

Chapter 4 Two-Phase Flow

I. Basic Definitions

(1) Void Fractions

: the time averaged volumetric fraction of vapor in a two phase mixture.

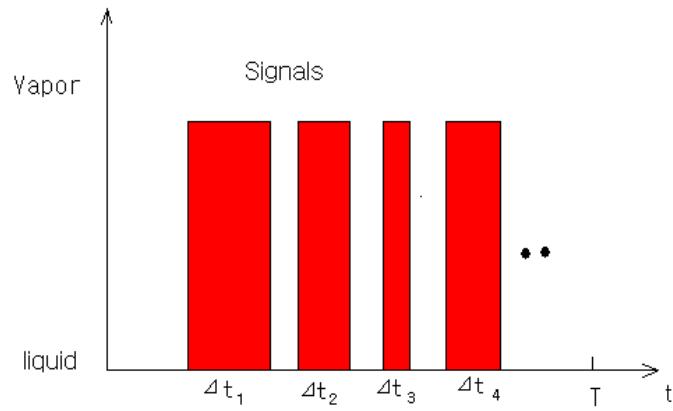
$$\alpha = \int \int \int_{V_i} dV / \int \int \int_{V_i} dV = V_v / (V_l + V_v) \quad (4.1.1)$$

o Measurement

1) Needle probe method

- electrical probe
- optical probe

2) X-ray absorption method (Jones and Zuber)



$$\alpha = \frac{\sum_{i=1}^N \Delta t_i}{T}$$

For differential length element

$$\begin{aligned} \alpha &= \Delta z \int \int_{V_i} dA / (\Delta z \int \int_{V_i} dA) \\ &= A_g / (A_f - A_g) \end{aligned} \quad (4.1.2)$$

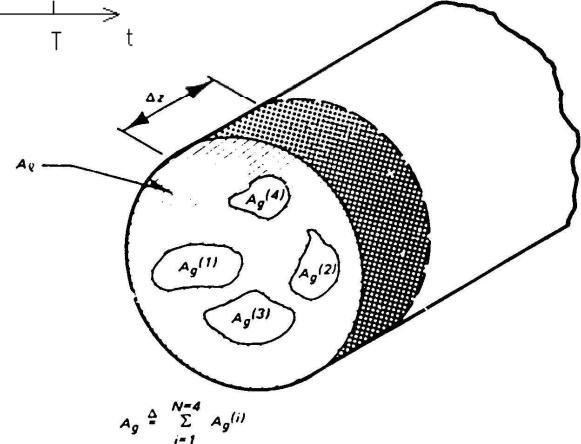


Fig. 4.1 Representation of void fraction.

(2) Phase Velocity

: the volumetric flow rate of the phase through its cross sectional flow area

$$u_i = Q_i / A_i$$

$$u_v = Q_v / A_v \quad (4.1.3)$$

where Q_i is the volumetric flow rate, $\text{ft}^3/\text{sec.}$, of phase i

(3) Volumetric flux or superficial velocity

: the volumetric flow rate of the phase divided by the total cross sectional flow area.

$$j_i = Q_i / (A_i + A_v)$$

$$j_v = Q_v / (A_i + A_v) \quad (4.1.4)$$

Now, using the relations given in Eqs.(4.1.2) and (4.1.3), we can write

$$j_i = u_i(1-\alpha)$$

$$j_v = u_v\alpha \quad (4.1.5)$$

And the total volumetric flux of the mixture, j, is

$$j = j_i + j_v \quad (4.1.6)$$

(4) Relative velocity, u_r

: velocity difference between the vapor and the liquid phase

$$u_r = u_v - u_i \quad (4.1.7)$$

For homogeneous flow

$$u_r = 0$$

Thus, from Eq. (4.1.5)

$$j_v/\alpha = j_i/(1-\alpha)$$

By solving for the void fraction

$$\alpha = j_v / (j_i + j_v) \quad (4.1.8)$$

(5) Slip ratio

: the ratio of the phase velocity

$$S = u_v / u_i \quad (4.1.9)$$

$$= \frac{x}{(1-x)} \cdot \frac{(1-\alpha)}{\alpha} \cdot \frac{\rho_i}{\rho_v}$$

For homogeneous flow, $S=1$, by definition.

(6) Quality

1. Thermal equilibrium quality : the flow fraction of vapor in the thermal equilibrium condition

$$x_e = (h - h_f)/h_{fg} \quad (4.1.10)$$

2. Flow quality : the true flow fraction of vapor regardless of whether thermal equilibrium exists or not

$$\begin{aligned} x &= \rho_v u_v A_v / (\rho_v u_v A_v + \rho_a u_a A_a) \\ &= u_v / u \end{aligned} \quad (4.1.11)$$

3. Static quality : the mass fraction of vapor

$$x_s = \rho_v A_v / (\rho_v A_v + \rho_a A_a) \quad (4.1.12)$$

From the definitions given in Eqs.(4.1.11) and (4.1.12), x and x_s can be related through the following identity

$$x/(1-x) = (u_v/u_f)x_s/(1-x_s) = Sx_s/(1-x_s) \quad (4.1.13)$$

(7) Cross-sectional Average Notation, $\langle \cdot \rangle$

$$\begin{aligned} \langle \zeta \rangle &= \frac{1}{A_{\perp,\perp}} \iint_A \zeta \, dA \\ \langle \zeta \rangle_f &= \frac{1}{A_{\perp,\perp}(1-\langle \alpha \rangle)} \iint_A \zeta / (1-\alpha) \, dA = \frac{\langle (1-\alpha)\zeta_f \rangle}{\langle 1-\alpha \rangle} \\ \langle \zeta \rangle_{\perp} &= \frac{1}{A_{\perp,\perp}\langle \alpha \rangle} \iint_A \zeta \alpha \, dA = \frac{\langle \alpha \zeta_{\perp} \rangle}{\langle \alpha \rangle} \end{aligned} \quad (4.1.14)$$

Note : $\langle j^2 \rangle / \langle j \rangle^2$

(8) Mass flux, G

$$G = w/A_{\perp,\perp} = \rho_f \langle u \rangle_f (1-\langle \alpha \rangle) / (1-\langle x \rangle) = \rho_v \langle u_v \rangle_f \langle \alpha \rangle / \langle x \rangle \quad (4.1.15)$$

where the last two equalities come from the identities

$$G(1-\langle x \rangle)A_{\perp,\perp} = \rho_f \langle u \rangle_f (1-\langle \alpha \rangle)A_{\perp,\perp}$$

and

$$G\langle x \rangle A_{\perp,\perp} = \rho_v \langle u_v \rangle_f \langle \alpha \rangle A_{\perp,\perp}$$

which also imply

$$\langle u \rangle_f = G(1-\langle x \rangle) / [\rho_f(1-\langle \alpha \rangle)] \quad (4.1.16)$$

$$\langle u_v \rangle_f = G\langle x \rangle / (\rho_v \langle \alpha \rangle) \quad (4.1.17)$$

(9) The fundamental Void Quality Relation

The fundamental void quality relation is obtained as,

$$\langle \alpha \rangle = \langle x \rangle / [\langle x \rangle - S(\rho_f/\rho_g)(1-\langle x \rangle)] \quad (4.1.18)$$

With respect to the static quality, Eq.(4.1.18) will be

$$\langle \alpha \rangle = \langle x_g \rangle \rho_f / [\langle x_g \rangle (\rho_f - \rho_g) - \rho_g] \quad (4.1.19)$$

and, from $\langle \bar{\rho} \rangle = (1-\alpha)\rho_f + \alpha\rho_g$

$$\langle \bar{\rho} \rangle = \rho_f \rho_g / [\langle x_g \rangle \rho_f - (1-\langle x_g \rangle) \rho_g] = 1/(v_f - \langle x_g \rangle v_{fg}) \quad (4.1.20)$$

For homogeneous flow condition, $S=1$

$$\langle \rho_w \rangle = \frac{G}{\langle j \rangle} = \frac{w}{Q} = 1/(v_f + \langle x_g \rangle v_{fg}) \quad (4.1.21)$$

but for slip flow condition

$$\langle \bar{\rho} \rangle = 1/(v_f - \langle x_g \rangle v_{fg})$$

(10) Drift velocity, V_d

: the void weighted average velocity of vapor phase with respect to the velocity of the center of volume of the mixture.

$$V_d = \langle \alpha(u_v - j) \rangle / \langle \alpha \rangle \quad (4.1.22)$$

II. Two phase Model

o Number of unknowns for three dimensional flow

Single Phase		Two Phase	
Parameter	No	Parameter	No
Velocity Vector	3	Void Fraction	1
Pressure	1	Velocity Vector of Liquid	3
Temperature	1	Velocity Vector of Vapor	3
Density	1	Density of Liquid	1
		Density of Vapor	1
		Pressure	1
		Temperature of each Phase	2
Total number of unknowns	6	Total number of unknowns	12

o Number of governing equations

Single Phase		Two Phase	
Parameter	No	Parameter	No
Conservation of Mass	1	Mass Eq. for Liquid	1
Conservation of Momentum	3	Mass Eq. for Vapor	1
Conservation of Energy	1	Momentum Eq. for Liquid	3
Equation of State	1	Momentum Eq. for vapor	3
		Energy Eq. for each Phase	2
		Eq. of State for each Phase	2
Total number of equations	6	Total number of equations	12

o Constitutive Equations

- interfacial transfer (jump condition)
- various correlations dependent on the flow regime

o Various forms for two phase flow modeling

- Local Instantaneous Formula
- Instantaneous Space or Local Time Averaged Equations
- Homogeneous Model
- Two Fluid Model (Drift Flux Model)

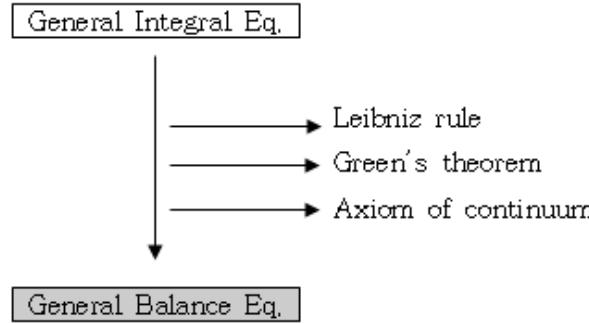
1. Local Instantaneous Formula

- balance equations for each phase
- interfacial condition(i.e. jump condition) as well as BC & IC

$$\bar{F} = \bar{F}(\underline{x}, t)$$

■ Fundamental Importances in the Local Instantaneous Formula

- direct application to study separated flow
: Δp calculation, instability analysis, bubble dynamics, etc
- fundamental base of all two-phase models
perform overall calculations using averaging techniques
void fraction calculation, development of drift flux model, etc



The general integral balance can be written as,

$$\frac{d}{dt} \int_{V_i} \rho_i \phi_i dV = - \int_{A_i} \bar{u}_i \cdot \bar{J}_i dA + \int_{V_i} \rho_i \Phi_i dV \quad (4.2.1)$$

Using the Reynolds transport theorem to obtain the differential form of the balance equation, i.e

$$\frac{d}{dt} \int_{V_i} F_i dV = \int_{V_i} \frac{\partial F_i}{\partial t} dV - \int_{A_i} F_i \bar{v}_i \cdot \bar{n} dA \quad (4.2.2)$$

the differential equation is,

$$-\frac{\partial \rho_i \phi_i}{\partial t} + \bar{v} \cdot (\bar{v}_i \rho_i \phi_i) = -\bar{v} \cdot \bar{J}_i + \rho_i \Phi_i \quad (4.2.3)$$

Balance	Ψ_k	\bar{J}_k	Φ_k
Mass	1	0	0
Momentum	v_k	$\rho_k \bar{I} - \bar{\tau}_k$	\bar{g}_k
Energy	$u_k + \frac{v_k^2}{2}$	$q_k - \bar{T}_k \cdot v_k$	$\bar{g}_k \cdot \bar{v}_k + \frac{Q_k}{\rho_k}$

(1) Interfacial balance (Jump condition)

Consider the interface as a singular surface across which the fluid density, energy and velocity suffer from jump discontinuity, although mass, momentum and energy must be conserved.

Mass Jump

$$\rho_1(\vec{v}_1 - \vec{v}_2) \cdot \vec{n}_1 + \rho_2(\vec{v}_2 - \vec{v}_1) \cdot \vec{n}_2 = 0 \quad (4.2.4)$$

or $\dot{m}_1 = \dot{m}_2 = 0$

If no mass transfer, $\dot{m}_1 = \dot{m}_2 = 0$

$$\therefore (\vec{v}_1 - \vec{v}_2) \cdot \vec{n} = 0 \quad \text{or} \quad \vec{v}_1 = \vec{v}_2$$

Linear Momentum Jump

$$\dot{m}_1 \vec{v}_1 + \dot{m}_2 \vec{v}_2 - \vec{n}_1 \cdot \vec{T}_1 - \vec{n}_2 \cdot \vec{T}_2 = 0 \quad (4.2.5)$$

From Eq. (4.2.5), $\dot{m}_1(\vec{v}_1 - \vec{v}_2) - \vec{n}_1 \cdot \vec{T}_1 - \vec{n}_2 \cdot \vec{T}_2 = 0$

a: for inviscid phase ($\vec{T} = \rho \vec{v}$)

$$\dot{m}_1(\vec{v}_1 - \vec{v}_2) - \vec{n}_1(\rho_1 - \rho_2) = 0$$

b: for viscous flow (1D flow which interface perpendicular to flow direction)

$$\dot{m}_1(\vec{v}_1 - \vec{v}_2) + \vec{n}_1 \cdot (\rho_1 - \rho_2) + \vec{n}_1 \cdot (\tau_2 - \tau_1) = 0$$

For 2 dimensional case with surface tension

$$\dot{m}_1 \vec{v}_1 - \dot{m}_2 \vec{v}_2 - \vec{n}_1 \cdot \vec{T}_1 - \vec{n}_2 \cdot \vec{T}_2 + \frac{d\sigma}{dl} \vec{\tau} - \frac{\sigma}{R} \vec{n}_1 = 0 \quad (4.2.6)$$

If no mass transfer and the phase are inviscid,

$$\vec{n}_1 \cdot (\rho_1 - \rho_2) + \frac{d\sigma}{dl} \vec{\tau} - \frac{\sigma}{R} \vec{n}_1 = 0 \quad (4.2.7)$$

If the surface tension is constant,

$$\rho_1 = \rho_2 = \frac{\sigma}{R} \quad (4.2.8)$$

Total Energy Jump

$$\tilde{m}_1(u_1 + \frac{1}{2}v_1^2) - \tilde{m}_2(u_2 + \frac{1}{2}v_2^2) + \vec{q}_1 \cdot \vec{u}_1 + \vec{q}_2 \cdot \vec{u}_2 - (\vec{u}_1 \cdot \vec{T}_1) \cdot \vec{v}_1 - (\vec{u}_2 \cdot \vec{T}_2) \cdot \vec{v}_2 = 0 \quad (4.2.9)$$

(2) Interfacial boundary conditions

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a) thermal boundary condition

$$T_1 = T_2 = T_i \quad (4.2.10)$$

b) mechanical boundary condition (No-slip condition)

$$\vec{v}_{i1} = \vec{v}_{i2} = \vec{v}_{ii} \quad (4.2.11)$$

c) phase change boundary condition

$$g_1 = g_2 = \left\{ \frac{|\vec{v}_2 - \vec{v}_i|^2}{2} - \frac{|\vec{v}_1 - \vec{v}_i|^2}{2} \right\} - \left\{ \frac{1}{\rho_2} (\vec{\tau}_2 \cdot \vec{u}_2) \cdot \vec{u}_2 - \frac{1}{\rho_1} (\vec{\tau}_1 \cdot \vec{u}_1) \cdot \vec{u}_1 \right\} \quad (4.2.12)$$

Note that in thermodynamic equilibrium,

$$T_1 = T_2 = 0$$

$$\rho_1 = \rho_2 = 0$$

$$g_1 = g_2 = 0$$