

2017 Spring

“Calculation and Applications Phase Equilibria”

Principles *of* Solidification

04.12.2017

Eun Soo Park

Office: 33-313

Telephone: 880-7221

Email: espark@snu.ac.kr

Office hours: by appointment

Contents for previous class

Solidification: Liquid → Solid

< Nucleation in Pure Metals >

* Homogeneous Nucleation

$$r^* = \frac{2\gamma_{SL}}{\Delta G_V} \quad \Delta G^* = \frac{16\pi\gamma_{SL}^3}{3(\Delta G_V)^2} = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3L_V^2} \right) \frac{1}{(\Delta T)^2}$$

r* & ΔG* ↓ as ΔT ↑

$$N_{hom} \approx f_0 C_o \exp\left\{-\frac{A}{(\Delta T)^2}\right\}$$

changes by orders of magnitude from zero to very high values over a very narrow temp. range

* Heterogeneous Nucleation

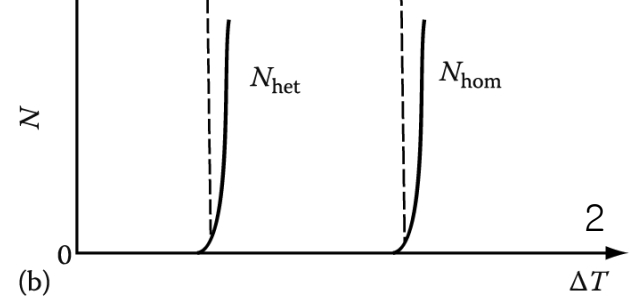
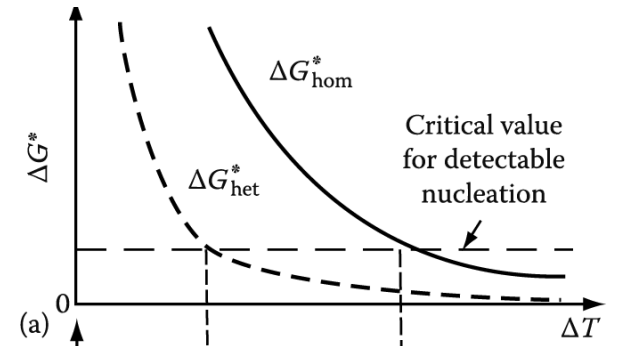
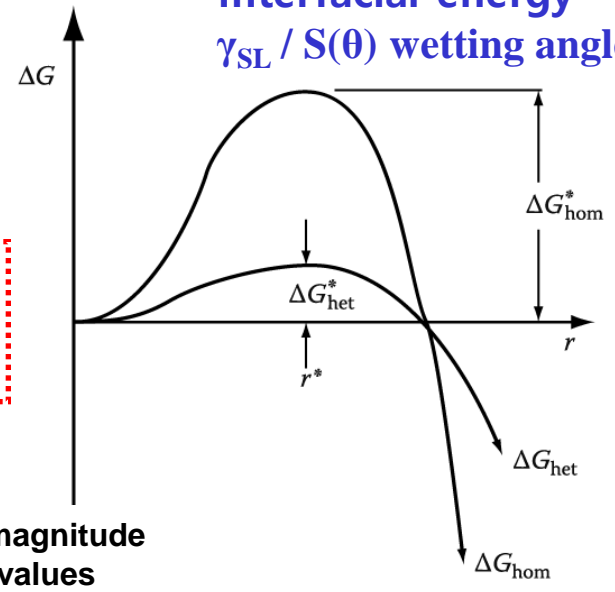
$$\Delta G_{het}^* = S(\theta)\Delta G_{hom}^*$$

$$\frac{V_A}{V_A + V_B} = \frac{2 - 3\cos\theta + \cos^3\theta}{4} = S(\theta)$$

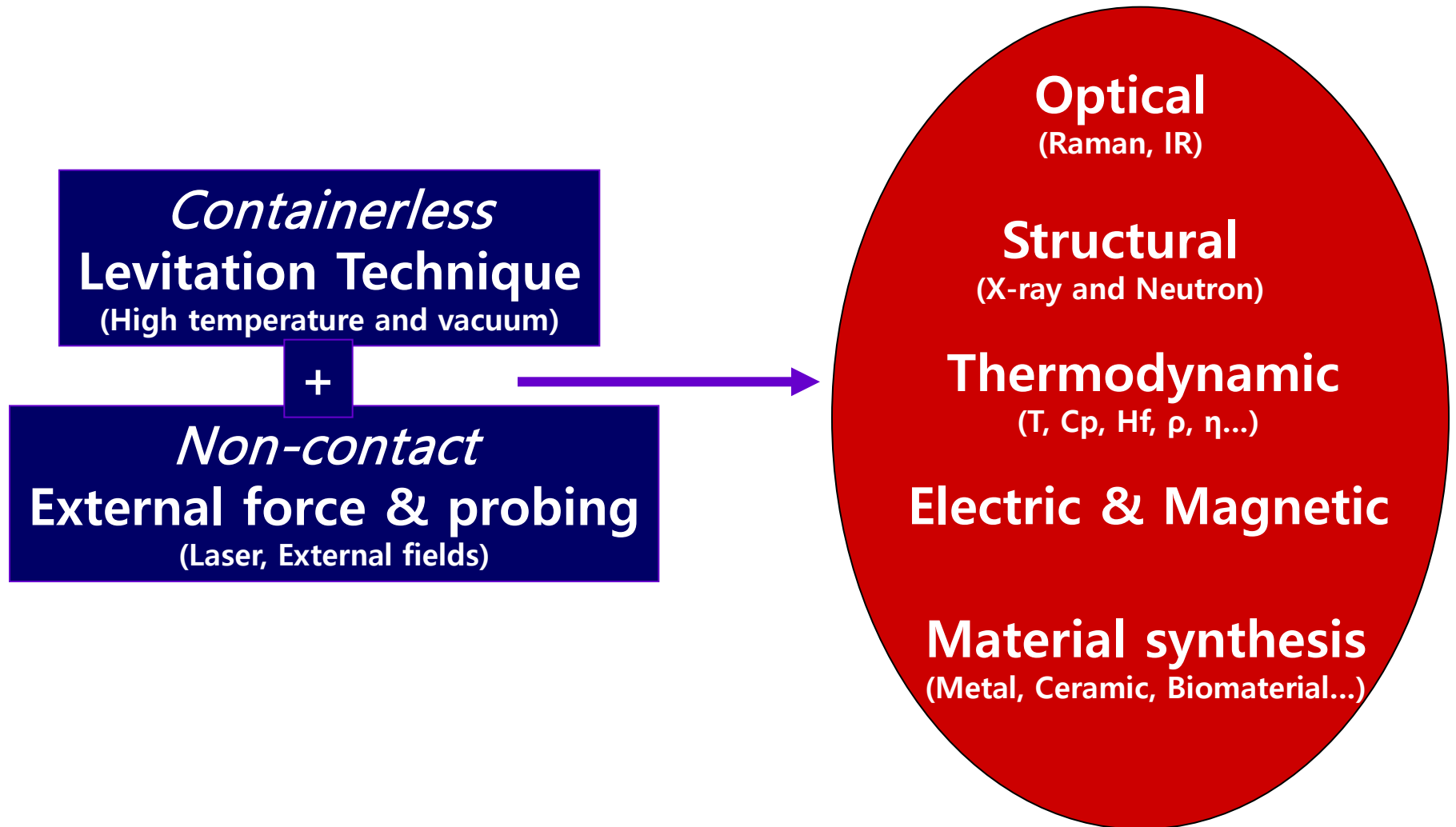
* Nucleation of melting

$$\gamma_{SL} + \gamma_{LV} < \gamma_{SV} \quad (\text{commonly})$$

- Undercooling ΔT
- Interfacial energy γ_{SL} / S(θ) wetting angle



Containerless and Contactless Measurement System

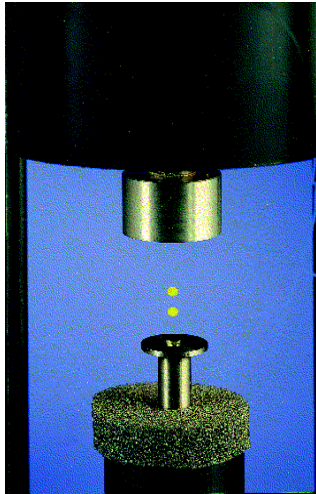
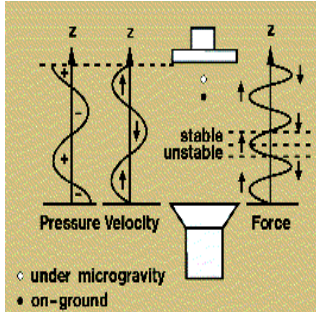


High Temperature Levitation

Ultra-high temperature > 3000 °C

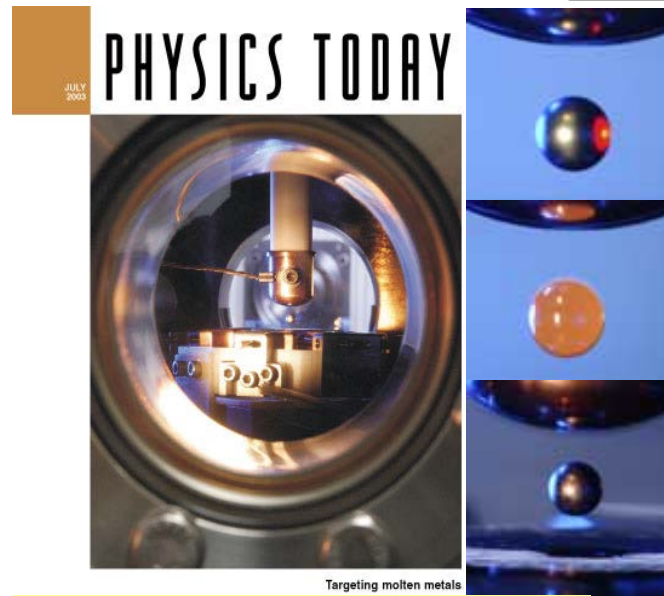
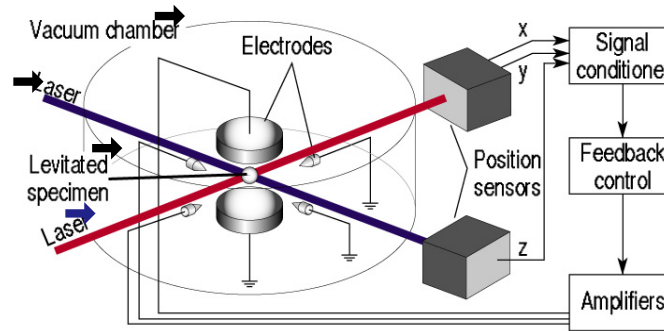
* Magnetic/diamagnetic/superconducting levitation → Only magnetic sample, below T_c

Acoustic



Requirement of acoustic media, Unstable at high T

Electrostatic



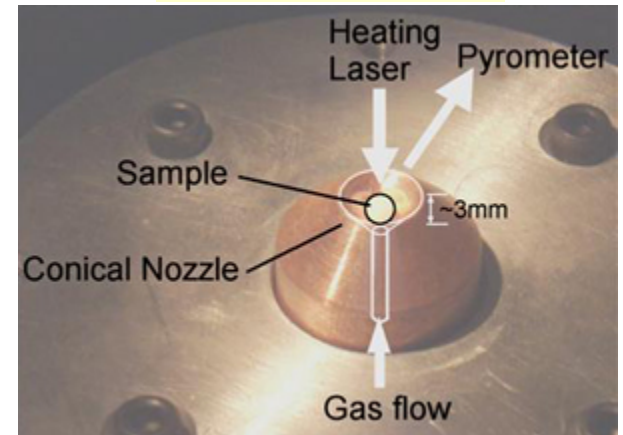
Physics Today, v56, p22, July 2003

All types of samples,
Suitable for sample heating

Electromagnetic



Only metallic & large mass sample
Aerodynamic



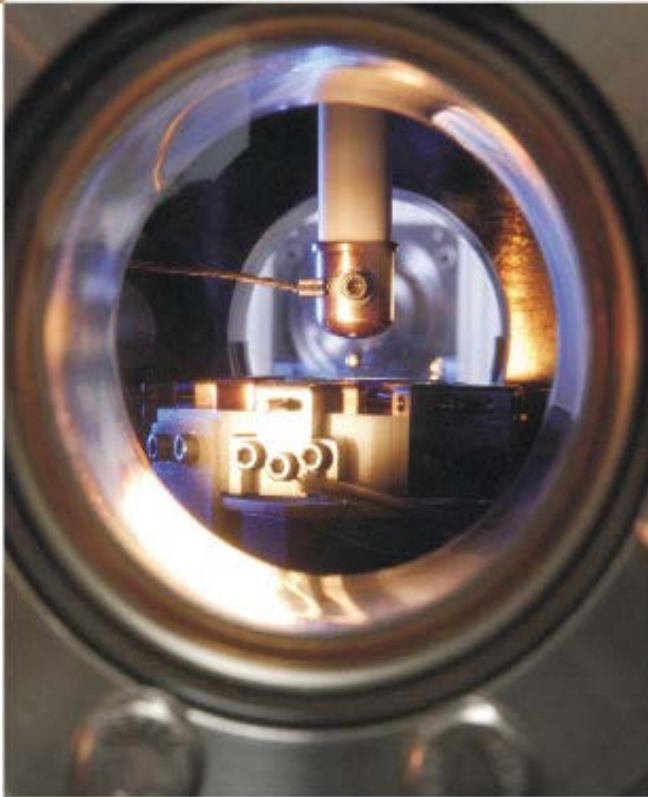
http://ec.europa.eu/research/industrial_technologies/articles/article_2288_en.html

Difficult to control rotation of sample,
Gas-sample reaction

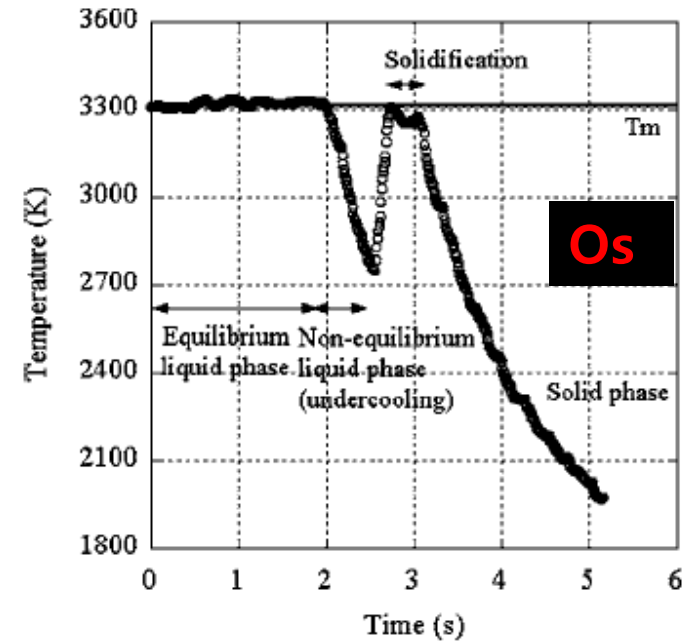
Electrostatic Levitation (NASA, MSFC (Huntsville))

JULY
2003

PHYSICS TODAY



Targeting molten metals

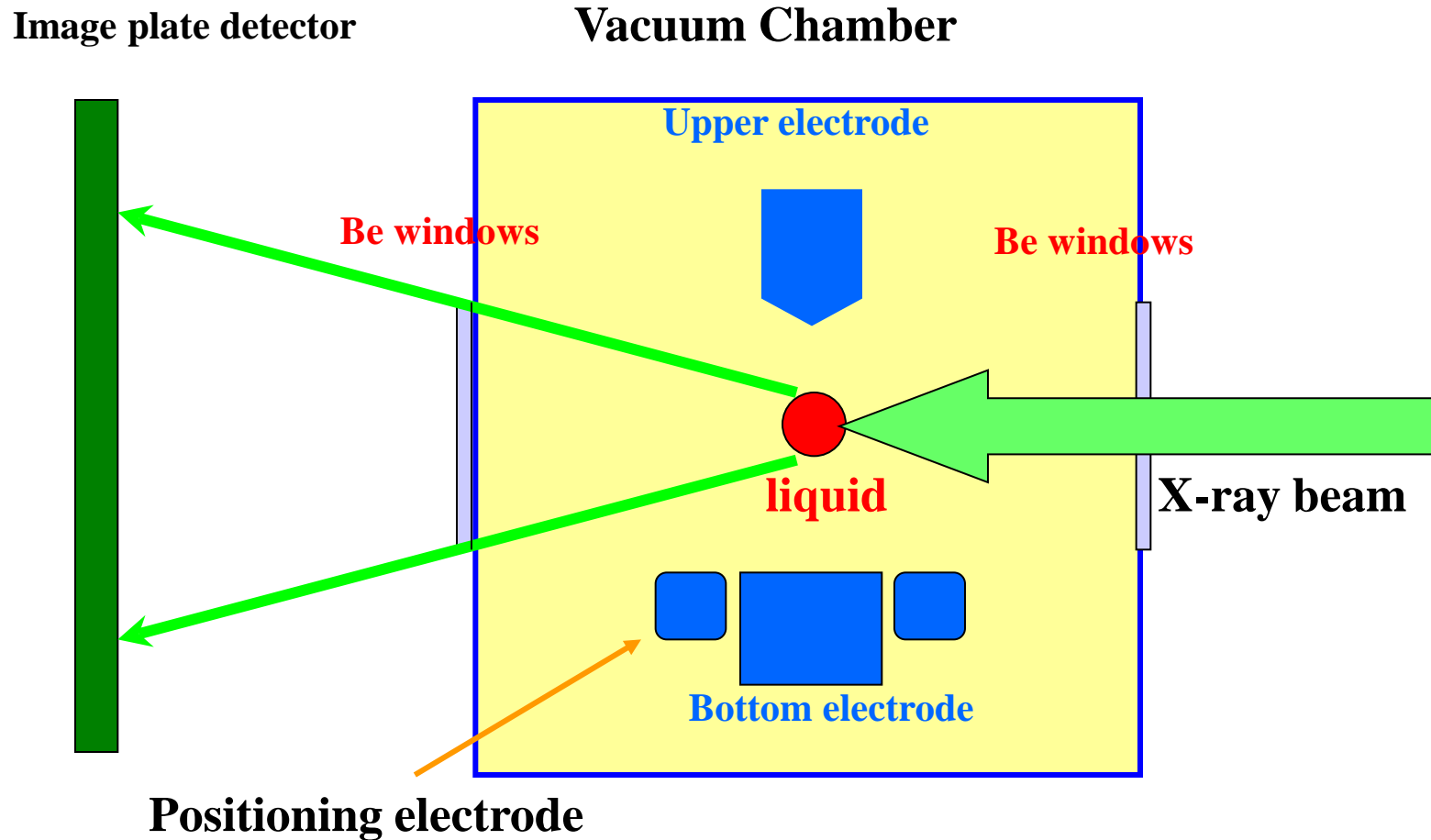


P.-F. Paradisa, et al., JOURNAL OF APPLIED PHYSICS 100, 103523 2006

- Specific heat,
- Emissivity,
- Density,
- Viscosity,
- Surface tension...

Physics Today, v56, p22, July 2003

Beam-ESL; High-energy synchrotron x-ray scattering experiment



Sample size : 30-100 mg,
Temperature : 300-2200 K,

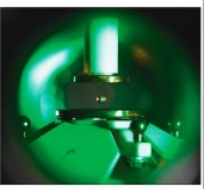
X-ray Energy : 125 keV,
Vacuum : 10^{-7} torr

Wavelength : 0.0988 \AA
Exposure time : 1 sec.

인공우주장치로 '건담 합금'도 가능할까

“직접 만들 한 방에 꺾을 수 있는 강력한 탄환을 만든다. 노르북이나 휴대전화의 배터리를 사용시간을 수십 배 늘린다.” 지구에서 직접으로 보내는 물질로는 이런 소재를 개발하기 어렵다. 지금까지 몇 배나 단단한 초강속, 전기저장량을 30배나 늘린 물질이 있어야 하기 때문이다. 과학자들이 대신 찾은 곳은 우주 공간이다. 정확히 말하면 진공 우주물 본론 ‘인공 우주’다.

● 국내 첫 인공우주실험실 국내에서 첫 인공우주실험실이 등장했다. 한국표준과학연구원 연구진이 4억 원에 달하는 인공우주 실험실 ‘인공우주 실험실’을 국내 최초로 개발했다. 우주공간과 비슷한 진공과 무중력의 공간을 만드는 장치다. 세계에서는 흔해져서 개발됐다. 진공을 만드는 것은 어렵지 않다. 펌프를 켜고 내면 된다. 그럴수록 공기를 빼내면 된다. 이 공기는 지구의 중력에 의해 땅에 떨어진다. 진공을 유지하는 데 성공했다. 공기와 진공이 상충하면서 장치 안에 공기를 넣으면 우주인처럼 공중에 뜰 수 있다. 50kg의 사람이면 0.000001바의 진공적인 환경을 필요로 한다. 이 장치는 고층에 레이저로 공중에 뜬 물질을 실험 4000시간까지 개발할 수 있다. 이공우 연구원은 “정확한 기체 원소 실험을 할 수 있게 된다. 더 깊이 있고 다양한 곳으로 뛰어나거나 지구 상에서보다 10배 이상 실험을 할 수 있다. 물체에 실을 때와 같은 미세하게 조절한 장치를 완성했다”고 말했다.



공중부양 실험을 위한 진공을 이룬 채 물체를 공중에 띄울 수 있다. 물체의 높이를 공중에서 띄우는 모습. 사진: 한국표준과학연구원 제공



세계 8번째 ‘진공-무중력 공간’ 구현 우주와 환경 유사해 각종 실험 가능
인공우주는 금속-저장량 늘린 배터리를 등 ‘연화성과 같은’ 신소재 개발 기대

장 큰 걸림돌 ‘배터리’가 문제일 전망이다. 공중이 있는 지구에서 실험할 때는 ‘이런저런’ 물체가 떨어진다. 무게는 공중에 띄워야 한다. 다른 물질이 떨어지거나 떨어지지 않게 해야 한다. 이공우 연구원은 “정확한 기체 원소 실험을 할 수 있게 된다. 더 깊이 있고 다양한 곳으로 뛰어나거나 지구 상에서보다 10배 이상 실험을 할 수 있다. 물체에 실을 때와 같은 미세하게 조절한 장치를 완성했다”고 말했다.

● 핵이동 한 스페이스 신소재 1970년 처음 발명된 일본 ‘메타-리온’ 기술은 ‘건담’의 주인공 코보트와 닮은 점이 많다. ‘메타-리온’이라는 금속으로 만든다는 장점을 갖고 있다. 대략 30cm 크기의 ‘메타-리온’을 갖고 있고, 고온이나 방사능도 견디는 등, 우주에서는 지구보다 더 뛰어난 소재를 만들 수 있을 것이라는 상상력이 바탕이 된 것이다. 과학자들은 인공우주실험실에서도 ‘연화성과 같은’ 신소재를 만들 수 있을 것으로 기대한다.

1970년 소개된 만화영화 ‘건담’의 주인공 코보트 300-78-2. 이 물체는 달에서 실험할 수 있는 핵이동 ‘메타-리온’으로 만든다. 원자력 연구원 제공

지는 이유는 물질이 결정 형태로 굳어 있기 때문이다. 비결정 금속은 결정을 이루지 않아 잘 깨지지 않는다. 이 기술은 핵의 정제와 정제, 재료정제 분야, 고온의 배터리를 더 다양한 충전물질을 만드는 데 응용되고 있다. 일본 우주항공연구개발기구(JAXA) 연구진은 2006년 비슷한 실험장치를 이용해 전기저장량을 30배나 늘린 배터리를 실험할 것을 개발했다.

이런 인공우주도 진짜 우주만큼 할 수는 없다. 그래서 국제우주정거장(ISS)을 거대한 우주실험실로 활용하는 방법은 모색되고 있다. NASA는 이미 우주정거장에 소재 개발을 위한 실험장치를 운영해오고 있으며 일본도 곧 비슷한 장치를 설치할 예정이다. 이공우 연구원은 “처음에는 국내에 이 장치를 어느 시설이 있어도 실험할 수 있는 일을 한다”는 냉소도 말했다. “앞으로 이 장치를 활용해 나노미터 크기의 초미세 반도체 소재를 연구할 계획”이라고 말했다.

대한=전승환 동아사이언스 기자 ehnkob@ncc.com



'공중부양 측정' 성공



우주환경 구현...공중부양장치 개발

한국표준과학연구원

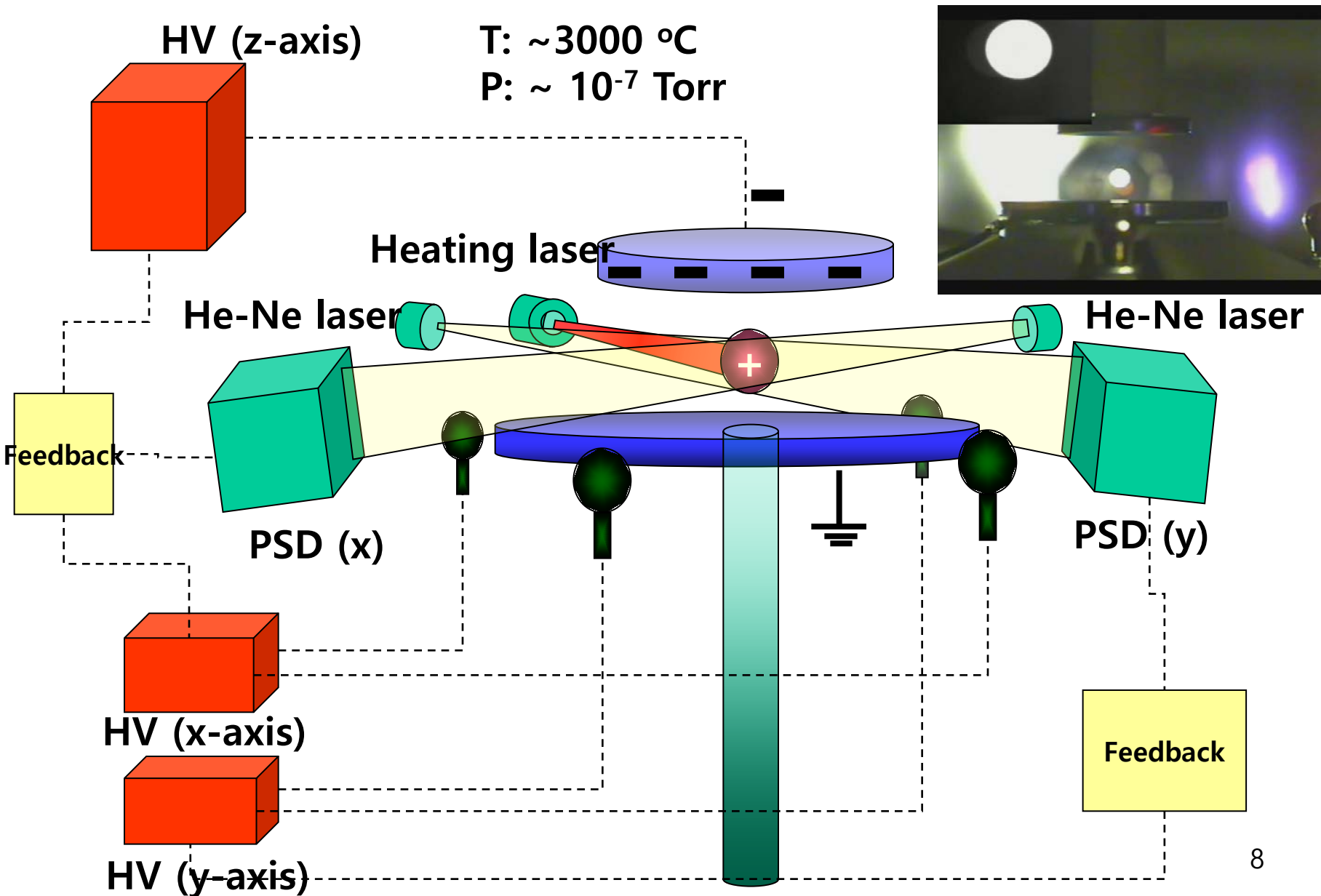


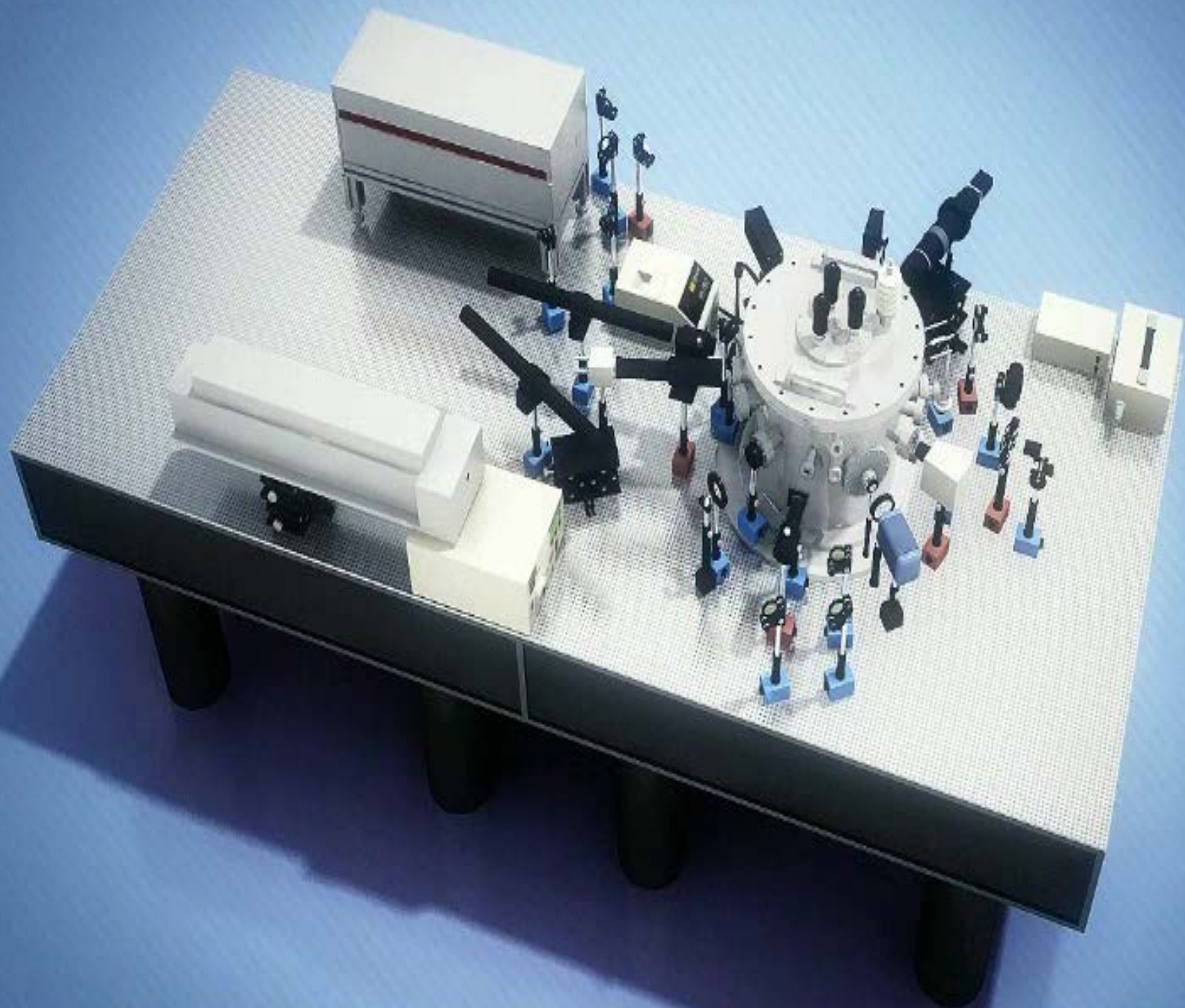
공중 부양 장치

무산 0.3%

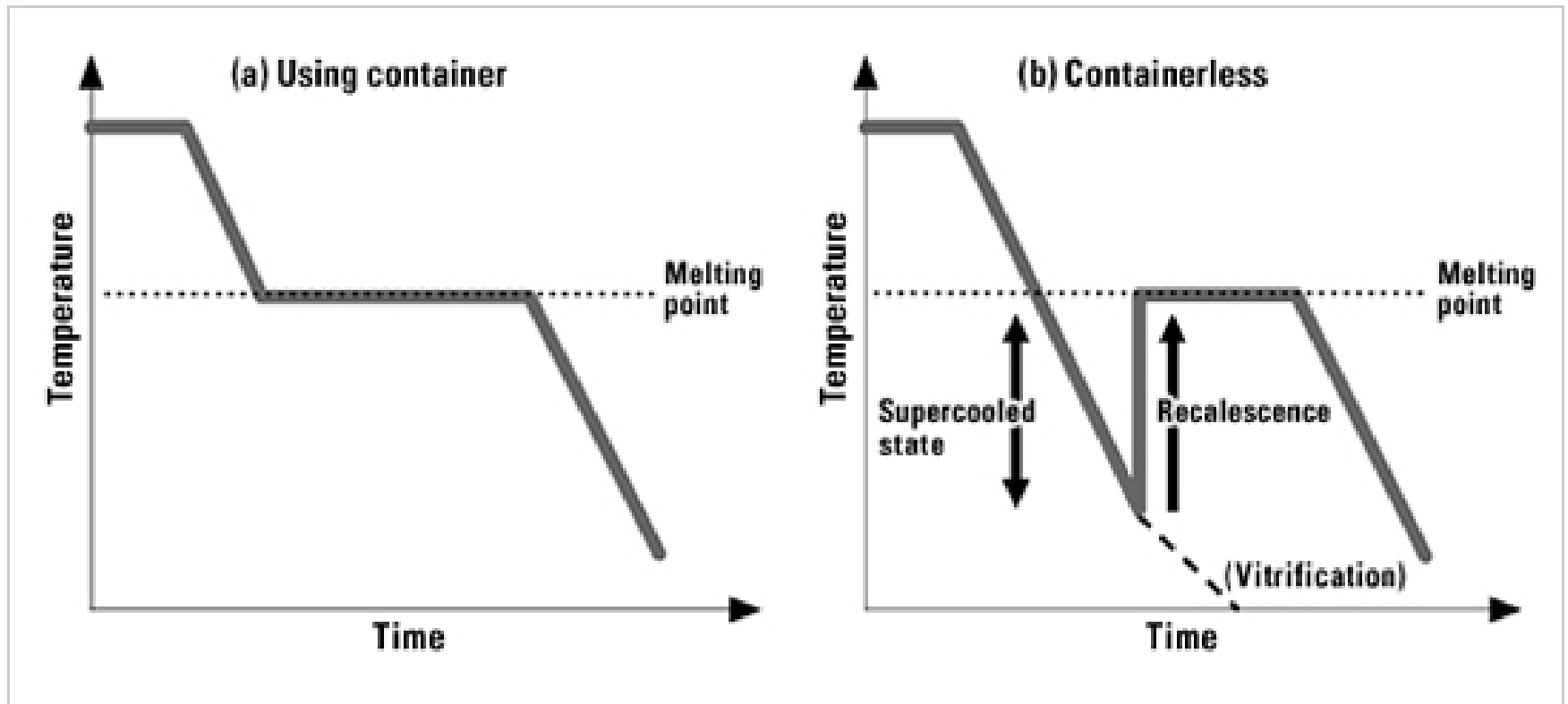
Electrostatic Levitation in KRISS

Containerless equipment: close to homogeneous nucleation

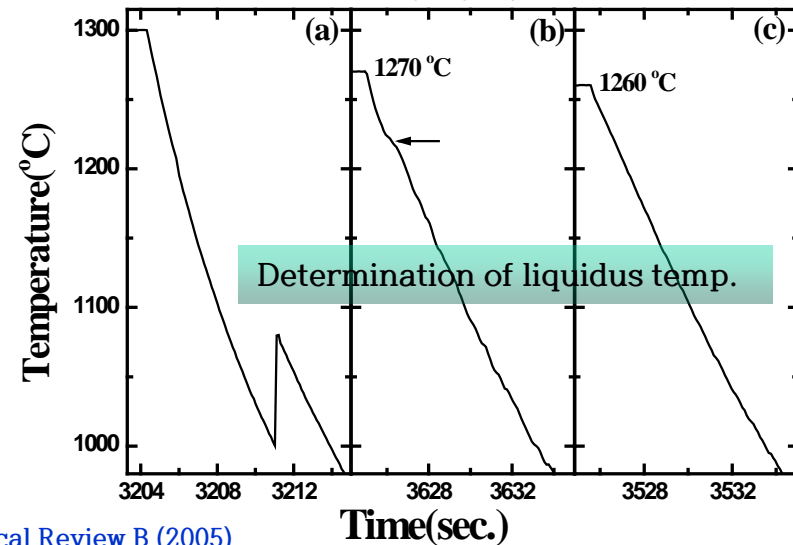
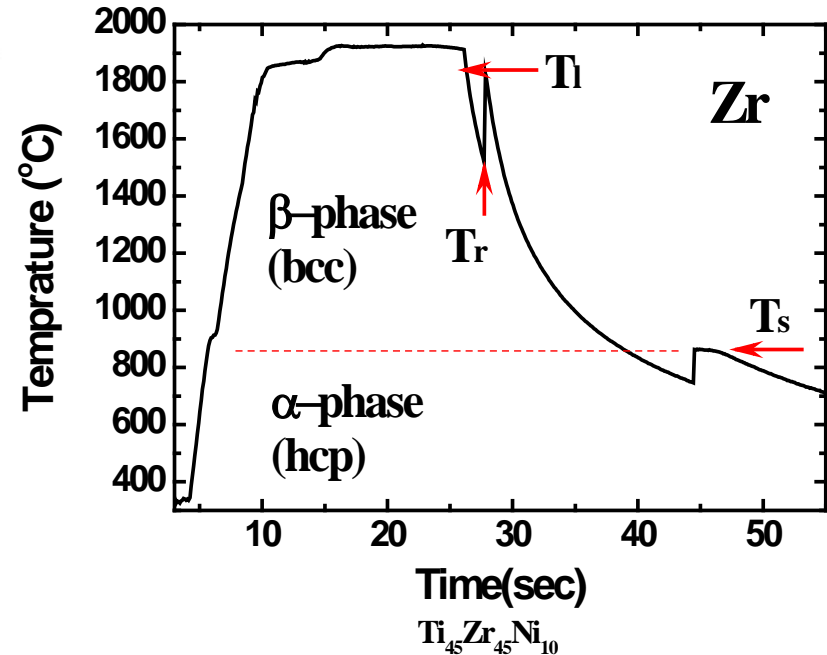
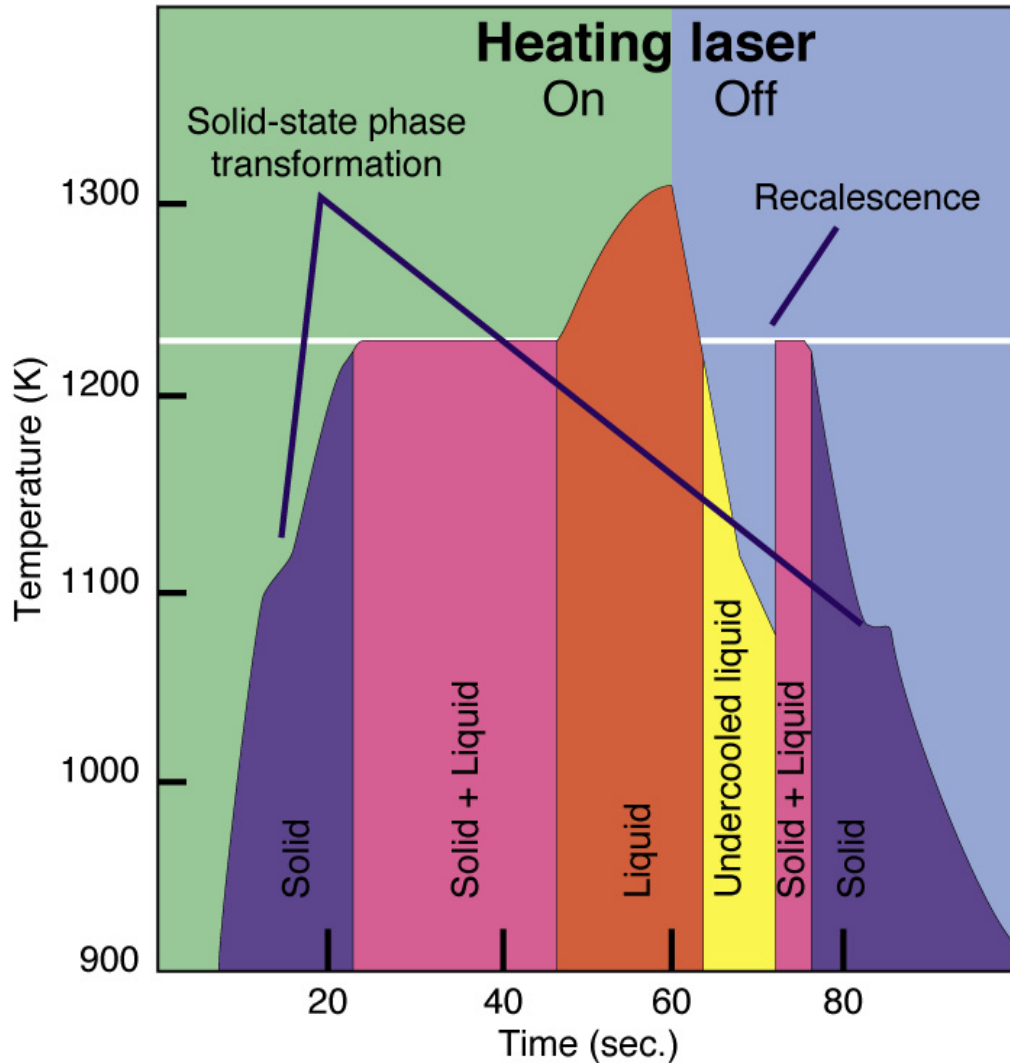


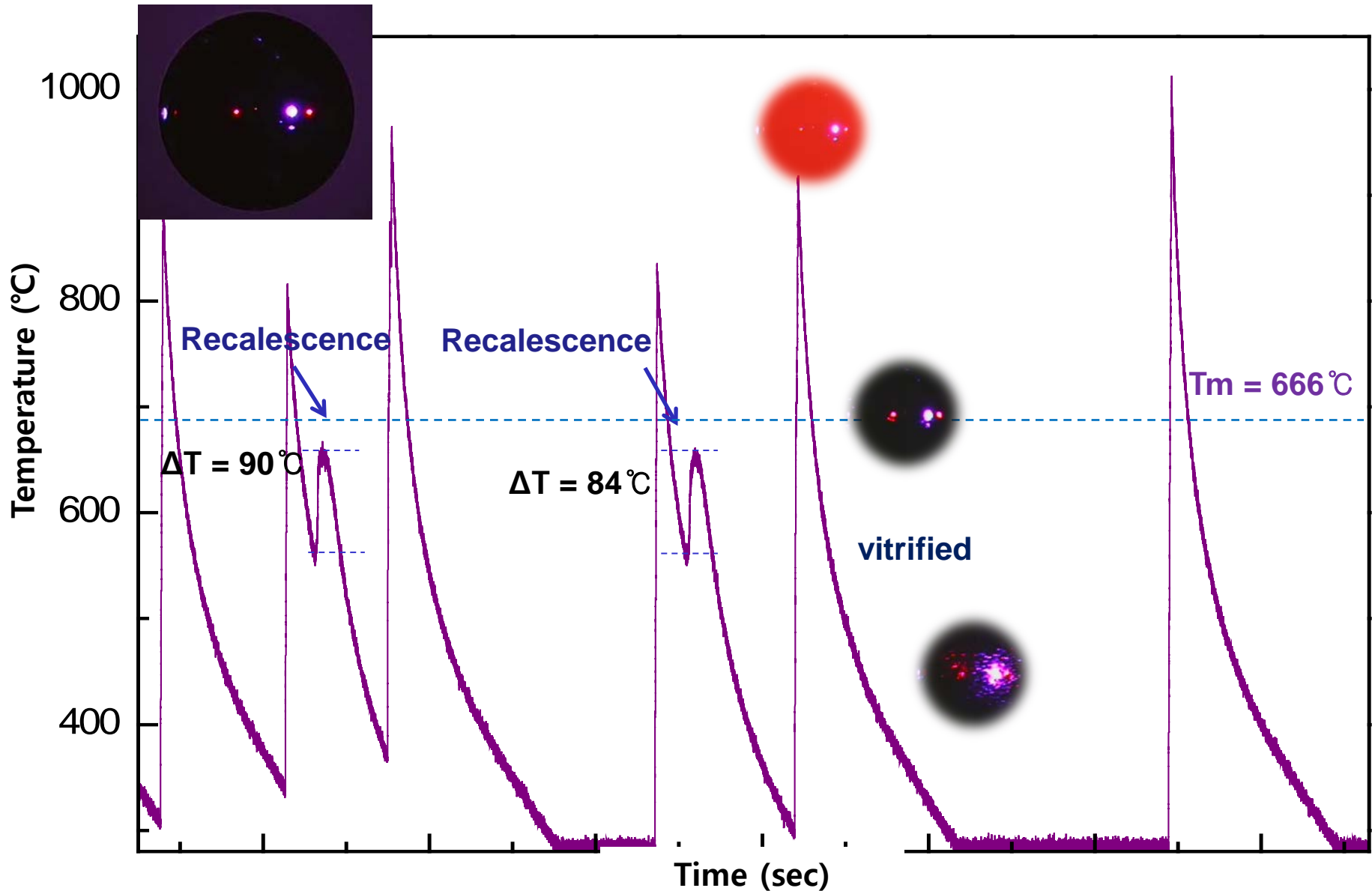


- ▶ No solid containers, No impurities from container
 - No heterogeneous nucleation site
 - Extremely large supercooling can be obtained ($\sim 100\text{ }^{\circ}\text{C}$), clear recalescence
 - Metallic glass can be formed through free cooling

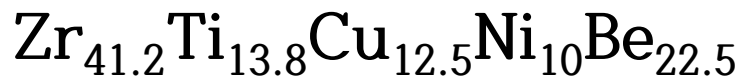


Melting and Freezing Using ESL

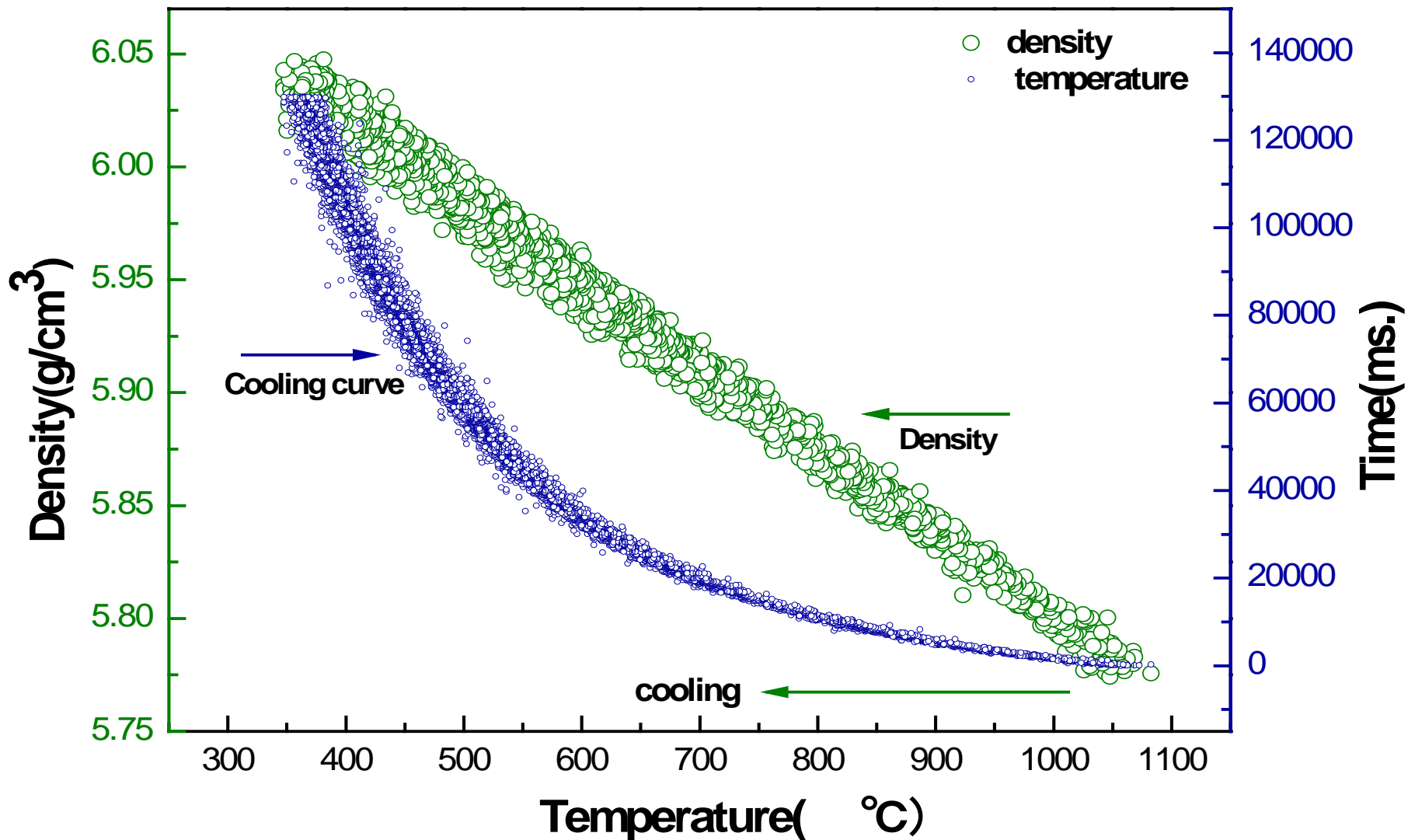




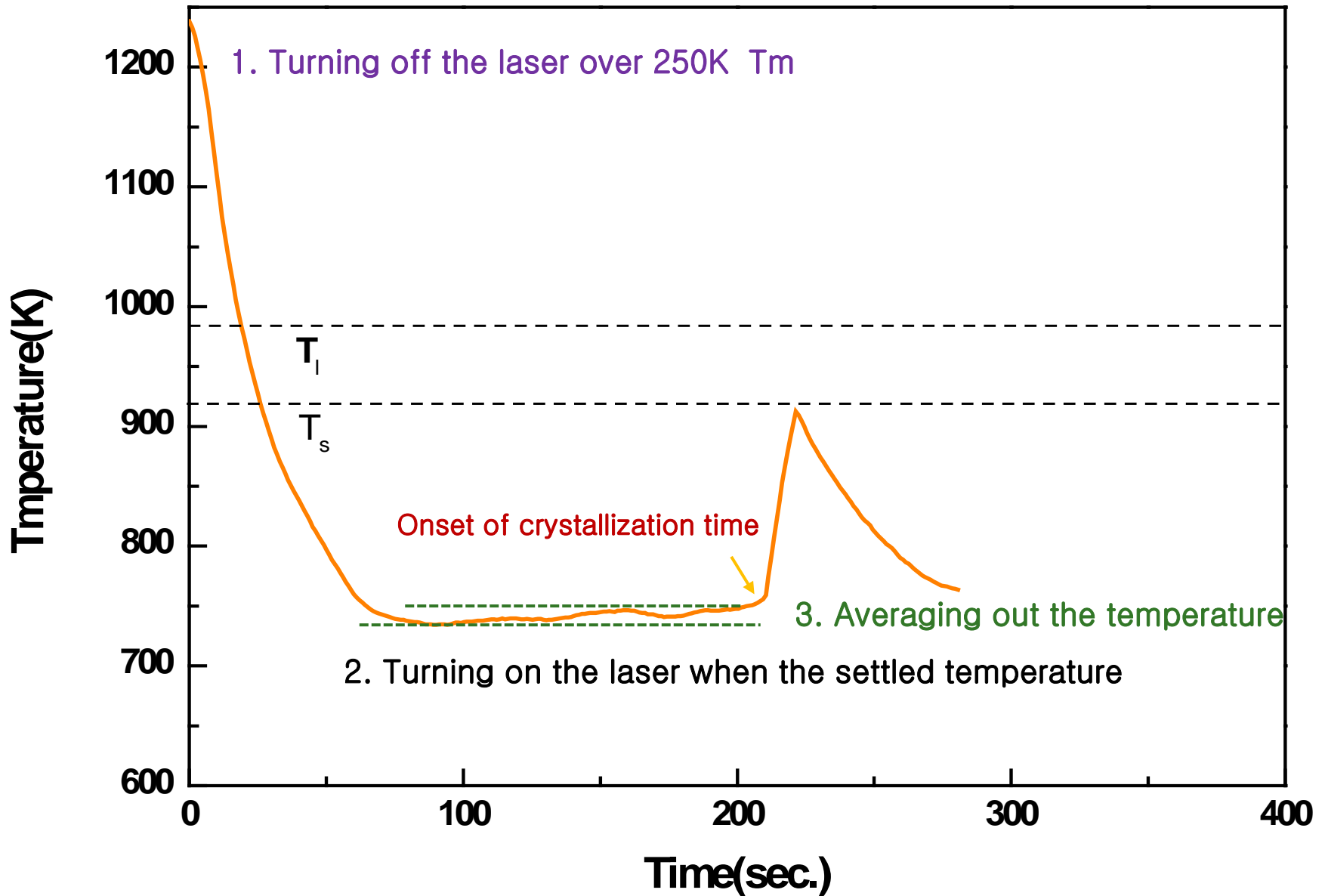
Cooling curve and density temperature profiles of



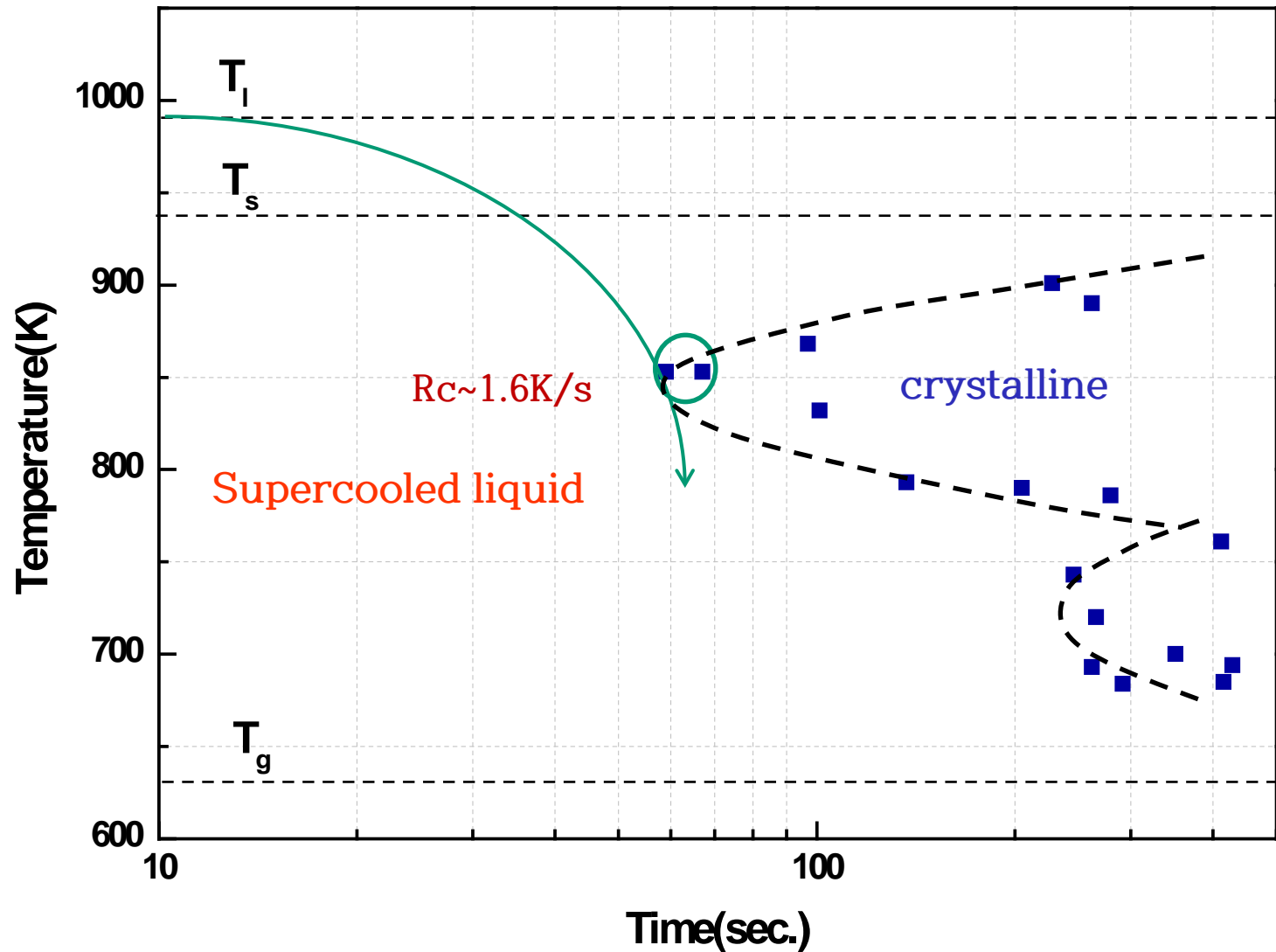
- Volume : CCD camera / Temperature measurement : pyrometer



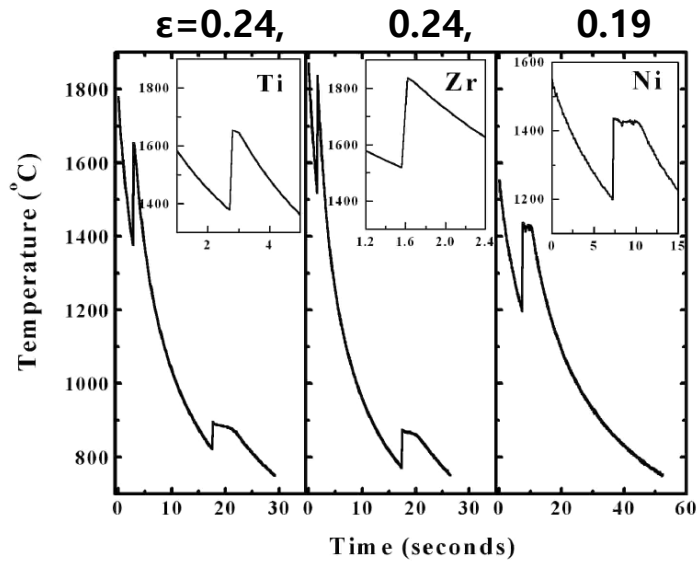
Measurement of TTT diagram - $Zr_{41.2}Ti_{13.8}Cu_{12.5}Ni_{10}Be_{22.5}$



Measurement of TTT diagram - $\text{Zr}_{41.2}\text{Ti}_{13.8}\text{Cu}_{12.5}\text{Ni}_{10}\text{Be}_{22.5}$



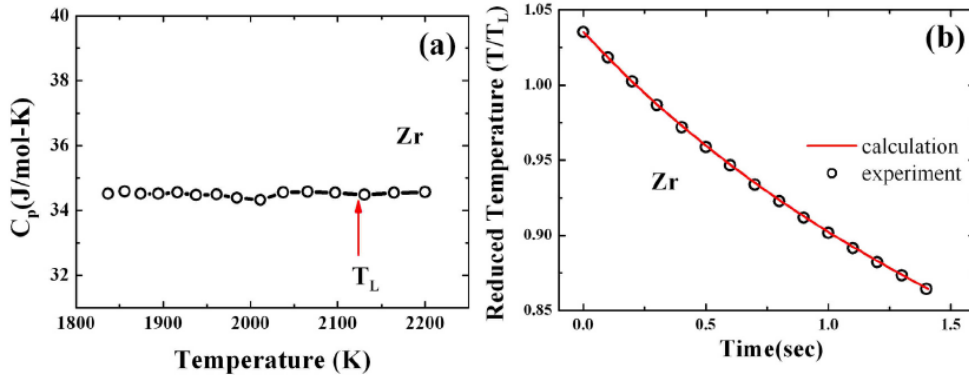
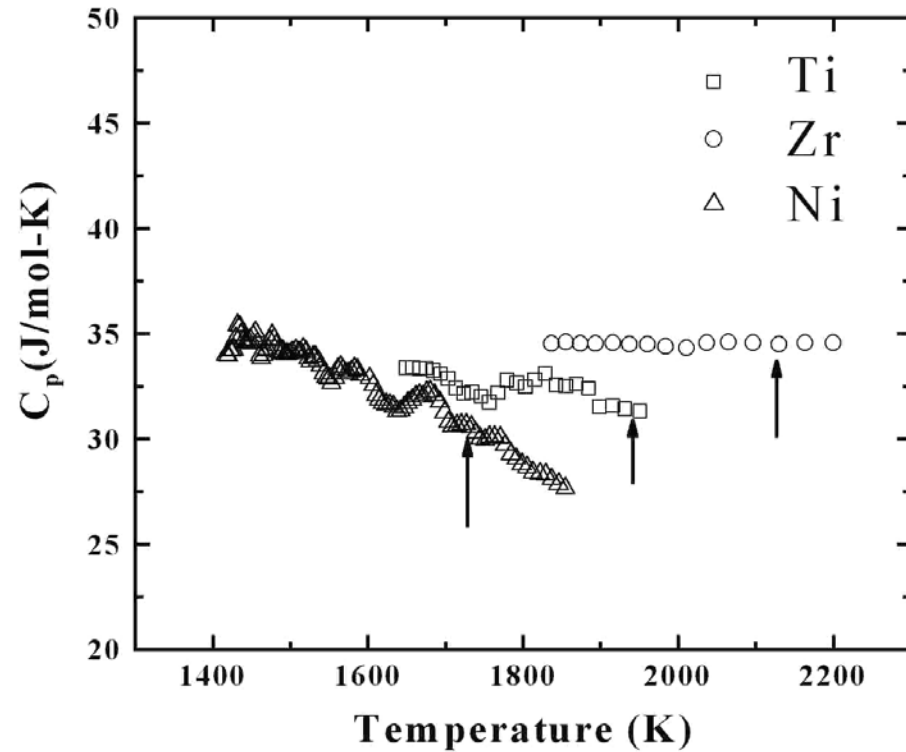
Specific heat capacity



$$mC_p \left(\frac{dT}{dt} \right) = 4\pi r^2 \sigma \epsilon (T^4 - T_s^4)$$

Specific heat of pure elements

Figure 6.1. Free cooling curves of Ti, Zr and Ni, measured in the ESL. Inset figures show recalescences at high temperature.

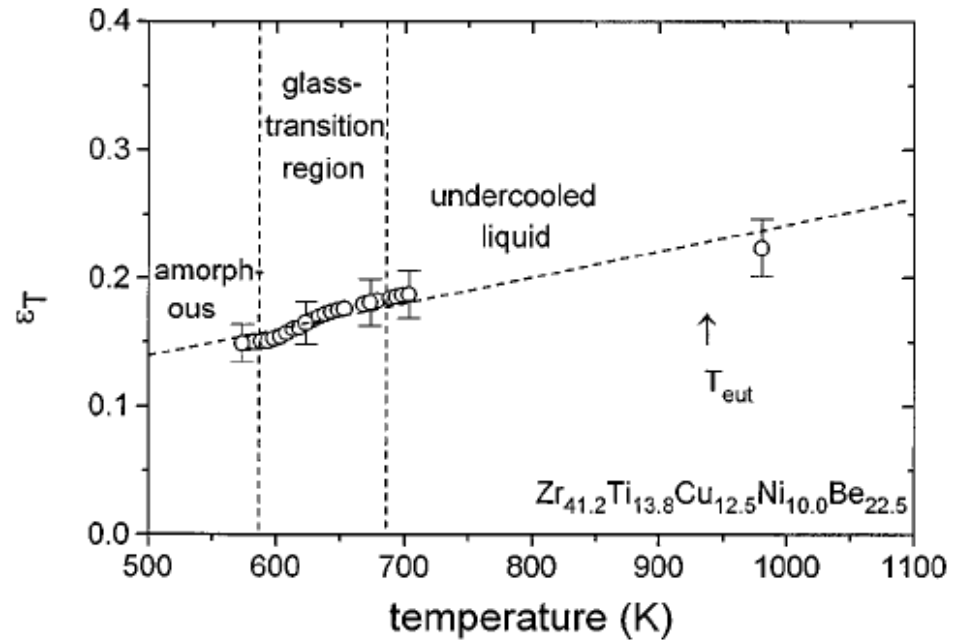
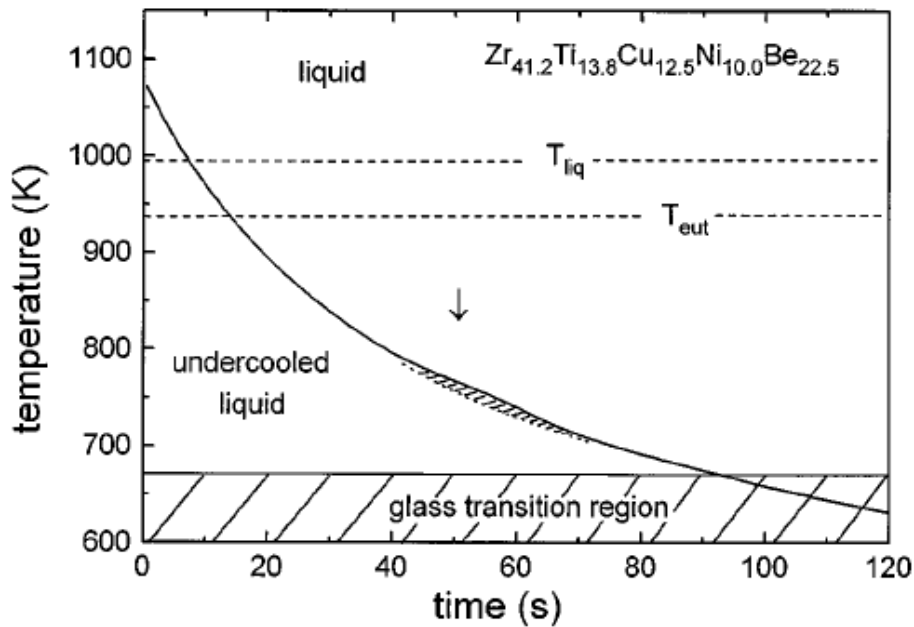


$$t = \frac{mC_p}{4\sigma\epsilon AT_r^3} \left(\ln \left(\frac{(T + T_r)(T_i - T_r)}{(T - T_r)(T_i + T_r)} \right) + 2 \tan^{-1} \left(\frac{T}{T_r} \right) - 2 \tan^{-1} \left(\frac{T_i}{T_r} \right) \right)$$

Emissivity

$$mC_p \left(\frac{dT}{dt} \right) + Power = 4\pi\sigma\epsilon(T^4 - T_o^4)$$

$$\text{Steady condition, } \left(\frac{dT}{dt} \right) = 0$$

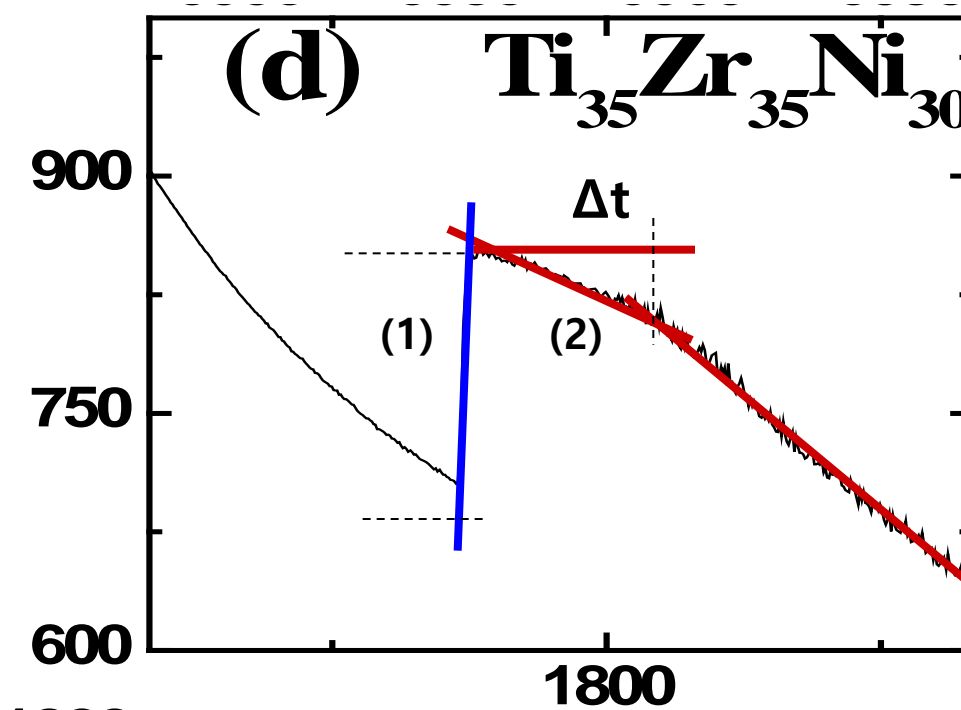


Kim, et.al., APL v68 (1995)

Fusion Enthalpy

$$\Delta H_f = (1) + (2) = C_p \Delta T_r + (4\pi\sigma\epsilon(T_P^4 - T_o^4)\Delta t - 4\pi\sigma\epsilon(T_P^4 - T_{end}^4)\Delta t)$$

Crystallization



Viscosity & Surface Tension: Oscillation

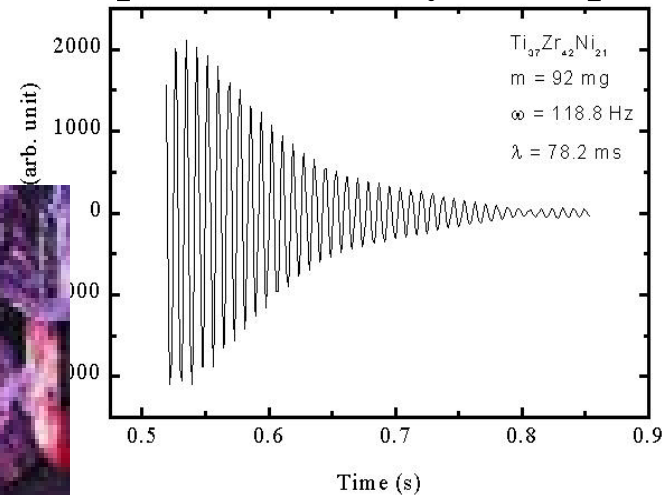
Snapshot of surface oscillations in a Ni droplet



Caution: No distortion from spherical symmetry
allowed multiple modes will be excited

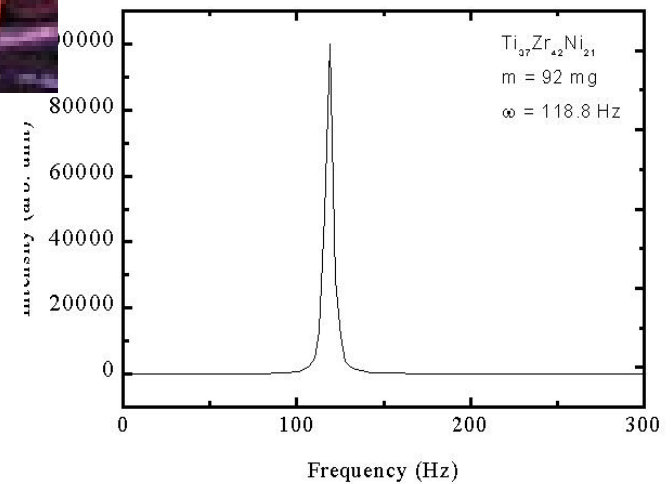
Damped oscillations

Exponential decay of amplitude



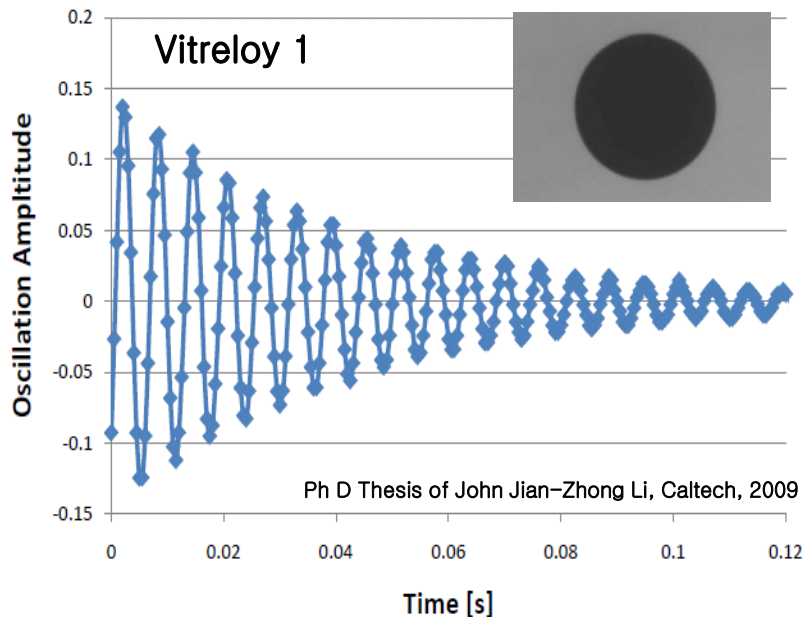
Power Spectrum

Single frequency



► Measurement of thermophysical properties

- Volume : CCD camera / Temperature measurement : pyrometer
- Surface tension & Viscosity : oscillating the sample by with a pulse of AC voltage



$$\omega_2^2 = \frac{8\gamma}{\rho r_0^3}$$

Oscillation frequency

Surface tension

Radius when melt adopts a spherical shape

Density

$$\frac{1}{\tau_2} = \frac{5\eta}{\rho r_0^2}$$

Decay time constant

Viscosity

- Specific heat & total hemispherical emissivity : $\frac{m}{M} C_P \frac{dT}{dt} = -\sigma_{SB} \epsilon_T A (T^4 - T_S^4)$
- Time- temperature-transformation curve : isothermal treatment

Oscillating drop 방식으로 고온에서 metal의 점도/ 표면 장력 측정 가능

Oscillating drop

- Induce surface oscillations in a levitating liquid droplet of radius 'R₀', mass 'm'
- Measure the frequency of oscillation (ω)
- Measure the damping constant (λ)

- Damped resonant oscillations:

$$R=R_0(1 + \delta \cos(\omega t)e^{-\lambda t})$$

Rayleigh (1879)

- Resonant frequency determined by surface tension:

$$\omega_l = \sqrt{\frac{l(l-1)(l+2)\gamma}{\rho R_o^3}}$$

- Damping determined by viscosity:

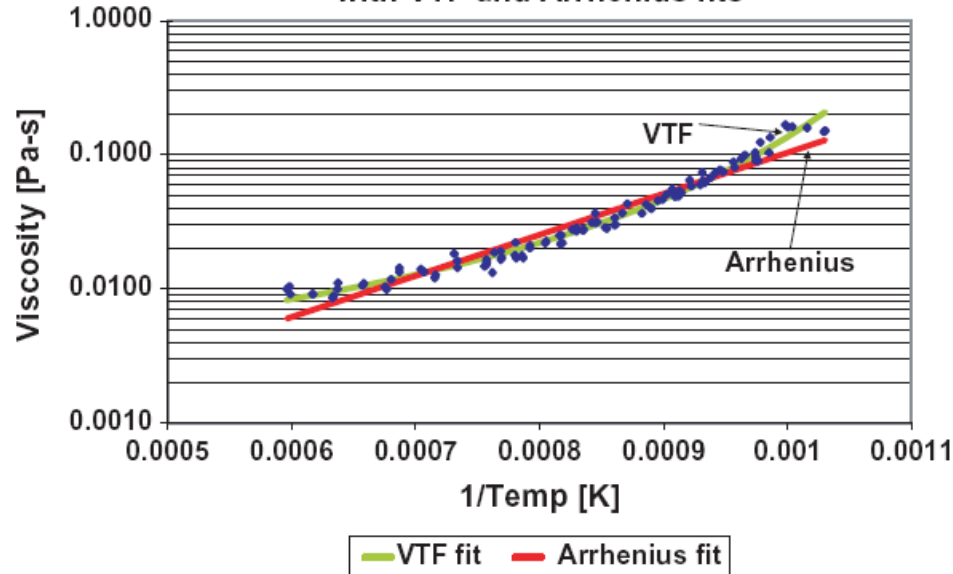
$$\lambda_l = \frac{(l-1)(2l+1)\eta}{\rho R_o^2}$$

Lamb (1881)

High Temp. Viscosity

JOURNAL OF APPLIED PHYSICS 100, 103523 2006

Ti_{39.5}Zr_{39.5}Ni₂₁ log Viscosity plot
with VTF and Arrhenius fits



Hyers, et.al., Philosophical Magazine Vol. 86, 2006(341-347)

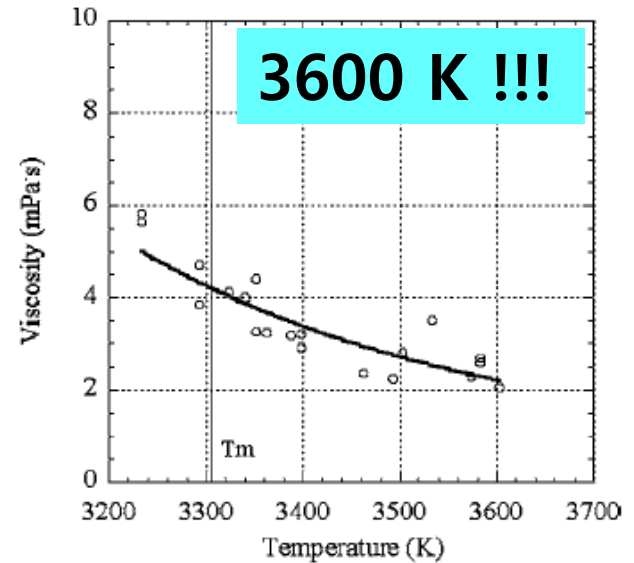


FIG. 5. Viscosity of equilibrium and nonequilibrium liquid Os as a function of temperature

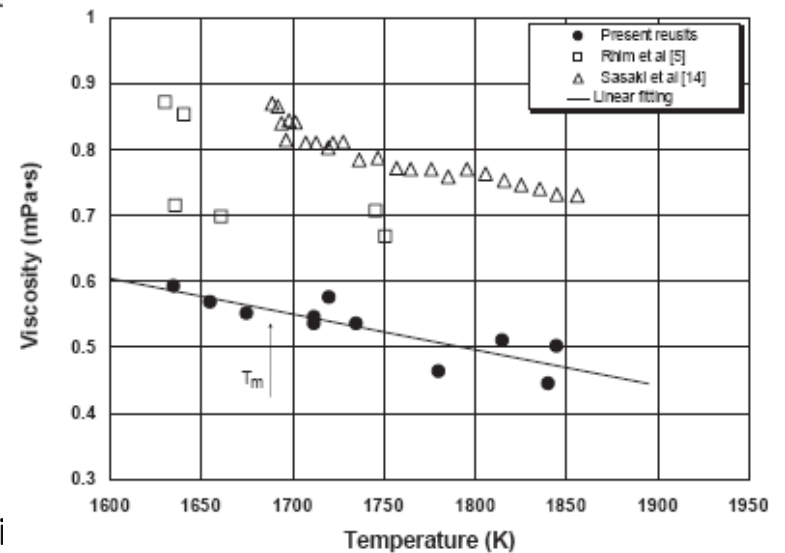


Fig. 8. Viscosity of molten silicon as a function of temperature.

High Temp. Surface Tension

Hyers, et.al., Philosophical Magazine Vol. 86, 2006(341–347)

JOURNAL OF APPLIED PHYSICS 100, 103523 2006

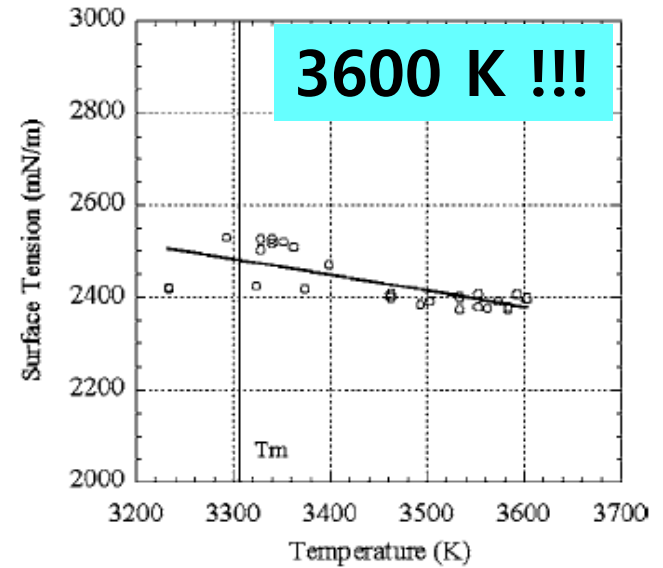
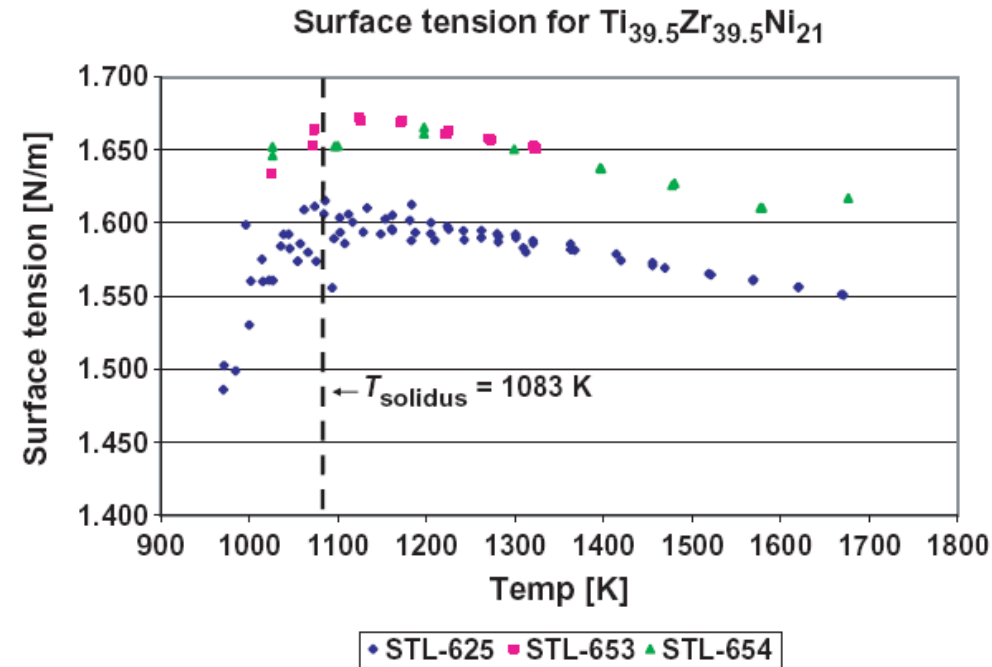


FIG. 4. Surface tension of equilibrium and nonequilibrium liquid Os as a function of temperature.

Density

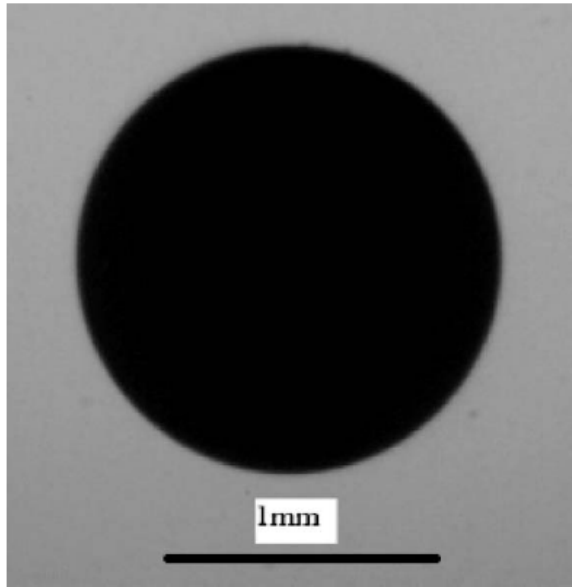


Fig. 2. A typical side view of a levitated molten silicon from which the density and the specific volume could be extracted.

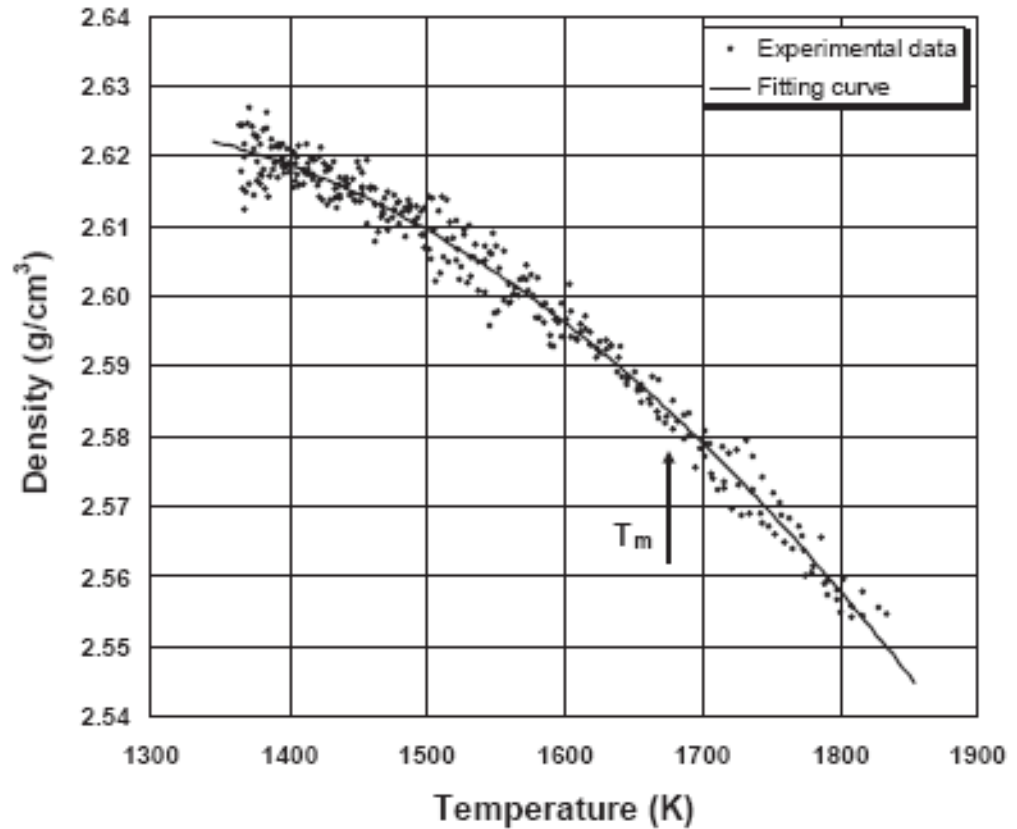


Fig. 4. Temperature dependent density of molten silicon.

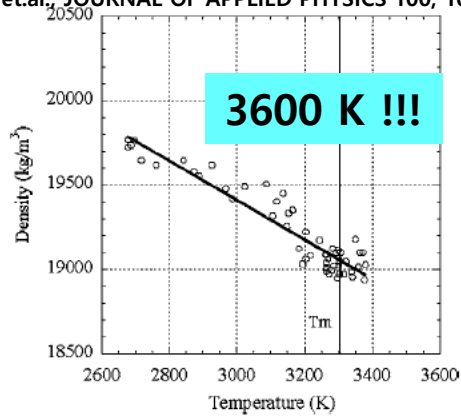


FIG. 3. Density of equilibrium and nonequilibrium liquid Os as a function of temperature.

Crystallization: Undercooling of Os

Paradis, et.al., JOURNAL OF APPLIED PHYSICS 100, 103523 2006

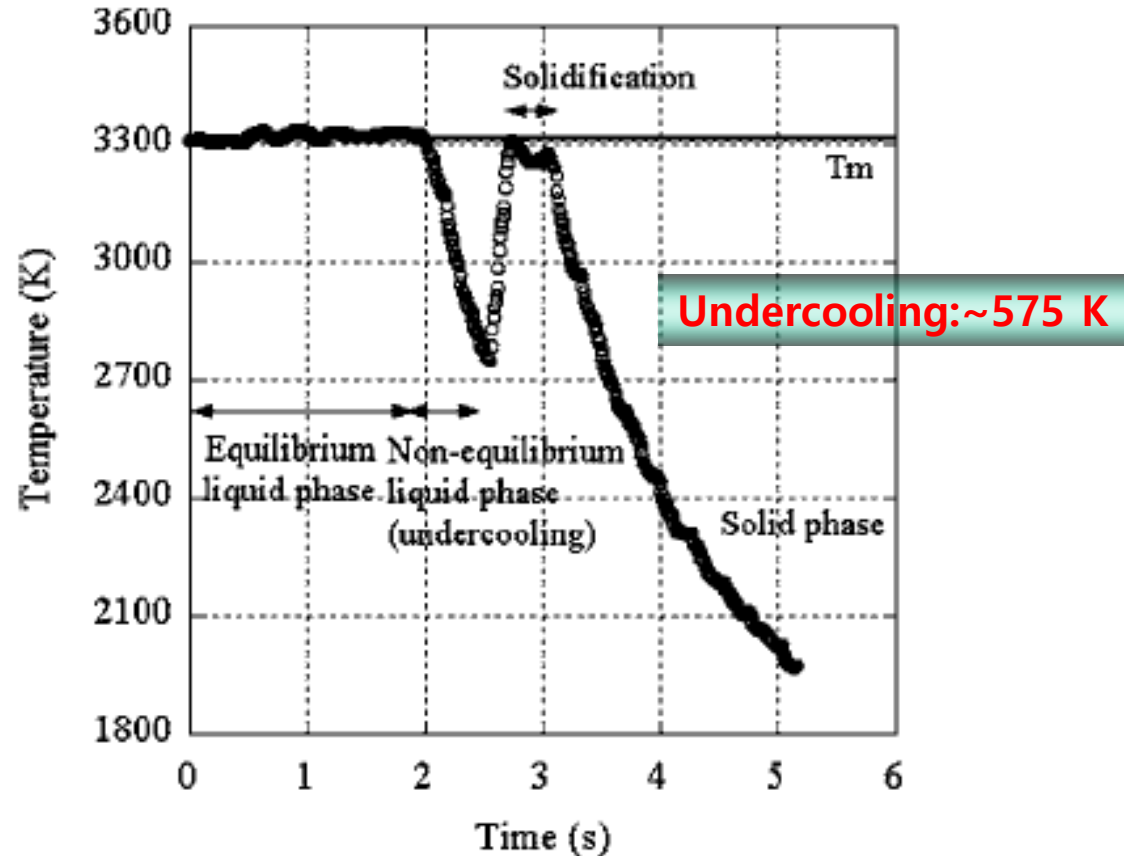
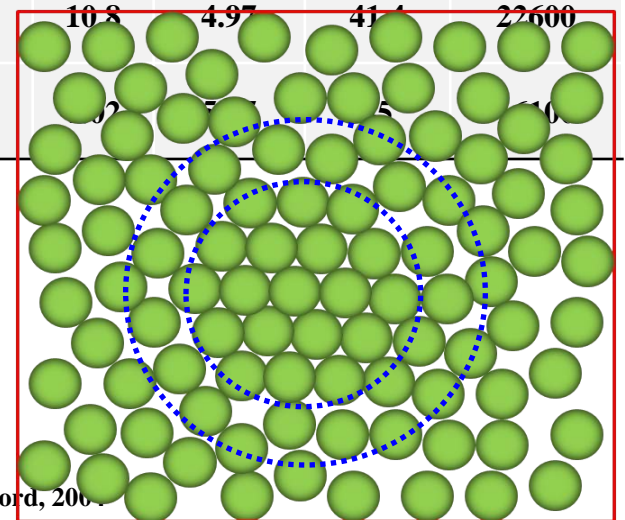


FIG. 2. Temperature history for an Os sample (diameter: ~ 1.76 mm; mass: 64.18 mg) showing heating slightly above the melting temperature (~ 3340 K), a radiative cooling rate of 1015 K/s, an undercooling of ~ 575 K, recalescence, and solidification.

Interfacial free energy of some elements

	$\frac{\Delta T_r}{\Delta T_{hyp}}$	σ ^[1]	σ	$\alpha = \frac{\sigma}{\Delta H_f}$	r^*	ρ	η_m ^[2]	C_p	ΔH_f ^[3]
	(K)	(J/m ²)	(J/m ²)		(nm)	(g/cm ³)	(10 ⁻³ Pa/s)	(J/mol·K)	(J/mol)
Ti	309 341	0.168	0.141 0.152	0.42 0.454	1.46 1.43	4.11	5.2	42.67	14550
Zr	332 345	0.158	0.154±0.009 0.159±0.010	0.410 0.423	1.54 1.52	6.08	4.67	42.5	19300
Hf	339 339	0.229	0.193±0.012	0.404	1.47	12.24	7.07	60.3	24070
Nb	443 563	0.262	0.258±0.016 0.303±0.024	0.394 0.462	1.33 1.23	7.63	4.94	52.0	29300
Rh	413 546	0.279	0.261±0.018 0.313±0.029	0.439 0.527	1.08 1.19				
Fe	195 357	0.269	0.158 0.228	0.33 0.478	1.45 1.15				

• Turnbull : $\alpha = 0.45$ for most metals



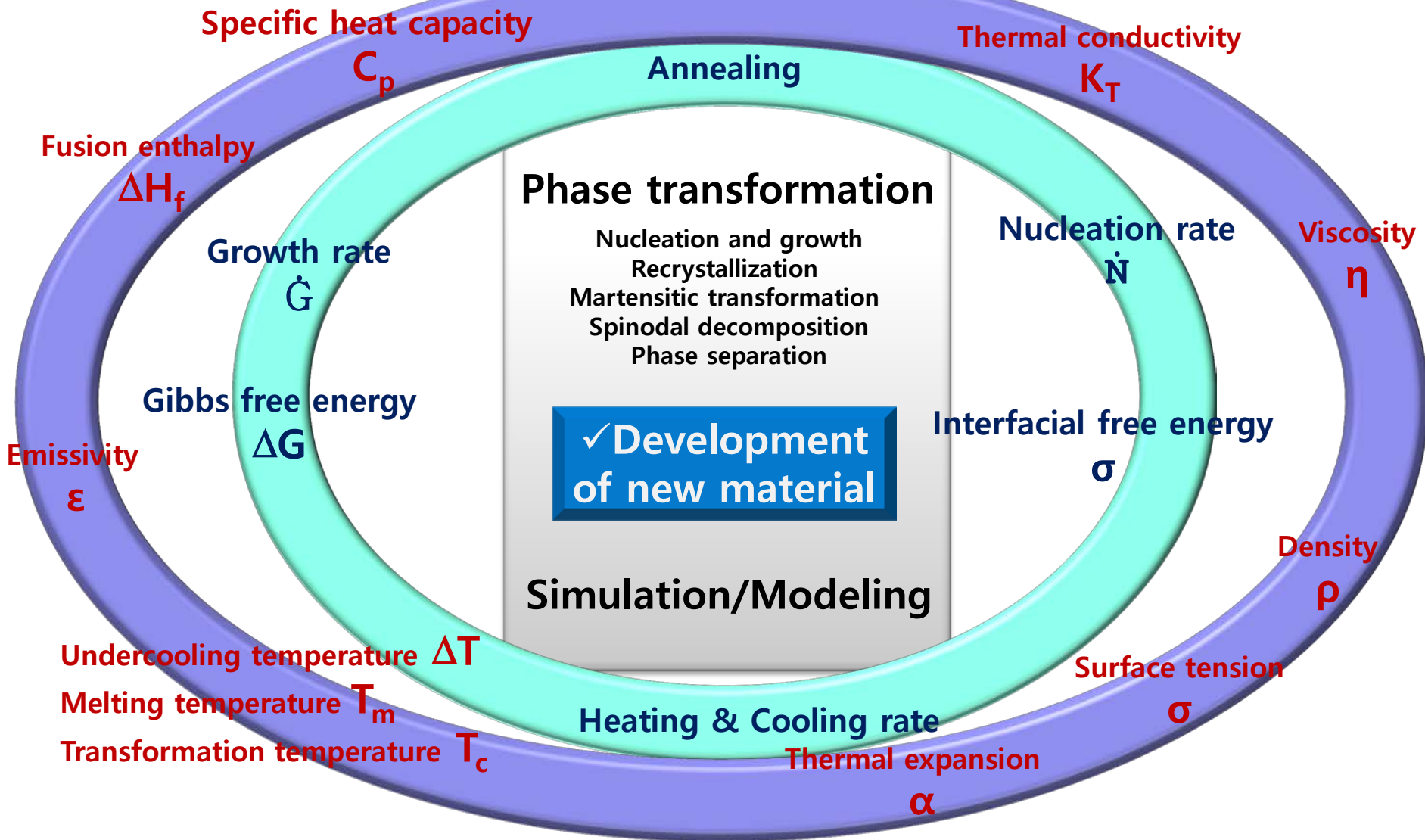
¹⁾ B. Vinet, L. Magnusson, H. Fredriksson, P. J. Desré, *J. Colloid Interf. Sci.* 255 (2002) 363

²⁾ T. Ishikawa, P.-F. Paradis, J. T. Okada, Y. Watanabe, *Meas. Sci. Technol.* 23 (2012) 025305

³⁾ W. F. Gale, T. C. Totemeier, in "Smithells Metals Reference Book", 8th ed. Butterworth-Heinemann, Oxford, 2000

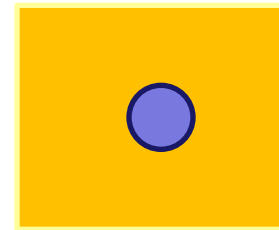
Development of extreme condition endurance materials

Thermophysical properties



Melting and Crystallization are Thermodynamic Transitions

Solidification: Liquid \rightarrow Solid



<Thermodynamic>

• Interfacial energy $\Rightarrow \Delta T_N$

Liquid

T_m Undercooled Liquid

Solid

No superheating required!

• Interfacial energy \Rightarrow No ΔT_N

$$\gamma_{SL} + \gamma_{LV} < \gamma_{SV}$$

vapor



Melting: Solid \leftarrow Liquid

Incentive Homework 2:

superheating이 일어나는 경우 정리 PPT 3 page 이내

Nucleation

* *Homogeneous Nucleation of crystal in supercooled liquid*

→ Well-defined by Turnbull and his coworker theoretically / experimentally.

* *Heterogeneous Nucleation*

→ detailed theory ~ less satisfactory

**Nucleation ~ a function of the temperature in liquids that are not in motion
but In practice, liquids are often exposed to dynamic conditions.**

< Two main type of dynamically stimulated nucleation >

1) completely metastable supercooled liquid containing no crystal

→ Nucleation by friction, ultrasonic vibration, pressure pulse , etc.

2) A phenomenon that the # of crystals is greatly increased by dynamic methods in solidifying liquid → It is difficult to conclude that it is not due to the fragmentation of pre-existing crystals.

* *Dynamically Stimulated Nucleation*

→ very poor understood

Chapter 4. microscopic Heat Flow Considerations

4.1 Qualitative Observation

Solidification: Liquid \longrightarrow **Solid**

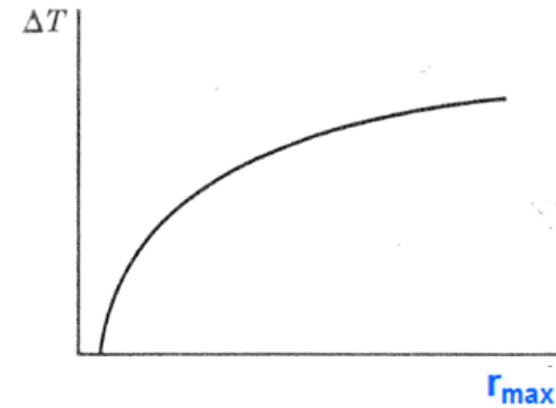
Presence of "Metastable supercooled liquid"

1) Atomic consideration

→ If it is curved, "escape angle" changes with curvature.

$$\therefore T_{E, \text{ small crystal}} < T_{E, \text{ large crystal}}$$

Thus, at any temperature below T_E , there is a radius of curvature at which the rates of melting and of freezing are equal. = **critical radius r^***



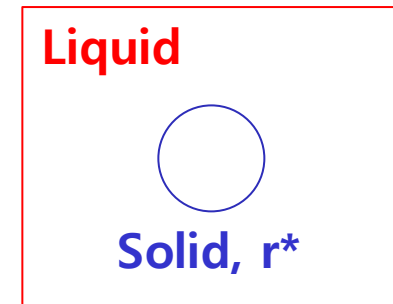
2) Thermodynamic treatment of equilibrium across a curved interface

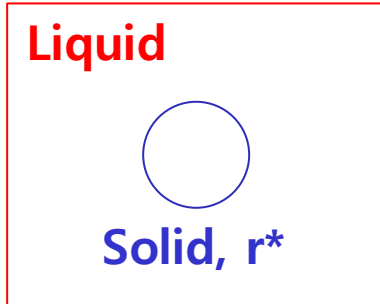
Extra pressure ΔP due to curvature

For incompressible solid,

$$\frac{L\Delta T}{T_E} = \frac{2\gamma_{SL}}{r^*}, \quad \text{or} \quad r^* = \frac{2\sigma T_E}{L\Delta T}$$

ΔH , ΔS : independent of temperature





$$T_E + \Delta T (-) \text{ 면, } R_M < R_F \rightarrow r \uparrow$$

$$\rightarrow T_E \rightarrow T_E'' \uparrow \rightarrow \Delta T (-) \uparrow \rightarrow R_M \ll R_F$$

For small departures from equilibrium, the rate is approximately proportional to the departure (ΔT); however, **the actual rate depends upon the crystallographic orientation of the interface.**

It should be emphasized that the foregoing remarks relate to the **actual temperature of the interface itself**; this may be different from the temperature of the liquid or solid at even a short distance from the interface because of the **“latent heat of fusion”** that is **generated at the interface during solidification or is absorbed there during melting.**

Who can explain the clear difference between two movies?

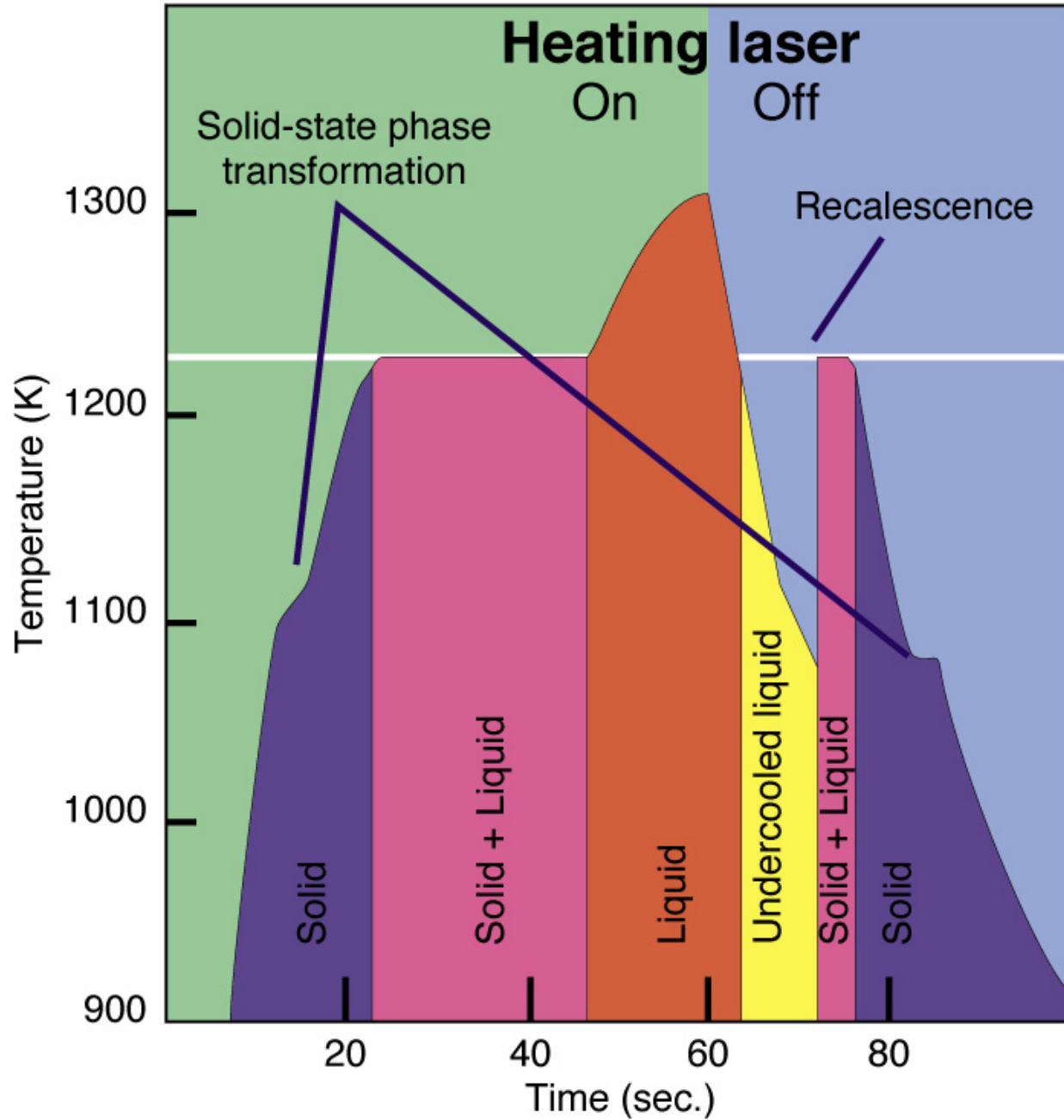


ZrCuAl alloy with non-purified Zr



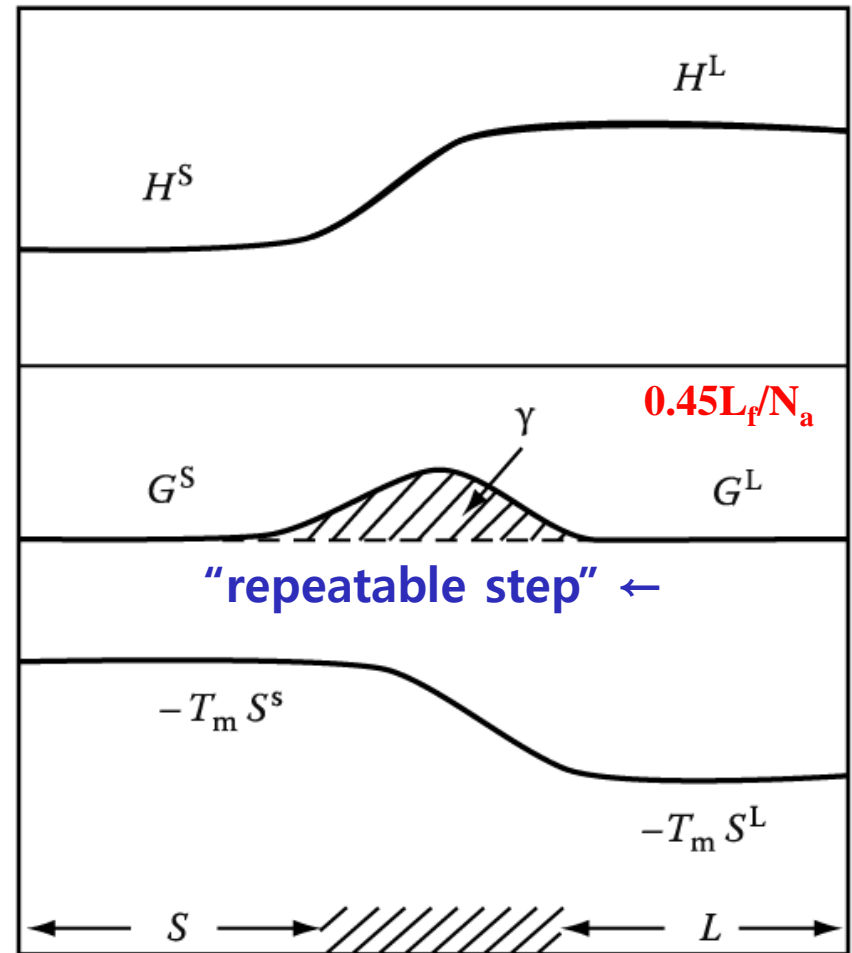
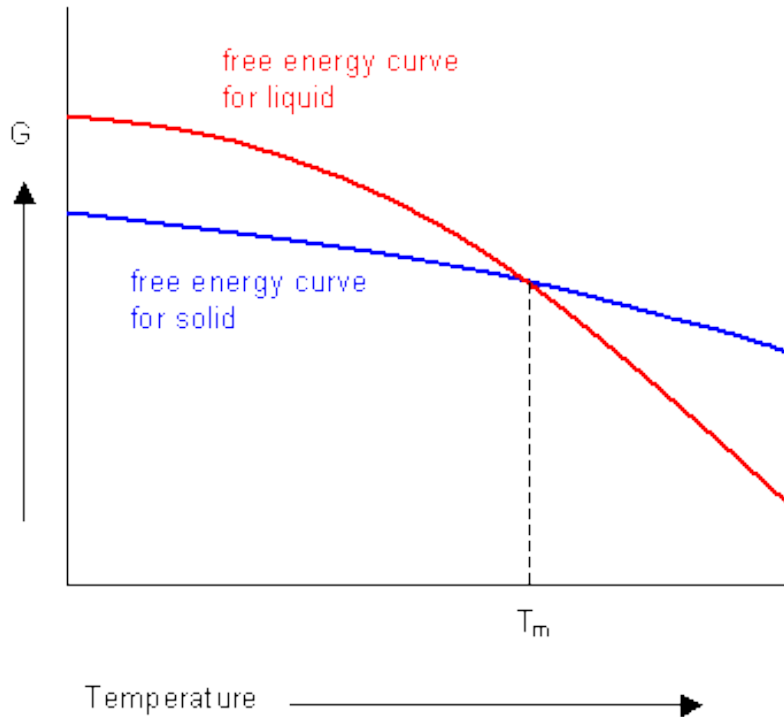
ZrCuAl alloy with purified Zr

Melting and Freezing Using ESL



* Broken bond model → calculation of the E of solid/ liquid interface

at equilibrium melting temp.



$\gamma_{SL} \approx 0.45 \gamma_b (= 0.15\gamma_{SV})$
for the most metals

$\gamma_{SV} > \gamma_{SL} + \gamma_{LV}$

Showing the origin of the solid/ liquid interfacial energy, γ 34

$T_{\text{interface}} < T_E \rightarrow \text{solidification} \uparrow \rightarrow \text{latent heat} \uparrow \rightarrow \Delta T \downarrow$

The “removal of latent heat” therefore controls the rate at which solidification can continue, and the interface temperature adjusts itself so that it corresponds to the rate of solidification determined by the externally imposed thermal conditions.

The local rate of growth at any point on the surface therefore depends on the thermal conditions and on the orientation of the surface, since this influences the relationship between temperature and rate of growth.

The interplay of the anisotropy of growth rate with the effects of the geometry of the surface on local heat flow is responsible for the very complicated morphology that may occur during solidification.