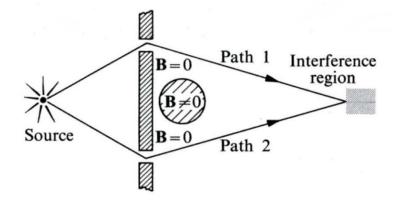


Aharonov-Bohm Effect

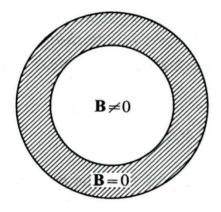
$$\frac{1}{2m} \left(-i\hbar \nabla - \frac{e\mathbf{A}}{c} \right) \psi + \mathbf{V} \psi = E \psi$$
$$\psi = \psi_1^{(0)} \exp\left[\frac{ie}{\hbar c} \int_{\text{Path 1}}^{\mathbf{s}(\mathbf{x})} \mathbf{A}(\mathbf{x}') \cdot d\mathbf{s}'\right] + \psi_2^{(0)} \exp\left[\frac{ie}{\hbar c} \int_{\text{Path 2}}^{\mathbf{s}(\mathbf{x})} \mathbf{A}(\mathbf{x}') \cdot d\mathbf{s}'\right]$$

1/2

$$\begin{cases} \cos \\ \sin \end{cases} \left[\frac{e}{\hbar c} \oint \mathbf{A} \cdot d\mathbf{s} \right] = \begin{cases} \cos \\ \sin \end{cases} \left[\frac{e}{\hbar c} \int \mathbf{B} \cdot \hat{n} \, dS \right] \\ = \begin{cases} \cos \\ \sin \end{cases} \frac{e\Phi}{\hbar c} \end{cases}$$







Center for Active Plasmonics



Atom-Radiation Interaction

$$\hat{H} = \frac{1}{2m} \left[\hat{\mathbf{p}} - \frac{e}{c} \mathbf{A}(\mathbf{r}, t) \right]^2 + V(\mathbf{r})$$

$$\hat{H} = \hat{H}^{(0)} + \hat{H}' + \hat{H}''$$

$$\hat{H}^{(0)} = \frac{\hat{p}^2}{2m} + V; \qquad \hat{H}' = -\frac{e}{2mc} [\hat{\mathbf{p}} \cdot \mathbf{A} + \mathbf{A} \cdot \hat{\mathbf{p}}]$$

$$H'' = \frac{e^2}{2mc^2} A^2$$

$$\langle n'|H'|n \rangle = -\frac{e}{2mc} \int \psi_{n'}^* (\hat{\mathbf{p}} \cdot \mathbf{A} + \mathbf{A} \cdot \hat{\mathbf{p}}) \psi_n \, d\mathbf{r}$$

$$\hat{\mathbf{p}} \cdot \mathbf{A} \psi_n = \frac{\hbar}{i} \nabla \cdot (\mathbf{A} \psi_n) = \frac{\hbar}{i} [\psi_n (\nabla \cdot \mathbf{A}) + (\mathbf{A} \cdot \nabla) \psi_n]$$

$$\langle n'|H'|n \rangle = -\frac{e}{mc} \int \psi_{n'}^* \mathbf{A} \cdot \hat{\mathbf{p}} \psi_n \, d\mathbf{r}$$

$$\langle n'|H'|n \rangle = -\frac{e}{mc} \int \psi_{n'}^* \hat{\mathbf{p}} \cdot \mathbf{A} \psi_n \, d\mathbf{r}$$

Center for Active Plasmonics





Atom-Radiation Interaction

$$\langle n'|\hat{H}'|n\rangle \rightarrow -\frac{e}{mc}\langle n';\hbar\omega|\hat{\mathbf{p}}|n\rangle$$
 (emission)

$$\langle n'|\hat{H}'|n\rangle \rightarrow -\frac{e}{mc}\langle n'|\hat{\mathbf{p}}|\hbar\omega;n\rangle$$
 (absorption)

$$\mathbf{A}(\mathbf{r},t) = \mathbf{a}A_0\cos(\mathbf{k}\cdot\mathbf{r}-\omega t)$$

Seoul National University

Center for Active Plasmonics



Atom-Radiation Interaction

$$\langle U \rangle = \frac{1}{4\pi} \langle \mathcal{B}^2 \rangle = \frac{1}{4\pi} \langle \mathcal{E}^2 \rangle$$

 $\boldsymbol{\mathcal{B}} = \boldsymbol{\nabla} \times \mathbf{A}$

$$\mathbf{A} = \frac{A_0}{2} \mathbf{a} [e^{i(\mathbf{k} \cdot \mathbf{r} - \omega t)} + e^{-i(\mathbf{k} \cdot \mathbf{r} - \omega t)}]$$

$$\mathbf{A}_{\pm} = c \left(\frac{2\pi\hbar}{\omega}\right)^{1/2} \mathbf{a} e^{\pm i \left(\mathbf{k} \cdot \mathbf{r} - \omega t\right)}$$

$$\hat{H}' = \hat{\mathbb{H}}_{\pm}(r) e^{\mp i \omega t}$$

$$\langle n' | \mathbb{H}_{\pm} | n \rangle = -\frac{e}{m} \left(\frac{2\pi\hbar}{\omega} \right)^{1/2} \langle n' | \mathbf{a} \cdot \hat{\mathbf{p}} e^{\pm i \mathbf{k} \cdot \mathbf{r}} | n \rangle$$

Seoul National University

Center for Active Plasmonics



The Dipole Approximation

$$\langle n' | \mathbb{H}_{\pm} | n \rangle = -\frac{e}{m} \left(\frac{2\pi\hbar}{\omega} \right)^{1/2} \langle n' | \mathbf{a} \cdot \hat{\mathbf{p}} | n \rangle$$

$$\langle n'|\mathbf{p}|n\rangle = -\frac{im}{\hbar} (E_n^{(0)} - E_{n'}^{(0)}) \langle n'|\mathbf{r}|n\rangle$$
$$\langle n'|\mathbf{p}|n\rangle = -im\omega \langle n'|\mathbf{r}|n\rangle$$

$$\langle n' | \mathbb{H}_+ | n, \mathbf{k} \rangle = \langle n', \mathbf{k} | \mathbb{H}_- | n \rangle$$

= $i e (2\pi \hbar \omega)^{1/2} \langle n' | \mathbf{a} \cdot \mathbf{r} | n \rangle$

Seoul National University

Center for Active Plasmonics



Quantum Mechanical Description of Vector Potential and its Results

$$\mathbf{A}(\mathbf{x},t) = \frac{1}{\sqrt{V}} \sum_{\mathbf{k}} \sum_{\alpha} c_{\sqrt{\frac{\hbar}{2\omega}}} [a_{\mathbf{k},\alpha}(0)\boldsymbol{\epsilon}^{(\alpha)} e^{i\mathbf{k}\cdot\mathbf{x}-i\omega t} + a_{\mathbf{k},\alpha}^{\dagger}(0)\boldsymbol{\epsilon}^{(\alpha)} e^{-i\mathbf{k}\cdot\mathbf{x}+i\omega t}]$$

$$\langle B; n_{\mathbf{k},\alpha} - 1 | H_{\text{int}} | A; n_{\mathbf{k},\alpha} \rangle$$

$$= -\frac{e}{mc} \langle B; n_{\mathbf{k},\alpha} - 1 | \sum_{i} c_{\sqrt{\frac{\hbar}{2\omega V}}} a_{\mathbf{k},\alpha}(0) e^{i\mathbf{k}\cdot\mathbf{x}_{i}-i\omega t} \mathbf{p}_{i} \cdot \boldsymbol{\epsilon}^{(\alpha)} | A; n_{\mathbf{k},\alpha} \rangle$$

$$= -\frac{e}{m} \sqrt{\frac{n_{\mathbf{k},\alpha}\hbar}{2\omega V}} \sum_{i} \langle B | e^{i\mathbf{k}\cdot\mathbf{x}_{i}} \mathbf{p}_{i} \cdot \boldsymbol{\epsilon}^{(\alpha)} | A \rangle e^{-i\omega t}$$

$$\langle B; n_{\mathbf{k},\alpha} + 1 | H_{\mathrm{int}} | A; n_{\mathbf{k},\alpha} \rangle = -\frac{e}{m} \sqrt{\frac{(n_{\mathbf{k},\alpha} + 1)\hbar}{2\omega V}} \sum_{i} \langle B | e^{-i\mathbf{k}\cdot\mathbf{x}_{i}} \mathbf{p}_{i} \cdot \boldsymbol{\epsilon}^{(\alpha)} | A \rangle e^{i\omega t}$$

absorption:
$$c\sqrt{\frac{n_{\mathbf{k},\alpha}\hbar}{2\omega V}} \epsilon^{(\alpha)} e^{i\mathbf{k}\cdot\mathbf{x}-i\omega t}$$
,
emission: $c\sqrt{\frac{(n_{\mathbf{k},\alpha}+1)\hbar}{2\omega V}} \epsilon^{(\alpha)} e^{-i\mathbf{k}\cdot\mathbf{x}+i\omega t}$

Center for Active Plasmonics

Application Systems

Seoul National University



Perturbation Theory

$$c_m^{(1)}(t) = \frac{1}{i\hbar} \int_0^t dt' \langle m | H_I(t') | l \rangle e^{i(E_m - E_l)t'/\hbar}$$

$$c_{m}^{(2)}(t) = \frac{1}{i\hbar} \sum_{n} \int_{0}^{t} dt'' \langle m | H_{I}(t'') | n \rangle e^{i(E_{m} - E_{n})t''/\hbar} c_{n}^{(1)}(t'')$$

$$= \frac{1}{(i\hbar)^{2}} \sum_{n} \int_{0}^{t} dt'' \int_{0}^{t''} dt' \langle m | H_{I}(t'') | n \rangle e^{i(E_{m} - E_{n})t''/\hbar} \langle n | H_{I}(t') | l \rangle e^{i(E_{n} - E_{l})t'/\hbar}$$

Seoul National University

Center for Active Plasmonics