

# Quantum Theory of Radiation (Sakurai, Advanced Quantum Mechanics, Ch. 2)

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# Gauge Transformation

$$\mathbf{E} = -\nabla \varphi - \frac{1}{c} \frac{\partial \mathbf{A}}{\partial t} , \quad \mathbf{B} = \nabla \times \mathbf{A}$$

$$\varphi \rightarrow \varphi - \frac{1}{c} \frac{\partial \Lambda}{\partial t} , \quad \mathbf{A} \rightarrow \mathbf{A} + \nabla \Lambda$$



# Hamiltonian

- Hamiltonian for a particle of electric charge  $e$  ( $e<0$  for the electron) subjected to the electromagnetic field.

$$H = \frac{1}{2m} \left( p - \frac{eA}{c} \right)^2 + e\varphi$$

$$\left( p - \frac{eA}{c} \right)^2 \rightarrow p^2 - \left( \frac{e}{c} \right) (p \cdot A + A \cdot p) + \left( \frac{e}{c} \right)^2 A^2$$

$p$ : canonical momentum

$$\Pi \equiv m \frac{dx}{dt} = p - \frac{eA}{c}$$

kinematical (or mechanical) momentum



# Transversality Condition

$$\nabla \cdot \mathbf{A} = 0$$

$$\mathbf{A} = \mathbf{A}_\perp + \mathbf{A}_{\parallel}, \quad \nabla \cdot \mathbf{A}_\perp = 0$$

$$\nabla \times \mathbf{A}_{\parallel} = 0$$

$\mathbf{A}_\perp$  : Transverse component of  $\mathbf{A}$

$\mathbf{A}_{\parallel}$  : Longitudinal component of  $\mathbf{A}$

$$H = \sum_j \frac{1}{2m_j} \left[ \mathbf{p}^{(j)} - e_j \frac{\mathbf{A}_\perp(\mathbf{x}^{(j)})}{c} \right]^2 + \sum_{i>j} \frac{e_i e_j}{4\pi |\mathbf{x}^{(i)} - \mathbf{x}^{(j)}|} + H_{\text{rad}}$$



# Fourier Decomposition

$$\begin{aligned} A(\mathbf{x}, t) &= \frac{1}{\sqrt{V}} \sum_{\mathbf{k}} \sum_{\alpha} \left( c_{\mathbf{k}, \alpha}(t) \varepsilon^{(\alpha)} e^{i \mathbf{k} \cdot \mathbf{x}} + c_{\mathbf{k}, \alpha}^*(t) \varepsilon^{(\alpha)} e^{-i \mathbf{k} \cdot \mathbf{x}} \right) \\ &= \frac{1}{\sqrt{V}} \sum_{\mathbf{k}} \sum_{\alpha} \left( c_{\mathbf{k}, \alpha}(0) \varepsilon^{(\alpha)} e^{i k \cdot x} + c_{\mathbf{k}, \alpha}^*(0) \varepsilon^{(\alpha)} e^{-i k \cdot x} \right) \end{aligned}$$

$$c_{\mathbf{k}, \alpha}(t) = c_{\mathbf{k}, \alpha}(0) e^{-i \omega t} \quad c_{\mathbf{k}, \alpha}^*(t) = c_{\mathbf{k}, \alpha}^*(0) e^{i \omega t}$$

$$\omega = |\mathbf{k}| c$$

$$k \cdot x = \mathbf{k} \cdot \mathbf{x} - \omega t = \mathbf{k} \cdot \mathbf{x} - |\mathbf{k}| c t$$



# Hamiltonian of Field

$$H = \frac{1}{2} \int \left( |B|^2 + |E|^2 \right) d^3x = \frac{1}{2} \int \left[ |\nabla \times A|^2 + \left| (1/c)(\partial A / \partial t) \right|^2 \right] d^3x$$

$$H = \sum_{\mathbf{k}} \sum_{\alpha} 2(\omega/c)^2 c_{\mathbf{k},\alpha}^* c_{\mathbf{k},\alpha}$$

$$\ddot{c}_{\mathbf{k},\alpha} = -\omega^2 c_{\mathbf{k},\alpha}$$

$$Q_{\mathbf{k},\alpha} = \frac{1}{c} (c_{\mathbf{k},\alpha} + c_{\mathbf{k},\alpha}^*) , \quad P_{\mathbf{k},\alpha} = -\frac{i\omega}{c} (c_{\mathbf{k},\alpha} - c_{\mathbf{k},\alpha}^*)$$

$$H = \sum_{\mathbf{k}} \sum_{\alpha} 2 \left( \frac{\omega}{c} \right)^2 \left[ \frac{c(\omega Q_{\mathbf{k},\alpha} - iP_{\mathbf{k},\alpha})}{2\omega} \right] \left[ \frac{c(\omega Q_{\mathbf{k},\alpha} + iP_{\mathbf{k},\alpha})}{2\omega} \right] = \sum_{\mathbf{k}} \sum_{\alpha} \frac{1}{2} \left( P_{\mathbf{k},\alpha}^2 + \omega^2 Q_{\mathbf{k},\alpha}^2 \right)$$



# Quantization of Radiation Oscillators

$$[Q_{\mathbf{k},\alpha}, P_{\mathbf{k}',\alpha'}] = i\hbar \delta_{\mathbf{kk}'} \delta_{\alpha\alpha'}$$

$$[Q_{\mathbf{k},\alpha}, Q_{\mathbf{k}',\alpha'}] = 0$$

$$[P_{\mathbf{k},\alpha}, P_{\mathbf{k}',\alpha'}] = 0$$

$$a_{\mathbf{k},\alpha} = \left(1/\sqrt{2\hbar\omega}\right) (\omega Q_{\mathbf{k},\alpha} + iP_{\mathbf{k},\alpha})$$

$$a_{\mathbf{k},\alpha}^\dagger = \left(1/\sqrt{2\hbar\omega}\right) (\omega Q_{\mathbf{k},\alpha} - iP_{\mathbf{k},\alpha})$$

$$[a_{\mathbf{k},\alpha}, a_{\mathbf{k}',\alpha'}^\dagger] = -\frac{i}{2\hbar} [Q_{\mathbf{k},\alpha}, P_{\mathbf{k}',\alpha'}] + \frac{i}{2\hbar} [P_{\mathbf{k},\alpha}, Q_{\mathbf{k}',\alpha'}] = \delta_{\mathbf{kk}'} \delta_{\alpha\alpha'}$$

$$[a_{\mathbf{k},\alpha}, a_{\mathbf{k}',\alpha'}] = [a_{\mathbf{k},\alpha}^\dagger, a_{\mathbf{k}',\alpha'}^\dagger] = 0 \quad N_{k,\alpha} = a_{\mathbf{k},\alpha}^\dagger a_{\mathbf{k},\alpha}$$



# Photon Number State

$$|n_{k_1,\alpha_1}, n_{k_2,\alpha_2}, \dots, n_{k_l,\alpha_l}, \dots\rangle = |n_{k_1,\alpha_1}\rangle |n_{k_2,\alpha_2}\rangle \dots |n_{k_l,\alpha_l}\rangle \dots$$

$$|n_{k_1,\alpha_1}, n_{k_2,\alpha_2}, \dots\rangle = \prod_{k_i, \alpha_i} \frac{\left(a_{k_i, \alpha_i}^\dagger\right)^{n_{k_i, \alpha_i}}}{\sqrt{n_{k_i, \alpha_i}!}} |0\rangle$$



# Quantized Radiation Field I

$$A(x, t) = \left(1/\sqrt{V}\right) \sum_{\mathbf{k}} \sum_{\alpha} c \sqrt{\hbar/2\omega} [a_{\mathbf{k}, \alpha}(t) \epsilon^{(\alpha)} e^{i\mathbf{k} \cdot \mathbf{x}} + a_{\mathbf{k}, \alpha}^\dagger(t) \epsilon^{(\alpha)} e^{-i\mathbf{k} \cdot \mathbf{x}}]$$

$$H = \frac{1}{2} \sum_{\mathbf{k}} \sum_{\alpha} \hbar\omega (a_{\mathbf{k}, \alpha}^\dagger a_{\mathbf{k}, \alpha} + a_{\mathbf{k}, \alpha} a_{\mathbf{k}, \alpha}^\dagger) = \sum_{\mathbf{k}} \sum_{\alpha} \left(N_{\mathbf{k}, \alpha} + \frac{1}{2}\right) \hbar\omega$$

$$H = \sum_{\mathbf{k}} \sum_{\alpha} \hbar\omega N_{\mathbf{k}, \alpha}$$



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# Quantized Radiation Field II

$$a_{k,\alpha} = a_{k,\alpha}(0) e^{-i\omega t}, \quad a_{k,\alpha}^\dagger = a_{k,\alpha}^\dagger(0) e^{i\omega t}$$

$$A(x,t) = \frac{1}{\sqrt{V}} \sum_k \sum_\alpha c \sqrt{\frac{\hbar}{2\omega}} [a_{k,\alpha}(0) \epsilon^{(\alpha)} e^{ik \cdot x - i\omega t} + a_{k,\alpha}^\dagger(0) \epsilon^{(\alpha)} e^{-ik \cdot x + i\omega t}]$$



# Vacuum Fluctuation

$$\langle 0 | \mathbf{E} \cdot \mathbf{E} | 0 \rangle - \langle \langle 0 | \mathbf{E} | 0 \rangle \rangle^2 = \langle 0 | \mathbf{E} \cdot \mathbf{E} | 0 \rangle = \infty$$

$$\overline{\mathbf{E}} = (1/\Delta V) \int_{\Delta V} \mathbf{E} d^3x$$

$$\langle 0 | \overline{\mathbf{E}} \cdot \overline{\mathbf{E}} | 0 \rangle \sim \hbar c / (\Delta l)^4 \quad \text{Casimir effect}$$



# Uncertainty Relation I

$$[E_x(\mathbf{x}, t), B_y(\mathbf{x}', t)] = i\hbar \frac{\partial}{\partial z} \delta^3(\mathbf{x} - \mathbf{x}')$$

$$[E_x(\mathbf{x}, t), B_y(\mathbf{x}', t')] = 0 \quad \text{for} \quad (\mathbf{x} - \mathbf{x}')^2 - (1/c^2)(t - t')^2 \neq 0$$



# Uncertainty Relation II

$$a_{k,\alpha} = e^{i(\phi_{k,\alpha} - \omega t)} \sqrt{N_{k,\alpha}}$$

$$a_{k,\alpha}^\dagger = \sqrt{N_{k,\alpha}} e^{i(\phi_{k,\alpha} - \omega t)}$$

$$[N, \phi] = i$$

$$\Delta N \Delta \phi \gtrsim 1$$



# Atom-Photon Interaction I

$$H_{\text{int}} = \sum_t \left[ -\frac{e}{2mc} (\mathbf{p}_i \cdot \mathbf{A}(\mathbf{x}_i, t) + \mathbf{A}(\mathbf{x}_i, t) \cdot \mathbf{p}_i) + \frac{e^2}{2mc^2} \mathbf{A}(\mathbf{x}_i, t) \cdot \mathbf{A}(\mathbf{x}_i, t) \right]$$

- Photon Absorption

$$\begin{aligned} \langle B; n_{\mathbf{k},\alpha} - 1 | H_{\text{int}} | A; n_{\mathbf{k},\alpha} \rangle &= -\frac{e}{mc} \langle B; n_{\mathbf{k},\alpha} - 1 | \sum_i c \sqrt{\frac{\hbar}{2\omega V}} a_{\mathbf{k},\alpha}(0) e^{i\mathbf{k} \cdot \mathbf{x}_i - i\omega t} \mathbf{p}_i \boldsymbol{\varepsilon}^{(\alpha)} | A; n_{\mathbf{k},\alpha} \rangle \\ &= -\frac{e}{mc} \sqrt{\frac{n_{\mathbf{k},\alpha} \hbar}{2\omega V}} \sum_i \langle B | e^{i\mathbf{k} \cdot \mathbf{x}_i} \mathbf{p}_i \cdot \boldsymbol{\varepsilon}^{(\alpha)} | A \rangle e^{-i\omega t} \end{aligned}$$

- Photon Emission

$$\langle B; n_{\mathbf{k},\alpha} + 1 | H_{\text{int}} | A; n_{\mathbf{k},\alpha} \rangle = -\frac{e}{m} \sqrt{\frac{(n_{\mathbf{k},\alpha} + 1) \hbar}{2\omega V}} \sum_i \langle B | e^{i\mathbf{k} \cdot \mathbf{x}_i} \mathbf{p}_i \cdot \boldsymbol{\varepsilon}^{(\alpha)} | A \rangle e^{i\omega t}$$



# Atom-Photon Interaction II

**absorption** :  $c\sqrt{\frac{n_{\mathbf{k},\alpha}\hbar}{2\omega V}}\mathcal{E}^{(\alpha)}e^{i\mathbf{k}\cdot\mathbf{x}-i\omega t}$

**emission** :  $c\sqrt{\frac{(n_{\mathbf{k},\alpha}+1)\hbar}{2\omega V}}\mathcal{E}^{(\alpha)}e^{-i\mathbf{k}\cdot\mathbf{x}+i\omega t}$



# Time-Dependent Perturbation Theory I

$$\psi = \sum_k c_k(t) u_k(\mathbf{x}) e^{-iE_k t/\hbar}$$

$$H_0 u_k(\mathbf{x}) = E_k u_k(\mathbf{x})$$

$$(H_0 + H_I)\psi = i\hbar (\partial\psi/\partial t) = i\hbar \sum_k (\dot{c}_k u_k e^{-iE_k t/\hbar} - i(E_k/\hbar) c_k u_k e^{-iE_k t/\hbar})$$

$$\sum_k H_I c_k u_k e^{-iE_k t/\hbar} = i\hbar \sum_l \dot{c}_l u_l e^{-iE_l t/\hbar}$$



# Time-Dependent Perturbation Theory II

$$\dot{c}_m = \sum_k (1/i\hbar) \langle m | H_I(t) | k \rangle e^{i(E_m - E_k)t/\hbar} c_k(t)$$

$$c_k(0) = \delta_{kl}$$

$$c_m^{(1)}(t) = \frac{1}{i\hbar} \int_0^t dt' \langle m | H_I(t') | l \rangle e^{i(E_m - E_l)t'/\hbar}$$

$$c_m \simeq c_m^{(1)} + c_m^{(2)}$$

$$c_m^{(2)}(t) = \frac{1}{(i\hbar)^2} \sum_n \int_0^t dt'' \int_0^{t''} dt' \langle m | H_I(t'') | n \rangle e^{i(E_m - E_n)t''/\hbar} \langle n | H_I(t') | l \rangle e^{i(E_n - E_l)t'/\hbar}$$



# Time-Dependent Perturbation Theory III

$$H_I(t) = H'_I e^{\mp i\omega t} \quad \text{for} \quad \begin{cases} \text{absorption} \\ \text{emission} \end{cases}$$

$$c_m^{(1)} = \frac{1}{i\hbar} \langle m | H'_I | l \rangle \int_0^t dt' e^{i(E_m - E_l \mp \hbar\omega)t'/\hbar}$$

$$|c_m^{(1)}(t)|^2 = (2\pi/\hbar) |\langle m | H'_I | l \rangle|^2 t \delta(E_m - E_l \mp \hbar\omega)$$



# Fermi's Golden Rule

$$\rho_{\hbar\omega,d\Omega} = \frac{V|k|^2}{(2\pi)^3} \frac{d|k|d\Omega}{d(\hbar\omega)} = \frac{V\omega^2}{(2\pi)^3} \frac{d\Omega}{\hbar c^3}$$

$$w_{d\Omega} = \int \left( |c_m^{(1)}|^2 / t \right) \rho_{\hbar\omega,d\Omega} d(\hbar\omega) = (2\pi/\hbar) |\langle m | H'_I | l \rangle|^2 \rho_{\hbar\omega,d\Omega}$$

Fermi's Golden Rule

$$E_m - E_l + \hbar\omega = 0$$



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# Spontaneous Emission in the Dipole Approximation I

$$w_{d\Omega} = \frac{2\pi}{\hbar} \frac{e^2 \hbar}{2m^2 \omega V} \left| \sum_i \langle B | e^{-i\mathbf{k} \cdot \mathbf{x}_i} \boldsymbol{\varepsilon}^{(\alpha)} \cdot \mathbf{p}_i | A \rangle \right|^2 \frac{V \omega^3 d\Omega}{(2\pi)^3 \hbar c^3}$$

$$\aleph_{\text{photon}} = 1/|k\rangle\rangle \gamma_{\text{atom}}$$

$$e^{-i\mathbf{k} \cdot \mathbf{x}_i} = 1 - i\mathbf{k} \cdot \mathbf{x}_i - (\mathbf{k} \cdot \mathbf{x}_i)^2/2 + \dots \approx 1$$

$$w_{d\Omega} = \frac{e^2 \omega}{8\pi^2 m^2 \hbar c^3} \left| \langle B | \mathbf{p} | A \rangle \cdot \boldsymbol{\varepsilon}^{(\alpha)} \right|^2 d\Omega$$



# Spontaneous Emission in the Dipole Approximation II

$$[\mathbf{p}^2, \mathbf{x}] = -2i\hbar p$$

$$\langle B | \mathbf{p} | A \rangle = \left\langle B \left| \frac{im}{\hbar} [H_0, \mathbf{x}] \right| A \right\rangle = \frac{im(E_B - E_A)}{\hbar} \langle B | \mathbf{x} | A \rangle = -im\omega \mathbf{x}_{BA}$$

$$-\frac{e\mathbf{A} \cdot \mathbf{p}}{mc} \rightarrow \frac{e}{c} \mathbf{x} \cdot \frac{\partial \mathbf{A}}{\partial t} = -e\mathbf{x} \cdot \mathbf{E}$$



# Spontaneous Emission in the Dipole Approximation III

$$w_{d\Omega} = \frac{e^2 \omega^3}{8\pi^2 \hbar c^3} |\mathbf{x}_{BA}|^2 \cos^2 \Theta^{(\alpha)} d\Omega$$

$$\cos \Theta^{(\alpha)} = \left| \mathbf{x}_{BA} \cdot \boldsymbol{\epsilon}^{(\alpha)} \right| / |\mathbf{x}_{BA}|$$

$$\cos \Theta^{(1)} = \sin \theta \cos \phi, \quad \cos \Theta^{(2)} = \sin \theta \sin \phi$$

$$2\pi \int_{-1}^1 \sin^2 \theta d(\cos \theta) = 8\pi/3$$

$$W = \frac{e^2 \omega^3}{3\pi \hbar c^3} |\mathbf{x}_{BA}|^2 = \left( \frac{e^2}{4\pi \hbar c} \right) \frac{4}{3} \frac{\omega^3}{c^2} |\mathbf{x}_{BA}|^2$$

