



Chapter 11 Diffraction

Reading Assignment: 1. D. Sherwood, Crystals, X-rays, and Proteins-chapter 5 and 7







- geometric arrangement

diffraction pattern amplitude $F(\vec{k})$

$$F(\vec{k}) = \int_{\text{all } \vec{r}} f(\vec{r}) e^{i\vec{k}\cdot\vec{r}} d\vec{r}$$

 $f(\vec{r})$: amplitude function

xz plane obstacle- along x-axis $\vec{\mathbf{r}} = (x, 0, 0)$ $f(\vec{r}) \rightarrow f(x)$ $\vec{k} = (k_x, 0, k_z)$ $\vec{k} \cdot \vec{\mathbf{r}} = k_x x = kx \sin \theta$





- geometric arrangement

$$F(\vec{k}) = \int_{\text{all } \vec{r}} f(\vec{r}) e^{i\vec{k}\cdot\vec{r}} d\vec{r} = \int_{-\infty}^{\infty} f(x) e^{ikx\sin\theta} dx$$
$$\rightarrow F(\sin\theta) = \int_{-\infty}^{\infty} f(x) e^{ikx\sin\theta} dx \quad (\because k: \text{ constant})$$

- in general, $F(\sin \theta)$ is complex

intensity of diffraction pattern $|F(\sin \theta)|^2$





- one narrow slit
 - an infinite opaque sheet along the *x*-axis containing narrow slit at the origin
- narrowness- compared to the wavelength
- amplitude function- δ function





$$-f(x) = \delta(x)$$

- $F(\sin\theta) = \int_{-\infty}^{\infty} f(x)e^{ikx\sin\theta} dx = \int_{-\infty}^{\infty} \delta(x)e^{ikx\sin\theta} dx$
= $\left[e^{ikx\sin\theta}\right]_{x=0} = 1$
- intensity ~ $\left|F(\sin\theta)\right|^2 = 1$

- intensity is uniform at all angles
- a single narrow slit \equiv an active point source







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Interference











- three narrow slits





- N narrow slits

equally spaced by a distance $x_0 \rightarrow N \delta$ function

N = 2p + 1 (odd number)

$$-f(x) = \sum_{n=-p}^{n=p} \delta(x - nx_0)$$



Diffraction by 1-D Obstacles

$$-F(\sin \theta) = \int_{-\infty}^{\infty} f(x)e^{ikx\sin\theta} dx$$

$$= e^{-ikpx_0\sin\theta} + e^{-ik(p-1)x_0\sin\theta} + \dots + 1 + \dots + e^{ikpx_0\sin\theta}$$

$$= e^{-ikpx_0\sin\theta} (1 + e^{ikx_0\sin\theta} + \dots + e^{ik^2px_0\sin\theta})$$

$$= e^{-ikpx_0\sin\theta} (\frac{1 - e^{ik(2p+1)x_0\sin\theta}}{1 - e^{ikx_0\sin\theta}}) = e^{-ikpx_0\sin\theta} (\frac{1 - e^{ikNx_0\sin\theta}}{1 - e^{ikx_0\sin\theta}})$$

$$= \frac{\sin \frac{Nkx_0\sin\theta}{2}}{\sin \frac{kx_0\sin\theta}{2}} \qquad \text{*main peak}$$

$$= \frac{\sin (1 + e^{ikx_0\sin\theta})}{\sin \frac{kx_0\sin\theta}{2}} \qquad \text{*main peak}$$

$$= \frac{\sin (1 + e^{ikx_0\sin\theta})}{\sin \frac{kx_0\sin\theta}{2}} \qquad \text{*main peak}$$

$$= \frac{\sin^2 \frac{Nkx_0\sin\theta}{2}}{\sin^2 \frac{kx_0\sin\theta}{2}} \qquad \text{*nain peak}$$



- infinite number of narrow slits separated by a distance $x_0 \rightarrow$ infinite array of δ function

$$-f(x) = \sum_{n=-\infty}^{\infty} \delta(x - nx_0)$$

$$-F(\sin\theta) = \sum_{n=-\infty}^{\infty} \delta(\sin\theta - \frac{2n\pi}{kx_0})$$

*infinitely sharp peak

$$\Delta(\sin\theta) = \frac{2\pi}{kx_0}$$







- one wide slit

an opaque screen containing a slit which is wide

 $\sim 100\lambda$ f (x) -f(x) = 0 if $-\infty < x < -X_0$ 1 if $-X_0 < x < +X_0$ 0 $if + X_0 < x < \infty$ - F(sin θ)= $\int_{0}^{\infty} f(x)e^{ikx\sin\theta}dx$ $X_0^{\dagger}X_0$ $=\int_{-X_0}^{X_0} f(x)e^{ikx\sin\theta} dx = \left[\frac{e^{ikx\sin\theta}}{ik\sin\theta}\right]_{-X_0}^{X_0}$

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Single-Slit Diffraction



http://www.phys.hawaii.edu/~teb/optics/java/slitdiffr/







D. Halliday, Fundamentals of Physics





Convolution

$$-c(\vec{u}) = f(\vec{r}) * g(\vec{r}) = \int_{all \ \vec{r}} f(\vec{r})g(\vec{u} - \vec{r})d\vec{r} = \int_{all \ \vec{r}} f(\vec{u} - \vec{r})g(\vec{r})d\vec{r}$$

- integrand is a function of \vec{u} and \vec{r}

integration is taken over $\vec{r} \rightarrow$ function of \vec{u}

- $g(\vec{r})$ vs. $g(\vec{u} - \vec{r})$ or g(x) vs. g(u - x) for 1-D

reflection+displacement

- example

f(x): δ function, g(x): arbitrary function

$$c(u) = \int_{-\infty}^{\infty} f(x)g(u-x)dx = \int_{-\infty}^{\infty} \delta(x+x_o)g(u-x)dx$$
$$+ \int_{-\infty}^{\infty} \delta(x-x_0)g(u-x)dx$$
$$= g(u+x_0) + g(u-x_0)$$
http://www.jhu.edu/~signals/convolve/index.html



















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The convolution of a rectangular and a triangular pulse. <u>http://eent3.sbu.ac.uk/staff/baoyb/foct</u>

-two wide slit

slits, each of width $2X_0$, centered at $x = -x_0$ and $x = x_0$ -f(x) = 0 if $-\infty < x < -(x_0 + X_0)$ 1 if $-(x_0 + X_0) < x < -(x_0 - X_0)$ 0 if $-(x_0 - X_0) < x < (x_0 - X_0)$ 1 if $(x_0 - X_0) < x < (x_0 + X_0)$ 0 if $(x_0 + X_0) < x < \infty$

-f(x) is convolution of one wide slit and two narrow slit





- Fourier transform

Tf(two wide slit) = T[f(one wide slit)*f(two narrow slit)]

 Fourier transform of a convolution is the product of the individual Fourier transforms
 Tf (two wide alit) = Tf (one wide alit). Tf (two warmous)

 $Tf(two wide slit) = Tf(one wide slit) \cdot Tf(two narrow slit)]$ - one wide slit

$$F(\sin\theta) = 2X_0 \frac{\sin(kX_0\sin\theta)}{kX_0\sin\theta}$$

- two narrow slit

 $F(\sin\theta) = 2\cos(kx_0\sin\theta)$





- two wide slit

$$F(\sin\theta) = 4X_0 \frac{\sin(kX_0\sin\theta)}{kX_0\sin\theta}\cos(kx_0\sin\theta)$$

-intensity





- two wide slit
 - form of diffraction pattern is a series of \cos^2 fringes modulated by $(\sin \alpha / \alpha)^2$
- \cos^2 function- first zero $kx_0 \sin \theta_1 = \pi/2$ $\sin \theta_1 = \pi/(2kx_0)$
- $(\sin \alpha / \alpha)^2$ function- first zero $kX_0 \sin \theta_2 = \pi$ $\sin \theta_2 = \pi / (kX_0)$ - $X_0 \gg X_0 \rightarrow \theta_2 > \theta_1$





- three wide slit
- amplitude function

 $f(three \ wide \ slit) = f(one \ wide \ slit) * f(three \ narrow \ slits)$ $Tf(three \ wide \ slit) = T[f(one \ wide \ slit) * f(three \ narrow \ slits)]$ $Tf(three \ wide \ slit) = Tf(one \ wide \ slit) * Tf(three \ narrow \ slits)$

$$-F(\sin\theta) = 2X_0 \frac{\sin(kX_0\sin\theta)}{kX_0\sin\theta} [1 + 2\cos(kx_0\sin\theta)]$$



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- three wide slits
- intensity

$$\left|F(\sin\theta)\right|^2 = 4X_0^2 \frac{\sin^2(kX_0\sin\theta)}{(kX_0\sin\theta)^2} [1 + 2\cos(kx_0\sin\theta)]^2$$



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- N wide slits

width of $2X_0$, centered on a δ function

- amplitude function

 $f(N \text{ wide slit}) = f(\text{one wide slit})^* f(N \text{ narrow slits})$ $Tf(N \text{ wide slit}) = T[f(\text{one wide slit})^* f(N \text{ narrow slits})]$ $Tf(N \text{ wide slit}) = Tf(\text{one wide slit})^* Tf(N \text{ narrow slits})$





- N wide slits
- intensity

$$\left|F(\sin\theta)\right|^{2} = 4X_{0}^{2} \frac{\sin^{2}(kX_{0}\sin\theta)}{\left(kX_{0}\sin\theta\right)^{2}} \frac{\sin^{2}\frac{Nkx_{0}\sin\theta}{2}}{\sin^{2}\frac{2}{1}}$$





Experiments

















Slit spacing s and slit width w





Diffraction patterns from gratings (a) and (b).















The zebra





- infinite wide slits
- amplitude function

 $f(\infty wide \ slit) = f(one \ wide \ slit)^* f(\infty \ narrow \ slits)$ $Tf(\infty \ wide \ slit) = T[f(one \ wide \ slit)^* f(\infty \ narrow \ slits)]$ $Tf(\infty \ wide \ slit) = Tf(one \ wide \ slit)^* Tf(\infty \ narrow \ slits)$



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- infinite number of wide slits
- intensity

$$\left|F(\sin\theta)\right|^{2} = 4X_{0}^{2} \frac{\sin^{2}(kX_{0}\sin\theta)}{(kX_{0}\sin\theta)^{2}} \left|\sum_{n=-\infty}^{n=\infty} \delta(\sin\theta - \frac{2n\pi}{kx_{0}})\right|^{2}$$



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- narrow slits
 - As the number of slits increases, the main peaks become sharper and narrower, whilest the subsidiary peaks become rapidly less intense
 - 2. The position and separation of the main peaks is constant independent of the number of peaks. The man peaks are separated by an angular deflection given by

$$\Delta(\sin\theta) = \frac{2\pi}{kx_0}$$

The distance is determined soly by λ and separation x_0 of any two neighboring narrow slits.





- narrow slits
- The position of the main peak in a diffraction pattern is determined solely by the lattice spacing in an obstacle.
- The shape of the main peak is determined by the overall shape of the obstacle.







- wide slits
- The effect of the motif (one wide slit) is to alter the intensity of each main peak, but the position of the main peaks are unchanged
- intensity envolope
 - \rightarrow structure of motif





- another way of looking at N wide slits
- shape function

is zero everywhere outside an obstacle corresponds to the macroscopic shape of the obstacle within the obstacle







- finite lattice
 - $f(finite \ lattice) = f(infinite \ lattice) \cdot f(shape \ function)$
 - $f(obstacle) = f(motif) * f(finite \ lattice)$
 - $f(obstacle) = f(motif) * [f(infinite \ lattice) \cdot f(shape \ function)]$
- diffraction pattern
 - $F(\sin\theta) = Tf(obstacle)$
 - $=T{f(motif)*[f(infinite \ lattice)•f(shape \ function)]}$
 - $= Tf(motif) \bullet T[f(infinite \ lattice) \bullet f(shape \ function)]$
 - = *Tf* (*motif*)•[*Tf* (*infinite lattice*) **Tf* (*shape function*)]





- N wide slits
 - f(N wide slits) = f(one wide slit) * f(N narrow slits)
 - f(N wide slits) = f(one wide slit)*

[$f(\infty \text{ narrow slits}) \bullet f(\text{shape function})$]





- diffraction pattern

 $Tf(N wide slits) = T\{f(one wide slit)*$

 $[f(\infty \text{ narrow slits}) \bullet f(\text{shape function})]\}$ = Tf(one wide slit) \u03c4 T[f(\u03c5 narrow slits) \u03c4 f(\text{shape function})] = Tf(one wide slit) \u03c4 Tf(\u03c5 narrow slits) \u03c4 Tf(\text{shape function})

- three types of structural information
 - that concerning the lattice
 - that concerning the motif
 - that concerning the shape of the entire crystal







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Summary



- the position of the main peaks gives information on the lattice
- the shape of each main peak gives information on the overall object shape
- the set of intensities of all the main peaks gives information on the structure of the motif

