



**Part I Fundamentals Electron Theory : Matter Waves** Chap. 1 Introduction Chap. 2 The Wave-Particle Duality **Chap. 3 The Schördinger Equation** Chap. 4 Solution of the Schördinger Equation for **Four Specific Problems** Chap. 5 Energy Bands in Crystals Chap. 6 Electrons in a Crystal **Electromagnetic Theory : Maxwell Equations** Chap. 4 Light Waves (Electrons in Solids, 3<sup>rd</sup> Ed., R. H. Bube)





## 6.1 Fermi Energy and Fermi Surface

The Fermi energy,  $E_{\rm F}$ :

- An important part of an electron band diagram
- Defined as "the highest energy that the electrons assume at T = 0 K"
- Fermi energy for Al and Cu : see Fig 5.21, 5.22

Fermi energy for semiconductor:

- The above definition can occasionally be misleading, particularly

when dealing with semiconductors

- Fermi function at  $E_F$ ,  $F(E_F) = \frac{1}{2}$  : see Section 6.2 for more accurate definition

Fermi surface (in 3-d k-space) for Cu : see Fig 5.27

## 6.2 Fermi Distribution Function

Fermi function, F(E): The probability that a certain energy level is occupied by electrons

$$F(E) = \frac{1}{\exp\left(\frac{E - E_F}{k_B T}\right) + 1}$$

Fermi distribution for T = 0 K (Fig 6.1)<sub>Figure 6.1. Fermi distribution function, F(E), versus energy, E, for T = 0. and for higher  $T(T \neq 0$  K) (Fig 6.2)</sub>

At high energy  $(E >> E_F)$ , F(E) is approximated by classical Boltzmann distribution

$$F(E) \approx \exp\left[-\left(\frac{E-E_F}{k_BT}\right)\right]$$

"Boltzmann tail"

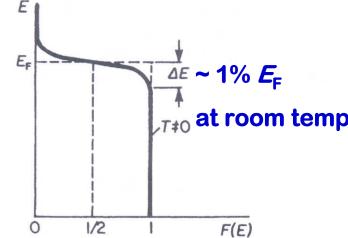
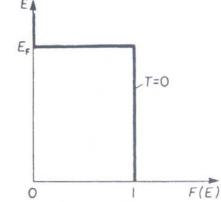


Figure 6.2. Fermi distribution function for  $T \neq 0$ .







## **6.3 Density of States**

### "How energy levels are distributed over a band?"

Assume free electrons are confined in a square potential well of crystal. Similar to the case in Sec. 4.2, by using B.C., the solution of the Schrödinger equation

$$E_{n} = \frac{\pi^{2}\hbar^{2}}{2ma^{2}}(n_{x}^{2} + n_{y}^{2} + n_{z}^{2})$$

$$n^{2} = n_{x}^{2} + n_{y}^{2} + n_{z}^{2}$$
where  $n_{x} \cdot n_{y} \cdot n_{z}$  are principal quantum numbers,  
*a* is the length of the crystal

Figure 6.3. Representation of an energy state in quantum number space.





## **6.3 Density of States**

A specific energy level,  $E_n$  for each set of  $n_x n_y n_z$ , called "energy state"

- Equal values of the energy,  $E_n$  lie on the surface of sphere with radius n

- All points within the sphere represent quantum states with energy smaller

than  $E_n$ 

- The **# of quantum state**,  $\eta$ , with an energy equal to or smaller than  $E_n$ , is proportional to the volume of the sphere

(*n* values can be defined in positive octant of the *n*-space)

In a one-eighth of the volume of the sphere with radius *n* (# of energy state  $\eta$  ),  $\eta = \frac{1}{8} \cdot \frac{4}{3} \pi n^3 = \frac{\pi}{6} \left( \frac{2\pi a^2}{\pi^2 \hbar^2} \right)^{3/2} E^{3/2}$ 







**Density of state**, Z(E): # of energy states per unit energy in the energy interval dE

 $(a^{3} =$ volume that the electrons can occupy)

Z(E): differentiation of  $\eta$  with respect to the energy, E

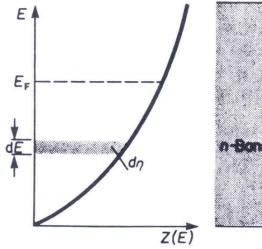


Figure 6.4. Density of states Z(E) within a band. The electrons in this band are considered to be free.

$$\frac{d\eta}{dE} = Z(E) = \frac{\pi}{4} \left(\frac{2ma^2}{\pi^2 \hbar^2}\right)^{3/2} E^{1/2} = \frac{V}{4\pi^2} \left(\frac{2m}{\hbar^2}\right)^{3/2} E^{1/2}$$

 $d\eta = Z(E) \cdot dE$ 



### **6.4 Population Density**

**Pauli principle :** each energy state can be occupied by one electron of positive spin and one of negative spin

**Population density**  $N(E) = 2 \cdot Z(E) \cdot F(E)$ 

$$N(E) = \frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2}\right)^{3/2} E^{1/2} \frac{1}{\exp\left(\frac{E - E_F}{k_B T}\right) + 1}$$

For  $T \rightarrow 0$  and  $E < E_F \rightarrow N(E) = 2 \cdot Z(E)$ , F(E) = 1For  $T \neq 0$ ,  $E \approx E_F$ ,  $\rightarrow$  the Fermi distribution function causes a smearing out of N(E)(Fig 6.5)

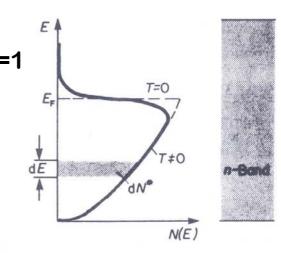


Figure 6.5. Population density N(E) within a band for free electrons.  $dN^*$  is the number of electrons in the energy interval dE.



### 6.4 Population Density

# of electrons  $N^*$ , that have an energy equal to or smaller than the energy  $E_n$  (The area within the curve in Fig 6.5) For an energy interval between E and E + dE  $dN^* = N(E)dE$ 

From (6.8) and (6.9) and consider simple case  $T \to 0$  and  $E < E_F$ , F(E) = 1  $N^* = \int_0^{E_F} N(E) dE = \int_0^{E_F} \frac{V}{2\pi^2} \left(\frac{2m}{\hbar^2}\right)^{3/2} E^{1/2} dE = \frac{V}{3\pi^2} \left(\frac{2m}{\hbar^2}\right)^{3/2} E_F^{-3/2}$  $E_F = \left(3\pi^2 \frac{N^*}{V}\right)^{2/3} \frac{\hbar^2}{2m}$ 

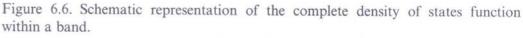
If we define #of electrons per unit volume as N' = N\*/V,  $E_F = (3\pi^2 N')^{2/3} \frac{\hbar^2}{2m}$ 



# 6.5 Complete Density of States Function Within a Band

#### Evs. Z(E) in actual crystals

- Low energy : free-electronlike
- Higher energy : fewer energy state available (Fig 5.26)
- $\rightarrow$  Z(E) decrease with increasing E



- The corners of the BZ : Z(*E*) dropped to zero

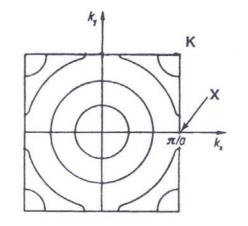
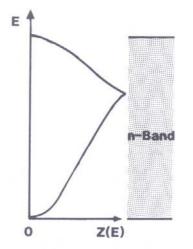


Figure 5.26. Curves of equal energy inserted into the first Brillouin zone for a twodimensional square lattice.



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## 6.6 Consequences of the Band Model

**Insulators** : solids in which the highest filled band is completely occupied by electron

Alkali metal: the valence band is essentially half-filled, electrons can drift under external field

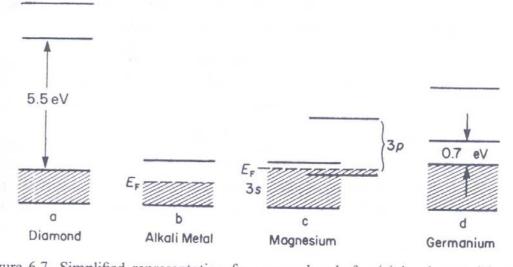


Figure 6.7. Simplified representation for energy bands for (a) insulators, (b) alkali metals, (c) bivalent metals, and (d) intrinsic semiconductors.

**Bivalent metals**: upper band partially overlapped ; weak binding forces of the valence electrons on atom

Semiconductors: valence band is completely filled with electron; relatively narrow band gap (intrinsic semiconductors); a sufficiently large energy can excite electron from valence band to conduction band  $\rightarrow$  some electron conduction





## 6.7 Effective Mass

#### Effective mass (of electron), m\*

experimentally determined electron mass

- Deviation of  $m^*$  from free electron mass  $m_0$ : usually attributed to interaction between drifting electrons and atoms in a crystal

### - For example,

Electron accelerated in an electric field might be slowed down slightly due to "collisions" with some atom  $\rightarrow$  ratio  $m^*/m_0 > 1$ 

The electron wave in another crystal might have just the right phase in order that the response to an external electric field is enhanced $\rightarrow m^*/m_0 < 1$ 

- Derivation of effective mass  $U_g = \frac{d\omega}{dk} = \frac{d(2\pi v)}{dk} = \frac{d(2\pi E/h)}{dk} = \frac{1}{\hbar} \frac{dE}{dk}$ 



## **6.7 Effective Mass**

(continued) 
$$a = \frac{d\upsilon_g}{dt} = \frac{1}{\hbar} \frac{d^2 E}{dk^2} \frac{dk}{dt}$$
  
(4.7) 
$$p = \hbar k \quad \rightarrow \quad \frac{dp}{dt} = \hbar \frac{dk}{dt}$$
$$a = \frac{1}{\hbar^2} \frac{d^2 E}{dk^2} \frac{dp}{dt} = \frac{1}{\hbar^2} \cdot \frac{d^2 E}{dk^2} \cdot \frac{d(m\upsilon)}{dt} = \frac{1}{\hbar^2} \frac{d^2 E}{dk^2} F$$
$$a = \frac{F}{m} \qquad m^* = \hbar^2 \left(\frac{d^2 E}{dk^2}\right)^{-1}$$



## **6.7 Effective Mass**

$$m^* = \hbar^2 \left(\frac{d^2 E}{dk^2}\right)^{-1}$$

Effective mass is inversely proportional to the curvature of and electron band.

In Fig 6.8,  $m^*$  is small and positive near the center of BZ

Negative *m*<sup>\*</sup> (upper part of the band in Fig 6.8): "particle travels in the opposite direction to an applied electric force ( and opposite to an electron"; called "**electron hole**"

cf) exciton : an electron/hole pair

