

# Four-wave mixing I

Dr Yoonchan Jeong

School of Electrical Engineering, Seoul National University

Tel: +82 (0)2 880 1623, Fax: +82 (0)2 873 9953

Email: [yunchan@snu.ac.kr](mailto:yunchan@snu.ac.kr)

# Tensor notation for the third-order susceptibility

Constitutive relations for E-field:

$$\mathbf{D} = \varepsilon \mathbf{E} = \varepsilon_0 \mathbf{E} + \mathbf{P},$$

$$\mathbf{P} = \varepsilon_0 \chi \mathbf{E} = \varepsilon_0 \left( \chi^{(1)} \mathbf{E} + \chi^{(2)} \mathbf{E}\mathbf{E} + \underline{\chi^{(3)}} \mathbf{E}\mathbf{E}\mathbf{E} + \dots \right)$$

Tensor notation for  $\chi^{(3)}$ :

$$\chi^{(3)} \rightarrow \chi_{ijkl}^{(3)}(\omega_4 = \omega_1 + \omega_2 + \omega_3)$$

$$P_{i,NL} \leftarrow \sum_{j,k,l} \chi_{ijkl}^{(3)} E_j E_k E_l$$

*Note that we normally omit summation notation.*

# Tensor nature of the third-order susceptibility

How many elements for  $\chi_{ijkl}^{(3)}$ ? *81 elements*

For isotropic media:

$$\chi_{1111}^{(3)} = \chi_{2222}^{(3)} = \chi_{3333}^{(3)},$$

$$\chi_{1122}^{(3)} = \chi_{1133}^{(3)} = \chi_{2211}^{(3)} = \chi_{2233}^{(3)} = \chi_{3311}^{(3)} = \chi_{3322}^{(3)},$$

$$\chi_{1212}^{(3)} = \chi_{1313}^{(3)} = \chi_{2323}^{(3)} = \chi_{2121}^{(3)} = \chi_{3131}^{(3)} = \chi_{3232}^{(3)},$$

$$\chi_{1221}^{(3)} = \chi_{1331}^{(3)} = \chi_{2112}^{(3)} = \chi_{2332}^{(3)} = \chi_{3113}^{(3)} = \chi_{3223}^{(3)}.$$

*→ 21 nonzero elements*

In addition:

$$\chi_{1111}^{(3)} = \chi_{1122}^{(3)} + \chi_{1212}^{(3)} + \chi_{1221}^{(3)} \quad \rightarrow \textit{Why?}$$

*→ In the compact form:*  $\chi_{ijkl}^{(3)} = \chi_{1122}^{(3)} \delta_{ij} \delta_{kl} + \chi_{1212}^{(3)} \delta_{ik} \delta_{jl} + \chi_{1221}^{(3)} \delta_{il} \delta_{jk}$

# Cross-phase modulation (XPM)

Suppose that there are two monochromatic waves:

$$\mathbf{E}(r, t) = \frac{1}{2} \hat{x} [E_1 \exp(-i\omega_1 t) + E_2 \exp(-i\omega_2 t)] + c.c.$$

Induced nonlinear polarisation:

$$\begin{aligned} \mathbf{P}_{NL}(r, t) &= \varepsilon_0 \chi^{(3)} \mathbf{E} \mathbf{E} \mathbf{E} = \frac{1}{2} \hat{x} \{ P_{NL}(\omega_1) \exp(-i\omega_1 t) + P_{NL}(\omega_2) \exp(-i\omega_2 t) \\ &+ P_{NL}(2\omega_1 - \omega_2) \exp[-i(2\omega_1 - \omega_2)t] \\ &+ P_{NL}(2\omega_2 - \omega_1) \exp[-i(2\omega_2 - \omega_1)t] \} + c.c \\ &+ \dots \end{aligned}$$

where

$$P_{NL}(\omega_1) = \chi_{eff} (|E_1|^2 + 2|E_2|^2) E_1,$$

$$P_{NL}(\omega_2) = \chi_{eff} (|E_2|^2 + 2|E_1|^2) E_2,$$

$$\chi_{eff} = \frac{3\varepsilon_0 \chi_{xxxx}^{(3)}}{4}.$$

# Nonlinear pulse propagation with XPM

**NLSE:**

$$\frac{\partial A}{\partial z} + \beta_1 \frac{\partial A}{\partial t} + i \frac{\beta_2}{2} \frac{\partial^2 A}{\partial t^2} + \frac{\alpha}{2} A = i\gamma |A|^2 A \quad \text{for } \exp(-i\omega t)$$

**Coupled NLSEs:**

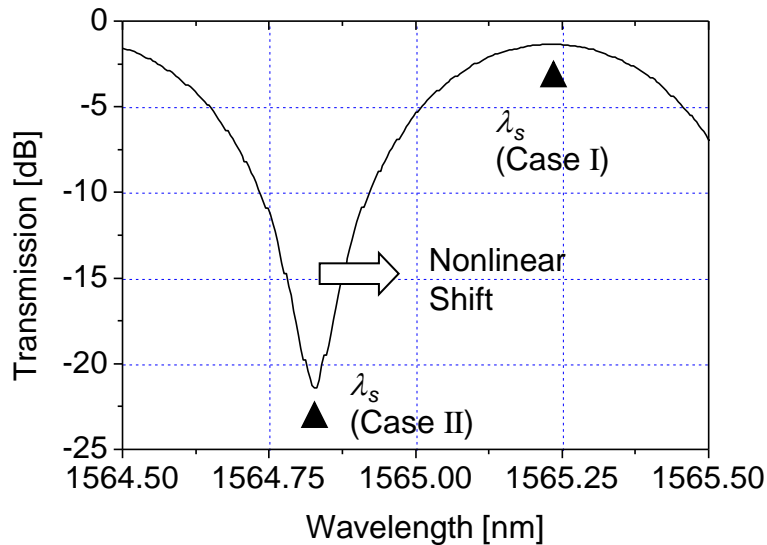
$$\begin{aligned} \frac{\partial A_1}{\partial z} + \beta_{11} \frac{\partial A_1}{\partial t} + i \frac{\beta_{21}}{2} \frac{\partial^2 A_1}{\partial t^2} + \frac{\alpha_1}{2} A_1 &= i\gamma_1 (|A_1|^2 + 2|A_2|^2) A_1, \\ \rightarrow \frac{\partial A_2}{\partial z} + \beta_{12} \frac{\partial A_2}{\partial t} + i \frac{\beta_{22}}{2} \frac{\partial^2 A_2}{\partial t^2} + \frac{\alpha_2}{2} A_2 &= i\gamma_2 (|A_2|^2 + 2|A_1|^2) A_2. \end{aligned}$$

***Modulation instability?***

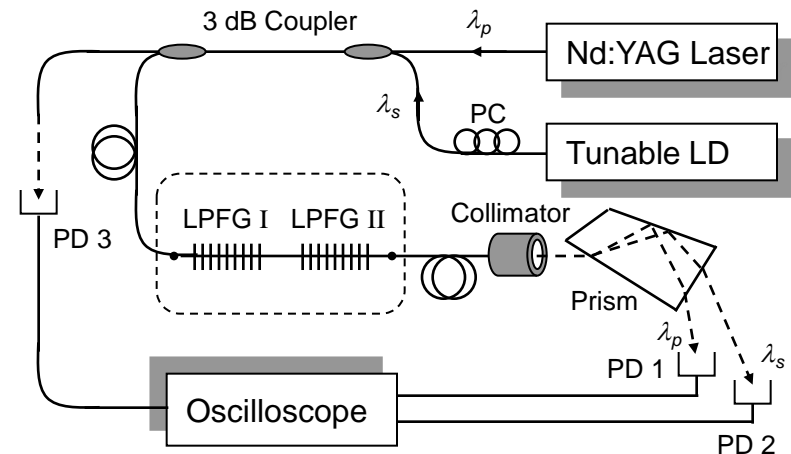
***Optical solitons?***

# All-optical switching via XPM

## Transmission spectra



## Experimental arrangement



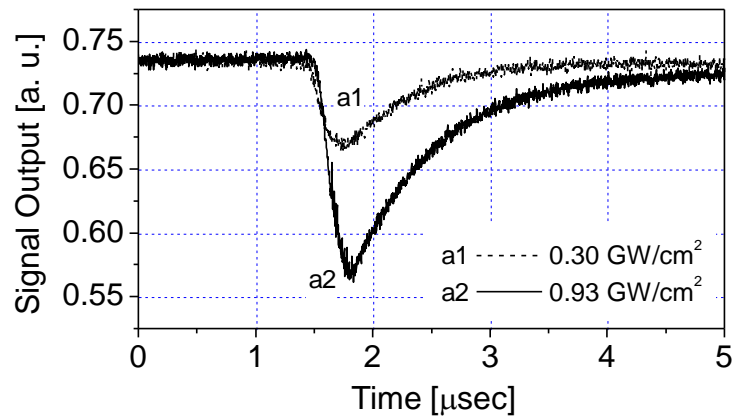
## Experiments:

Signal wave: Tunable LD @ 1565.2 nm (Case I), @ 1564.8 nm (Case II)

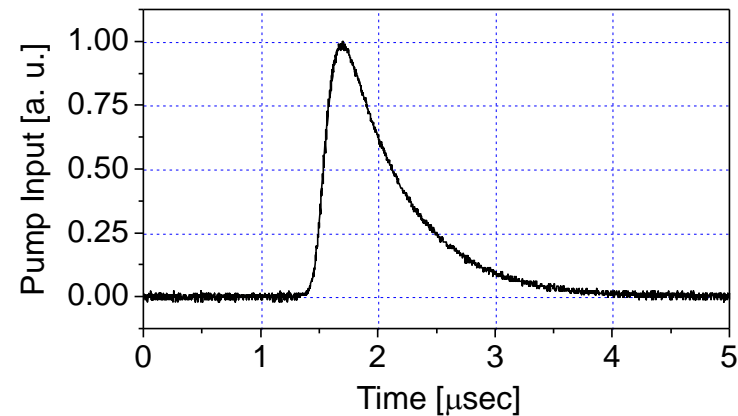
Pump wave: Q-switched Nd:YAG laser (1 kHz)

# All-optical switching via XPM

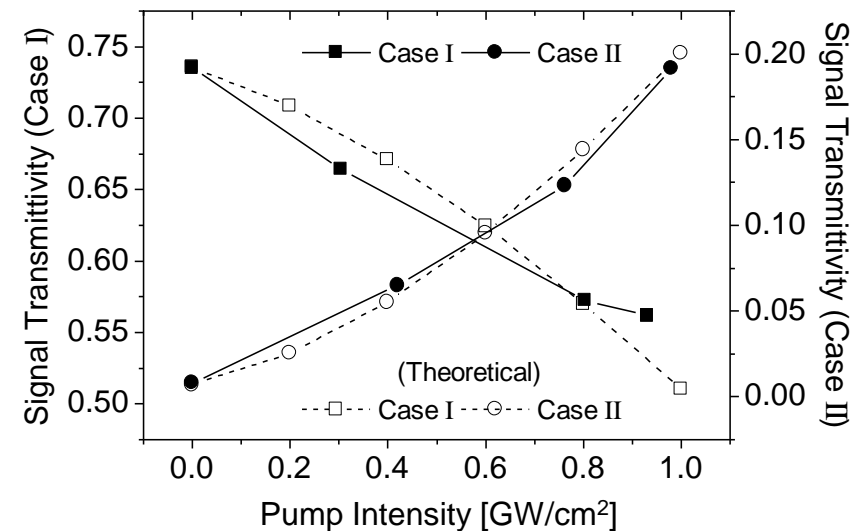
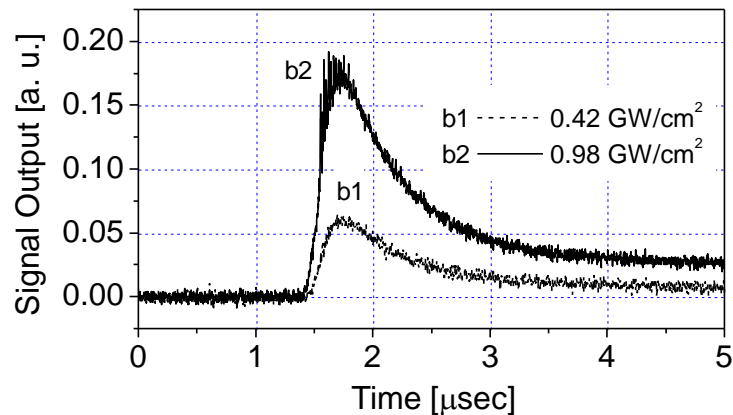
## Case I



## Pump



## Case II



$$\Delta\lambda_s \sim 0.12 \text{ nm}/(\text{GW}/\text{cm}^2)$$