

**Introduction to
Nuclear Fusion
(409.308A, 3 Credits)**

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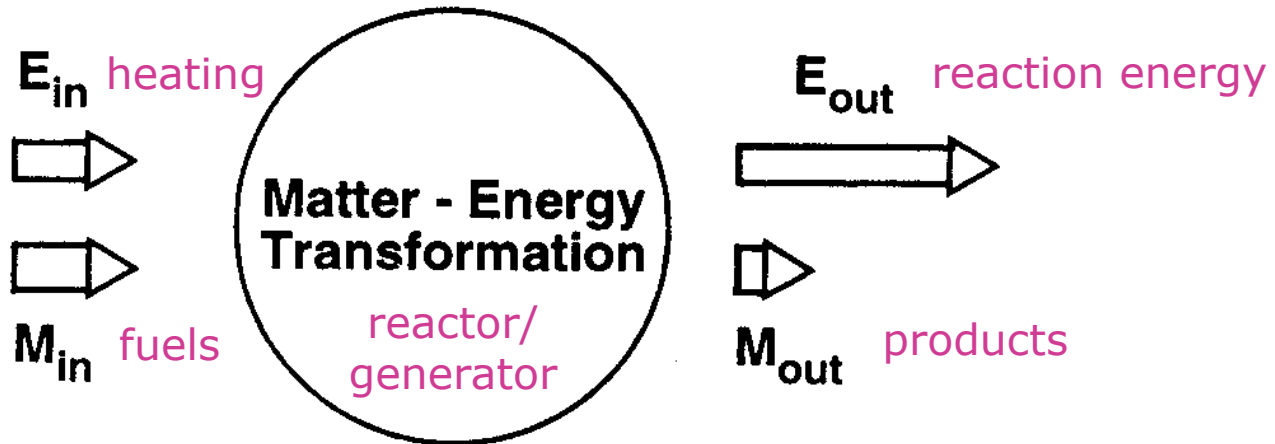
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Week 5. Inertial Confinement

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Matter and Energy



- Hydro-electric process
- Chemical reactions (combustion)
- Fission process
- Fusion process

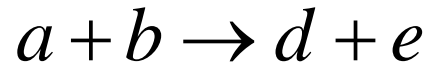
- Total energy conservation including rest mass energy

$$E_{in} + M_{in} \rightarrow E_{out} + M_{out}$$

- If $\Delta m = M_{out} - M_{in} < 0$, then we can get $E_{out} > E_{in}$.

Mass Defect Energy of Nuclear Reaction

reactants products



$$\Delta m_{ab} = (m_d + m_e) - (m_a + m_b)$$

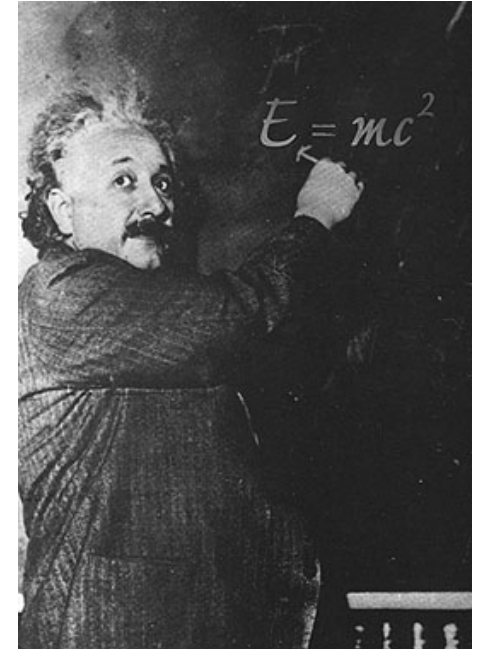
$\Delta m_{ab} < 0$: exothermic or exoergic

$\Delta m_{ab} > 0$: endothermic or endoergic

$$Q_{ab} = (-\Delta m)_{ab} c^2 \quad \text{Einstein's mass-energy relation}$$

$$E_{\text{before}}^* = E_{\text{after}}^*$$

$$(E_{k,a} + m_a c^2) + (E_{k,b} + m_b c^2) = (E_{k,d} + m_d c^2) + (E_{k,e} + m_e c^2)$$



Mass Defect Energy of Nuclear Reaction

For $E_{k,a} + E_{k,b} \ll Q_{ab}$

$$Q_{ab} \approx E_{k,d} + E_{k,e} = \frac{1}{2}m_d v_d^2 + \frac{1}{2}m_e v_e^2$$

Momentum conservation for reactions with CM at rest

$$m_d v_d = m_e v_e$$

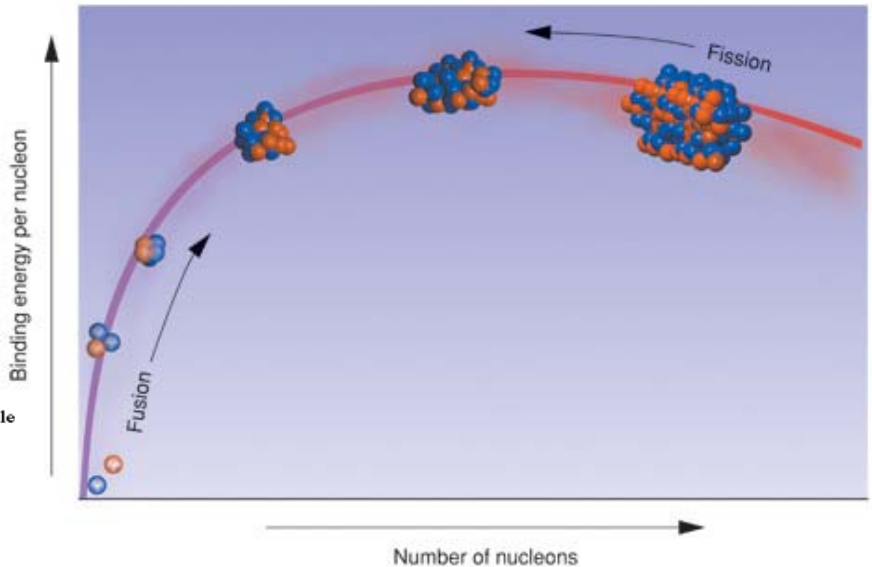
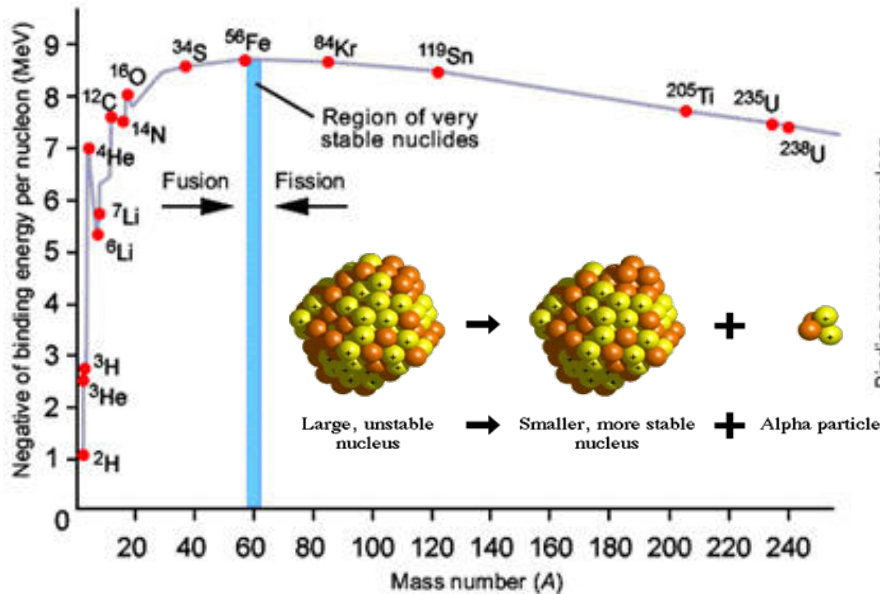
$$E_{k,d} \approx \left(\frac{m_e}{m_d + m_e} \right) Q_{ab}, \quad E_{k,e} \approx \left(\frac{m_d}{m_d + m_e} \right) Q_{ab}$$

Ex) d-t fusion reaction $d + t \rightarrow n + \alpha + 17.6 \text{ MeV}$

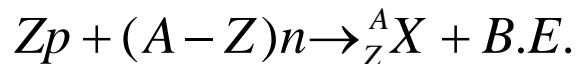
$$Q_{dt} = 17.6 \text{ MeV}$$

$$E_{k,n} \approx \frac{4}{5} Q_{dt} \approx 14.1 \text{ MeV}, \quad E_{k,\alpha} \approx \frac{1}{5} Q_{dt} \approx 3.5 \text{ MeV}$$

Binding Energy for an Assembled Nucleus



- The amount of energy released when a particular isotope is formed.
- The strength of the bonding is measured by the binding energy per nucleon where "nucleon" is a collective name for neutrons and protons, sometimes called the mass defect per nucleon.
- The difference in mass is equivalent to the energy released in forming the nucleus.



$$B.E. \equiv -[(m_X - Zm_p + (A - Z)m_n)]c^2 = -\Delta mc^2$$

$\Delta m < 0$: released energy
(exothermic or exoergic)

Fusion in Nature

- **Fusion reactions by which stars convert hydrogen to helium**
 - The PP (proton-proton) chain: in stars the mass of the Sun and less
 - The CNO cycle (Bethe-Weizsäcker-cycle): in more massive stars

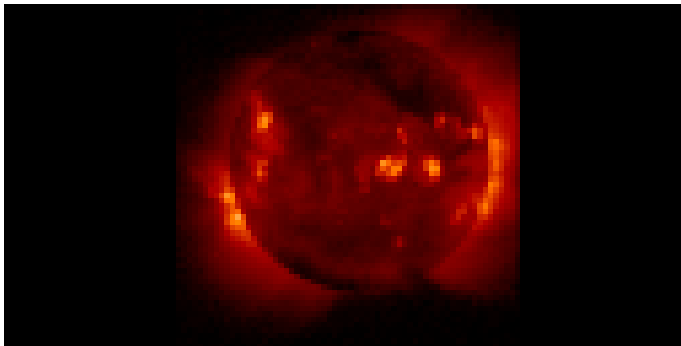
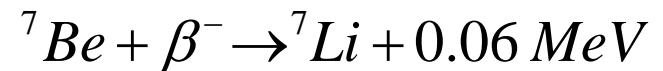
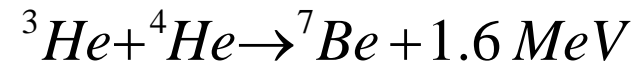
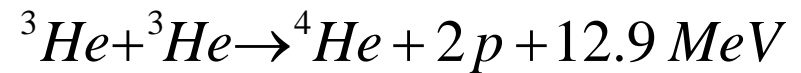
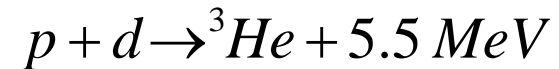
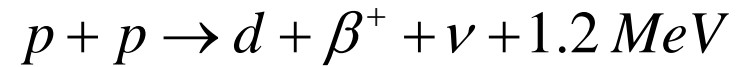
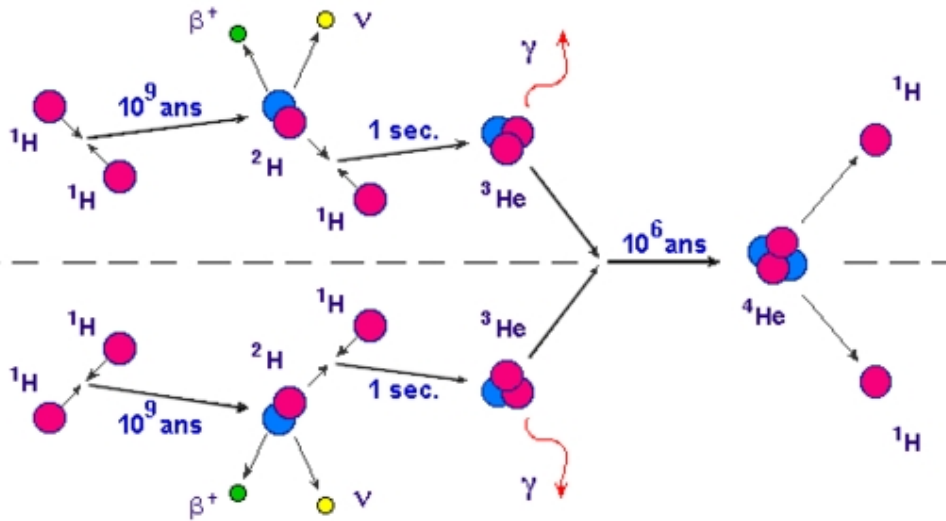


Nobel prize in physics 1967
“for his contribution to the
theory of nuclear reactions,
especially his discoveries
concerning the energy
production in stars”

Hans Albrecht Bethe
(1906. 7. 2 – 2005. 3. 6)

Fusion in Nature

• The PP (Proton-Proton) Chain

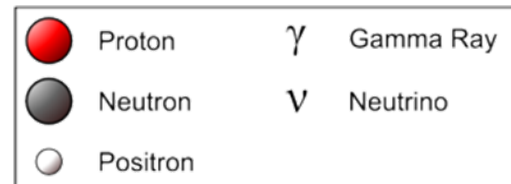
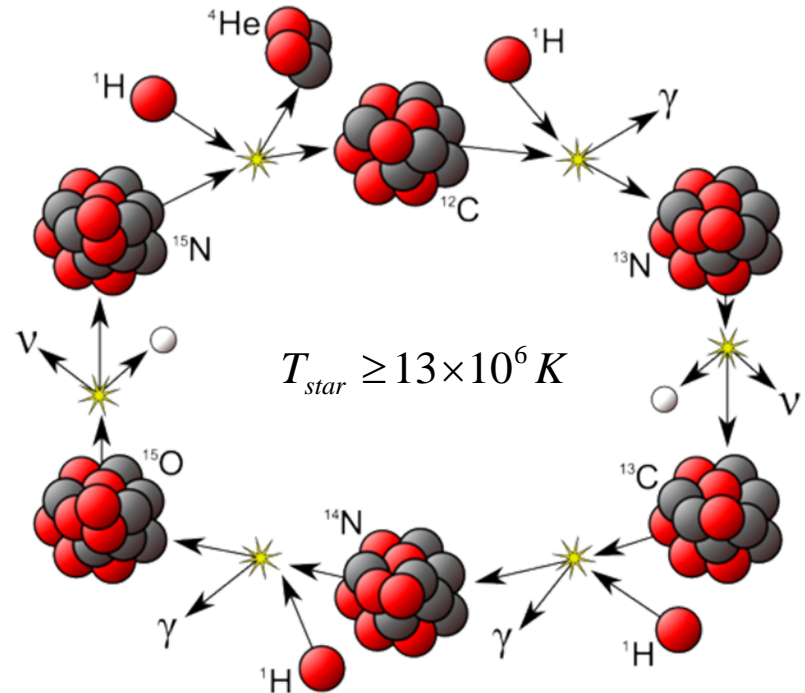
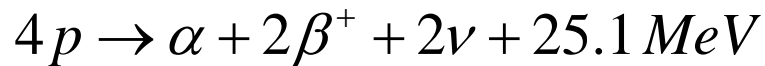
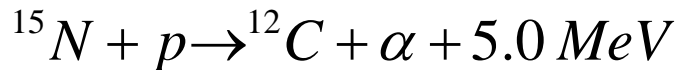
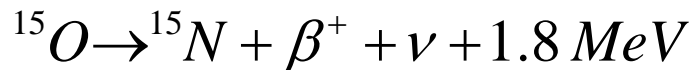
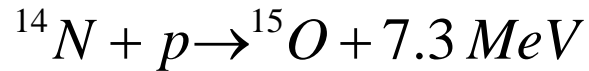
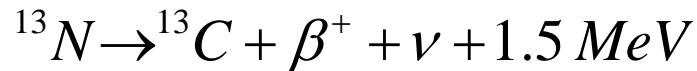
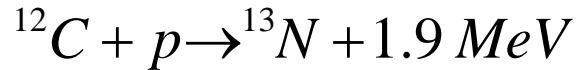


Sun $\leq 15 \times 10^6 \text{ K}$

- Only 1.7% of ${}^4\text{He}$ nuclei being produced in the Sun are born in the CNO cycle

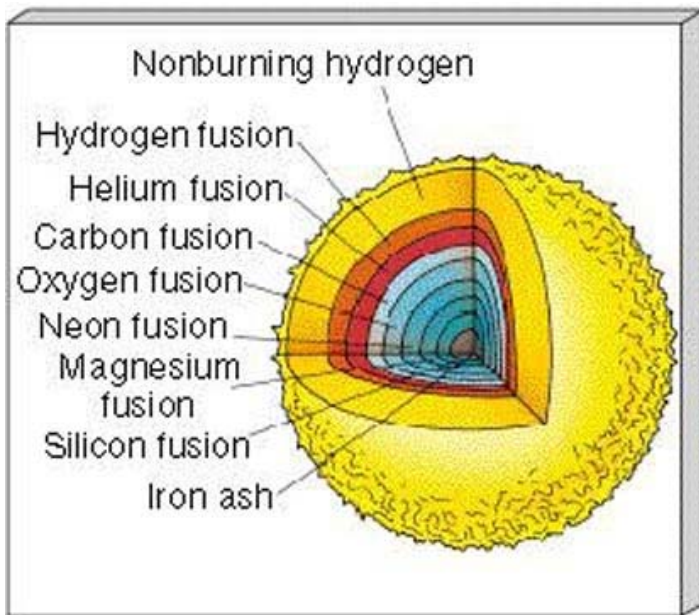
Fusion in Nature

• CNO (Carbon-Nitrogen-Oxygen) Cycle

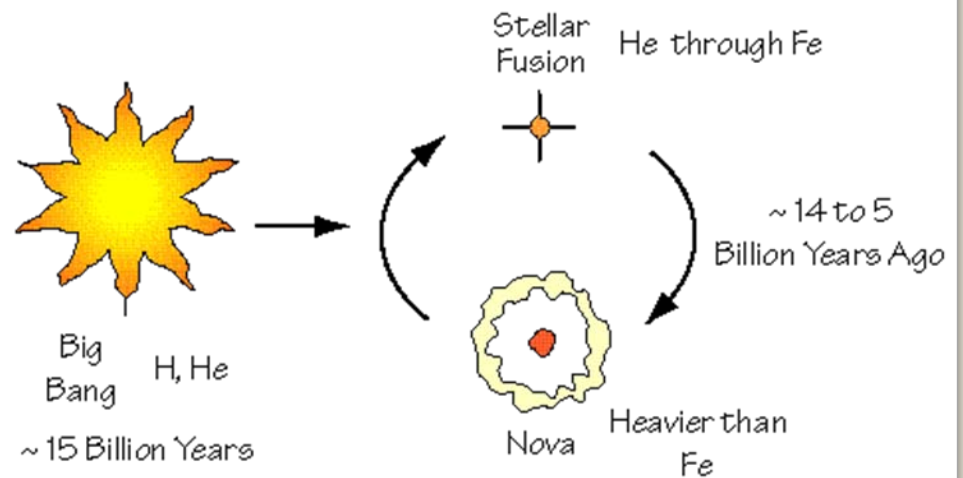


Fusion in Nature

Layers of Fusion in a Star



■ The Universe and the Formation of the Elements



<http://jconwell.wordpress.com/2009/07/20/formation-of-the-elements/>

http://eqseis.geosc.psu.edu/~cammon/HTML/Classes/IntroQuakes/Notes/earth_origin_lecture.html

Fusion in Nature

Galaxy Evolution and Merger

from 16 million to 13.7 billion years old

AMR simulation, 250 million light year region

scientific simulation

Brian O'Shea

Michigan State University

Michael Norman

University of California, San Diego

Dynamic universe birth of the milky way galaxy

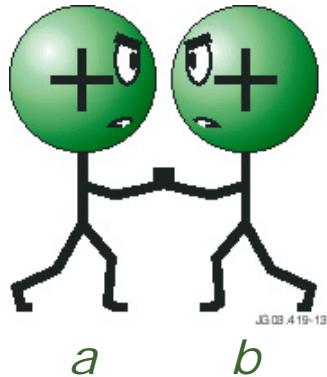
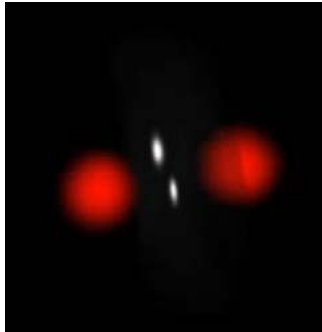
Physical Characterisation of Fusion Reaction



- (ab) : a complex short-lived dynamic state which disintegrates into products d and e .

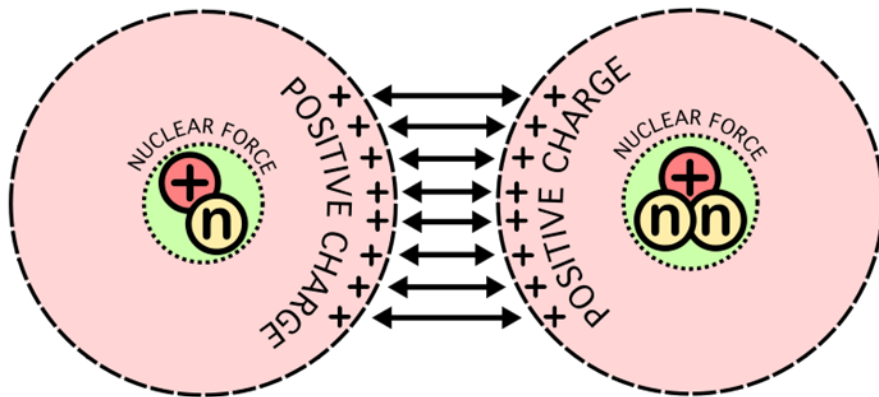
→ The energetics are determined according to nucleon kinetics analysis, with nuclear excitation and subsequent gamma ray emission known to play a comparatively small role in fusion processes at the energies of interest envisaged for fusion reactors.

Physical Characterisation of Fusion Reaction



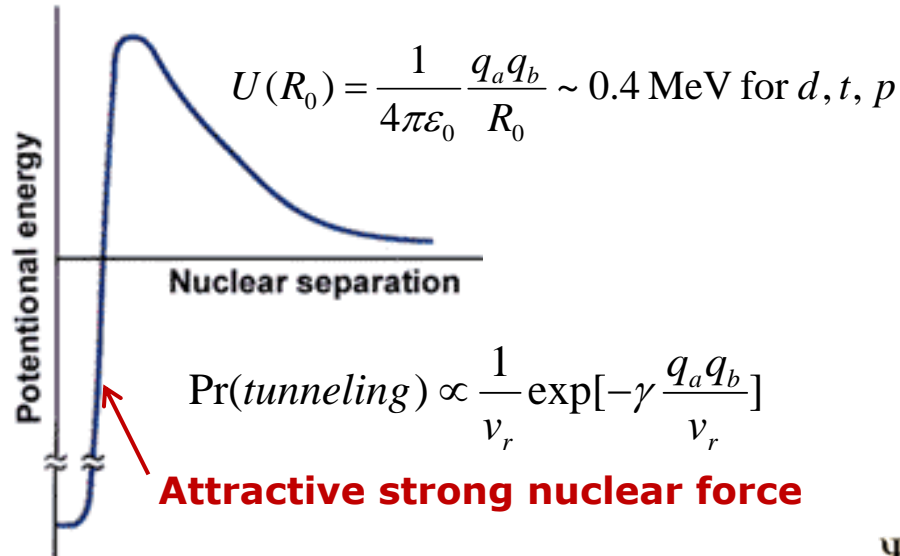
$$F_{g,a} = -G \frac{m_a m_b}{r^3} \vec{r}$$

$$F_{c,a} = \frac{1}{4\pi\epsilon_0} \frac{q_a q_b}{r^3} \vec{r}$$

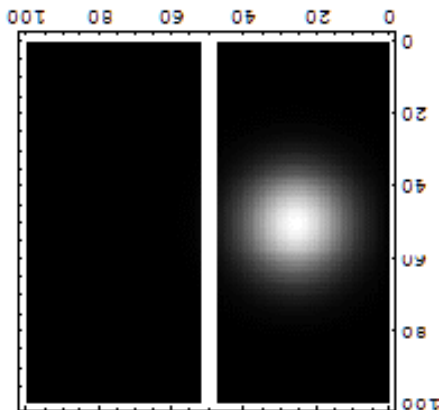
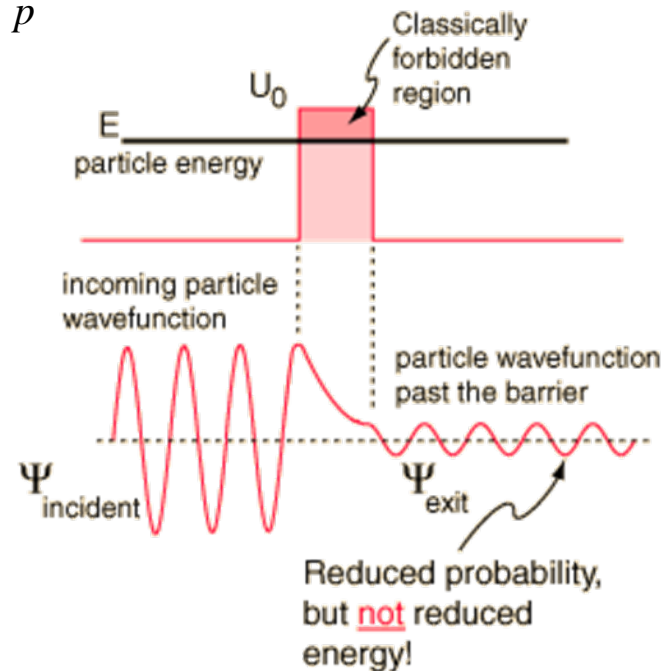


- The electrostatic force caused by positively charged nuclei is very strong over long distances, but at short distances the nuclear force is stronger.
- As such, the main technical difficulty for fusion is getting the nuclei close enough to fuse. Distances not to scale.

Physical Characterisation of Fusion Reaction



Finite potential barrier



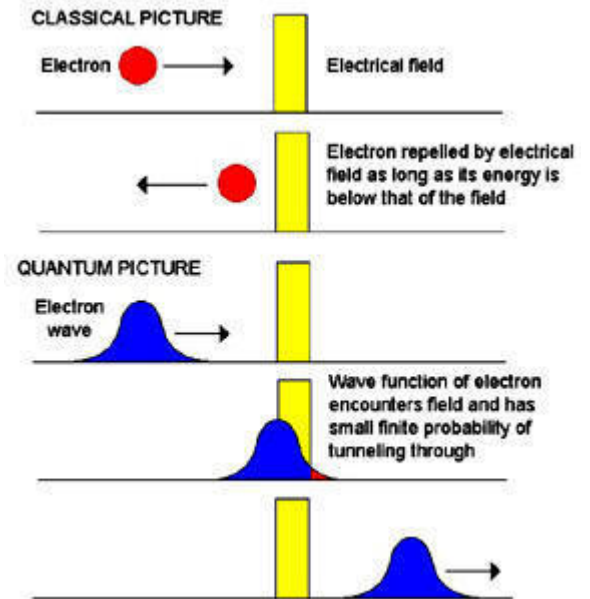
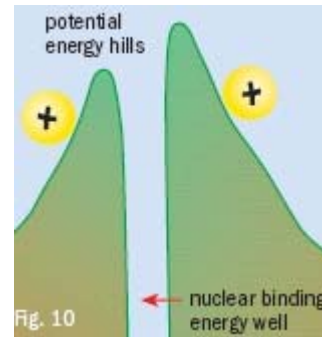
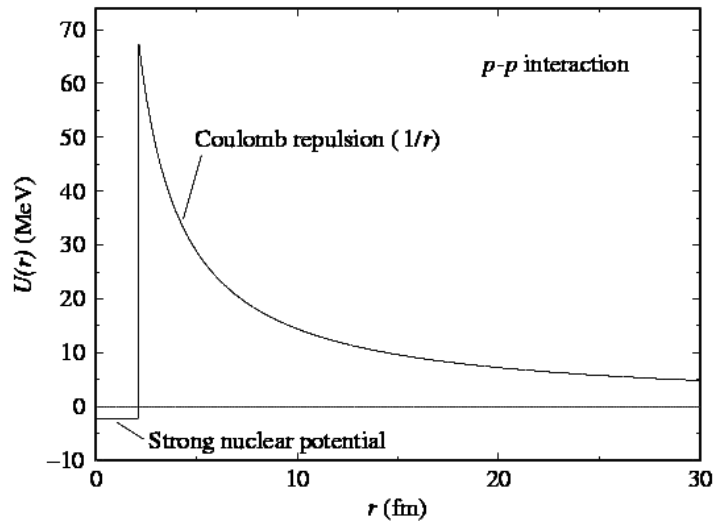
$$E \ll V_0$$

$$|T|^2 = \frac{(2k_0/\kappa)^2}{(2k_0/\kappa)^2 + (\kappa^2 + k_0^2)^2 \sinh^2 \kappa a}$$

$$\cong \left(\frac{4k_0\kappa}{\kappa^2 + k_0^2} \right)^2 e^{-2\kappa a} \neq 0$$

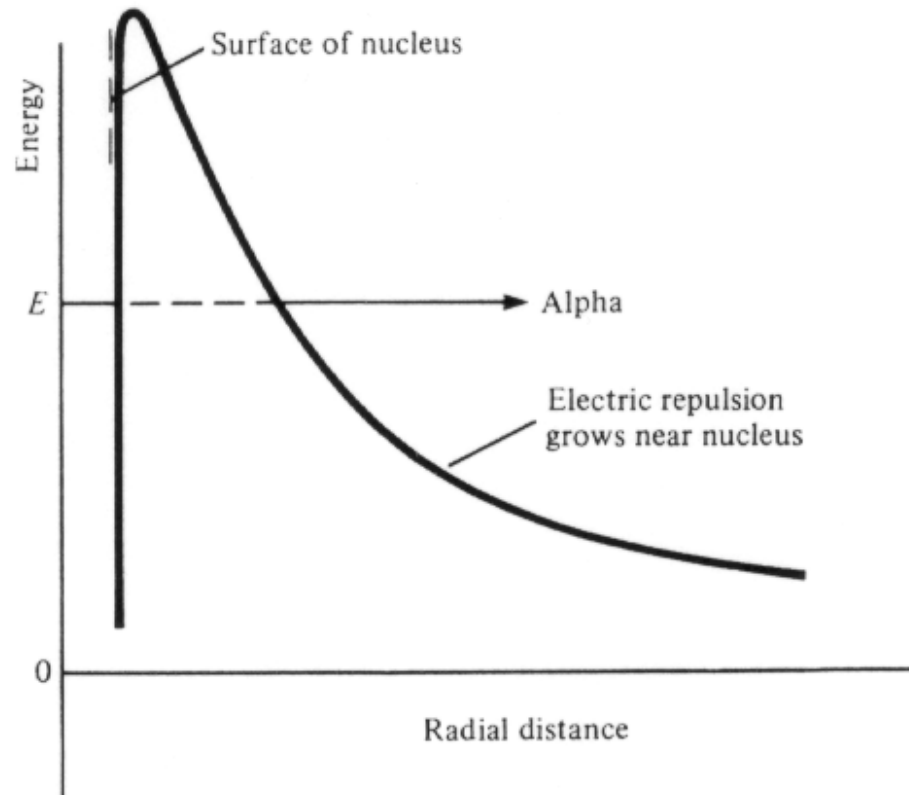
Reflection and tunneling of an electron wavepacket directed at a potential barrier

Physical Characterisation of Fusion Reaction

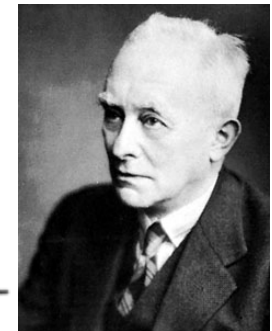


Physical Characterisation of Fusion Reaction

Potential barrier around a uranium nucleus presented to an alpha particle. The central well is due to the average nuclear attraction of all the nucleons and the hill is due to the electric repulsion of the protons. Alpha particles with energy E trapped inside the nuclear well may still escape to become alpha rays, by quantum mechanically tunnelling through the barrier.



George Gamow
(1904-1968)

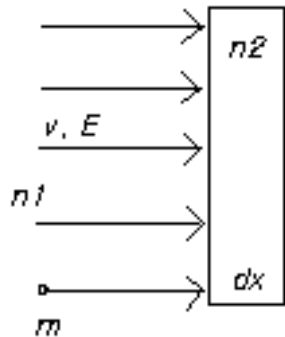


Max Born
(1882-1970)

- By 1928, George Gamow had solved the theory of the alpha decay of a nucleus via tunneling. After attending a seminar by Gamow, Max Born recognized the generality of quantum-mechanical tunneling. (Max Born, Nobel Prize in Physics 1954)

Fusion Reaction Cross Sections

- Beam-target collisions (Binary interactions)



- For fixed target

$$m = m_1, \quad v = v_1, \quad E = m_1 v_1^2 / 2$$

- For moving target

$$m = m_r, \quad v = |v_1 - v_2|, \quad E = E_{CM}$$

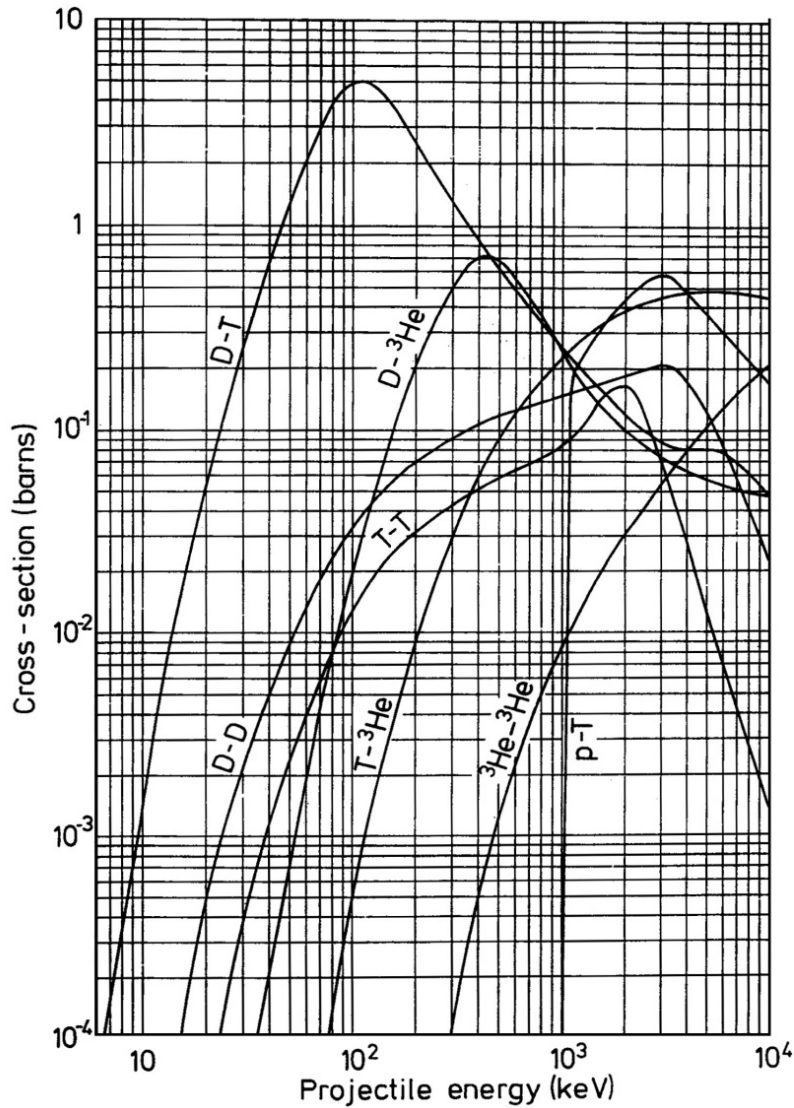
$$dn_1 = -\sigma_{12}(E)n_1n_2dx$$

- Fusion cross section for low energy $E_{CM} < U(R_0)$ by quantum mechanical tunneling process:

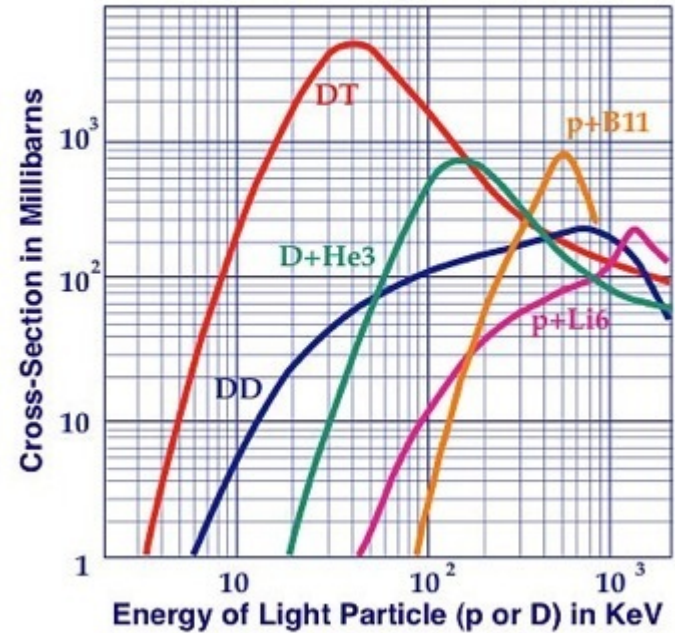
$$\sigma_{12}(E) = \frac{A}{E} e^{-B/\sqrt{E}} \quad \text{Gamow theory (1938)}$$

$$A = \text{const.}, \quad B = 2^{-1/2} \pi m_r^{1/2} Z_1 Z_2 e^2 / h \epsilon_0$$

Fusion Reaction Cross Sections



Fusion Reaction Cross-Sections
Particles Have Equal Momentum



Fusion Reaction Rate Parameter (Reactivity or σ - v Parameter)

- σ - v parameter

$$\langle \sigma v \rangle_{ab} = \int \int_{v_a v_b} \sigma_{ab}(|v_a - v_b|) |v_a - v_b| F_a(v_a) F_b(v_b) d^3 v_a d^3 v_b$$

$$dn_1 = -\sigma_{12}(E) n_1 n_2 dx$$

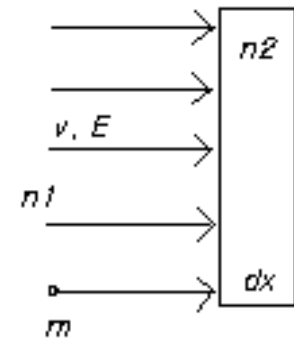
- Fusion reaction rate density

$$dR_{fu} \equiv \frac{d}{dt}(-dn) = dn_1 dn_2 \sigma_{12}(v) v$$

$$dn_1 = n_1 f_1(v_1) d^3 v_1$$

$$dn_2 = n_2 f_2(v_2) d^3 v_2$$

$f_{1,2}$: normalised distribution function



$$R_{fu} = \int \int_{v_a v_b} \sigma_{fu}(|v_a - v_b|) |v_a - v_b| N_a F_a(v_a) N_b F_b(v_b) d^3 v_a d^3 v_b$$

$$= N_a N_b \int \int_{v_a v_b} \sigma_{fu}(|v_a - v_b|) |v_a - v_b| F_a(v_a) F_b(v_b) d^3 v_a d^3 v_b$$

$$F_x(v_x) \rightarrow M_x(v_x)$$

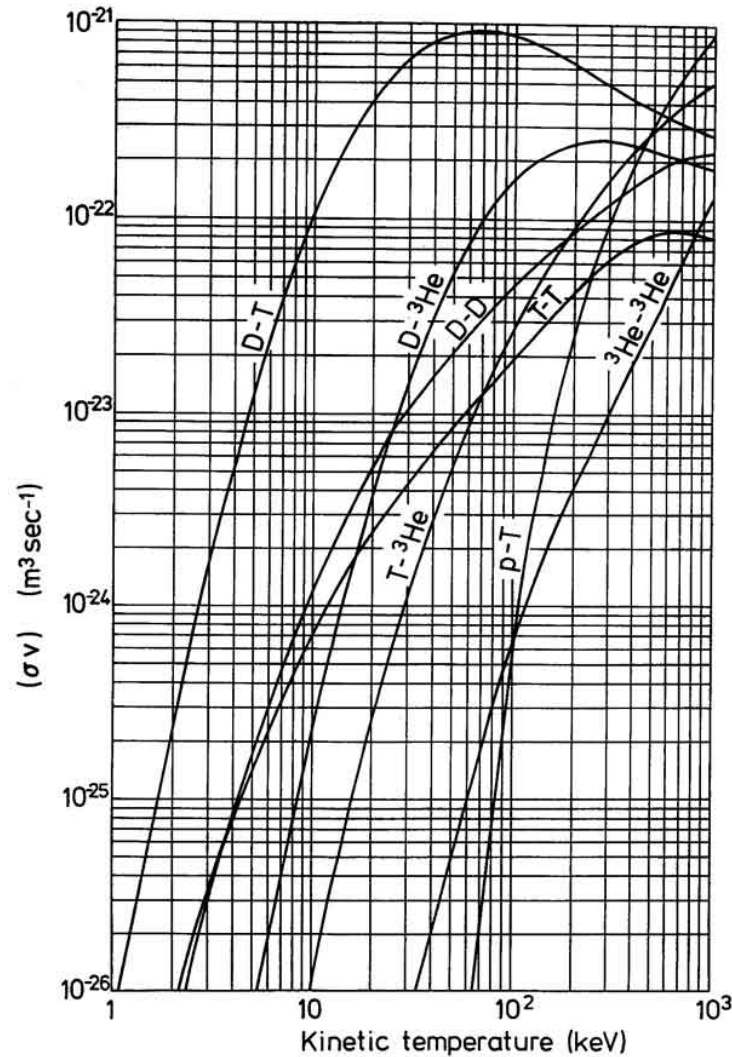
Thermodynamic equilibrium

$$R_{fu} = N_a N_b \langle \sigma v \rangle_{ab}$$

- Fusion power density

$$P_{fu} = R_{fu} Q_{fu} = N_a N_b \langle \sigma v \rangle_{ab} Q_{fu}$$

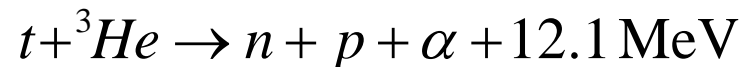
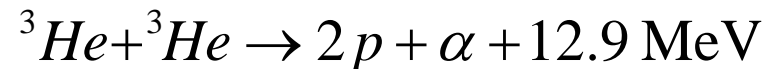
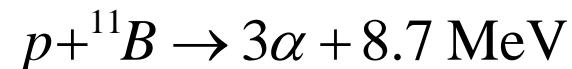
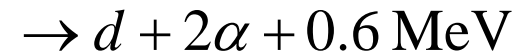
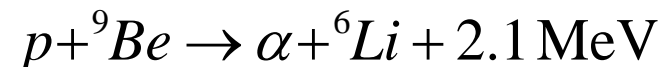
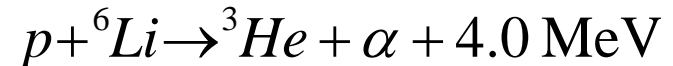
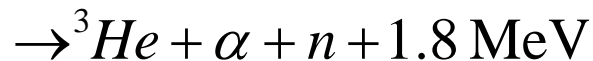
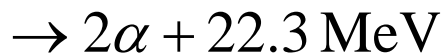
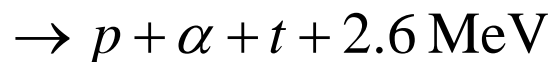
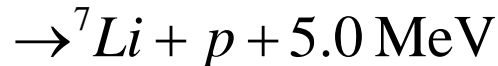
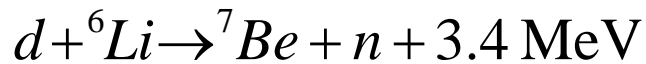
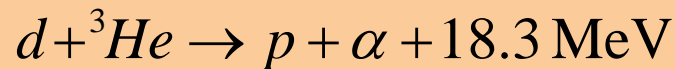
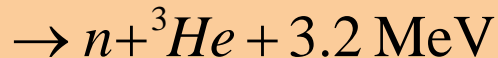
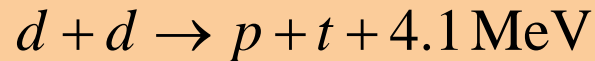
Fusion Reaction Rate Parameter (Reactivity or $\sigma\text{-}v$ Parameter)



Both species at the same temperatures

Fusion Fuels

- Possible fusion reactions



Fusion Fuels

- **Choice of a fusion reaction as a fuel in a fusion reactor**
 - Availability of fusion fuels
 - Requirements for attaining a sufficient reaction rate density
- **D-T reaction: 1st generation** $d + t \rightarrow n + \alpha + 17.6 \text{ MeV}$
 - Considered for the first generation of fusion reactors
 - Ample supply of deuterium: $d/(p+d) \sim 1/6700$ in the world's oceans, fresh water lakes, rivers
 - Scarce of tritium: radioactive β^- decay with a half life of 12.3 years.
total steady state atmospheric and oceanic
quantity produced by cosmic radiation $\sim 50 \text{ kg}$

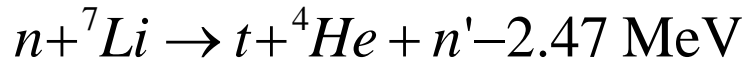
Fusion Fuels

- **D-T reaction: 1st generation**

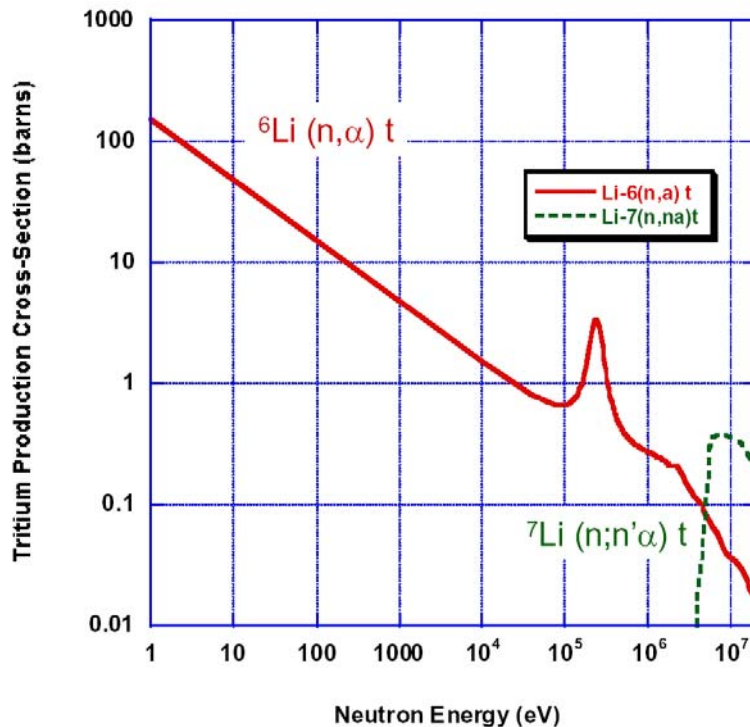
- Tritium breeding



7.42% of natural Li



92.58% of natural Li



The ${}^7\text{Li}(n,n'\alpha)t$ reaction is a threshold reaction and requires an incident neutron energy in excess of 2.8 MeV.

D-T Burn

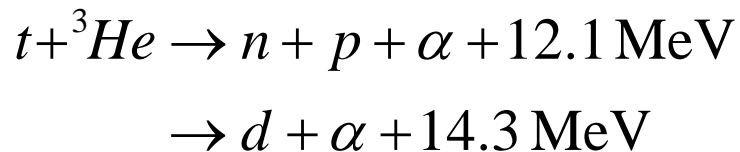
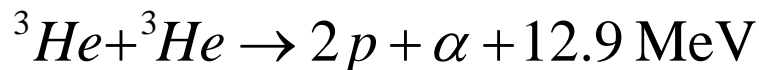
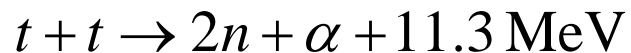
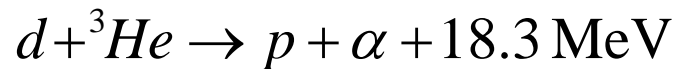
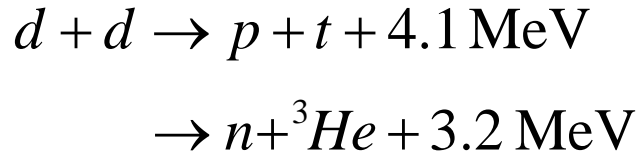
- $d+t \rightarrow n+\alpha$, $Q_{dt}=17.6$ MeV

$$\begin{aligned}R_{dt}(\vec{r}, t) &= \int \int N_d(\vec{r}, \vec{v}_d, t) N_t(\vec{r}, \vec{v}_t, t) \sigma_{dt}(|\vec{v}_d - \vec{v}_t|) |\vec{v}_d - \vec{v}_t| d^3 v_d d^3 v_t \\&= N_d^* N_t^* \int \int f_d(\vec{r}, \vec{v}_d, t) f_t(\vec{r}, \vec{v}_t, t) \sigma_{dt}(|\vec{v}_d - \vec{v}_t|) |\vec{v}_d - \vec{v}_t| d^3 v_d d^3 v_t \\&= N_d(\vec{r}, t) N_t(\vec{r}, t) \frac{\int \int f_d(\vec{r}, \vec{v}_d, t) f_t(\vec{r}, \vec{v}_t, t) \sigma_{dt}(|\vec{v}_d - \vec{v}_t|) |\vec{v}_d - \vec{v}_t| d^3 v_d d^3 v_t}{\int \int f_d(\vec{r}, \vec{v}_d, t) f_t(\vec{r}, \vec{v}_t, t) d^3 v_d d^3 v_t} \\&= N_d(\vec{r}, t) N_t(\vec{r}, t) \langle \sigma v \rangle_{dt}(\vec{r}, t)\end{aligned}$$

$$P_{dt}(\vec{r}, t) = R_{dt}(\vec{r}, t) Q_{dt} = N_d(\vec{r}, t) N_t(\vec{r}, t) \langle \sigma v \rangle_{dt}(\vec{r}, t) Q_{dt}$$

D-D Burn Modes

- D-D reactions and side reactions



D-D Burn Modes

- PURE-D Mode



$$R_{dd,t} = \frac{N_d^2}{2} \langle \sigma v \rangle_{dd,t}$$

$$R_{dd,{}^3\text{He}} = \frac{N_d^2}{2} \langle \sigma v \rangle_{dd,{}^3\text{He}}$$

$$\langle \sigma v \rangle_{dd} = \langle \sigma v \rangle_{dd,t} + \langle \sigma v \rangle_{dd,{}^3\text{He}}$$

$$\langle \sigma v \rangle_{dd,t} \approx \langle \sigma v \rangle_{dd,{}^3\text{He}} \approx \frac{1}{2} \langle \sigma v \rangle_{dd}$$

At temperatures
of common interest

D-D Burn Modes

	a_1	a_2	...	a_x	...	a_{Na}
b_1						
b_2						
...						
b_y				(a_x, b_y)		
...						
b_{Nb}						

Interaction between N_a a-type and N_b b-type particles

$$R_{ab} = N_a N_b \langle \sigma v \rangle_{ab}$$

	a_1	a_2	...	a_x	...	a_{Na}
a_1						
a_2						
...						
a_x				(a_x, a_x)		
...						
a_{Na}						

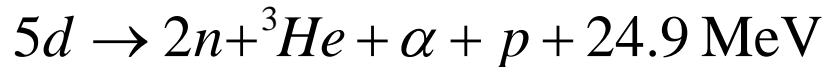
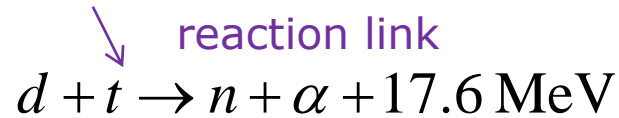
Interaction between N_a a-type particles

$$R_{aa} = \frac{N_a(N_a - 1)}{2} \langle \sigma v \rangle_{aa}$$

$$\approx \frac{N_a^2}{2} \langle \sigma v \rangle_{aa}$$

D-D Burn Modes

- Semi-Catalyzed-D cycle (SCAT-D Mode)



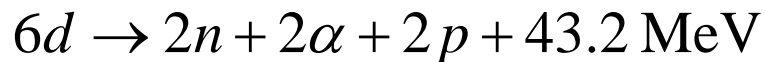
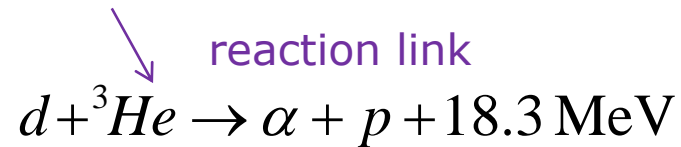
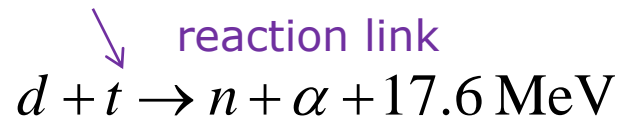
Providing $R_{dd,t} = R_{dt}$

$$\frac{N_d^2}{2} \langle \sigma v \rangle_{dd,t} = N_d N_t \langle \sigma v \rangle_{dt}$$

$$\frac{N_t}{N_d} = \frac{1 \langle \sigma v \rangle_{dd,t}}{2 \langle \sigma v \rangle_{dt}} \approx \frac{1 \langle \sigma v \rangle_{dd}}{4 \langle \sigma v \rangle_{dt}}$$

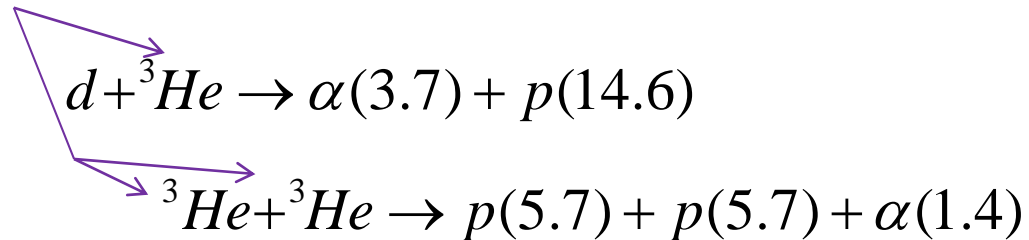
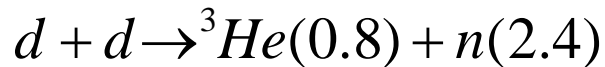
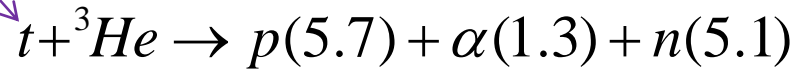
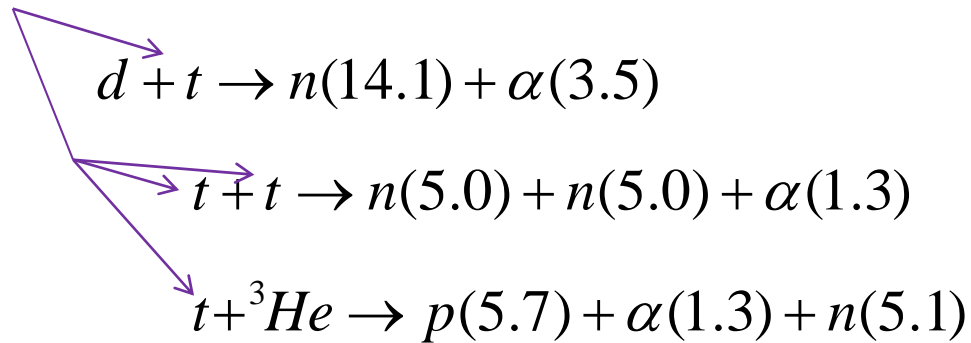
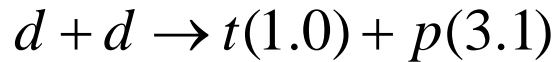
D-D Burn Modes

- Catalyzed-D cycle (CAT-D Mode)

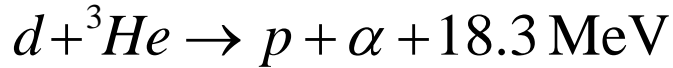


D-D Burn Modes

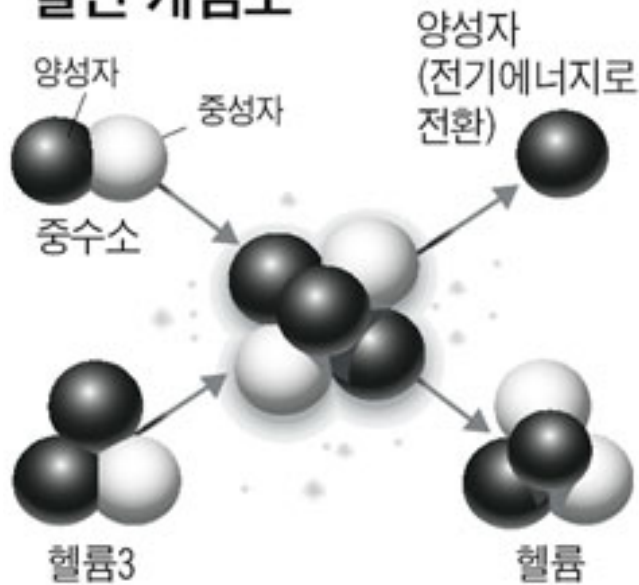
- General D-D initiated fusion linkage processes



D-³He Fusion






헬륨3를 이용한 핵융합 발전 개념도



헬륨3 차세대 핵융합 발전의 연료.

헬륨3의 원자는 양성자 2개와 중성자 1개로 이뤄져 있으며, 중수소(양성자 1개 중성자 1개)와 핵융합을 하면 정상적인 헬륨 원자(양성자 2개, 중성자 2개)가 되면서 강한 에너지를 가진 양성자를 방출한다.

1000MW급 발전소 가동을 위한
연료별 소모량 단위:kg/day

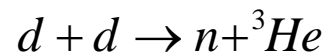
	석탄	8,640,000
	우라늄(235)	3
	헬륨3	0.2

D-³He Fusion

- An attainable “clean” fusion reaction, direct energy conversion
 - Tritium, neutron: problems of radiological safety, first wall endurance, shielding and induced radioactivity
- Higher reaction temperature required
- More severe Bremsstrahlung radiation
- Scarce ³He: ${}^3\text{He}/({}^3\text{He}+{}^4\text{He})\sim 10^{-6}$

cf) Lunar Rock

$$t \rightarrow {}^3\text{He} + \beta^-, \quad \tau_{1/2} = 12.3 \text{ years}$$



$$d + {}^3\text{He} \rightarrow p + \alpha + 18.3 \text{ MeV}, \quad R_{d^3\text{He}} = \langle \sigma v \rangle_{d^3\text{He}} N_d N_{^3\text{He}}$$

D-³He Fusion

SCIENCE

Mining The Moon

An Apollo astronaut argues that with its vast stores of nonpolluting nuclear fuel, our lunar neighbor holds the key to Earth's future.

BY HARRISON H. SCHMITT
ILLUSTRATION BY PAUL DIMARE

Apollo 17 astronaut Harrison Schmitt left the moon 32 years ago with 244 pounds of rocks and an abiding desire to see humankind continue its exploration of space. Now, in an exclusive essay for POPULAR MECHANICS, Schmitt explains why the time is right for America to return.

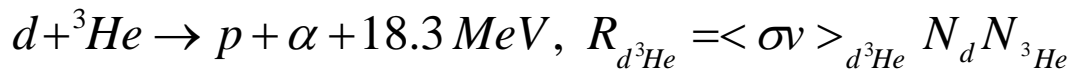
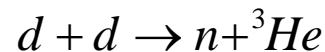
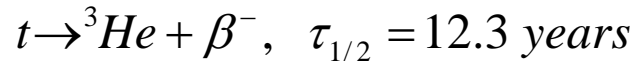


FUTURE MINERS: Robotic equipment would scrape and refine lunar soil. Helium-3 would be sent to Earth aboard a future space shuttle or perhaps be shot from an electric rail gun.

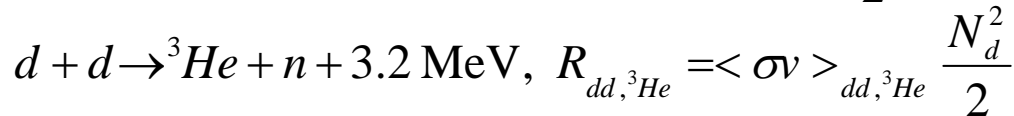
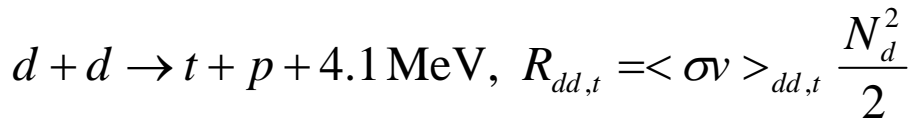
D-³He Fusion

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cf) Lunar Rock



unclean side reactions



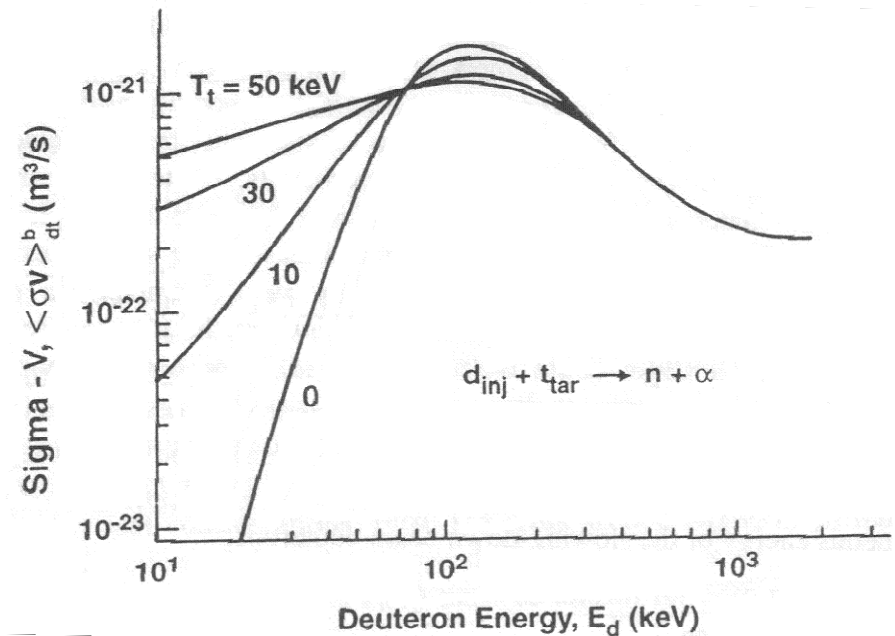
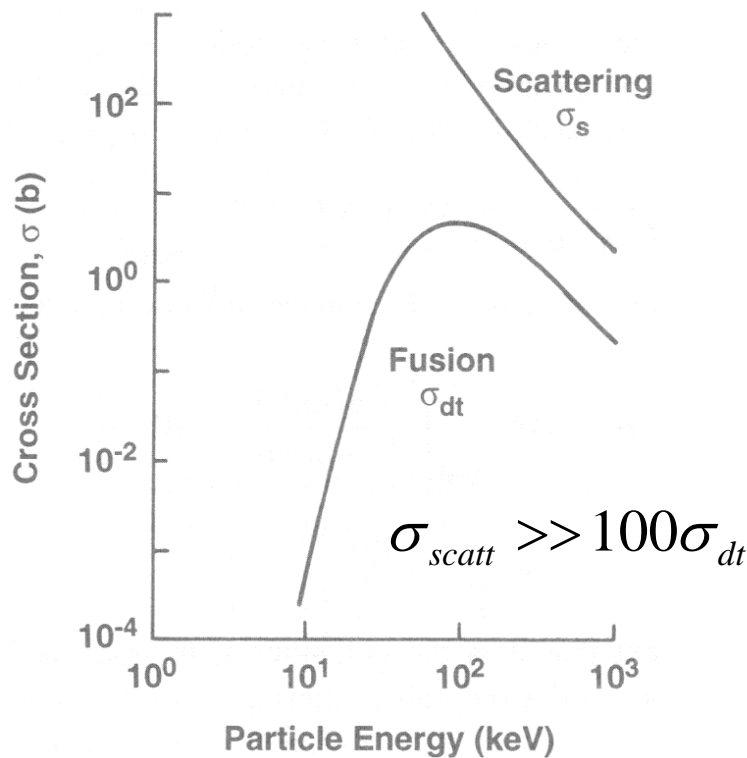
Control on high temperature and ³He and d fuel ions

$$\frac{R_{d^3\text{He}}}{R_{dd,t}} = 2 \frac{\langle \sigma v \rangle_{d^3\text{He}} N_{{}^3\text{He}}}{\langle \sigma v \rangle_{dd,t} N_d}$$

$$\frac{R_{d^3\text{He}}}{R_{dd,{}^3\text{He}}} = 2 \frac{\langle \sigma v \rangle_{d^3\text{He}} N_{{}^3\text{He}}}{\langle \sigma v \rangle_{dd,{}^3\text{He}} N_d}$$

Beam-target Fusion

- Beam-target collisions (Binary interactions)



- loss energy \gg fusion energy
- Fusion by beam-target collisions are not proper for practical energy-producing fusion reactors

Confinement needed!