# 2. Wave Propagation on Transmission Lines

# A. General Solutions of Transmission-Line Equations

### 1) Wave solutions in the phasor domain

For uniform transmission lines with time-harmonic variation  $e^{j\omega t}$ , Transmission-line equations:

$$\frac{d^2 V(z)}{dz^2} - \gamma^2 V(z) = 0$$
(8-10)

$$\frac{d^2 I(z)}{dz^2} - \gamma^2 I(z) = 0 \tag{8-11}$$

where 
$$\gamma = \alpha + j\beta = \sqrt{(R + j\omega L)(G + j\omega C)} = \sqrt{ZY}$$
 (m<sup>-1</sup>) (8-12)

$$\alpha = Re[\sqrt{(R+j\omega L)(G+j\omega C)}] \qquad (Np/m) \qquad (8-12a)$$

$$\beta = Im[\sqrt{(R+j\omega L)(G+j\omega C)}] \qquad (rad/m) \qquad (8-12b)$$

General solutions of (8-10, 11):

$$V(z) = V_o^+ e^{-\gamma z} + V_o^- e^{+\gamma z} \quad (\forall)$$
(8-33, 62)

$$I(z) = I_o^+ e^{-\gamma z} + I_o^- e^{+\gamma z} \quad (A)$$
(8-34, 63)
  
FTW BTW(reflected wave)

where unknown amplitudes (  $V_o^+, \, V_o^-, \, I_o^+, \, I_o^-$  ) are to be determined by BCs. Generally,  $V_o^+$ ,  $V_o^-$ ,  $I_o^+$ ,  $I_o^-$  are complex quantities, like  $V_o^{\pm} = |V_o^{\pm}| e^{j\phi^{\pm}}$ (8-33) in (8-8):

$$I(z) = \frac{\gamma}{(R+j\omega L)} \left[ V_o^+ e^{-\gamma z} - V_o^- e^{+\gamma z} \right]$$
(8-34)\*

### 2) Characteristic impedance

Comparison of (8-34)\* with (8-34) leads to

$$\frac{V_o^+}{I_o^+} = -\frac{V_o^-}{I_o^-} = \frac{R + j\omega L}{\gamma} \equiv Z_o = R_o + jX_o$$
(8-35, 64)

Define the Characteristic Impedance  $Z_o$  of the transmission line by

$$Z_o = \frac{R + j\omega L}{\gamma} = \frac{\gamma}{G + j\omega C} = \sqrt{\frac{R + j\omega L}{G + j\omega C}} = \sqrt{\frac{Z}{Y}} \quad (\Omega) \quad (8-38)$$

Notes)

i)  $Z_o \ {
m and} \ \gamma$  are independent of z and the length of the line, but

depends only on distributed parameters(R, L, G, C) and frequency ( $\omega$ ). ii) Phasor solution in terms of  $Z_o$  from (8-34)\*:

$$I(z) = \frac{V_o^+}{Z_o} e^{-\gamma z} - \frac{V_o^-}{Z_o} e^{+\gamma z}$$
(8-34)\*\*

#### B. Wave Characteristics on an Infinite Transmission Line

#### 1) Wave solutions

For an infinite uniform transmission line,  $\exists$  no reflection waves (BTW). Then,  $V(z) = V^+(z) = V_o^+ e^{-\gamma z}$  (8-36)

$$I(z) = I^{+}(z) = I_{o}^{+} e^{-\gamma z} = \frac{V_{o}^{+}}{Z_{o}} e^{-\gamma z}$$
(8-37)

or

$$v(z,t) = |V_o^+| e^{-\alpha z} \cos(\omega t - \beta z + \phi^+)$$
(8-36)\*

$$i(z,t) = \frac{|V_o^+|}{|Z_o|} e^{-\alpha z} \cos\left(\omega t - \beta z + \phi^+ - \phi_{Z_o}\right)$$
(8-37)\*

 $\Rightarrow$  The transmission line is characterized by two fundamental properties,  $\gamma$  and  $Z_o$  which are specified by R, L, G, C, and  $\omega$ .

#### 2) Characteristics in the lossless line

For lossless (R = 0, G = 0) or high frequency  $(\omega L \gg R, \omega C \gg G)$ ,

#### a) Propagation constant

$$(8-12) \implies \gamma = \alpha + j\beta = j\omega \sqrt{LC}$$
(8-39)

i.e., 
$$\alpha = 0$$
 (no attenuation),  $\beta = \omega \sqrt{LC}$  (8-40, 41)

(cf) For lossless unbounded medium,  $\gamma = jk$ ,  $k = \beta = \omega \sqrt{\mu \epsilon}$  (7-4, 42)

#### b) Characteristic impedance

$$(8-38) \implies Z_o = R_o + jX_o = \sqrt{L/C}$$

$$(8-43)$$

i.e., 
$$R_o = \sqrt{L/C}$$
 (constant),  $X_o = 0$  (8-44, 45)

(cf) For lossless unbounded medium,  $\eta = \sqrt{\mu/\epsilon}$  (7-14)

#### c) Phase velocity

(8-41) in (7-50) 
$$\Rightarrow u_p = \frac{\omega}{\beta} = \frac{1}{\sqrt{LC}}$$
 (constant: ind. of  $f$ ) (8-42)

 $\Rightarrow$  Distortionless (or nondispersive) line

(cf) For lossless unbounded medium, 
$$u_p = \frac{1}{\sqrt{\mu\epsilon}} = \frac{c}{\sqrt{\mu_r\epsilon_r}}$$

# 3) Characteristics in the distortionless lossy line

For the (distortionless) condition of 
$$\frac{R}{L} = \frac{G}{C}$$
 (8-46)

# a) Propagation constant

(8-46) in (8-12) 
$$\Rightarrow \gamma = \alpha + j\beta = \sqrt{C/L} (R + j\omega L)$$
 (8-47)

i.e., 
$$\alpha = R\sqrt{C/L}$$
 (attenuation),  $\beta = \omega\sqrt{LC}$  (8-48, 49)

#### b) Characteristic impedance

$$(8-46) \text{ in}(8-38) \implies Z_o = R_o + jX_o = \sqrt{L/C}$$
(8-51)

i.e., 
$$R_o = \sqrt{L/C}$$
 (constant),  $X_o = 0$  (8-52, 53)

# c) Phase velocity

(8-49) in (7-50) 
$$\Rightarrow$$
  $u_p = \frac{\omega}{\beta} = \frac{1}{\sqrt{LC}}$  (constant: ind. of  $f$ ) (8-42)

 $\Rightarrow$  Distortionless (or nondispersive) lossy line

# 4) Characteristics in the lossy line

For the lossy transmission line,



Distortionless (or nondispersive) line



Distorted (or dispersive) short line



Distorted (or dispersive) long line

For small losses ( $\omega L \gg R$ ,  $\omega C \gg G$ )

In most practical transmission lines of good conductors and very low leakage dielectrics,

$$rac{G}{C} < rac{R}{L}$$
 in (8-51)\*  $\Rightarrow$   $X_o < 0$  (capacitive reactance)

Therefore, from (8-36) and (8-37) ,

$$V(z) = Z_o I(z) = (R_o - j | X_o|) I(z)$$
  

$$\Rightarrow V(z) \text{ lags behind } I(z) \text{ by } \phi_{Z_o} = \tan^{-1} \frac{|X_o|}{R_o} V$$

#### Attenuation constant from power relation

From (8-36) and (8-37) ,

$$V(z) = V_o e^{-(\alpha + j\beta)z}, \qquad I(z) = \frac{V_o}{Z_o} e^{-(\alpha + j\beta)z}$$
 (8-54, 55)

Time-average power along the line (like time-ave Poynting vector 7-79):

$$P(z) = \mathscr{P}_{av}(z) = \frac{1}{2} Re[V(z)I^{*}(z)] = \frac{V_{o}^{2}e^{-2\alpha z}}{2} Re\left[\frac{1}{Z_{o}}\right]$$
$$= \frac{V_{o}^{2}e^{-2\alpha z}}{2} Re\left[\frac{1}{R_{o}+jX_{o}}\right] = \frac{V_{o}^{2}e^{-2\alpha z}}{2} Re\left[\frac{R_{o}-jX_{o}}{R_{o}^{2}+X_{o}^{2}}\right] = \frac{V_{o}^{2}R_{o}e^{-2\alpha z}}{2|Z_{o}|^{2}} (8-56)$$

From energy conservation law,

Decrease rate of P(z) along z = Time-ave. power loss per length  $\partial P(z) = P(z)$ 

$$-\frac{\partial P(z)}{\partial z} = P_L(z) \quad \Rightarrow \quad 2\alpha P(z) = P_L(z),$$

from which the attenuation constant can be found by

$$\alpha = \frac{P_L(z)}{2P(z)} \quad (Np/m)$$
(8-57)

For lossy line,  $P_L(z) = (1/2)(I^2R + V^2G) = (V_o^2/2|Z_o|^2)(R + G|Z_o|^2)e^{-2\alpha z}$ (8-58)

(8-58) in (8-57) : 
$$\alpha = \frac{1}{2R_o} (R + G |Z_o|^2)$$
 (8-59)

For a low loss line with  $Z_{o}\cong R_{o}=\sqrt{L/C}\,\text{,}$ 

(8-59) becomes 
$$\alpha \simeq \frac{1}{2} (R \sqrt{C/L} + G \sqrt{L/C})$$
 (cf) (7-47) (8-60)

For a distortionless lossy line with  $Z_{\!_o} = R_{\!_o} = \sqrt{L/C}$  using (8-46),

(8-60) yields 
$$\alpha = R \sqrt{\frac{C}{L}}$$
 (8-61) = (8-48)



### 1) General solutions

FTW

For finite uniform transmission lines,  $\exists$  reflection waves (BTW). General solutions of (8-10, 11):

$$V(z) = V_o^+ e^{-\gamma z} + V_o^- e^{+\gamma z}$$
(8-33, 62)

$$I(z) = I_o^+ e^{-\gamma z} + I_o^- e^{+\gamma z} = Z_o^{-1} [V_o^+ e^{-\gamma z} - V_o^- e^{+\gamma z}]$$
(8-34, 63)

where  $\frac{V_o^+}{I_o^+} = -\frac{V_o^-}{I_o^-} \equiv Z_o = R_o + jX_o$ : characteristic impedance (8-64)

BCs : (8-62, 63) at the load end (z = l) using (8-64),

$$V_{L} = V_{o}^{+} e^{-\gamma l} + V_{o}^{-} e^{+\gamma l} , \qquad I_{L} = \frac{V_{o}^{+}}{Z_{o}} e^{-\gamma l} - \frac{V_{o}^{-}}{Z_{o}} e^{+\gamma l}$$
(8-66, 67)

Solution of (8-66, 67) in (8-62, 63) with change of variable z' = l - z,

$$V(z') = \frac{I_L}{2} [(Z_L + Z_o)e^{\gamma z'} + (Z_L - Z_o)e^{-\gamma z'}] = \frac{I_L}{2} (Z_L + Z_o)e^{\gamma z'} [1 + \Gamma e^{-2\gamma z'}] (8-72, 87)$$
$$= I_L (Z_L \cosh \gamma z' + Z_o \sinh \gamma z')$$
(8-74)

$$I(z') = \frac{I_L}{2Z_o} [(Z_L + Z_o)e^{\gamma z'} - (Z_L - Z_o)e^{-\gamma z'}] = \frac{I_L}{2Z_o} (Z_L + Z_o)e^{\gamma z'} [1 - \Gamma e^{-2\gamma z'}] (8-73, 89)$$
$$= \frac{I_L}{Z_o} (Z_L \sinh \gamma z' + Z_o \cosh \gamma z')$$
(8-75)

where  $\Gamma \equiv \frac{V_o^-}{V_o^+} = \frac{Z_L - Z_o}{Z_L + Z_o} = |\Gamma| e^{j\theta_{\Gamma}} = \text{(voltage) reflection coeff. of } Z_L \quad (8-88)$ : complex value with  $|\Gamma| \le 1$ 

Notes)

i) Current reflection coeff. 
$$\equiv \frac{I_o^-}{I_o^+} = -\Gamma$$
 (out of phase)

ii) For  $Z_L = Z_o$ ,  $\Gamma = 0$  and  $V_o^- = 0$  (no reflection wave)

 $\Rightarrow$  The transmission line is said to be matched to the load.

- iii) For an open-circuit line  $(Z_L \rightarrow \infty)$ ,  $\Gamma = 1$  and  $V_o^- = V_o^+$  (in phase)
- iv) For a short-circuit line  $(Z_L = 0)$ ,  $\Gamma = -1$  and  $V_o^- = -V_o^+$  (out of phase)
- v) For  $Z_L \neq Z_o$ ,  $\exists$  standing voltage and current waves along the line,

standing-wave ratio (SWR): 
$$S = \frac{|V_{\text{max}}|}{|V_{\text{min}}|} = \frac{1+|\Gamma|}{1-|\Gamma|}$$
 or  $20 \log_{10} S$  in (dB) (8-90)  

$$\Rightarrow \quad |\Gamma| = \frac{S-1}{S+1} \qquad (8-91)$$

$$|\Gamma| = 0: matched \qquad |\Gamma| = 1: o.c. \text{ or } s.c.$$

vi) For a lossless(  $\alpha = 0, X_o = 0; \gamma = j\beta, Z_o = R_o$ ) line, (8-87, 89) become

$$V(z') = \frac{I_L}{2} (Z_L + R_o) e^{j\beta z'} \left[ 1 + |\Gamma| e^{j(\theta_{\Gamma} - 2\beta z')} \right]$$
(8-92)

$$I(z') = \frac{I_L}{2R_o} (Z_L + R_o) e^{j\beta z'} \left[ 1 - |\Gamma| e^{j(\theta_{\Gamma} - 2\beta z')} \right]$$
(8-93)

### 2) Input impedance

Impedance Z(z') looking toward the load end at z' from the load:

$$Z(z') \equiv \frac{V(z')}{I(z')} = Z_o \frac{Z_L + Z_o \tanh \gamma z'}{Z_o + Z_L \tanh \gamma z'} = Z_o \frac{1 + \Gamma e^{-2\gamma z'}}{1 - \Gamma e^{-2\gamma z'}} \quad (\Omega)$$

$$(8-77)$$

$$\frac{(8-74)}{(8-75)} \qquad \frac{(8-87)}{(8-89)}$$

Input impedance  $Z_i$  looking into the line from the source at z' = l:

$$Z_{i} \equiv \frac{V_{i}}{I_{i}} = (Z)_{\substack{z'=l\\z=0}}$$

$$= Z_{o} \frac{Z_{L} + Z_{o} \tanh \gamma l}{Z_{o} + Z_{L} \tanh \gamma l} = Z_{o} \frac{1 + \Gamma e^{-2\gamma l}}{1 - \Gamma e^{-2\gamma l}} \quad (8-78)$$

Note) When  $Z_L = Z_o$ ,  $Z_i = Z_o$  irrespective of the length l

 $\Rightarrow$  The transmission line is matched

For a lossless(  $\alpha=0,~X_{o}=0~;~\gamma=j\beta,~Z_{o}=R_{o}$  ) line, (8-78) become

$$Z_{i} = R_{o} \frac{Z_{L} + jR_{o} \tan\beta l}{R_{o} + jZ_{L} \tan\beta l}$$
(8-79)

From the standpoint of the generator circuit,

$$V_{i} = Z_{i}I_{i} = \frac{Z_{i}V_{g}}{Z_{g} + Z_{i}} = V_{g} - I_{i}Z_{g}$$
(8-94)

If  $Z_L \neq Z_o$  but  $Z_g = Z_o$ , reflected at the load and ending at the generator If  $Z_L \neq Z_o$  but  $Z_g \neq Z_o$ , reflected at both the load and generator repeating indefinitely

## 3) Standing waves

For lossless lines ( $\gamma\!=\!j\beta),$  (8–62, 63) with (8–88) becomes

$$V(z') = V_o^+ (e^{j\beta z'} + \Gamma e^{-j\beta z'})$$
(8-62)\*

$$I(z') = \frac{V_o^+}{Z_o} \left( e^{j\beta z'} - \Gamma e^{-j\beta z'} \right)$$
(8-63)\*

Polar expression of (8-88) in (8-62)\* using  $|V(z')| = [V(z')V^*(z')]^{1/2}$  gives

$$\begin{aligned} |V(z')| &= \{ [V_o^+(e^{j\beta z'} + |\Gamma|e^{j\theta_{\Gamma}}e^{-j\beta z'})] \cdot [(V_o^+)^* e^{-j\beta z'} + |\Gamma|e^{-j\theta_{\Gamma}}e^{j\beta z'})] \}^{-1/2} \\ &= |V_o^+|[1 + |\Gamma|^2 + |\Gamma|(e^{j(2\beta z' - \theta_{\Gamma})} + e^{-j(2\beta z' - \theta_{\Gamma})})]^{1/2} \\ &= |V_o^+|[1 + |\Gamma|^2 + 2|\Gamma|\cos\left(2\beta z' - \theta_{\Gamma}\right)]^{1/2} \quad (8-62)^{\star\star} \end{aligned}$$

⇒ Standing-wave pattern resulted from interference of incid. and reflec. waves



 $|V|_{\max} = |V_o^+|(1+|\Gamma|) \tag{3}$ when  $\cos(2\beta z'_{\max} - \theta_{\Gamma}) = 1 \implies 2\beta z'_{\max} - \theta_{\Gamma} = 2n\pi$   $\left(\begin{array}{c} \theta_{\Gamma} & n \end{array}\right) \qquad (n = 1, 2, \dots, \text{ for } \theta_{\Gamma} < 0$ 

$$\Rightarrow z'_{\max} = (\theta_{\Gamma} + 2n\pi)/2\beta = \left(\frac{\delta T}{4\pi} + \frac{\pi}{2}\right)\lambda \text{ for } \begin{cases} n = 1, 2, \dots, \text{ for } \theta_{\Gamma} < 0 \\ n = 0, 1, 2, \dots, \text{ for } \theta_{\Gamma} > 0 \end{cases}$$
(4)  
$$V|_{\min} = |V_{o}^{+}|(1 - |\Gamma|) \qquad -\pi \le \theta_{\Gamma} \le \pi, \quad \beta = 2\pi/\lambda$$
(5)  
when  $\cos\left(2\beta z'_{\min} - \theta_{\Gamma}\right) = -1 \Rightarrow 2\beta z'_{\min} - \theta_{\Gamma} = (2n+1)\pi$ 

$$\begin{array}{l} \text{men } \cos\left(2\beta z'_{\min} - \theta_{\Gamma}\right) = -1 \implies 2\beta z'_{\min} - \theta_{\Gamma} = (2n+1)\pi \\ \Rightarrow z'_{\min} = \left(\frac{\theta_{\Gamma}}{4\pi} + \frac{2n+1}{4}\right)\lambda \qquad n = 0, 1, 2, \dots \end{array}$$

$$\tag{6}$$

1st minimum position 
$$(n=0)$$
:  $l_{\min} = \left(\frac{\theta_{\Gamma}}{\pi} + 1\right) \frac{\lambda}{4}$  (7)

1st maximum position(n = 0 or 1):  $l_{\max} = \frac{\theta_{\Gamma} \lambda}{4\pi} \text{ or } \frac{\theta_{\Gamma} \lambda}{4\pi} + \frac{\lambda}{2}$  (8) Using  $|z'_{\min} - z'_{\max}| = \lambda/4$ ,

$$l_{\min} = \begin{cases} l_{\max} + \lambda/4 & \text{for } l_{\max} < \lambda/4 \\ l_{\max} - \lambda/4 & \text{for } l_{\max} \ge \lambda/4 \end{cases}$$
(9)



Determination SWR,  $|\Gamma|$ ,  $\theta_{\Gamma}$  and  $Z_L$  by a slotted-line probe:



Measurements of  $~|V_{\rm max}|,~|V_{\rm min}|,~{\rm and}~l_{\rm min}$ 

$$\Rightarrow \quad \text{Determine } S \text{ and } |\Gamma| \text{ from } S = \frac{|V_{\text{max}}|}{|V_{\text{min}}|} (8-90) \text{,} \quad |\Gamma| = \frac{S-1}{S+1} (8-91)$$
  
$$\theta_{\Gamma} \text{ from } l_{\text{min}} = \left(\frac{\theta_{\Gamma}}{\pi} + 1\right) \frac{\lambda}{4} (7)$$
  
$$\Gamma \text{ and } Z_{L} \text{ from } \Gamma = |\Gamma| e^{j\theta_{\Gamma}} = \frac{Z_{L} - Z_{o}}{Z_{L} + Z_{o}} (8-88)$$

(e.g. 8-5) For a lossless terminated transmission line,

given 
$$Z_0 = R_o = 50 (\Omega)$$
,  $S = 3$ ,  $l_{\min} = 5$  (cm), volt. min. dist.= 20 (cm)  
a)  $\Gamma = ?$  and b)  $Z_I = ?$ 

a) 
$$\lambda = 2 \times 0.2 = 0.4$$
,  $|\Gamma| = (3-1)/(3+1) = 0.5$   
 $l_{\min} = \left(\frac{\theta_{\Gamma}}{\pi} + 1\right) \frac{\lambda}{4} \implies \theta_{\Gamma} = \pi (4l_{\min}/\lambda - 1) = \pi (4 \times 0.05/0.4 - 1) = -\pi/2$   
 $\therefore \underline{\Gamma} = |\Gamma| e^{j\theta_{\Gamma}} = 0.5 e^{-j\pi/2} \underline{= -j0.5}$   
b)  $\Gamma = \frac{Z_L - Z_o}{Z_L + Z_o} \implies -j0.5 = \frac{Z_L - 50}{Z_L + 50} \implies \underline{Z_L} = 30 - j40$  ( $\Omega$ )

- 4) Characteristics in lossless finite lines ( $\alpha = 0, X_o = 0$ ;  $\gamma = j\beta, Z_o = R_o$ )
- a) Open-circuited line ( $Z_L \rightarrow \infty$ ,  $\Gamma = 1, S \rightarrow \infty$ )

$$Z_{ioc} \longrightarrow Z_{o} = R_{o} \qquad Z_{L} \rightarrow \infty \text{ (practically for } f < f_{UHF} \text{ )}$$

$$Q_{L} \longrightarrow Q_{L} \rightarrow \infty \text{ (practically for } f < f_{UHF} \text{ )}$$

$$(8-62)\star \implies V_{oc}(z') = V_o^+(e^{j\beta z'} + e^{-j\beta z'}) = 2V_o^+\cos\beta z' \qquad (8-62)\star_{oc}$$

$$(8-63)\star \implies I_{oc}(z') = \frac{V_o^+}{R_o} \left( e^{j\beta z'} - e^{-j\beta z'} \right) = \frac{2jV_o^+}{R_o} \sin\beta z' \tag{8-63}\star_{oc}$$

$$(8-79) \implies Z_{ioc} = \frac{V_{oc}(l)}{I_{oc}(l)} = jX_{ioc} = \frac{-jR_o}{\tan\beta l} = -jR_o \cot\beta l \text{ (purely reactive)} \quad (8-80)$$

For very shot line ( $\beta l = 2\pi l/\lambda \ll 1$ ),

$$Z_{ioc} = jX_{ioc} \simeq \frac{-jR_o}{\beta l} = -j\frac{\sqrt{L/C}}{\omega\sqrt{LC}l} = \frac{1}{2\omega}\frac{1}{\omega cl}$$
: capacitively reactive (8-81)



b) Short-circuited line  $(Z_L = 0, \Gamma = -1, S \rightarrow \infty)$ 

$$Z_{ioc} \longrightarrow Z_{o} = R_{o} \qquad Z_{L} = 0$$
short circuit
$$z' \leftarrow 0$$

$$(8-62)\star \implies V_{oc}(z') = V_o^+ (e^{j\beta z'} - e^{-j\beta z'}) = 2jV_o^+ \sin\beta z'$$
(8-62)\*<sub>sc</sub>

$$(8-63)^{\star} \implies I_{oc}(z') = \frac{V_o^+}{R_o} \left( e^{j\beta z'} + e^{-j\beta z'} \right) = \frac{2V_o^+}{R_o} \cos\beta z' \qquad (8-63)^{\star} \mathrm{sc}$$

$$(8-79) \implies Z_{isc} = \frac{V_{sc}(l)}{I_{sc}(l)} = jX_{isc} = jR_o \tan\beta l = jR_o \tan\left(\frac{2\pi l}{\lambda}\right)$$

$$(8-82)$$

 $\Rightarrow$  purely reactive (inductive or capacitive depending on aneta l)

 $\Rightarrow$  Proper choice of l of s.c. line can substitute for inductors and capacitors. For very shot line ( $\beta l=2\pi l/\lambda\ll 1),$ 

$$\begin{split} Z_{isc} &= j X_{isc} \cong j R_o \beta l = j \sqrt{L/C} \omega \sqrt{LC} l = j \omega L \, l \ : \ \text{inductively reactive} \ \ \text{(8-83)} \\ \text{For} \ \beta l = \pi/2, \ i.e., \ l = \lambda/4 \ , \ \ Z_{isc} \to \infty \end{split}$$

 $\Rightarrow$  A s.c. quarter-wavelength line is effectively an o.c. line.



Load condition	General case ( $\alpha \neq 0$ )	Lossless case ( $\alpha = 0$ )
Any value of load $Z_L$	$Z_{t} = Z_{0} \frac{Z_{L} + Z_{0} \tanh \gamma x}{Z_{0} + Z_{L} \tanh \gamma x}$	$Z_{\rm I} = Z_0 \frac{Z_L + jZ_0 \tan \beta x}{Z_0 + jZ_L \tan \beta x}$
Open-circuited line $(Z_L = \infty)$	$Z_{\rm f} = Z_0 \coth \gamma x$	$Z_{i} = -jZ_{0} \cot \beta x$
Short-circuited line ( $Z_L = 0$ )	$Z_{0} = Z_{0} \tanh \gamma x$	$Z_{1} = jZ_{0} \tan \beta x$

Input impedance of terminated transmission line<sup>†</sup>

<sup>†</sup> $\gamma = \alpha + j\beta$ , where  $\alpha$  = attenuation constant in nepers per meter,  $\beta = 2\pi/\lambda$  = phase constant in radians per meter, and  $\lambda$  = wavelength.

## c) Half-wavelength lossless line $(l = n\lambda/2)$

For 
$$l = n\lambda/2$$
  $(n = 1, 2, 3, ....)$ ,  $\tan\beta l = \tan\left(\frac{2\pi l}{\lambda}\right) = \tan n\pi = 0$  in (8-79):  
 $Z_i = Z_L$  for  $l = n\lambda/2$   $(n = 1, 2, 3, ....)$  (10)

- ⇒ A half-wave lossless line transfers a load impedance to the generator end without change.
- $\Rightarrow$  The generator induce the same V and I across the load as when the line does not exist there.

### d) Quater-wavelength lossless line $(l = \lambda/4 + n\lambda/2)$

For  $l = \lambda/4 + n\lambda/2$  (n = 0, 1, 2, 3, ....),

$$\tan\beta l = \tan\left(\frac{2\pi l}{\lambda}\right) = \tan\left(\pi/2 + n\pi\right) \to \infty \quad \text{in (8-79):}$$
$$Z_{i} = \frac{Z_{o}^{2}}{Z_{L}} \quad \text{for} \quad l = \lambda/4 + n\lambda/2 \quad (n = 0, 1, 2, 3, ....) \tag{8-111}$$

 $\Rightarrow$  Quater-wave transformer to eliminate reflections at the load terminal.



If  $Z_{\!i} = Z_{\!01}$  , no reflections at the terminal AA'.

By (8-111),  $Z_{01} = Z_o^2/Z_L \implies Z_o = \sqrt{Z_{01}Z_L}$  (11) Therefore, if a quarter-wave lossless line having a characteristic impedance of  $Z_o = \sqrt{Z_{01}Z_L}$  is inserted between the feedline and the load, there are no reflections at the terminal and all the incident power is transferred into the load. e) Determination of  $Z_o \ {
m and} \ \gamma$  by input impedance measurements

From (8-78) 
$$Z_i(l) = Z_o \frac{Z_L + Z_o \tanh \gamma l}{Z_o + Z_L \tanh \gamma l}$$
,  
 $Z_{ioc} = Z_o \coth \gamma l \text{ for } Z_L \rightarrow \infty \text{ and } Z_{isc} = Z_o \tanh \gamma l \text{ for } Z_L = 0$  (8-84a, b)  
 $\Rightarrow Z_o = \sqrt{Z_{ioc} Z_{isc}}$  (Q) (8-85)

$$\gamma = \frac{1}{l} \tanh^{-1} \sqrt{\frac{Z_{isc}}{Z_{ioc}}} \quad (m^{-1})$$
(8-86)

### d) Power flow along the transmission lines

Time-average power flow along the line by analogy with (7-79),

$$P(z') = \mathscr{P}_{av}(z') = \frac{1}{2} \operatorname{Re}[V I^*]$$
(12)

For lossless lines  $(\gamma = j\beta)$ ,

$$V(z') = V_o^+ (e^{j\beta z'} + \Gamma e^{-j\beta z'})$$
(8-62)\*

$$I(z') = \frac{V_o^+}{Z_o} \left( e^{j\beta z'} - \Gamma e^{-j\beta z'} \right)$$
(8-63)\*

At the load (z'=0), the incident and reflected waves are

$$V_i(0) = V_o^+, \qquad I_i(0) = V_o^+ / Z_o$$
 (13)

$$V_r(0) = \Gamma V_o^+, \quad I_i(0) = -\Gamma V_o^+ / Z_o$$
(14)

(13), (14) in (12) :

$$P_{i} = \frac{1}{2} Re[V_{o}^{+}V_{o}^{+*}/Z_{o}] = \frac{|V_{o}^{+}|^{2}}{2Z_{o}}$$
(12)

$$P_{r} = \frac{1}{2} Re \left[ \Gamma V_{o}^{+} \left( -\Gamma^{*} V_{o}^{+*} / Z_{o} \right) \right] = -|\Gamma|^{2} \frac{|V_{o}^{+}|^{2}}{2Z_{o}}$$
(12)

Net average power delivered to the load :

$$P = P_i + P_r = \frac{|V_o^+|^2}{2Z_o} (1 - |\Gamma|^2) \qquad (cf) \ (7-104) \tag{15}$$