (Ch.1)

Linear potential sweep with a sweep rate  $v$  (in V/s)

$$E = vt$$

$$E = E_R + E_C = iR_s + q/C_d$$

$$vt = R_s(dq/dt) + q/C_d$$

If  $q = 0$  at  $t = 0$ ,

$$i = vC_d[1 - \exp(-t/R_sC_d)]$$

- Current rises from 0 and attains a steady-state value ( $vC_d$ ): measure  $C_d$

## Effect of double-layer capacitance & uncompensated resistance

Charging current at potential sweep

$$|i_c| = AC_d v$$

Faradaic current measured with baseline of  $i_c$

$i_p$  varies with  $v^{1/2}$ ,  $i_c$  varies with  $v \rightarrow$   $i_c$  more important at faster  $v$

$$|i_c|/i_p = [C_d v^{1/2} (10^{-5})] / [2.69 n^{3/2} D_O^{1/2} C_O^*]$$

At high  $v$  & low  $C_O^*$   $\rightarrow$  severe distortion of the LSV wave

$R_u$  cause  $E_p$  to be a function of  $v$

## Totally irreversible systems

### Solution of the boundary value problem

Totally irreversible one-step, one-electron reaction:  $O + e \xrightarrow{k_f} R$

$$i/FA = D_O(\partial C_O(x, t)/\partial x)_{x=0} = k_f(t)C_O(0, t)$$

Where  $k_f = k^0 e^{-\alpha f(E(t) - E^{0'})}$ ,  $E(t) = E_i - vt$

$$\rightarrow k_f(t)C_O(0, t) = k_{fi}C_O(0, t)e^{bt}$$

Where  $b = \alpha fv$  &  $k_{fi} = k^0 \exp[-\alpha f(E_i - E^{0'})]$

$$i = FAC_O^* D_O^{1/2} v^{1/2} (\alpha F/RT)^{1/2} \chi(bt)$$

$\chi(bt)$  (Table 6.3.1).  $i$  varies with  $v^{1/2}$  and  $C_O^*$

For spherical electrodes

$$i = i(\text{plane}) + FAD_O C_O^* \phi(bt)/r_0$$



## Peak current and potential

Maximum  $\chi(bt)$  at  $\pi^{1/2}\chi(bt) = 0.4958$

Peak current

$$i_p = (2.99 \times 10^5) \alpha^{1/2} A C_O^* D_O^{1/2} v^{1/2}$$

n-electron process with RDS: n in right side

Peak potential

$$\alpha(E_p - E^{0'}) + (RT/F) \ln[(\pi D_O b)^{1/2}/k^0] = -0.21(RT/F) = -5.34 \text{ mV at } 25^\circ\text{C}$$

Or

$$E_p = E^{0'} - (RT/\alpha F)[0.780 + \ln(D_O^{1/2}/k^0) + \ln(\alpha F v/RT)^{1/2}]$$
$$|E_p - E_{p/2}| = 1.857RT/\alpha F = 47.7/\alpha \text{ mV at } 25^\circ\text{C}$$

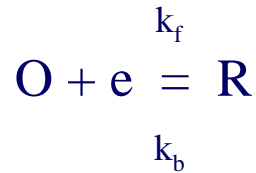
$E_p$ : ftn of  $v \rightarrow$  for reduction,  $1.15RT/\alpha F$  (or  $30/\alpha$  mV at  $25^\circ\text{C}$ ) negative shift for tenfold increase in  $v$

$$i_p = 0.227FAC_O^* k^0 \exp[-\alpha f(E_p - E^{0'})]$$

$\rightarrow i_p$  vs.  $E_p - E^{0'}$  plot at different  $v$ : slope of  $-\alpha f$  and intercept proportional to  $k^0$   
n-electron process with RDS: n in right side

## Quasireversible systems

For one-step, one-electron system



For the quasireversible one-step, one-electron case (5.5.3, p. 191)

$$i/FA = D_O(\partial C_O(x, t)/\partial x)_{x=0} = k_f C_O(0, t) - k_b C_R(0, t)$$

Where  $k_f = k^0 e^{-\alpha f(E - E_0')}$  &  $k_b = k^0 e^{(1 - \alpha)f(E - E_0')}$ ,  $f = F/RT$

The shape of peak & peak parameters  $\rightarrow$  ftns of  $\alpha$  &  $\Lambda$

$$\Lambda = k^0 / (D_O^{1-\alpha} D_R^\alpha f v)^{1/2}$$

Or for  $D_O = D_R = D$

$$\Lambda = k^0 / (D f v)^{1/2}$$

Current

$$i = F A D_O^{1/2} C_O^* f^{1/2} v^{1/2} \Psi(E)$$

$\Psi(E)$  (Fig. 6.4.1):  $\Lambda > 10 \rightarrow$  approach to the reversible

## Cyclic voltammetry

$$(0 < t \leq \lambda) \quad E = E_i - vt$$

$$(t > \lambda) \quad E = E_i - 2v\lambda + vt$$

## **Nernstian systems**

i-t curve at different  $E_\lambda$

i-E curve (CV) at different  $E_\lambda$

(1)  $E_\lambda$  (1)  $E_{1/2} - 90/n$ , (2)  $E_{1/2} - 130/n$ , (3)  $E_{1/2} - 200/n$  mV, (4) after  $i_{pc} \rightarrow 0$

$i_{pa}/i_{pc} = 1$  for nernstian regardless of scan rate,  $E_\lambda (> 35/n$  mV past  $E_{pc})$ , D

$i_{pa}/i_{pc} \rightarrow$  kinetic information

If actual baseline cannot be determined,

$$i_{pa}/i_{pc} = (i_{pa})_0/i_{pc} + 0.485(i_{sp})_0/i_{pc} + 0.086$$

Reversal charging current is same as forward scan, but opposite sign

$$\Delta E_p = E_{pa} - E_{pc} \sim 2.3RT/nF \text{ (or } 59/n \text{ mV at } 25^\circ\text{C)}$$

## Quasireversible systems

Wave shape &  $\Delta E_p \rightarrow$  ftns of  $v$ ,  $k^0$ ,  $\alpha$  &  $E_\lambda$

If  $E_\lambda > 90/n$  mV beyond cathodic peak  $\rightarrow$  small  $E_\lambda$  effect

$$\Psi = \Lambda \pi^{-1/2} = [k^0(D_O/D_R)^\alpha]/(\pi D_O f v)^{1/2}$$

(1)  $\Psi = 0.5$ ,  $\alpha = 0.7$ , (2)  $\Psi = 0.5$ ,  $\alpha = 0.3$ , (3)  $\Psi = 7$ ,  $\alpha = 0.5$ , (4)  $\Psi = 0.25$ ,  $\alpha = 0.5$

For  $0.3 < \alpha < 0.7 \rightarrow \Delta E_p$  independent of  $\alpha$ ; depend only on  $\Psi$   
 $\rightarrow$  estimating  $k^0$  in quasireversible systems

$\Delta E_p$  vs.  $v \rightarrow \Delta E_p$  vs  $\Psi$

# Multicomponent systems & Multistep charge transfers

O & O' system



Method for obtaining baselines

Constant E after 1

Sweep stop beyond  $E_{p1}$

In vivo applications of LSV & CV  
e.g., rat brain