

# Ch. 5 Three-dimensional potential flows

노트 제목

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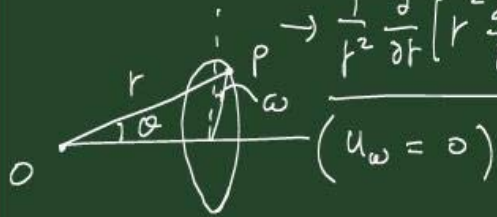
(axisymmetric flows)

$\psi$  does not exist.  $(r, \theta)$  coord.

- Velocity potential

$$\nabla \times \underline{u} = 0 \rightarrow \underline{u} = \nabla \phi \rightarrow u_r = \frac{\partial \phi}{\partial r}, \quad u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta}$$

$$\nabla \cdot \underline{u} = 0 \rightarrow \nabla^2 \phi = 0$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left( r^2 \frac{\partial \phi}{\partial r} \right) + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left( \sin \theta \frac{\partial \phi}{\partial \theta} \right) = 0$$


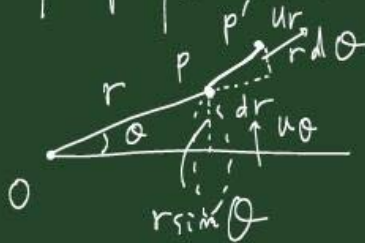
$(u_\theta = 0)$

- Stokes stream ft.

$$\nabla \cdot \underline{u} = 0 \rightarrow \frac{1}{r^2} \frac{\partial}{\partial r} (r^2 u_r) + \frac{1}{r \sin \theta} \frac{\partial}{\partial \theta} (u_\theta \sin \theta) = 0$$

$$\rightarrow u_r = \frac{1}{r^2 \sin \theta} \frac{\partial \psi}{\partial \theta}, \quad u_\theta = -\frac{1}{r \sin \theta} \frac{\partial \psi}{\partial r}$$

$\psi$  physical meaning



$$dQ = 2\pi r \sin \theta (u_r r d\theta - u_\theta dr)$$

(P-P')

$$= 2\pi r \sin \theta \left( \frac{1}{r^2 \sin \theta} \frac{\partial \psi}{\partial \theta} r d\theta \right)$$

$$+ \frac{1}{r \sin \theta} \frac{\partial \psi}{\partial r} dr$$

$$= 2\pi \left( \frac{\partial \psi}{\partial \theta} d\theta + \frac{\partial \psi}{\partial r} dr \right) = 2\pi d\psi$$

$\therefore 2\pi r d\psi$  corresponds to the volume of fluid crossing the surface of revolution which is formed by rotating the position vector  $pp'$  around the reference axis.



- Solution of the potential equation  $\nabla^2\phi = 0$   
 Separation of variables  $\phi(r, \theta) = R(r)T(\theta)$   
 $\rightarrow \frac{T}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) + \frac{R}{r^2 \sin\theta} \frac{d}{d\theta} \left( \sin\theta \frac{dT}{d\theta} \right) = 0$

$$\rightarrow \underbrace{\frac{1}{R} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right)}_{\text{fct. } r} = - \underbrace{\frac{1}{T \sin\theta} \frac{d}{d\theta} \left( \sin\theta \frac{dT}{d\theta} \right)}_{\text{fct. } \theta} = \text{const} = l(l+1)$$

$$\frac{1}{r^2} \frac{d}{dr} \left( r^2 \frac{dR}{dr} \right) = l(l+1) : \text{Euler-Cauchy eq.}$$

$$\rightarrow R_l(r) = A_l r^l + B_l r^{-(l+1)}$$

$$- \frac{1}{T \sin\theta} \frac{d}{d\theta} \left( \sin\theta \frac{dT}{d\theta} \right) = l(l+1) : \text{Legendre eq.}$$

$$\text{Let } x = \cos\theta$$

$$\rightarrow \frac{d}{dx} \left[ (1-x^2) \frac{dT}{dx} \right] + l(l+1)T = 0$$

$$\rightarrow T_l(\theta) = \sqrt{C_l} P_l(\cos\theta) + D_l Q_l(\cos\theta)$$

$\uparrow$  1st kind                       $\uparrow$  2nd kind  
 for all  $l$ ,  $Q_l(\pm 1) \rightarrow \infty$   
 $P_l(\pm 1)$  diverges unless  
 $l$  is an integer.

$\therefore \phi_l(r, \theta) = (A_l r^l + B_l r^{-(l+1)}) P_l(\cos \theta)$

$\therefore \phi(r, \theta) = \sum_{l=0}^{\infty} (A_l r^l + \frac{B_l}{r^{l+1}}) P_l(\cos \theta)$

Legendre ft.  $P_l(x) = \frac{1}{2^l l!} \frac{d^l}{dx^l} (x^2 - 1)^l$

$P_0(x) = 1$ ,  $P_1(x) = x$ ,  $P_2(x) = \frac{1}{2}(3x^2 - 1)$ , ...

We use this sol. for some specific potential flows.

- Uniform flow

$B_l = 0$  and  $A_l = \begin{cases} 0 & \text{for } l \neq 1 \\ U & \text{for } l = 1 \end{cases}$        $P_1(\cos \theta) = \cos \theta$

$\rightarrow \phi(r, \theta) = U r \cos \theta$

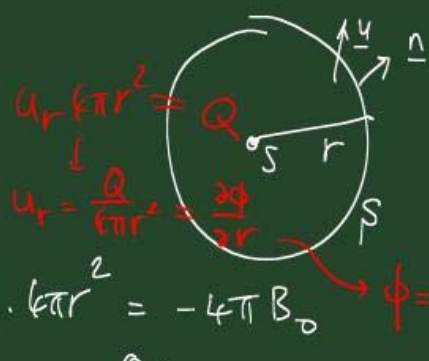
$u_r = \frac{\partial \phi}{\partial r} = U \cos \theta = \frac{1}{r^2 \sin \theta} \frac{\partial \psi}{\partial \theta}$

$u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -U \sin \theta = -\frac{1}{r \sin \theta} \frac{\partial \psi}{\partial r}$

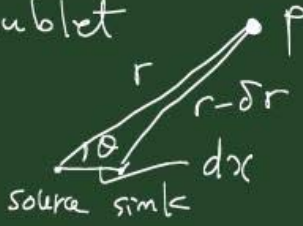
$\psi = \frac{1}{2} U r^2 \sin \theta$
- Source and sink

$A_l = 0$  and  $B_l = \begin{cases} 0 & \text{for } l \neq 0 \\ B_0 & \text{for } l = 0 \end{cases}$ ,  $P_0(\cos \theta) = 1$

$\rightarrow \phi(r, \theta) = \frac{B_0}{r}$   
 $u_r = -\frac{B_0}{r^2}$   
 $u_\theta = 0$   
 $Q = \int_S \mathbf{u} \cdot \mathbf{n} \, ds = -\frac{B_0}{r^2} \cdot 4\pi r^2 = -4\pi B_0$   
 $\therefore \phi(r, \theta) = -\frac{Q}{4\pi r}$   
 $\rightarrow \psi = -\frac{Q}{4\pi} (1 + \cos\theta)$



$u_r \cdot 4\pi r^2 = Q$   
 $u_r = \frac{Q}{4\pi r^2} = \frac{\partial \phi}{\partial r}$   
 $\phi = -\frac{Q}{4\pi r}$   
 $Q > 0$  source  
 $Q < 0$  sink

• Doublet 
 $\delta r \approx \delta z \cos\theta$

$\phi(r, \theta) = -\frac{Q}{4\pi r} + \frac{Q}{4\pi(r-\delta r)}$   
 $= -\frac{Q}{4\pi} \left( \frac{1}{r} - \frac{1}{r-\delta r} \right) \approx +\frac{Q\delta r}{4\pi r^2}$   
 $\approx \frac{Q}{4\pi r^2} \delta z \cos\theta = \frac{Q\delta z}{4\pi r^2} \cos\theta$   
 $= \frac{\mu}{4\pi r^2} \cos\theta$

$\left( \begin{array}{l} Q\delta z \rightarrow \text{finite} \\ \text{as } \delta z \rightarrow 0 \end{array} \right)$

$u_r = \frac{\partial \phi}{\partial r} = -\frac{\mu}{2\pi r^3} \cos\theta = \frac{1}{r^2 \sin\theta} \frac{\partial \psi}{\partial r}$   
 $u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = \frac{-\mu}{4\pi r^3} \sin\theta = -\frac{1}{r \sin\theta} \frac{\partial \psi}{\partial r}$

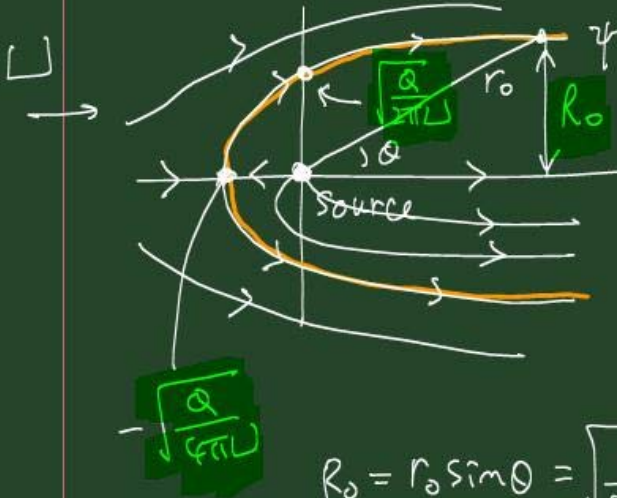
$\rightarrow \psi = -\frac{\mu}{4\pi r} \sin^2\theta$



• Flow near a blunt nose : uniform flow + source

$$\psi = \frac{1}{2} U r^2 \sin^2 \theta - \frac{Q}{4\pi} (1 + \cos \theta)$$

$$\begin{aligned} (1 + \cos \theta) &= 2 \cos^2 \frac{\theta}{2} \\ \sin \theta &= 2 \sin \frac{\theta}{2} \cos \frac{\theta}{2} \end{aligned}$$



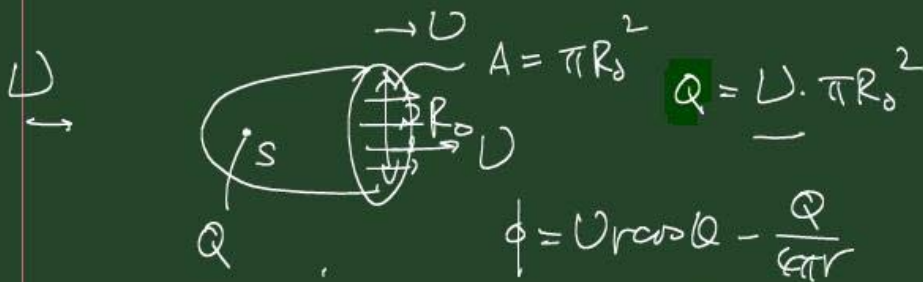
$$r_0 = \sqrt{\frac{Q}{4\pi U} \frac{1}{\sin^2 \frac{\theta}{2}}}$$

$$\theta = 0 : r_0 \rightarrow \infty$$

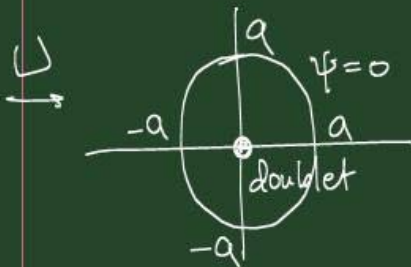
$$\theta = \frac{\pi}{2} : r_0 = \sqrt{\frac{Q}{2\pi U}}$$

$$\theta = \pi : r_0 = \sqrt{\frac{Q}{4\pi U}}$$

$$R_0 = r_0 \sin \theta = \sqrt{\frac{Q}{4\pi U}} \frac{\sin \theta}{\sin \frac{\theta}{2}} \rightarrow \text{as } \theta \rightarrow 0, R_0 = \sqrt{Q/\pi U}$$



• Flow around a sphere : uniform flow + doublet



$$\psi = \frac{1}{2} U r^2 \sin^2 \theta - \frac{\mu}{4\pi r} \sin^2 \theta$$

$$\psi = 0 : 0 = \frac{1}{2} U r_0^2 \sin^2 \theta - \frac{\mu}{4\pi r_0} \sin^2 \theta$$

$$\rightarrow r_0 = \left(\frac{\mu}{2\pi U}\right)^{\frac{1}{3}}$$

$$\text{Let } \mu = 2\pi U a^3 \rightarrow r_0 = a : \text{radius of sphere.}$$

$$\text{Then, } \psi = \frac{1}{2} U \left( r^2 - \frac{a^2}{r} \right) \sin^2 \theta$$
$$\phi = U \left( r + \frac{1}{2} \frac{a^3}{r^2} \right) \cos \theta .$$