### Computer Aided Ship Design Part.3 Grillage Analysis of Midship Cargo Hold

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# **Summary**

*u*: Axial Displacement *G*: Shear Modulus *A*: Sectional Area *G*: Shear Modulus *E*: Young's Modulus *w*: Vertical Displacement  $\theta$ : Angle of Twist *l*: Length *J*: Polar Moment of Inertia *I*: Moment of Inertia



Beam Theory : Sign Convention, Deflection of Beam

Elasticity : Displacement, Strain, Stress, Force Equilibrium, Compatibility, Constitutive Equation



**Chapter 2. Element : Beam** 







### Element : Beam - Differential Eqn.

$$\sigma_{x} = \sigma \mathbf{i} \ , \ \mathbf{\epsilon}_{x} = \varepsilon \mathbf{i} \ , \ \mathbf{\theta} = \theta \mathbf{k}, \ \mathbf{y} = y \mathbf{j} \qquad \sigma = E\varepsilon$$

$$(\mathbf{p}) \quad \mathbf{\rho} \cdot d\theta = ds \implies \frac{d\theta}{ds} = \frac{1}{\rho}$$
(1) strain at  $y$  in x-direction :

 $\mathbf{\varepsilon}_{x} = \varepsilon \mathbf{i}$  $\varepsilon = \frac{(\rho - y) \cdot d\theta - \rho d\theta}{ds} = -y \frac{d\theta}{ds} = -\frac{y}{\rho} , ds: \text{initial length}$  $\rho$  $d\theta$ ds' neutra surface V ds y σ X \* neutral surface : Longitudinal lines on the lower part of the beam are elongated, whereas those on the upper part are shortened. Thus the lower part of the beam is in tension and the upper part is in compression. Somewhere between the top and bottom of the beam is a dxsurface in which longitudinal lines do not change in length. This surface is called neutral surface 2009 Fall, Computer Aided Ship Design, Part 3. Grillage 02-Beam Element Seoul National **SDAL** Advanced Ship Design Automation Lab. http://asdal.snu.ac.kr 4/37

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### **Element : Beam - Differential Eqn.**

$$\boldsymbol{\sigma}_{x} = \boldsymbol{\sigma} \mathbf{i} \ , \ \boldsymbol{\varepsilon}_{x} = \varepsilon \mathbf{i} \ , \ \boldsymbol{\theta} = \boldsymbol{\theta} \mathbf{k}, \ \mathbf{y} = y \mathbf{j} \qquad \boldsymbol{\sigma} = E\varepsilon$$

(1) strain at y in x-direction :

$$\boldsymbol{\varepsilon}_{x} = \boldsymbol{\varepsilon} \mathbf{i}$$
$$\boldsymbol{\varepsilon} = \frac{(\rho - y) \cdot d\theta - \rho d\theta}{ds} = -y \frac{d\theta}{ds} = -\frac{y}{\rho}$$

, ds : initial length ,  $y \cdot d\theta$ : elongated length

y  
x 
$$\rho$$
  
d $\theta$   
neutral  
surface  
y ds  
ds  
dx

② stress at 
$$\mathcal{Y}$$
 in x-direction :  $\mathbf{\sigma}_x = \sigma \mathbf{i} = E \cdot \varepsilon \mathbf{i}$  , where  $\varepsilon = -\frac{y}{\rho}$   $\therefore \mathbf{\sigma}_x = \sigma \mathbf{i} = -E\frac{y}{\rho}\mathbf{i}$ 

**(3)** force acting on dA in x-direction  $: d\mathbf{F}_x = \mathbf{\sigma}_x dA = (\sigma \mathbf{i}) dA = \sigma dA \mathbf{i}$   $\therefore d\mathbf{F} = -E \frac{y}{\rho} dA \mathbf{i}$ 

(a) moment about z-axis :  $d\mathbf{M} = \mathbf{y} \times d\mathbf{F} = (y\mathbf{j}) \times (-E\frac{y}{\rho}dA\mathbf{i}) = E\frac{y^2}{\rho}dA\mathbf{k}$ 





### **Element : Beam - Differential Eqn.**

$$\sigma_{x} = \sigma \mathbf{i} , \ \mathbf{\epsilon}_{x} = \varepsilon \mathbf{i} , \ \mathbf{\theta} = \theta \mathbf{k}, \ \mathbf{y} = y \mathbf{j} \qquad \sigma = E\varepsilon$$

$$(\mathbf{p}, \mathbf{b}) \quad \mathbf{h} + d\theta = ds \implies \frac{d\theta}{ds} = \frac{1}{\rho}$$

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$$(\mathbf{p}, \mathbf{b}) \quad \mathbf{h} + d\theta = ds \implies \frac{d\theta}{ds} = -E \frac{y}{ds}$$

$$(\mathbf{p}, \mathbf{b}) \quad \mathbf{h} + d\theta = ds \implies \mathbf{h} + d\theta = -E \frac{y}{\rho} d\mathbf{A} \mathbf{i}$$

$$(\mathbf{p}, \mathbf{b}) \quad \mathbf{h} + d\theta = ds \implies \mathbf{h} + d\theta = (y\mathbf{j}) \times (-E \frac{y}{\rho} d\mathbf{A} \mathbf{i}) = E \frac{y^{2}}{\rho} d\mathbf{A} \mathbf{k} \quad \therefore \mathbf{M} = \int_{A} d\mathbf{M} = \int_{A} E \frac{y^{2}}{\rho} d\mathbf{A} \mathbf{k}$$

$$(\mathbf{p}, \mathbf{h}) = \frac{1}{\rho} \mathbf{h}, \ \mathbf{M} = \frac{EI}{\rho} \mathbf{h}, \ \mathbf{M} = \frac{EI}{\rho}$$

$$(\mathbf{p}, \mathbf{h}) = \frac{1}{\rho} \mathbf{h}, \ \mathbf{M} = \frac{EI}{\rho} \mathbf{h}, \ \mathbf{M} = EI \frac{d\theta}{ds} \mathbf{k}$$

$$(\mathbf{M}) = EI \frac{d^{2}y}{dx^{2}} \mathbf{k} \quad \mathbf{M} = EI \frac{d^{2}y}{dx^{2}}$$

$$(\mathbf{M}) = EI \frac{d\theta}{ds} \mathbf{k}$$

$$(\mathbf{M}) = EI \frac{d^{2}y}{dx^{2}} \mathbf{k} \quad \mathbf{h} = EI \frac{d^{2}y}{dx^{2}}$$

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$$(\mathbf{M}) = EI \frac{d$$

6/37

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# Element : Beam - Differential Eqn

$$\boldsymbol{\sigma}_x = \boldsymbol{\sigma} \mathbf{i}$$
,  $\boldsymbol{\varepsilon}_x = \boldsymbol{\varepsilon} \mathbf{i}$ ,  $\boldsymbol{\theta} = \boldsymbol{\theta} \mathbf{k}$ ,  $\mathbf{y} = y\mathbf{j}$   $\boldsymbol{\sigma} = E\boldsymbol{\varepsilon}$ 

strain at Y in x-direction : ε = (ρ-y)·dθ-ρ·dθ/ds = -y dθ/ds
 ε<sub>x</sub> = ε**i** , dθ : initial length, y·dθ : elongated length
 stress at Y in x-direction : σ<sub>x</sub> = σ**i** = -E Y/ρ i
 force acting on dA in x-direction : d**F** = -E y/ρ dA**i** moment about z-axis :



7/37

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$$d\mathbf{M} = \mathbf{y} \times d\mathbf{F} = (y\mathbf{j}) \times (-E\frac{y}{\rho}dA\mathbf{i}) = E\frac{y^2}{\rho}dA\mathbf{k} \quad \therefore \mathbf{M} = \int_A d\mathbf{M} = \int_A E\frac{y^2}{\rho}dA\mathbf{k} = \frac{EI}{\rho}\mathbf{k} \quad , I = \int_A y^2 dA$$
  
(5) assume  $ds \approx dx, \ \theta \approx \tan(\theta) = \frac{dy}{dx} \Rightarrow \frac{d\theta}{ds} = \frac{d^2y}{dx^2} \Rightarrow \mathbf{M} = \frac{EI}{\rho}\mathbf{k} = EI\frac{d\theta}{ds}\mathbf{k} \Rightarrow \mathbf{M} = EI\frac{d^2y}{dx^2}\mathbf{k} \quad , M = EI\frac{d^2y}{$ 

**(6)** relationships between loads, shear forces, and bending moments

200

$$\mathbf{V}_{1} = V\mathbf{j}, \quad \mathbf{V}_{2} = -\left(V + \frac{\partial V}{\partial x}dx\right)\mathbf{j}, \quad \mathbf{M}_{1} = -M\mathbf{k}, \quad \mathbf{M}_{2} = \left(M + \frac{\partial M}{\partial x}dx\right)\mathbf{k}$$
  
•force equilibrium 
$$\sum_{\mathbf{v}} \mathbf{F}_{\mathbf{v}} = \mathbf{V}_{1} + \mathbf{V}_{2} + \mathbf{f}(x) = 0$$
  

$$(V_{1}\mathbf{j}) + \left(-\left(V_{1} + \frac{\partial V_{1}}{\partial x}dx\right)\mathbf{j}\right) + (-f(x)dx\mathbf{j}) = 0$$
  

$$\left(V_{1} - V_{1} - \frac{\partial V_{1}}{\partial x}dx - f(x)dx\right)\mathbf{j} = 0$$
  

$$\therefore \quad \frac{dV}{dx} = -f(x)$$
  
9 Fall. Computer Aided Ship Design. Part 3. Grillage 02-Beam Element

# Element : Beam - Differential Eqn

$$\boldsymbol{\sigma}_x = \boldsymbol{\sigma} \mathbf{i}$$
,  $\boldsymbol{\varepsilon}_x = \boldsymbol{\varepsilon} \mathbf{i}$ ,  $\boldsymbol{\theta} = \boldsymbol{\theta} \mathbf{k}$ ,  $\mathbf{y} = y\mathbf{j}$   $\boldsymbol{\sigma} = E\boldsymbol{\varepsilon}$ 

(1) strain at y in x-direction :  $\varepsilon = \frac{(\rho - y) \cdot d\theta - \rho \cdot d\theta}{ds} = -y \frac{d\theta}{ds}$   $\varepsilon_x = \varepsilon \mathbf{i}$ ,  $d\theta$ : initial length,  $y \cdot d\theta$ : elongated length (2) stress at y in x-direction :  $\sigma_x = \sigma \mathbf{i} = -E \frac{y}{\rho} \mathbf{i}$ (3) force acting on dA in x-direction :  $d\mathbf{F} = -E \frac{y}{\rho} dA\mathbf{i}$ (4) moment about z-axis :



) = 0

8/37

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$$d\mathbf{M} = \mathbf{y} \times d\mathbf{F} = (y\mathbf{j}) \times (-E\frac{y}{\rho}dA\mathbf{i}) = E\frac{y^2}{\rho}dA\mathbf{k} \quad \therefore \mathbf{M} = \int_A d\mathbf{M} = \int_A E\frac{y^2}{\rho}dA\mathbf{k} = \frac{EI}{\rho}\mathbf{k} \quad , I = \int_A y^2 dA$$
  
(5) assume  $ds \approx dx, \ \theta \approx \tan(\theta) = \frac{dy}{dx} \Rightarrow \frac{d\theta}{ds} = \frac{d^2y}{dx^2} \Rightarrow \mathbf{M} = \frac{EI}{\rho}\mathbf{k} = EI\frac{d\theta}{ds}\mathbf{k} \Rightarrow \mathbf{M} = EI\frac{d^2y}{dx^2}\mathbf{k} \quad , M = EI\frac{d^2y}{$ 

**6** relationships between loads, shear forces, and bending moments

$$\mathbf{V}_{1} = V\mathbf{j}, \quad \mathbf{V}_{2} = -\left(V + \frac{\partial V}{\partial x}dx\right)\mathbf{j}, \quad \mathbf{M}_{1} = -M\mathbf{k}, \quad \mathbf{M}_{2} = \left(M + \frac{\partial M}{\partial x}dx\right)\mathbf{k}$$
  
•force equilibrium  
•moment equilibrium  

$$\sum \mathbf{M}_{z} = \mathbf{M}_{1} + \mathbf{M}_{2} + d\mathbf{x} \times \mathbf{V}_{2} + \frac{1}{2}d\mathbf{x} \times (\mathbf{f}(x) \cdot dx)$$

 $-M\mathbf{k} + \left(M + \frac{\partial M}{\partial x}dx\right)\mathbf{k} + \left(dx\mathbf{i}\right) \times \left(-\left(V + \frac{\partial V}{\partial x}dx\right)\mathbf{j}\right) + \left(\frac{1}{2}dx\mathbf{i}\right) \times \left(-f(x)dx\mathbf{j}\right) = 0 \qquad \therefore \frac{dM}{dx}$ 

# Element : Beam - Differential Eqn

$$\boldsymbol{\sigma}_x = \boldsymbol{\sigma} \mathbf{i}$$
,  $\boldsymbol{\varepsilon}_x = \boldsymbol{\varepsilon} \mathbf{i}$ ,  $\boldsymbol{\theta} = \boldsymbol{\theta} \mathbf{k}$ ,  $\mathbf{y} = y\mathbf{j}$   $\boldsymbol{\sigma} = E\boldsymbol{\varepsilon}$ 

strain at Y in x-direction : ε = (ρ-y)·dθ-ρ·dθ/ds = -y dθ/ds
 ε<sub>x</sub> = ε**i** , dθ : initial length, y·dθ : elongated length
 stress at Y in x-direction : σ<sub>x</sub> = σ**i** = -E Y/ρ i
 force acting on dA in x-direction : d**F** = -E y/ρ dA**i** moment about z-axis :



9/37

$$d\mathbf{M} = \mathbf{y} \times d\mathbf{F} = (y\mathbf{j}) \times (-E\frac{y}{\rho}dA\mathbf{i}) = E\frac{y^2}{\rho}dA\mathbf{k} \quad \therefore \mathbf{M} = \int_A d\mathbf{M} = \int_A E\frac{y^2}{\rho}dA\mathbf{k} = \frac{EI}{\rho}\mathbf{k} \quad , I = \int_A y^2 dA$$
  
(5) assume  $ds \approx dx, \ \theta \approx \tan(\theta) = \frac{dy}{dx} \Rightarrow \frac{d\theta}{ds} = \frac{d^2y}{dx^2} \Rightarrow \mathbf{M} = \frac{EI}{\rho}\mathbf{k} = EI\frac{d\theta}{ds}\mathbf{k} \Rightarrow \mathbf{M} = EI\frac{d^2y}{dx^2}\mathbf{k} \quad , M = EI\frac{d^2y}{$ 

**(6)** relationships between loads, shear forces, and bending moments

$$\mathbf{V}_{1} = V\mathbf{j}, \quad \mathbf{V}_{2} = -\left(V + \frac{\partial V}{\partial x}dx\right)\mathbf{j}, \quad \mathbf{M}_{1} = -M\mathbf{k}, \quad \mathbf{M}_{2} = \left(M + \frac{\partial M}{\partial x}dx\right)\mathbf{k}$$
  
•force equilibrium  $\frac{dV}{dx} = -f(x)$  •moment equilibrium  $\frac{dM}{dx} = V(x)$   
 $\frac{d^{2}y}{dx^{2}} = \frac{M}{EI} \rightarrow \frac{d^{3}y}{dx^{3}} = \frac{1}{EI} \cdot \frac{dM}{dx} = \frac{1}{EI} \cdot V(x) \rightarrow \frac{d^{4}y}{dx^{4}} = \frac{1}{EI} \cdot \frac{dV}{dx} = -\frac{1}{EI} \cdot f(x)$   
 $\therefore EI \frac{d^{4}y}{dx^{4}} = -f(x)$ 

### **Element : Beam - Variational Method**

#### multiply by $\delta u$ and integrate

$$\int_0^l \left( EI \frac{d^4 v}{dx^4} + f \right) \delta v \, dx = 0$$

L.H.S:

$$\int_0^l \left( EI \frac{d^4 v}{dx^4} \delta v + f \, \delta v \right) dx$$

integration by part

$$= EI\left[\frac{d^{3}v}{dx^{3}}\delta v\right]_{0}^{l} - \int_{0}^{l} \left(EI\frac{d^{3}v}{dx^{3}}\frac{d(\delta v)}{dx}\right) dx + \int_{0}^{l} (f\delta v) dx$$
$$= -\int_{0}^{l} \left(EI\frac{d^{3}v}{dx^{3}}\frac{d(\delta v)}{dx}\right) dx + \int_{0}^{l} (f\delta v) dx$$
$$= -EI\left[\frac{d^{2}v}{dx^{2}}\frac{d\delta v}{dx}\right]_{0}^{l} + \int_{0}^{l} EI\frac{d^{2}v}{dx^{2}}\frac{d^{2}\delta v}{dx^{2}} dx + \int_{0}^{l} (f\delta v) dx$$

$$= EI \int_0^l \frac{d^2 v}{dx^2} \frac{d^2 \delta v}{dx^2} dx + \int_0^l (f \,\delta v) dx$$

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### **Differential Equation**

$$EI\frac{d^4v}{dx^4} + f(x) = 0$$

### **Boundary condition**

$$v|_{x=0} = 0 , v|_{x=l} = 0$$
  
,  $EI \frac{d^2 v}{dx^2}|_{x=0} = 0, EI \frac{d^2 v}{dx^2}|_{x=l} = 0$ 

 $\delta$  operation

• 
$$f \,\delta v = \delta (fv)$$
  
•  $v \,\delta v = \delta \left(\frac{1}{2}v^2\right)$   
•  $\frac{\delta^2 v}{\delta x^2} \delta \frac{\delta^2 v}{\delta x^2} = \frac{1}{2} \delta \left(\frac{\delta^2 v}{\delta x^2}\right)^2$ 

• 
$$\frac{d}{dx}\delta v = \delta \frac{d}{dx}v$$

• 
$$\delta \int_{a}^{b} h(x) dx = \int_{a}^{b} \delta h(x) dx$$



10/37

### **Element : Beam - Variational Method**

#### multiply by $\delta u$ and integrate

$$\int_0^l \left( EI \frac{d^4 v}{dx^4} + f \right) \delta v \, dx = 0$$

L.H.S:

$$\int_0^l \left( EI \frac{d^4 v}{dx^4} \delta v + f \, \delta v \right) dx$$

integration by part

$$= \int_0^l EI \frac{d^2 v}{dx^2} \frac{d^2 \delta v}{dx^2} dx + \int_0^l (f \,\delta v) dx$$

$$= \int_0^l \delta \frac{1}{2} EI \frac{d^2 v}{dx^2} \frac{d^2 v}{dx^2} dx + \int_0^l (\delta f v) dx$$
$$\delta \int_0^l \left[ \frac{EI}{2} \left( \frac{d^2 v}{dx^2} \right) + (f v) \right] dx$$

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### **Differential Equation**

$$EI\frac{d^4v}{dx^4} + f(x) = 0$$

### **Boundary Condition**

$$v|_{x=0} = 0$$
,  $v|_{x=l} = 0$   
,  $EI \frac{d^2 v}{dx^2}|_{x=0} = 0$ ,  $EI \frac{d^2 v}{dx^2}|_{x=l} = 0$ 

 $\delta$  operation

• 
$$f \,\delta v = \delta (fv)$$
  
•  $v \,\delta v = \delta \left(\frac{1}{2}v^2\right)$   
•  $\frac{\delta^2 v}{\delta x^2} \delta \frac{\delta^2 v}{\delta x^2} = \frac{1}{2} \delta \left(\frac{\delta^2 v}{\delta x^2}\right)^2$ 

• 
$$\frac{d}{dx}\delta v = \delta \frac{d}{dx}v$$

• 
$$\delta \int_{a}^{b} h(x) dx = \int_{a}^{b} \delta h(x) dx$$



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Variational Method  $\delta \int_{0}^{l} \left[ \frac{EI}{2} \left( \frac{d^{2}v}{dx^{2}} \right)^{2} + (fv) \right] dx$ 



discretization

finite element method  $\oint$  1 element , 2 nodes



assume: 
$$v(x) = c_0 + c_1 x + c_2 x^2 + c_3 x^3$$
,  $v(0) = v_1$ ,  $v(l) = v_2$   
,  $\frac{dv}{dx}(0) = \phi_1$ ,  $\frac{dv}{dx}(l) = \phi_2$ 



assume: 
$$v(x) = c_0 + c_1 x + c_2 x^2 + c_3 x^3$$
,  $v(0) = v_1$ ,  $v(l) = v_2$ ,  
 $\frac{dv}{dx}(0) = \phi_1$ ,  $\frac{dv}{dx}(l) = \phi_2$ 

$$v(0) = c_{0} \implies c_{0} = v_{1}$$

$$v(l) = c_{0} + c_{1}l + c_{2}l^{2} + c_{3}l^{3} = v_{2} - \frac{3}{l^{2}}(v_{1} - v_{2}) - \frac{1}{l}(2\phi_{1} + \phi_{2})$$

$$\frac{dv}{dx}(0) = c_{1} \implies c_{1} = \phi_{1}$$

$$\frac{dv}{dx}(l) = c_{1} + 2c_{2}l + 3c_{3}l^{2} = \phi_{2} - \frac{3}{l^{2}}(v_{1} - v_{2}) + \frac{1}{l^{2}}(\phi_{1} + \phi_{2})$$

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13/37



$$v(x) = \frac{1}{l^{3}}(2x^{3} - 3x^{2}l + l^{3})v_{1} + \frac{1}{l^{3}}(x^{3}l - 2x^{2}l^{2} + xl^{3})\phi_{1} + \frac{1}{l^{3}}(-2x^{3} + 3x^{2}l)v_{2} + \frac{1}{l^{3}}(x^{3}l - x^{2}l^{2})\phi_{2}$$

$$v(x) = \begin{bmatrix} N_{1} & N_{2} & N_{3} & N_{4} \end{bmatrix} \begin{bmatrix} v_{1} \\ \phi_{1} \\ v_{2} \\ \phi_{2} \end{bmatrix}$$

$$N_{1} = \frac{1}{l^{3}}(2x^{3} - 3x^{2}l + l^{3})$$

$$N_{2} = \frac{1}{l^{3}}(x^{3}l - 2x^{2}l^{2} + xl^{3})$$

$$N_{3} = \frac{1}{l^{3}}(-2x^{3} + 3x^{2}l)$$

$$N_{4} = \frac{1}{l^{3}}(x^{3}l - x^{2}l^{2})$$

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15/37

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$$v(x) = \frac{1}{l^3} (2x^3 - 3x^2l + l^3)v_1 + \frac{1}{l^3} (x^3l - 2x^2l^2 + xl^3)\phi_1 + \frac{1}{l^3} (-2x^3 + 3x^2l)v_2 + \frac{1}{l^3} (x^3l - x^2l^2)\phi_2$$

$$v(x) = \begin{bmatrix} N_1 \ N_2 \ N_3 \ N_4 \end{bmatrix} \begin{bmatrix} v_1 \\ \phi_1 \\ v_2 \\ \phi_2 \end{bmatrix}$$

$$N_1 = \frac{1}{l^3} (-2x^3 + 3x^2l) \qquad N_2 = \frac{1}{l^3} (x^3l - 2x^2l^2 + xl^3) \\ N_3 = \frac{1}{l^3} (-2x^3 + 3x^2l) \qquad N_4 = \frac{1}{l^3} (x^3l - x^2l^2) \\ N_4 = \frac{1}{l^3} (x^3l - x^3l^2) \\ N_4 = \frac{1}{l^3} (x^3l - x^3l^$$

 $igcup_{}$  differentiation with respect to  $~\mathcal{X}$  twice

$$\frac{d^2 v(x)}{dx^2} = \begin{bmatrix} B_1 & B_2 & B_3 & B_4 \end{bmatrix} \begin{vmatrix} v_1 \\ \phi_1 \\ v_2 \\ \phi_2 \end{vmatrix}$$

 $B_{1} = \frac{1}{l^{3}}(12x - 6l) \qquad B_{2} = \frac{1}{l^{3}}(6xl - 4l^{2})$  $B_{3} = \frac{1}{l^{3}}(-12x + 6l) \qquad B_{4} = \frac{1}{l^{3}}(6xl - 2l^{2})$ 

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16/37

 $\therefore v(x) = \mathbf{N}\mathbf{d}, \quad \frac{d^2v(x)}{dx^2} = \mathbf{B}\mathbf{d}$ 

![](_page_16_Figure_1.jpeg)

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# (derivation)

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$$\delta \int_{0}^{l} \left[ \frac{EA}{2} \left( \frac{d^{2}v}{dx^{2}} \right)^{2} + (fv) \right] dx \qquad f \operatorname{Nd} = (\operatorname{Nd})^{\mathrm{T}} f = \operatorname{d}^{\mathrm{T}} \operatorname{N}^{\mathrm{T}} f \quad \because \operatorname{Nd} : scalar$$

$$= \delta \left\{ \frac{EI}{2} \int_{0}^{l} \left( \operatorname{d}^{\mathrm{T}} \operatorname{B}^{\mathrm{T}} \operatorname{Bd} \right) dx + \int_{0}^{l} (f \operatorname{Nd}) dx \right\}$$

$$= \delta \left\{ \frac{EI}{2} \int_{0}^{l} \left( \operatorname{d}^{\mathrm{T}} \operatorname{B}^{\mathrm{T}} \operatorname{Bd} \right) dx + \int_{0}^{l} (\operatorname{d}^{\mathrm{T}} \operatorname{N}^{\mathrm{T}} f) dx \right\}$$

$$= \delta \left\{ \frac{1}{2} \operatorname{d}^{\mathrm{T}} \left[ \int_{0}^{l} EI \left( \operatorname{B}^{\mathrm{T}} \operatorname{B} \right) dx \right] \operatorname{d} + \operatorname{d}^{\mathrm{T}} \left[ \int_{0}^{l} (\operatorname{N}^{\mathrm{T}} f) dx \right] \right\}$$

$$= \delta \left\{ \frac{1}{2} \operatorname{d}^{\mathrm{T}} \left[ \operatorname{Kd} - \operatorname{d}^{\mathrm{T}} \operatorname{F} \right] \right\}$$

$$= \frac{1}{2} (\delta \operatorname{d})^{\mathrm{T}} \operatorname{Kd} + \frac{1}{2} (\delta \operatorname{d})^{\mathrm{T}} \operatorname{Kd} - (\delta \operatorname{d})^{\mathrm{T}} \operatorname{F} \quad \because (\delta \operatorname{d})^{\mathrm{T}} \operatorname{Kd} = \operatorname{d}^{\mathrm{T}} \operatorname{K} \delta \operatorname{d}$$

$$\mathbf{K} = EI \int_{0}^{l} (\mathbf{B}^{T} \mathbf{B}) dx$$
$$= \frac{EI}{l^{3}} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^{2} & -6l & 2l^{2} \\ -12 & -6l & 12 & -6l \\ 6l & 2l^{2} & -6l & 4l^{2} \end{bmatrix}$$

 $\Rightarrow$  **K** = **K**<sup>T</sup>

symmetry

 $\mathbf{F} = -\int_0^l (\mathbf{N}^{\mathrm{T}} f) dx$  $d = \begin{bmatrix} v_1 & \phi_1 & v_2 & \phi_2 \end{bmatrix}^T$ 

![](_page_17_Picture_7.jpeg)

![](_page_17_Picture_8.jpeg)

### (derivation)

 $\mathbf{K} = EI \int_0^l \left( \mathbf{B}^T \mathbf{B} \right) dx$ 

$$= EI \int_{0}^{1} \left[ \frac{\frac{1}{l^{3}}(12x-6l)\frac{1}{l^{3}}(12x-6l)}{\frac{1}{l^{3}}(12x-6l)\frac{1}{l^{3}}(6xl-4l^{2})}{\frac{1}{l^{3}}(6xl-4l^{2})} - \frac{\frac{1}{l^{3}}(12x-6l)\frac{1}{l^{3}}(-12x+6l)}{\frac{1}{l^{3}}(6xl-4l^{2})\frac{1}{l^{3}}(6xl-2l^{2})} \right] dx$$

$$= EI \int_{0}^{1} \left[ \frac{\frac{1}{l^{3}}(6xl-4l^{2})\frac{1}{l^{3}}(12x-6l)}{\frac{1}{l^{3}}(12x-6l)\frac{1}{l^{3}}(6xl-4l^{2})\frac{1}{l^{3}}(6xl-4l^{2})}{\frac{1}{l^{3}}(6xl-4l^{2})\frac{1}{l^{3}}(-12x+6l)\frac{1}{l^{3}}(-12x+6l)\frac{1}{l^{3}}(6xl-4l^{2})\frac{1}{l^{3}}(6xl-2l^{2})} \right] dx$$

$$= EI \int_{0}^{1} \left[ \frac{1}{l^{6}}(144x^{2}-144xl+36l^{2}) - \frac{1}{l^{6}}(72x^{2}l-84xl^{2}+24l^{3})}{\frac{1}{l^{6}}(36x^{2}l^{2}-48xl^{3}+16l^{4})} - \frac{1}{l^{6}}(-72x^{2}l+84xl^{2}-24l^{3})} - \frac{1}{l^{6}}(-72x^{2}l+84xl^{2}-24l^{3})} - \frac{1}{l^{6}}(-72x^{2}l+60xl^{2}-12l^{3})} \right] dx$$

$$\begin{bmatrix} l^{6} & l^{6} & l^{6} \\ \frac{1}{l^{6}}(72x^{2}l - 60xl^{2} + 12l^{3}) & \frac{1}{l^{6}}(36x^{2}l^{2} - 36xl^{3} + 8l^{4}) & \frac{1}{l^{6}}(-72x^{2}l + 60xl^{2} - 12l^{3}) & \frac{1}{l^{3}}(36x^{2}l^{2} - 24xl^{3} + 4l^{4}) \end{bmatrix}$$

Seoul National Univ

19/37

$$=\frac{EI}{l^{3}}\begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^{2} & -6l & 2l^{2} \\ -12 & -6l & 12 & -6l \\ 6l & 2l^{2} & -6l & 4l^{2} \end{bmatrix}$$

![](_page_19_Figure_1.jpeg)

$$\therefore \mathbf{Kd} = \mathbf{F} \quad \text{where, } \mathbf{K} = \frac{EI}{l^3} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \quad \mathbf{F} = -\int_0^l (\mathbf{N}^T f) dx$$

![](_page_19_Picture_4.jpeg)

![](_page_19_Picture_5.jpeg)

# **Element : Beam - Finite Element Method** $\int_{0}^{t} \left[ \frac{EI}{2} \left( \frac{d^2 v}{dx^2} \right)^2 + (fv) \right] dx$

equivalent nodal forces

#### Variational Method Element : Beam - Finite Element Metho $\delta \int_0^l \left[ \frac{EI}{2} \left( \frac{d^2 v}{dx^2} \right)^2 + (fv) \right] dx$

![](_page_21_Figure_1.jpeg)

# **Element : Beam - Finite Element Method** $S_{j_0}^{[I]} \left[ \frac{EI}{2} \left( \frac{d^2 v}{dx^2} \right)^2 + (fv) \right] dx$

equivalent nodal forces

$$f(x) = f: const$$

![](_page_22_Figure_3.jpeg)

$$v(x) = \frac{f}{24EI} (x^3 l - 2x^2 l^2 + x l^3) \phi_1 + \frac{1}{l^3} (x^3 l - x^2 l^2) \phi_2 , 0 \le x \le l$$

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displacement given : x find : v(x)

![](_page_22_Picture_7.jpeg)

23/37

![](_page_22_Picture_8.jpeg)

![](_page_22_Picture_9.jpeg)

### **Galerkin's Residual Method**

![](_page_23_Figure_1.jpeg)

Thus substituting the approximated solution into the differential equation results in a residual over the whole region of the problem as follows

 $\iiint_V R \, dV$ 

In the residual method, we require that a weighted value of the residual be a minimum over the whole region. The weighting functions allow the weighted integral of residuals to go to zero

$$\iiint_V R W dV = 0$$

weighting function or test function

*ref.*)  $\int_{0}^{1} (-u'v' + uv - xv) dx = 0$ 

### **Galerkin Method**

 $\overline{\mathbf{n}}$ 

the basis functions  $N_i$  are chosen to play the role of the weighting functions W

$$\iiint_{V} R N_{i} dV = 0 \qquad , (i = 1, 2)$$

![](_page_23_Picture_11.jpeg)

![](_page_23_Picture_12.jpeg)

![](_page_23_Picture_13.jpeg)

### Element : Beam - Galerkin's Residual Method

#### Beam - Galerkin's Residual Method

$$\int_{0}^{l} \left[ EI \frac{d^{4}v(x)}{dx^{4}} + f \right] N_{i} dx = 0 \quad , (i = 1, 2, 3, 4)$$
  
where,  $v(x) = \mathbf{N}\mathbf{d}$ 

 $N_{1} = \frac{1}{l^{3}}(2x^{3} - 3x^{2}l + l^{3})$   $N_{2} = \frac{1}{l^{3}}(x^{3}l - 2x^{2}l^{2} + xl^{3})$   $N_{3} = \frac{1}{l^{3}}(-2x^{3} + 3x^{2}l)$   $N_{4} = \frac{1}{l^{3}}(x^{3}l - x^{2}l^{2})$ 

integration by parts

$$\left[N_{i}EI\frac{d^{3}v}{dx^{3}}\right]_{0}^{l} - \int_{0}^{l}EI\frac{d^{3}v}{dx^{3}}\frac{dN_{i}}{dx}dx + \int_{0}^{l}f N_{i}dx = 0 \quad , (i = 1, 2, 3, 4)$$

#### integration by parts again

$$EI\int_{0}^{l} \frac{d^{2}N_{i}}{dx^{2}} \frac{d^{2}v}{dx^{2}} dx + EI\left[N_{i}\frac{d^{3}v}{dx^{3}} - \frac{dN_{i}}{dx}\frac{d^{2}v}{dx^{2}}\right]_{0}^{l} + \int_{0}^{l} f N_{i} dx = 0 \quad , (i = 1, 2, 3, 4)$$

$$EI \int_{0}^{l} \frac{d^{2}N_{i}}{dx^{2}} \mathbf{B} \, dx \, \mathbf{d} + EI \left[ N_{i}V - \frac{dN_{i}}{dx} m \right]_{0}^{l} + \int_{0}^{l} f \, N_{i} \, dx = 0 \quad , (i = 1, 2, 3, 4)$$

In matrix form,

$$EI\int_0^l \mathbf{B}^T \mathbf{B} \, dx \, \mathbf{d} = EI\left[\frac{d\mathbf{N}^T}{dx}m - \mathbf{N}^T V\right]_0^l - \int_0^l \mathbf{N}^T f \, dx = 0$$

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![](_page_24_Picture_12.jpeg)

### **Differential Equation**

![](_page_24_Picture_14.jpeg)

Recall,

$$\frac{d^3v}{dx^3} = V(x), \ \frac{d^2v}{dx^2} = m(x)$$

![](_page_24_Picture_17.jpeg)

### Element : Beam - Galerkin's Residual Method

### *ref.*) $\int_{-u'v'+uv-xv}^{-uv-xv} dx = 0$ Beam - Galerkin's Residual Method **Galerkin Method** $N_1 = \frac{1}{l^3} (2x^3 - 3x^2l + l^3)$ the test functions $N_i$ are chosen to play the role of $\int_{0}^{l} EI \frac{d^{4}v(x)}{dx^{4}} + f \left[ N_{i} dx = 0 , (i = 1, 2) \right]$ $N_2 = \frac{1}{l^3} (x^3 l - 2x^2 l^2 + x l^3)$ the weighting functions W $\iiint_{V} R N_{i} dV = 0 , (i = 1, 2)$ weighting function *where*, $v(x) = \mathbf{Nd}$ $N_3 = \frac{1}{l^3}(-2x^3 + 3x^2l)$ $N_4 = \frac{1}{l^3} (x^3 l - x^2 l^2)$ integration by parts **Differential Equation** $EI\int_{0}^{T} \mathbf{B}^{T} \mathbf{B} \, dx \, \mathbf{d} = EI \left| \frac{d\mathbf{N}^{T}}{dx} m - \mathbf{N}^{T} V \right|^{l} - \int_{0}^{l} \mathbf{N}^{T} f \, dx = 0$ $EI\frac{d^4v(x)}{dx^4} + f = 0$ $\mathbf{N}(x) = \frac{1}{l^3} \begin{bmatrix} 2x^3 - 3x^2l + l^3 & x^3l - 2x^2l^2 + xl^3 & -2x^3 + 3x^2l & x^3l - x^2l^2 \end{bmatrix}$ $\frac{d\mathbf{N}(x)}{dx} = \frac{1}{l^3} \begin{bmatrix} 6x^2 - 6xl & 3x^2l - 4xl^2 + l^3 & -6x^2 + 6xl & 3x^2l - 2xl^2 \end{bmatrix}$ $\mathbf{N}(0) = \begin{bmatrix} 1 & 0 & 0 & 0 \end{bmatrix} \quad , \mathbf{N}(l) = \begin{bmatrix} 0 & 0 & 1 & 0 \end{bmatrix}$ $\mathbf{N}(0) = \begin{bmatrix} 1 & 0 & 0 & 0 \end{bmatrix} \quad , \mathbf{N}(l) = \begin{bmatrix} 0 & 0 & 1 & 0 \end{bmatrix}$ $\mathbf{N}(0) = \begin{bmatrix} 0 & 1 & 0 & 0 \end{bmatrix} \quad , \mathbf{N}(l) = \begin{bmatrix} 0 & 0 & 1 & 0 \end{bmatrix}$ $\mathbf{K} = EI \int_{0}^{l} (\mathbf{B}^{T} \mathbf{B}) dx$ $=\frac{EI}{l^3} \begin{vmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{vmatrix}$ R.H.S

$$EI\left[\frac{d\mathbf{N}^{T}}{dx}m - \mathbf{N}^{T}V\right]_{0}^{l} - \int_{0}^{l}\mathbf{N}^{T}f \, dx = m(l)\frac{d\mathbf{N}}{dx}(l) - V(l)\mathbf{N}(l) - m(0)\frac{d\mathbf{N}^{T}}{dx}(0) + V(0)\mathbf{N}^{T}(0) - \int_{0}^{l}\mathbf{N}^{T}f \, dx$$

![](_page_25_Picture_4.jpeg)

### Element : Beam - Galerkin's Residual Method

![](_page_26_Figure_1.jpeg)

### Element : Beam - Galerkin's Residual Method *ref.*) $\int (-u'v' + uv - xv) dx = 0$

d

 $\therefore \mathbf{Kd} = \mathbf{F} \text{ where, } \mathbf{K} = \frac{EI}{l^3} \begin{vmatrix} 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{vmatrix}, \mathbf{F} = [f_1 \ f_2 \ f_3 \ f_4]^T$ 

#### Beam - Galerkin's Residual Method

$$\int_{0}^{l} \left[ EI \frac{d^{4}v(x)}{dx^{4}} + f \right] N_{i} dx = 0 \quad , (i = 1, 2)$$
where,  $v(x) = \mathbf{N}$ 

#### **Galerkin Method**

the test functions  $N_i$  are chosen to play the role of the weighting functions W

$$\iiint_{V} R N_{i} dV = 0 , (i = 1, 2)$$

$$\xrightarrow{V} \text{weighting function}$$

$$\xrightarrow{V} \text{residual (test function N used}$$

### **Differential Equation**

$$EI\frac{d^4v(x)}{dx^4} + f = 0$$

For simple support beam,

$$f(x) = f: const$$

$$V(0) = \frac{l}{2}f, m(0) = 0, V(l) = -\frac{l}{2}f, m(l) = 0$$

![](_page_27_Figure_11.jpeg)

$$f_1 = V(0) - \frac{l}{2}f, \quad f_2 = -m(0) - \frac{l}{12}f$$
$$f_3 = -V(l) - \frac{l}{2}f, \quad f_4 = m(l) + \frac{l^2}{12}f$$

 $\mathbf{F} = [f_1 \ f_2 \ f_3 \ f_4]^T = \begin{bmatrix} 0 \\ -\frac{l^2}{12}f \\ 0 \\ \frac{l^2}{-1}f \end{bmatrix}$ 

$$f_1 = V(0) - \frac{l}{2}f, \quad f_2 = -m(0) - \frac{l^2}{12}f$$
$$f_3 = -V(l) - \frac{l}{2}f, \quad f_4 = m(l) + \frac{l^2}{12}f$$

the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy  $\Pi$  is *defined* as the

sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

 $\Pi = \Pi_{in} + \Pi_{ext}$ 

![](_page_28_Picture_7.jpeg)

![](_page_28_Picture_8.jpeg)

the total potential energy  $\Pi$  is *defined* as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

 $\Pi = \Pi_{in} + \Pi_{ext}$ 

To evaluate the strain energy for a bar, we consider only the work done by the internal forces during deformation.

$$d\Pi_{in} = \int_{0}^{\varepsilon_{x}} \sigma \, d\varepsilon_{x} \, dx \, dy \, dz$$
$$= \int_{0}^{\varepsilon_{x}} E\varepsilon_{x} \, d\varepsilon_{x} \, dx \, dy \, dz = \frac{1}{2} E(\varepsilon_{x})^{2} d\varepsilon_{x} \, dx \, dy \, dz$$
$$= \frac{1}{2} \sigma \varepsilon_{x} \, dx \, dy \, dz$$

$$\Pi_{in} = \iiint_{V} d\Pi_{in} = \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV$$

the strain energy for one-dimensional stress.

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![](_page_29_Figure_7.jpeg)

Linear-elastic (Hooke's law)material

![](_page_29_Picture_9.jpeg)

![](_page_29_Picture_10.jpeg)

#### the total potential energy $\Pi$ is *defined* as the

sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

 $\Pi = \Pi_{in} + \Pi_{ext}$ 

The potential energy of the external forces, being opposite in sign from the external work expression because the potential energy of external forces is lost when the work is done by the external forces, is given by

$$\Pi_{ext} = -\iint_{S_1} T_y v_s dS - \sum_{i=1}^2 f_{iy} v_i - \sum_{i=1}^2 m_i \phi_i$$

- body forces  $X_b$  typically from the self-weight of the bar (in units of force per unit volume) moving through displacement function v
- surface loading or traction  $T_y$  typically from distributed loading acting along the surface of the element (in units of force per unit surface area) moving through displacements  $v_s$  where  $v_s$  are the displacements occurring over surface  $s_1$
- nodal concentrated force  $f_{iv}$  moving through nodal displacements  $v_i$

![](_page_30_Figure_9.jpeg)

![](_page_30_Picture_11.jpeg)

#### the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy II is defined as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

$$\Pi = \Pi_{in} + \Pi_{ext} \qquad \Pi_{in} = \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV \quad , \Pi_{ext} = - \iint_{S_{1}} T_{y} v_{s} dS - \sum_{i=1}^{2} f_{iy} v_{i} - \sum_{i=1}^{2} m_{i} \phi_{i}$$

Apply the following steps when using the principle of minimum potential energy to derive the finite element equations.

- 1. Formulate an expression for the total potential energy.
- 2. Assume the displacement pattern to vary with a finite set of undetermined parameters (here these are the nodal displacements  $V_i$ ), which are substituted into the expression for total potential energy.
- 3. Obtain a set of simultaneous equations minimizing the total potential energy with respect to these nodal parameters. These resulting equations represent the element equations.

![](_page_31_Picture_12.jpeg)

#### the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy  $\Pi$  is defined as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

$$\Pi = \Pi_{in} + \Pi_{ext} \qquad \Pi_{in} = \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV \quad , \Pi_{ext} = - \iint_{S} T_{y} v_{s} dS - \sum_{i=1} f_{iy} v_{i} - \sum_{i=1} m_{i} \phi_{i}$$

assume that there is no surface traction and body force and the sectional area A is constant

Apply the following steps when using the principle of minimum potential energy to derive the finite element equations.

The differential volume for the beam element

$$dV = dAdx$$

The differential area over which the surface loading acts is

$$dS = b dx$$
 ,  $b$ : Width of beam

$$: \Pi = \Pi_{in} + \Pi_{ext}$$

$$= \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV - \iint_{S_{1}} T_{y} v_{s} dS - \sum_{i=1}^{2} f_{iy} v_{i} - \sum_{i=1}^{2} m_{i} \phi_{i}$$

$$= \frac{1}{2} \iiint_{x} \sigma_{x} \varepsilon_{x} dA dx - \int_{0}^{l} bT_{y} v_{s} dx - \sum_{i=1}^{2} f_{iy} v_{i} - \sum_{i=1}^{2} m_{i} \phi_{i}$$

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![](_page_32_Figure_13.jpeg)

![](_page_32_Picture_14.jpeg)

![](_page_32_Picture_15.jpeg)

33/37

### the principle of minimum potential energy Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy the total potential energy $\Pi$ is defined as the sum of the internal strain energy $\Pi_{in}$ and the potential energy of the external forces II<sub>ext</sub> $\Pi = \Pi_{in} + \Pi_{ext} \qquad \Pi_{in} = \frac{1}{2} \iiint \sigma_x \varepsilon_x dV \quad , \Pi_{ext} = -\iint T_y v_s dS - \sum_{i=1}^2 f_{iy} v_i - \sum_{i=1}^2 m_i \phi_i$ assume that there is no surface traction and body force and the $y, v \wedge$ sectional area A is constant $\therefore \Pi = \frac{1}{2} \iiint_{xA} \sigma_x \varepsilon_x dA dx - \int_0^l bT_y v_s dx - \sum_{i=1}^2 f_{iy} v_i - \sum_{i=1}^2 m_i \phi_i$ $\varepsilon_x = \frac{du}{dx} = -y \frac{d^2 v}{dx^2}$ $=-v\mathbf{Bd}$ $\sigma_{\rm r} = E \varepsilon_{\rm r} = -y E \mathbf{B} \mathbf{d}$ dvdx $bT_y = f_y$ , $\int_0^l bT_y v_s dx = \int bT_y \mathbf{d}^T \mathbf{N}^T dx$ $f_{1y}$ $-\int_{0}^{l} f_{y} \mathbf{d}^{T} \mathbf{N}^{T} dx - \sum_{i=1}^{2} f_{iy} v_{i} - \sum_{i=1}^{2} m_{i} \phi_{i} = -\mathbf{d}^{T} \int_{0}^{l} \mathbf{N}^{T} f_{y} dx - [v_{1} \ \phi_{1} \ v_{2} \ \phi_{2}] \begin{vmatrix} v_{1y} \\ m_{1} \\ f_{2y} \end{vmatrix} = -\mathbf{d}^{T} \mathbf{F}$ $\Pi = \frac{EI}{2} \int_0^l \mathbf{d}^T \mathbf{B}^T \mathbf{B} \, \mathbf{d} dx - \mathbf{d}^T \mathbf{F}$

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![](_page_33_Picture_3.jpeg)

![](_page_33_Picture_4.jpeg)

 $u = -y \frac{dv}{dx}$ 

#### the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy  $\Pi$  is defined as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

$$\Pi = \Pi_{in} + \Pi_{ext} \qquad \Pi_{in} = \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV \quad , \Pi_{ext} = - \iint_{S} T_{y} v_{s} dS - \sum_{i=1} f_{iy} v_{i} - \sum_{i=1} m_{i} \phi_{i}$$

assume that there is no surface traction and body force and the sectional area A is constant

$$\Pi = \frac{1}{2} \iiint_{xA} \sigma_x \varepsilon_x dA dx - \int_0^l bT_y v_s dx - \sum_{i=1}^2 f_{iy} v_i - \sum_{i=1}^2 m_i \phi_i$$
$$\Box = \frac{EI}{2} \int_0^l \mathbf{d}^{\mathrm{T}} \mathbf{B}^{\mathrm{T}} \mathbf{B} \mathbf{d} dx - \mathbf{d}^{\mathrm{T}} \mathbf{F}$$

 $\begin{array}{c} y, v \land \qquad l \\ m_1 \\ f_{1y} \\ f_{1y} \\ f_{1y} \\ f_{1y} \\ f_{2y} \\ f_{2y} \end{array} >$ 

3. Obtain a set of simultaneous equations minimizing the total potential energy with respect to these nodal parameters. These resulting equations represent the element equations.

#### The minimization of $\Pi$ with respect to each nodal displacement requires that

$$\frac{\partial \Pi}{\partial v_1} = 0, \frac{\partial \Pi}{\partial \phi_1} = 0, \frac{\partial \Pi}{\partial v_2} = 0 \text{ and } \frac{\partial \Pi}{\partial \phi_2} = 0$$

![](_page_34_Picture_12.jpeg)

![](_page_34_Picture_13.jpeg)

#### the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy  $\Pi$  is defined as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{avt}$ 

![](_page_35_Figure_4.jpeg)

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36/37

#### the principle of minimum potential energy

Of all the geometrically possible shapes that a body can assume, the true one, corresponding to the safisfaction of stable equilibrium of the body, is identified by a minimum value of the total potential energy

the total potential energy  $\Pi$  is defined as the sum of the internal strain energy  $\Pi_{in}$  and the potential energy of the external forces  $\Pi_{ext}$ 

$$\Pi = \prod_{in} + \prod_{ext} \qquad \Pi_{in} = \frac{1}{2} \iiint_{V} \sigma_{x} \varepsilon_{x} dV \qquad , \Pi_{ext} = - \iiint_{V} X_{b} u dV - \iint_{S_{b}} T_{x} u_{s} dS - \sum_{i=1} f_{ix} u_{i}$$

$$\mathbf{Kd} = \mathbf{F} \qquad \text{where, } \mathbf{K} = \frac{EI}{l^3} \begin{bmatrix} 12 & 6l & -12 & 6l \\ 6l & 4l^2 & -6l & 2l^2 \\ -12 & -6l & 12 & -6l \\ 6l & 2l^2 & -6l & 4l^2 \end{bmatrix} \quad \mathbf{F} = \int_0^l \mathbf{N} f_y dx + \begin{bmatrix} f_{1y} \\ m_1 \\ f_{2y} \\ m_2 \end{bmatrix}$$

For simple support beam with uniform load f

$$x \qquad f_{y} = -f, \quad f_{1y} = \frac{f}{2} \\ m_{1} = 0 \\ f_{2y} = \frac{f}{2} \\ m_{2} = 0 \end{aligned} , \mathbf{F} = -\int_{0}^{l} \mathbf{N} f \, dx + \begin{bmatrix} \frac{f}{2} \\ 0 \\ \frac{f}{2} \\ 0 \end{bmatrix} = \begin{bmatrix} 2^{-j} + 2^{-j} \\ -\frac{l^{2}}{12} f + 0 \\ -\frac{l}{2} f + \frac{l}{2} f \\ \frac{l^{2}}{12} f + 0 \end{bmatrix} = \begin{bmatrix} 0 \\ -\frac{l^{2}}{12} f \\ 0 \\ \frac{l^{2}}{12} f \end{bmatrix}$$

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f(x) = f:const

![](_page_36_Picture_9.jpeg)

![](_page_36_Picture_10.jpeg)

 $\begin{bmatrix} l & l \\ -l & f + l \\ f \end{bmatrix}$