



Chapter 16.

MOS FUNDAMENTALS

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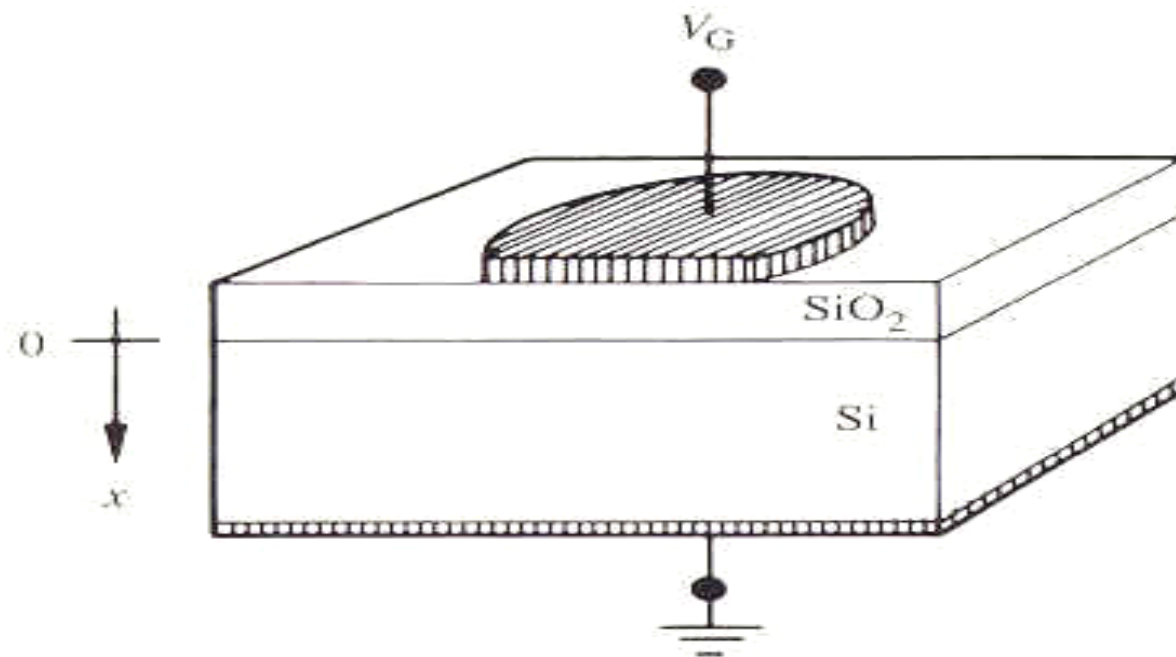
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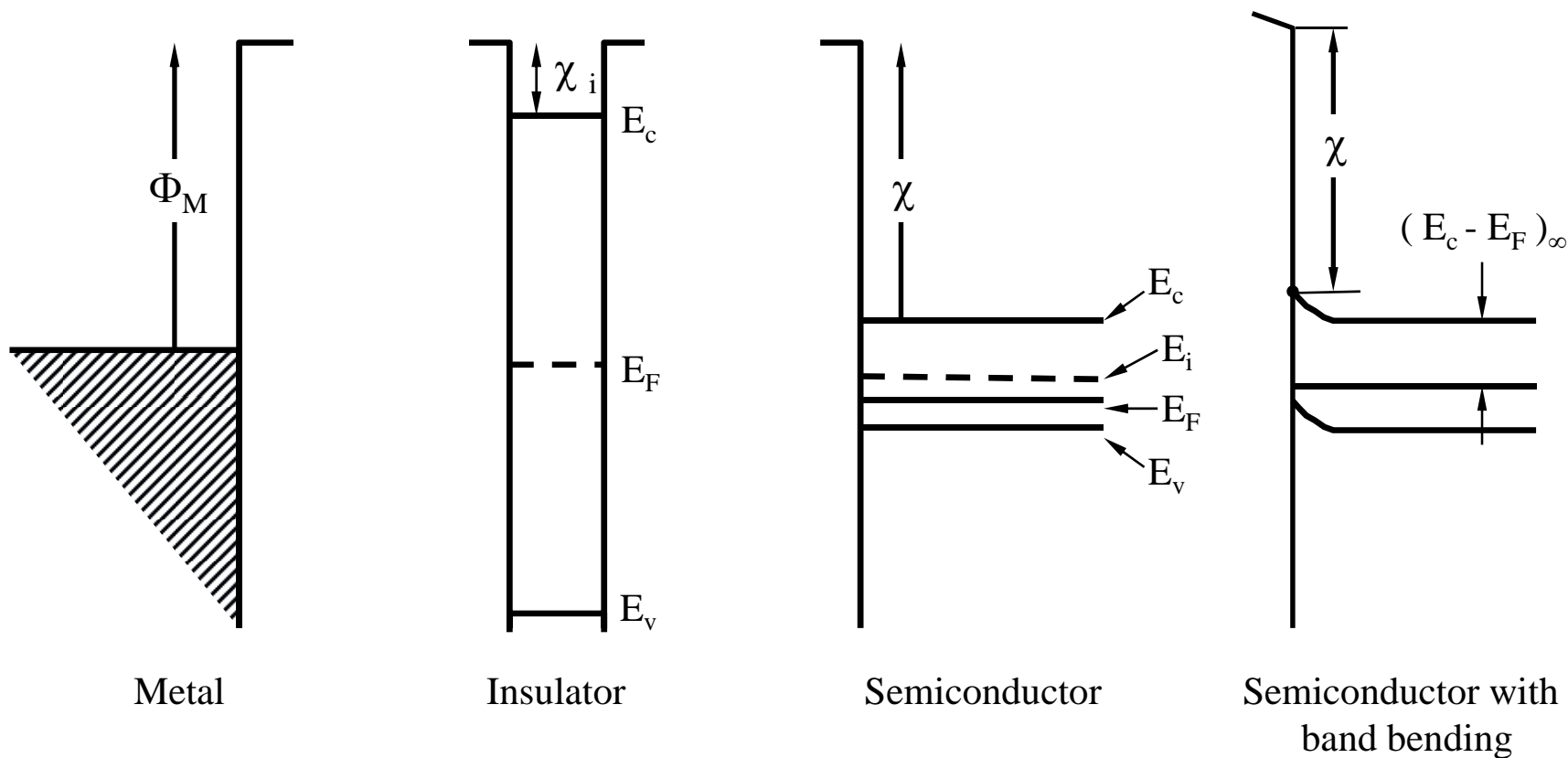
MOS Structure

Metal - Oxide (SiO_2) - Semiconductor (Si)



- The most common field plate (gate) materials are heavily doped polycrystalline silicon.
- The silicon-side terminal is called the back or substrate contact.
- The more general designation: metal-insulator-semiconductor (MIS)

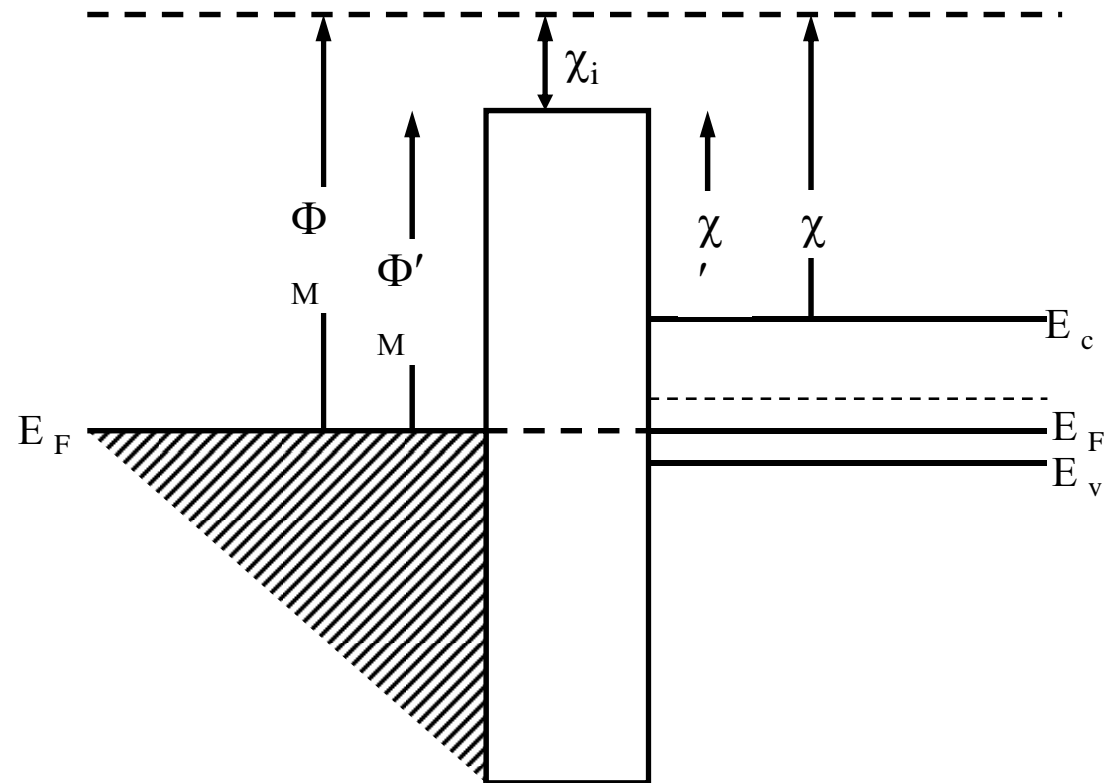
Individual energy band diagrams for the metal, insulator, and semiconductor components.



Ideal Structure Assumption

- (1) The metallic gate is sufficiently thick so that it can be considered an equipotential region.
- (2) The oxide is a perfect insulator.
- (3) No charge centers located in the oxide or at the interface.
- (4) Uniformly doped.

- (5) The semiconductor is sufficiently thick so that a field-free region(“bulk”) is encountered before reaching the back contact.
- (6) An ohmic contact between the semiconductor and the metal on the back side.
- (7) One-dimensional structure.
- (8) No work function difference between metal and semiconductor.



Energy band diagram of an ideal MOS structure with no bias.

Effect Of An Applied Bias - Qualitative description

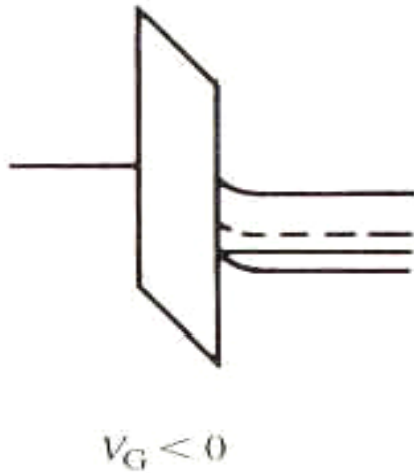
- General observations
 - $E_F(\textit{metal}) - E_F(\textit{semiconductor}) = -qV_G$
 - With $V_G \neq 0$, semiconductor Fermi energy is unaffected by the bias and remains invariant as a function of position because of the assumed zero current flow.



- Since the barrier heights are fixed quantities, the movement of the metal Fermi level leads to a band bending.
 - In the metal, no bend-bending.
 - In the oxide and semiconductor,
 - an upward slope when $V_G > 0$
 - a downward slope when $V_G < 0$.
 - With no oxide charges, the Poisson's equation yields a constant slope in the oxide.
 - Band bending in the semiconductor is somewhat more complex.

- Specific biasing regions

- Accumulation ($V_G < 0$)



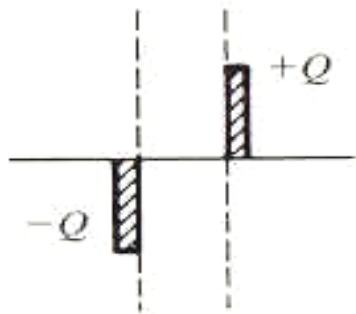
- $V_G < 0$ raises $E_F(\text{metal})$ and the hole concentration inside the semiconductor,

$$p = n_i \exp\left[-(E_F - E_i)/kT\right]$$

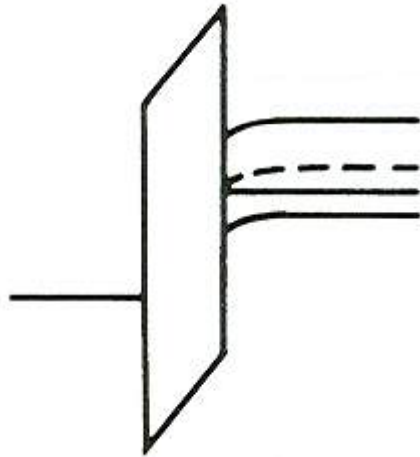
increases as one approaches the oxide-semiconductor interface.

- $V_G < 0$ places negative charges on the gate.

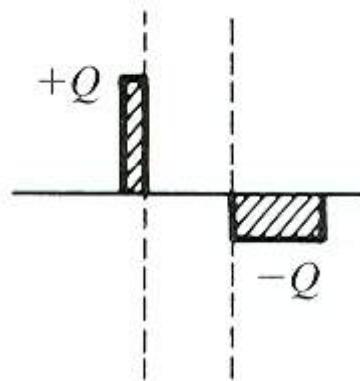
To maintain a balance of charge, positively charged holes must be drawn toward the Si-SiO₂ interface.



– Depletion ($0 < V_G < V_T$)



$$V_T > V_G > 0$$

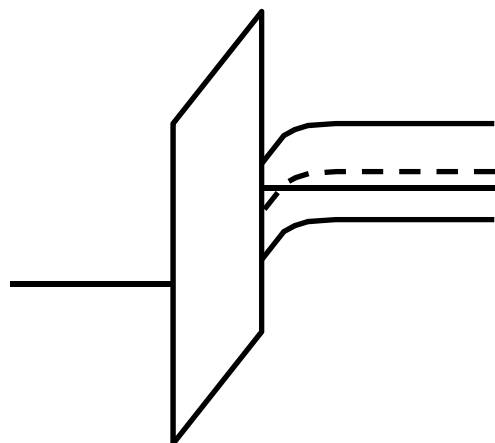


Depletion

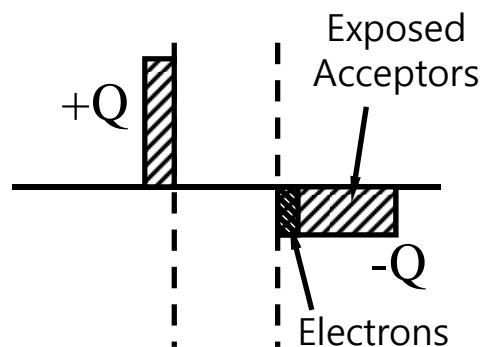
- $V_G > 0$ slightly lowers $E_F(\text{metal})$ and the hole concentration is decreased (depleted) in the vicinity of the Si-SiO₂ interface.
- $V_G > 0$ places positive charges on the gate, which in turn repels holes from the interface and exposes the negatively charged acceptor sites.



– Onset of Inversion ($V_G = V_T$)



$$V_G = V_T$$



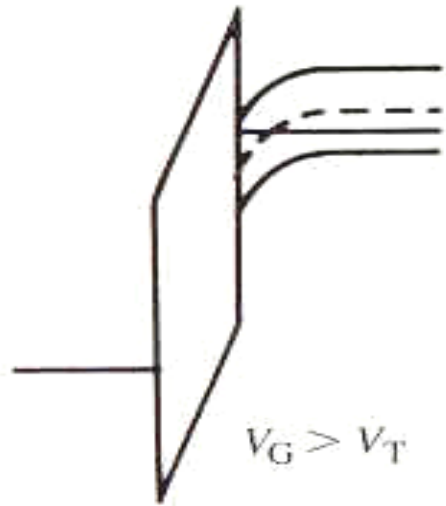
- As V_G is increased positively, the bands at the Si surface will bend down more and the electron concentration at the surface (n_s) will increase from less than n_i when E_i (surface) $>$ E_F , to n_i when E_i (surface) $=$ E_F , to greater than n_i when E_i (surface) $<$ E_F .
- At $V_G = V_T$,

$$E_i(\text{bulk}) - E_i(\text{surface}) = 2[E_i(\text{bulk}) - E_F]$$

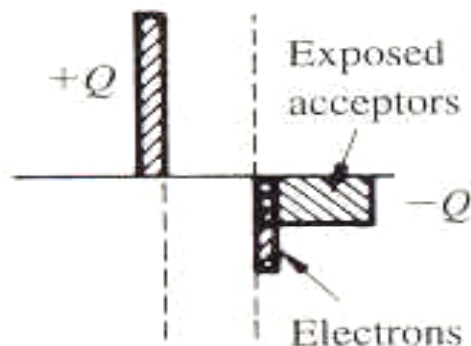
$$n_s = n_i \exp\left[\frac{E_F - E_i(\text{surface})}{kT}\right]$$

$$= n_i \exp\left[\frac{E_i(\text{bulk}) - E_F}{kT}\right] = p_{\text{bulk}} = N_A$$

- Inversion ($V_G > V_T$)

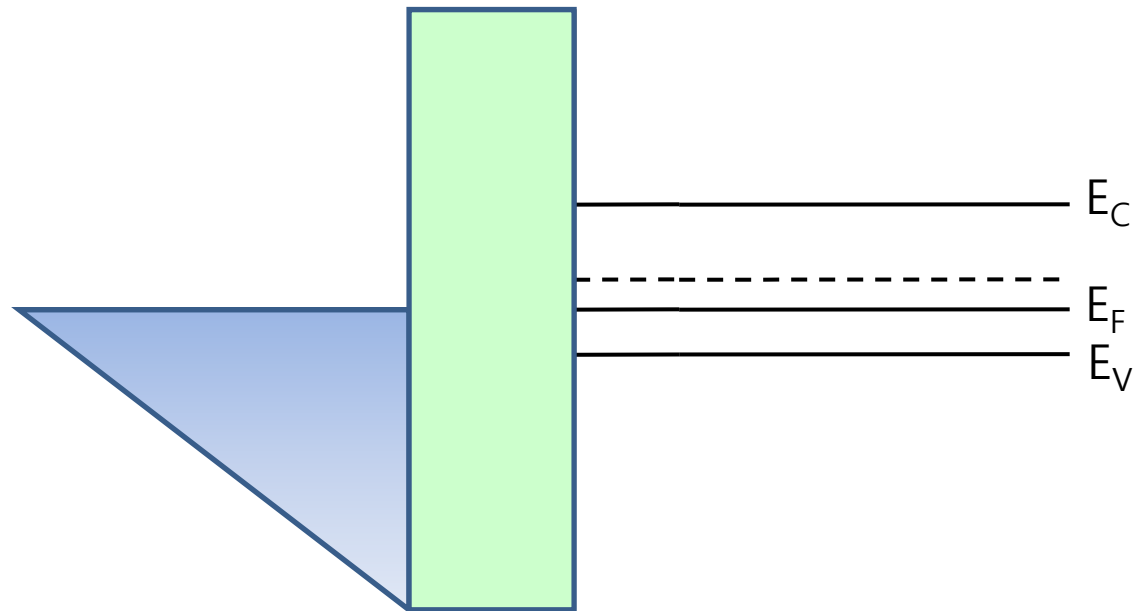


- For $V_G > V_T$, $n_s > N_A$.
- The surface region :p-type n-type
- In inversion, the depletion width changes little since



$$n_s \propto \exp\left[\frac{E_F - E_i(\text{surface})}{kT}\right]$$

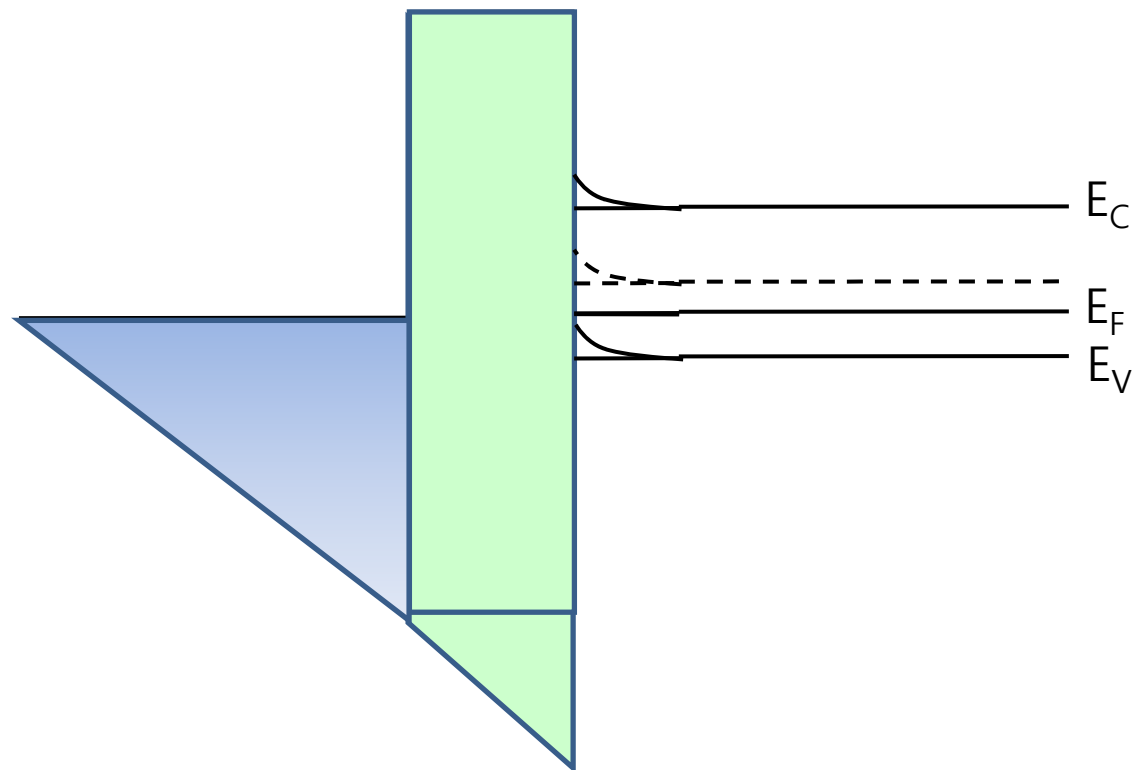
No Bias ($V_G = 0$)



Metal

Insulator

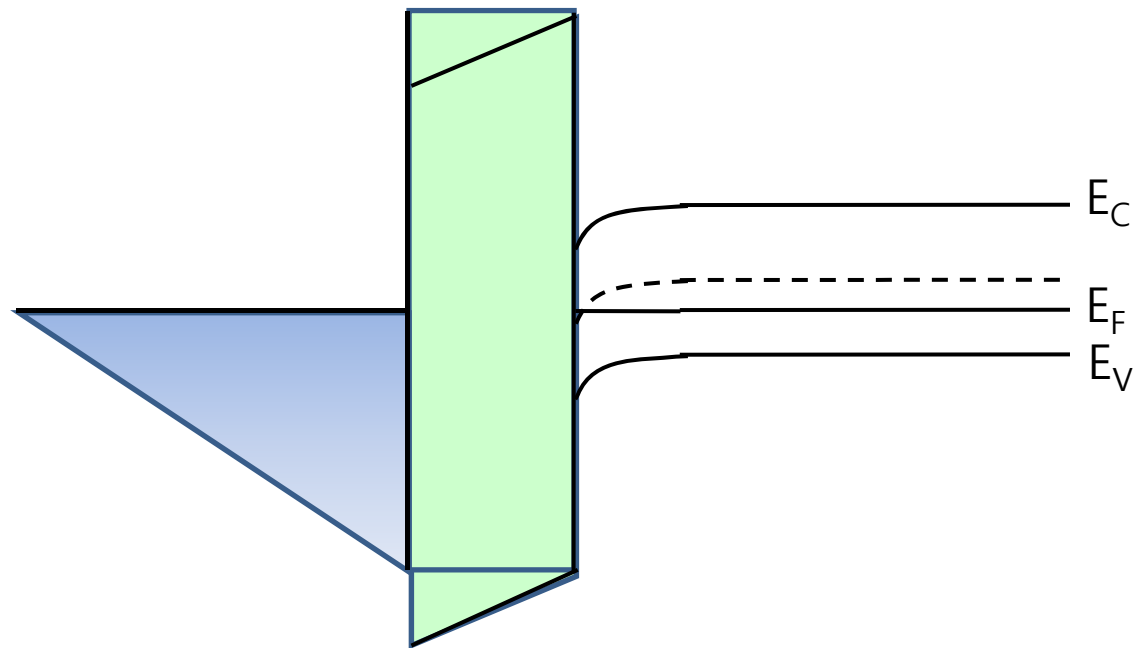
Si

Accumulation ($V_G < 0$)

Metal

Insulator

Si

Depletion ($0 < V_G < V_T$)

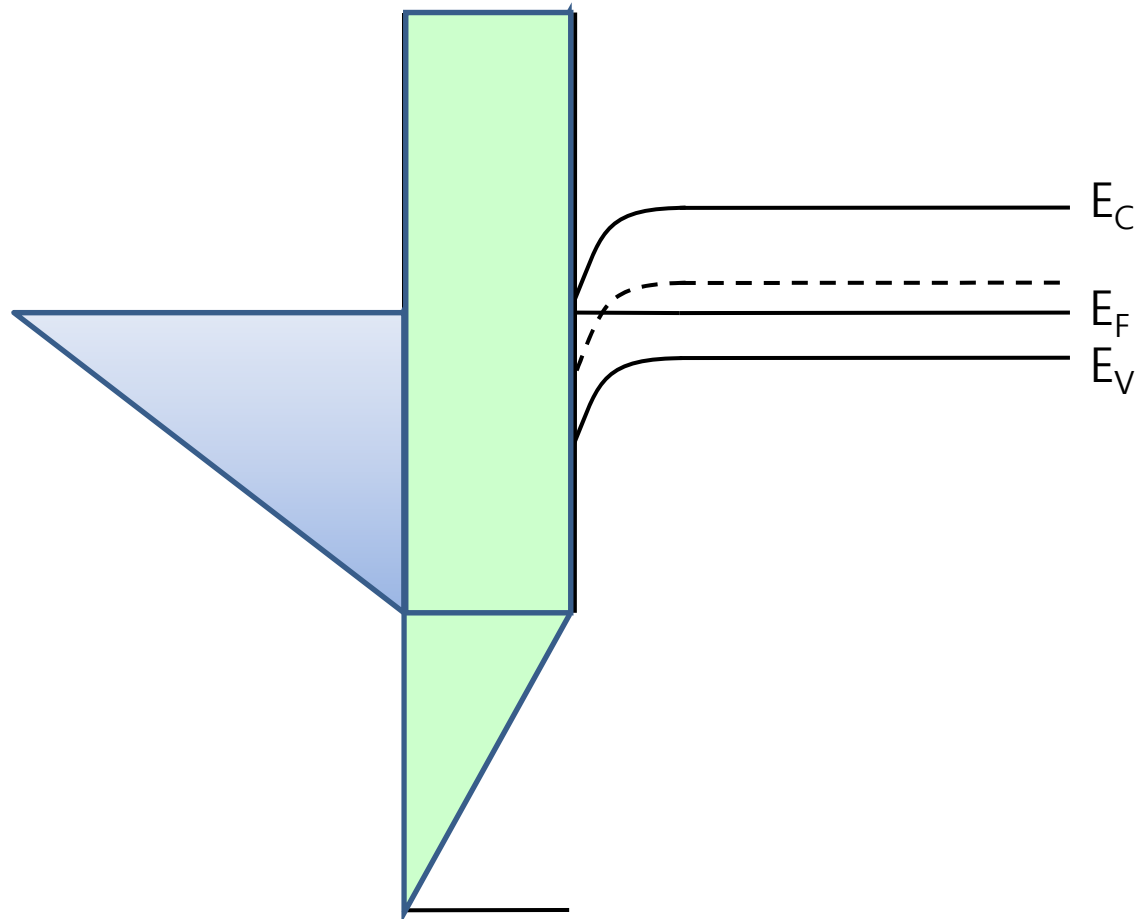
Metal

Insulator

Si



Inversion ($V_G > V_T$)



Metal

Insulator

Si

Effect Of An Applied Bias - Quantitative formulation

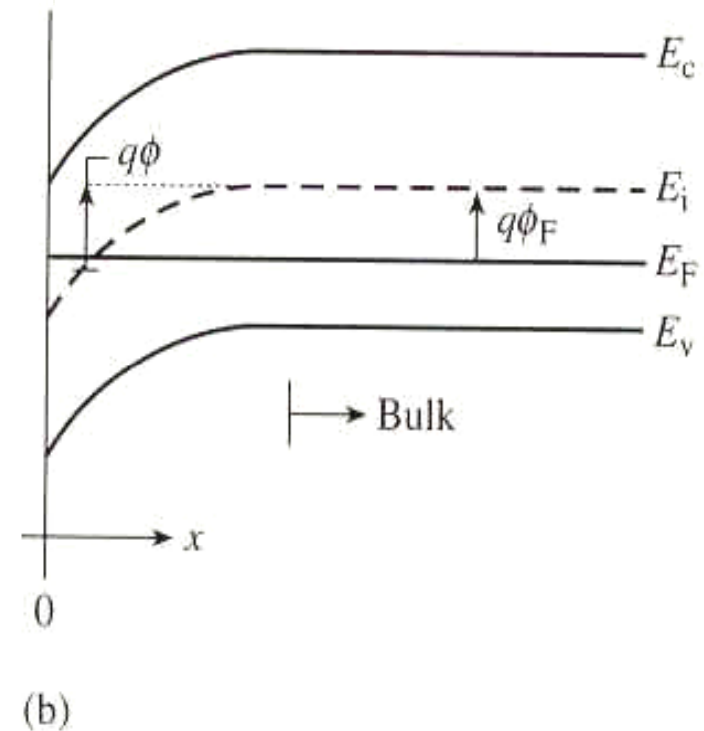
- Preparatory considerations

$$\phi(x) = \frac{1}{q} [E_i(\text{bulk}) - E_i(x)]$$

ϕ_S : the surface potential.

$$\phi_S = \frac{1}{q} [E_i(\text{bulk}) - E_i(\text{surface})]$$

$$\begin{aligned} \phi_F &= \frac{1}{q} [E_i(\text{bulk}) - E_F] \\ &= \frac{kT}{q} \ln(N_A/n_i) \end{aligned}$$



For an p-type semiconductor,

- Accumulation : $\phi_s < 0$
- Flat band : $\phi_s = 0$
- Depletion : $0 < \phi_s < 2\phi_F$
- Onset of inversion : $\phi_s = 2\phi_F$
- Inversion : $\phi_s > 2\phi_F$



- **Delta-depletion solution**

Delta-depletion assumption :

- The functional form of the accumulation charge & the inversion charge : δ - function.
- Because the depletion width increases only slightly once the semiconductor inverts, it is assumed the δ - function of charge added in inversion precisely balances the charge added to the gate.
- The actual depletion charge is replaced with a squared-off distribution



- Accumulation :
 - Because of the assumed δ - function, the electric field and potential are zero for all $x > 0$.
- Depletion :

For $0 \leq x \leq W$,

$$\frac{dE}{dx} = \frac{\rho}{K_S \epsilon_0} \cong -\frac{qN_A}{K_S \epsilon_0}$$

$$E(x) = \frac{qN_A}{K_S \epsilon_0} (W - x)$$

$$\phi(x) = \frac{qN_A}{2K_S \epsilon_0} (W - x)^2$$

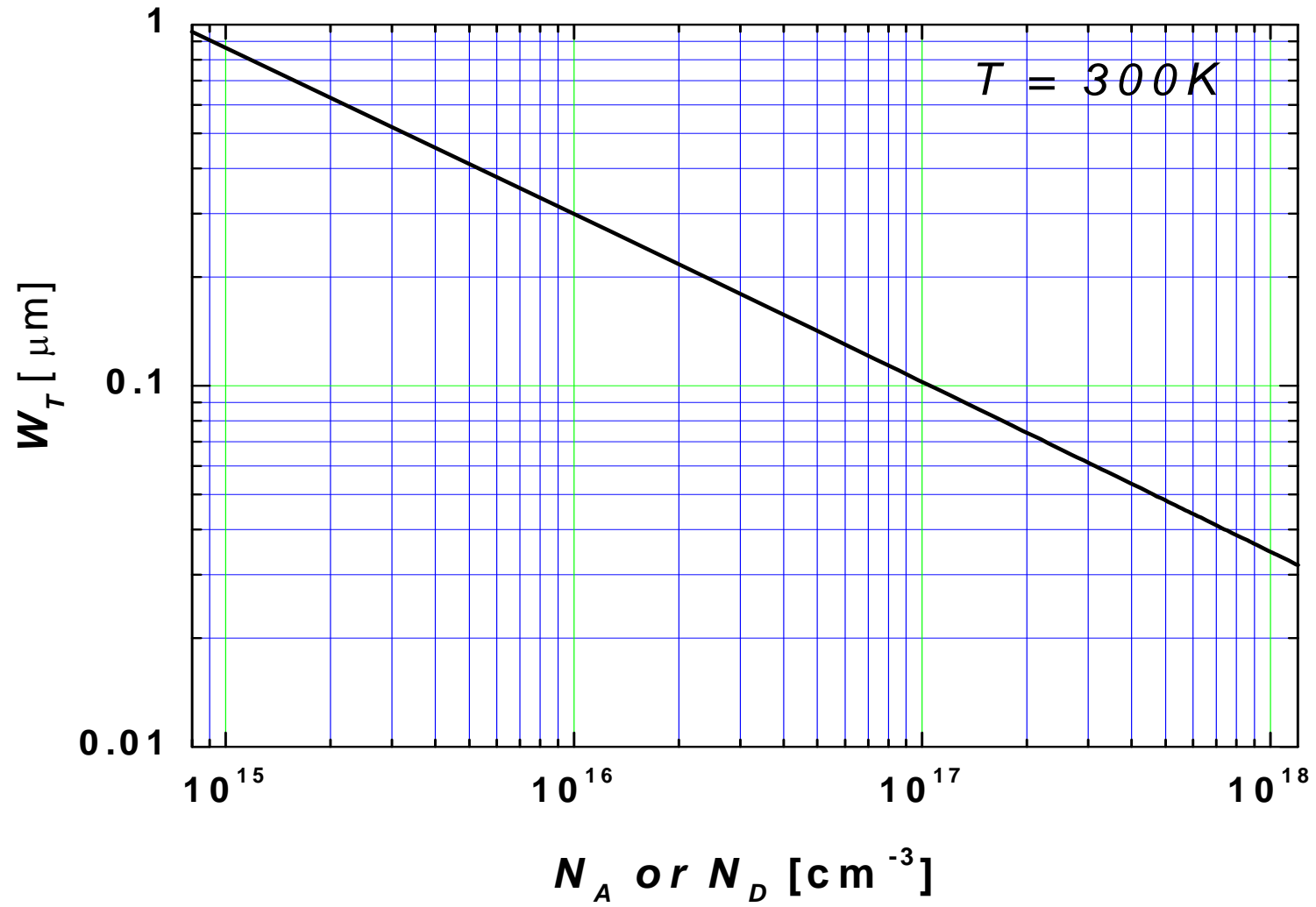
with $\phi = \phi_s$ at $x = 0$,

$$W = \left[\frac{2K_S \epsilon_0}{qN_A} \phi_s \right]^{1/2}$$

The maximum depletion width,

$$W_T = \left[\frac{2K_S \epsilon_0}{qN_A} (2\phi_F) \right]^{1/2}$$





– Inversion :

- The solution is established by merely adding a δ -function of surface charge to the solution existing at the end of depletion.
- The depletion charge, the $x > 0$ electric field, and the $x > 0$ potential remain fixed at their values.

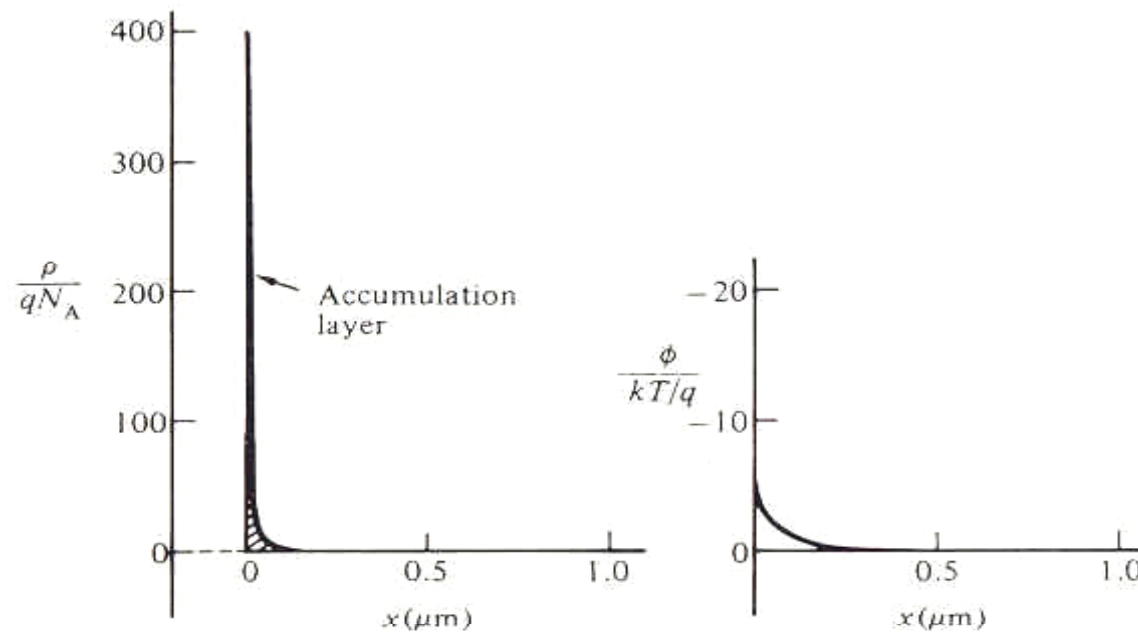
$$\phi_s = 2\phi_F$$



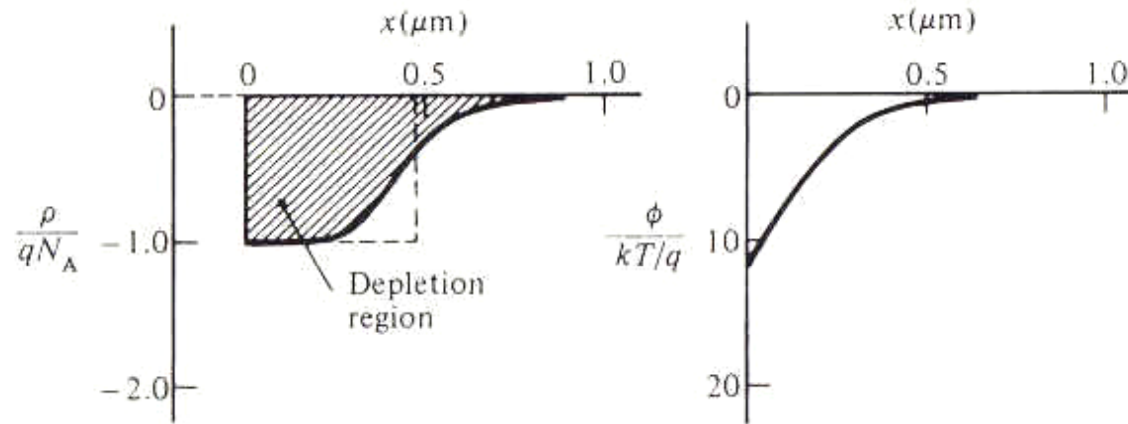
Exact solution for the charge density and potential
assuming

$$\phi_F = 12kT/q \text{ and } T = 300K$$

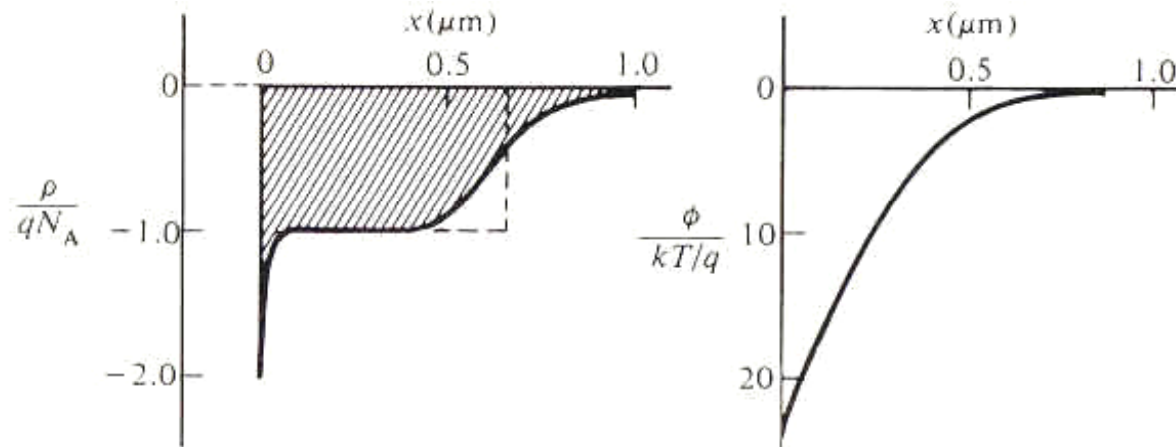
(a) Accumulation ($f_s = -6kT/q$)



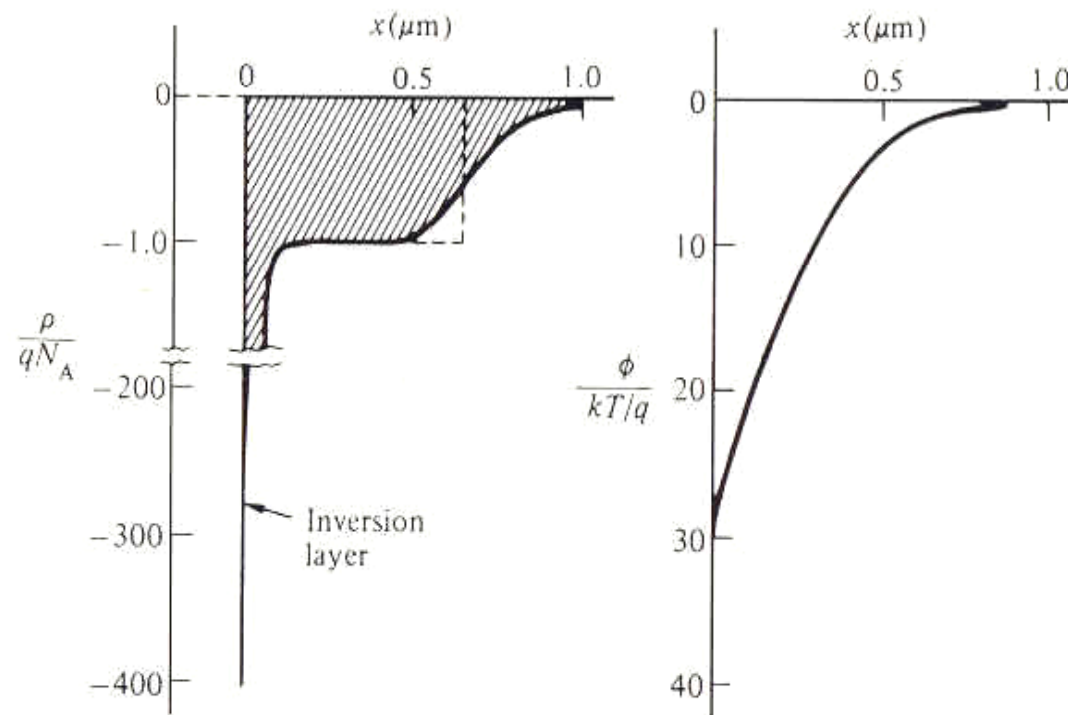
(b) Middle of depletion ($f_s = f_F = 12kT/q$)



(c) Onset of inversion ($f_s = 2f_F = 24kT/q$)



(d) Deep into inversion



- Gate voltage relationship (delta-depletion solution)

$$V_G = \Delta\phi_{\text{Semi}} + \Delta\phi_{\text{OX}}$$

Because ϵ_{OX} is constant in an ideal oxide with no charges,

$$\Delta\phi_{\text{OX}} = x_o \epsilon_{\text{OX}}$$

Since there is no charges at the interface, (in the depletion region)

$$D_{\text{ox}} = D_{\text{Semi}} \Big|_{x=0}$$

$$E_{\text{ox}} = \frac{K_s}{K_o} E_s$$

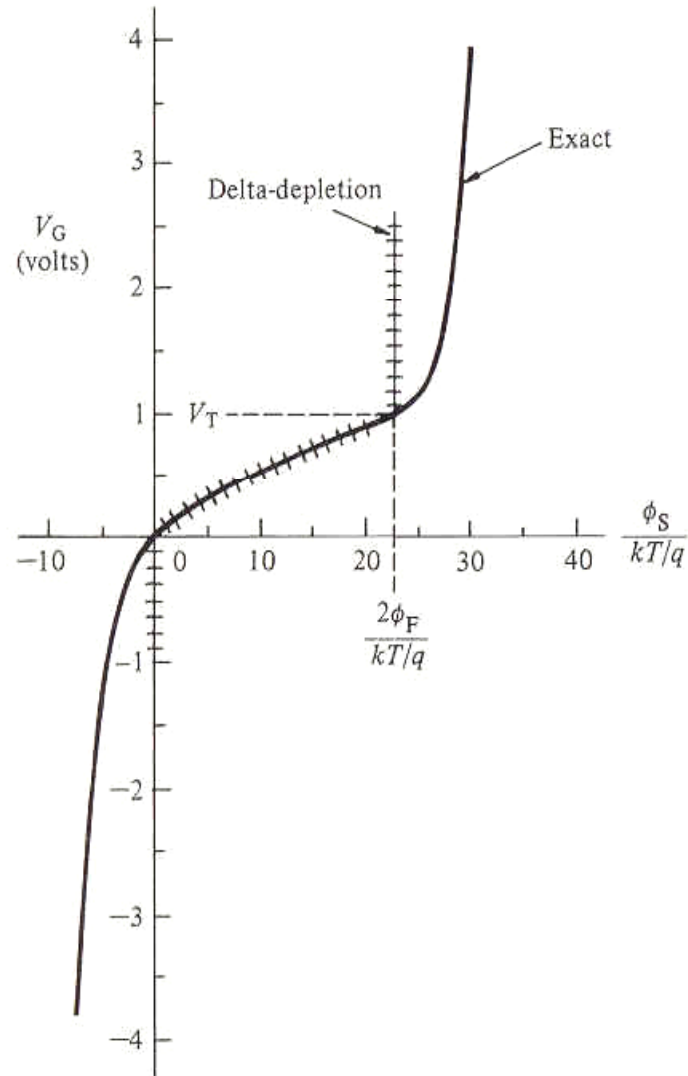
$$V_G = \phi_S + \frac{K_s}{K_o} x_o E_s = \phi_S + \frac{K_s}{K_o} x_o \sqrt{\frac{2qN_A}{K_s \epsilon_0}} \phi_S \quad (0 \leq \phi_S \leq 2\phi_F)$$



At threshold, $\phi_s = 2\phi_F$

$$V_T = 2\phi_F + \frac{K_s x_o}{K_o} \sqrt{\frac{4qN_A}{K_s \epsilon_0} \phi_F}$$





ϕ_s is a rather rapidly varying function of V_G when the device is in depletion. However, when it is accumulated or inverted, it takes a large change in V_G to produce a small change in ϕ_s .

; +++++ delta-depletion solution, ——— exact solution.

