

Ch. 5. EM Optics

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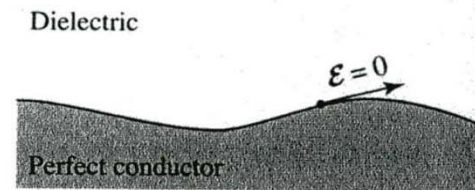
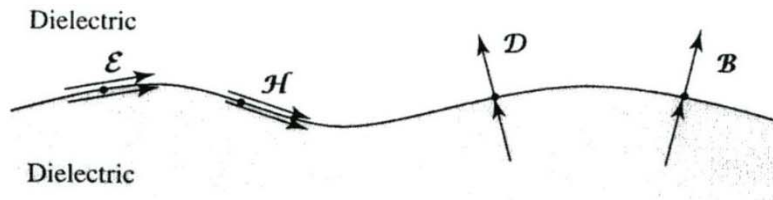


Maxwell's equations

In source-free medium

$$\begin{aligned}\nabla \times \mathcal{H} &= \frac{\partial \mathcal{D}}{\partial t} & \mathcal{D} &= \epsilon_0 \mathcal{E} + \mathcal{P} \\ \nabla \times \mathcal{E} &= -\frac{\partial \mathcal{B}}{\partial t} & \mathcal{B} &= \mu_0 \mathcal{H} + \mu_0 \mathcal{M} \\ \nabla \cdot \mathcal{D} &= 0 \\ \nabla \cdot \mathcal{B} &= 0.\end{aligned}$$

Boundary conditions





Poynting vector and Poynting theorem

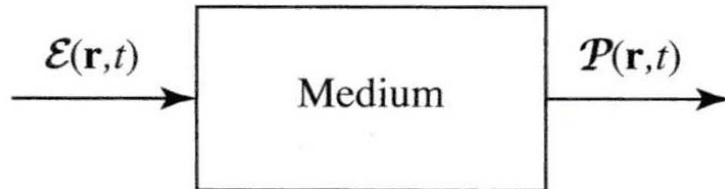
$$\mathbf{S} = \mathbf{E} \times \mathbf{H}$$

$$\nabla \cdot \mathbf{S} = -\frac{\partial}{\partial t} \left(\frac{1}{2} \epsilon_0 \mathbf{E}^2 + \frac{1}{2} \mu_0 \mathbf{H}^2 \right) + \mathbf{E} \cdot \frac{\partial \mathcal{P}}{\partial t} + \mu_0 \mathbf{H} \cdot \frac{\partial \mathcal{M}}{\partial t}$$





EM waves in media



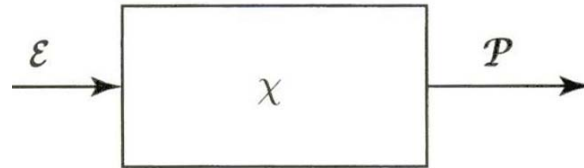
$$\mathcal{P} = \epsilon_0 \chi \mathcal{E}$$

Definitions

- A dielectric medium is said to be *linear* if the vector field $\mathcal{P}(\mathbf{r}, t)$ is linearly related to the vector field $\mathcal{E}(\mathbf{r}, t)$. The principle of superposition then applies.
- The medium is said to be *nondispersive* if its response is instantaneous, i.e., if \mathcal{P} at time t is determined by \mathcal{E} at the same time t and not by prior values of \mathcal{E} . Nondispersiveness is clearly an idealization since all physical systems, no matter how rapidly they may respond, do have a response time that is finite.
- The medium is said to be *homogeneous* if the relation between \mathcal{P} and \mathcal{E} is independent of the position \mathbf{r} .
- The medium is said to be *isotropic* if the relation between the vectors \mathcal{P} and \mathcal{E} is independent of the direction of the vector \mathcal{E} , so that the medium exhibits the same behavior from all directions. The vectors \mathcal{P} and \mathcal{E} must then be parallel.
- The medium is said to be *spatially nondispersive* if the relation between \mathcal{P} and \mathcal{E} is local, i.e., if \mathcal{P} at each position \mathbf{r} is influenced only by \mathcal{E} at the same position \mathbf{r} . The medium is assumed to be spatially nondispersive throughout this chapter (optically active media, considered in Sec. 6.4A, are spatially dispersive).



Linear, nondispersive, homogeneous and isotropic medium



$$\mathcal{P} = \epsilon_o \chi \mathcal{E}$$

$$\mathcal{D} = \epsilon \mathcal{E}$$

$$\epsilon = \epsilon_o (1 + \chi)$$

$$\mathcal{B} = \mu \mathcal{H}$$

$$\nabla \times \mathcal{H} = \epsilon \frac{\partial \mathcal{E}}{\partial t}$$

$$\nabla \times \mathcal{E} = -\mu \frac{\partial \mathcal{H}}{\partial t}$$

$$\nabla \cdot \mathcal{E} = 0$$

$$\nabla \cdot \mathcal{H} = 0.$$

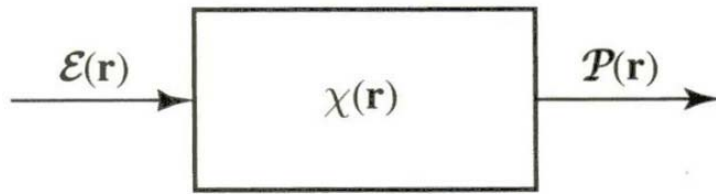
$$\nabla^2 u - \frac{1}{c^2} \frac{\partial^2 u}{\partial t^2} = 0$$

$$c = \frac{1}{\sqrt{\epsilon \mu}}$$

$$n = \sqrt{\frac{\epsilon}{\epsilon_o}} = \sqrt{1 + \chi}$$



Inhomogeneous media



$$\frac{\epsilon_o}{\epsilon} \nabla \times (\nabla \times \mathcal{E}) = -\frac{1}{c_o^2} \frac{\partial^2 \mathcal{E}}{\partial t^2}$$

$$\nabla \times \left(\frac{\epsilon_o}{\epsilon} \nabla \times \mathcal{H} \right) = -\frac{1}{c_o^2} \frac{\partial^2 \mathcal{H}}{\partial t^2}$$

$$\nabla^2 \mathcal{E} + \nabla \left(\frac{1}{\epsilon} \nabla \epsilon \cdot \mathcal{E} \right) - \mu_o \epsilon \frac{\partial^2 \mathcal{E}}{\partial t^2} = 0$$

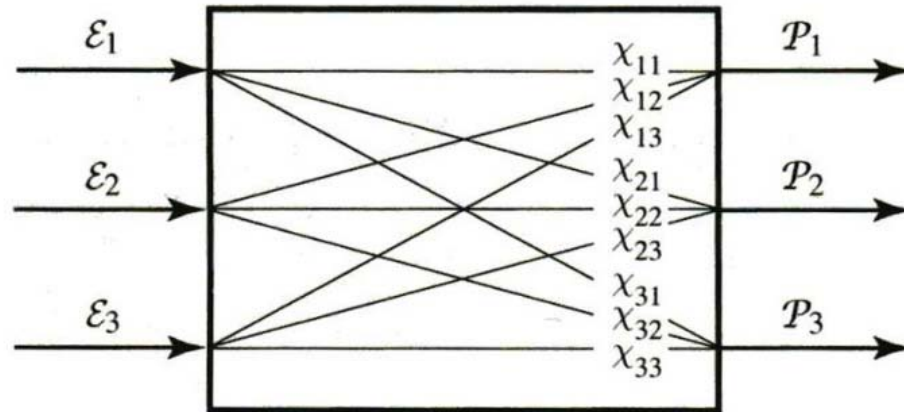
$$\nabla^2 \mathcal{E} - \frac{1}{c^2(\mathbf{r})} \frac{\partial^2 \mathcal{E}}{\partial t^2} \approx 0$$



Anisotropic media

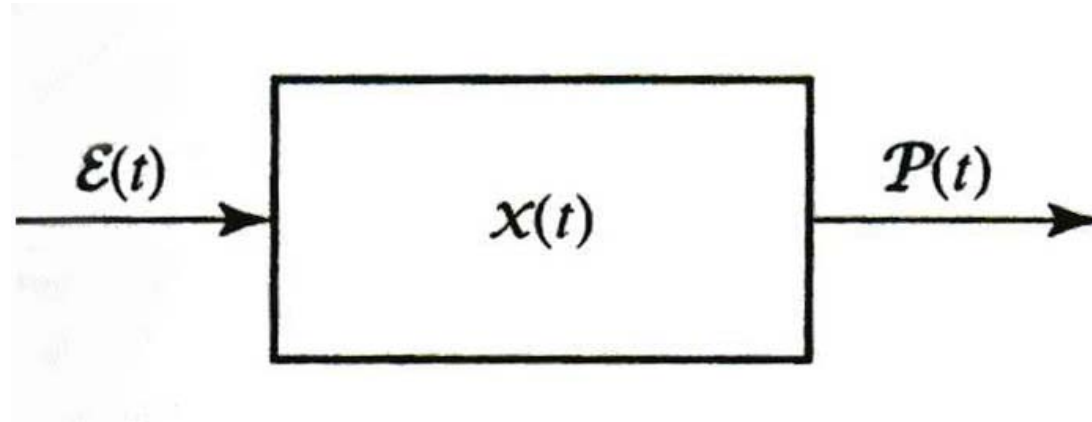
$$\mathcal{P}_i = \sum_j \epsilon_o \chi_{ij} \mathcal{E}_j$$

$$\mathcal{D}_i = \sum_j \epsilon_{ij} \mathcal{E}_j$$





Dispersive media



$$\mathcal{P}(t) = \epsilon_0 \int_{-\infty}^{\infty} \chi(t - t') \mathcal{E}(t') dt'$$



Nonlinear media

$$\nabla^2 \mathcal{E} - \frac{1}{c_0^2} \frac{\partial^2 \mathcal{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \mathcal{P}}{\partial t^2}$$

$$P = \Psi(E)$$

$$\nabla^2 \mathcal{E} - \frac{1}{c_0^2} \frac{\partial^2 \mathcal{E}}{\partial t^2} = \mu_0 \frac{\partial^2 \Psi(\mathcal{E})}{\partial t^2}$$





Monochromatic EM waves

$$\mathcal{E}(\mathbf{r}, t) = \text{Re}\{\mathbf{E}(\mathbf{r}) \exp(j\omega t)\}$$

$$\mathcal{H}(\mathbf{r}, t) = \text{Re}\{\mathbf{H}(\mathbf{r}) \exp(j\omega t)\}$$

$$\nabla \times \mathbf{H} = j\omega \mathbf{D}$$

$$\nabla \times \mathbf{E} = -j\omega \mathbf{B}$$

$$\nabla \cdot \mathbf{D} = 0$$

$$\nabla \cdot \mathbf{B} = 0$$

$$\mathbf{D} = \epsilon_0 \mathbf{E} + \mathbf{P}$$

$$\mathbf{B} = \mu_0 \mathbf{H} + \mu_0 \mathbf{M}$$

$$\mathbf{S} = \frac{1}{2} \mathbf{E} \times \mathbf{H}^*$$

$$\mathbf{D} = \epsilon \mathbf{E} \quad \text{and} \quad \mathbf{B} = \mu \mathbf{H}$$

$$\nabla \times \mathbf{H} = j\omega \epsilon \mathbf{E}$$

$$\nabla \times \mathbf{E} = -j\omega \mu \mathbf{H}$$

$$\nabla \cdot \mathbf{E} = 0$$

$$\nabla \cdot \mathbf{H} = 0$$

$$\nabla^2 U + k^2 U = 0$$

$$k = nk_0 = \omega \sqrt{\epsilon \mu}$$





Dispersive media

$$\mathbf{P} = \epsilon_o \chi(\nu) \mathbf{E}$$

$$\chi(\nu) = \int_{-\infty}^{\infty} x(t) \exp(-j2\pi\nu t) dt$$

$$\mathbf{D} = \epsilon(\nu) \mathbf{E}$$

$$\epsilon(\nu) = \epsilon_o [1 + \chi(\nu)]$$

$$k = \omega \sqrt{\epsilon(\nu) \mu_o}$$





TEM wave

$$\mathbf{H}(\mathbf{r}) = \mathbf{H}_0 \exp(-j\mathbf{k} \cdot \mathbf{r})$$

$$\mathbf{E}(\mathbf{r}) = \mathbf{E}_0 \exp(-j\mathbf{k} \cdot \mathbf{r})$$

$$\mathbf{k} \times \mathbf{H}_0 = -\omega \epsilon \mathbf{E}_0$$

$$\mathbf{k} \times \mathbf{E}_0 = \omega \mu \mathbf{H}_0$$

$$\eta = \frac{E_0}{H_0} = \sqrt{\frac{\mu}{\epsilon}}$$

$$\eta = \frac{\eta_0}{n}$$

$$I = \frac{|E_0|^2}{2\eta}$$

$$W = \frac{1}{2}\epsilon|E_0|^2$$

$$I = cW$$

