
Ch12. Filters and Tuned Amplifiers

Introduction

□ Passive LC filters

□ Electronic Filter → Active filter

– Active RC filters

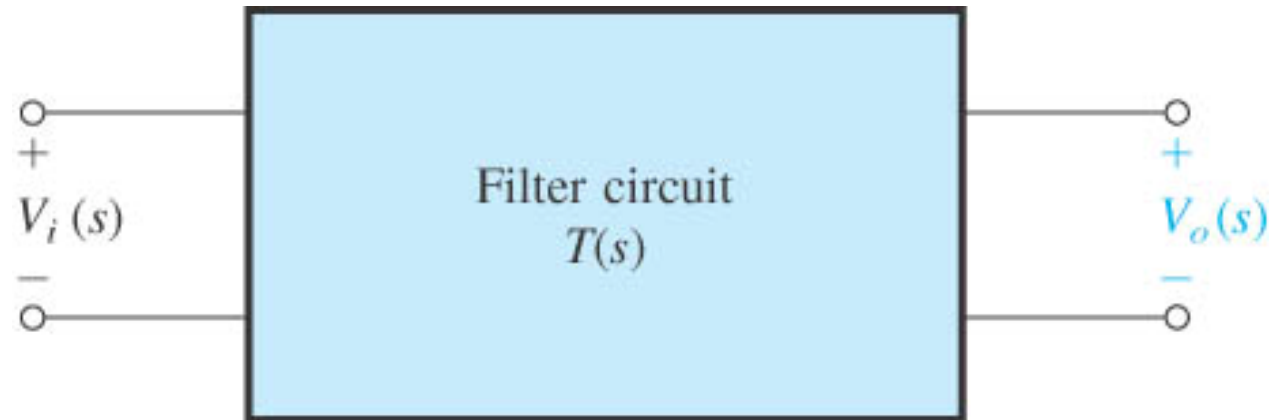
– Switched capacitor circuits

→ **Advantages: No inductors!**

Inductors are large and physically bulky for low frequency applications
(such as those used in passive LC filters)

Filter Transmission

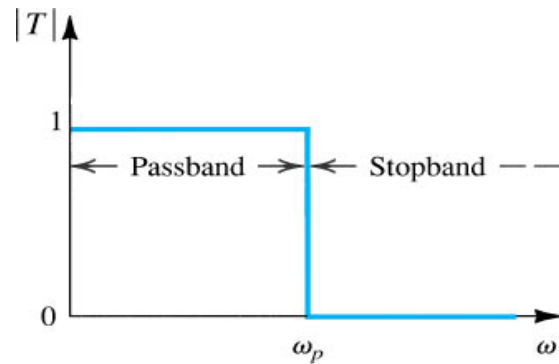
- Filter - a two port device



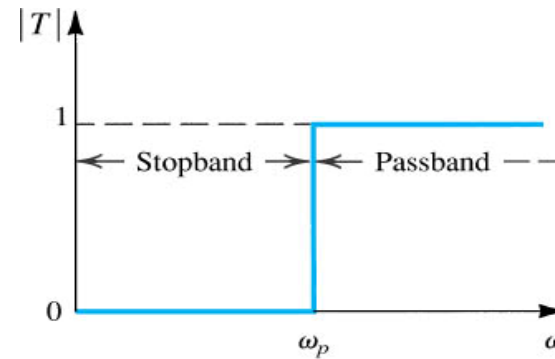
- Transfer function: $T(s) \equiv \frac{V_o(s)}{V_i(s)}$
- Transfer transmission: $T(j\omega) = \boxed{|T(j\omega)|} e^{j\phi(\omega)}$ (s = j\omega)
 ↑ magnitude ↑ phase
- Gain function: $G(\omega) \equiv 20 \log |T(j\omega)| \text{ dB}$
- Attenuation function: $A(\omega) \equiv -20 \log |T(j\omega)| \text{ dB}$
- Input Output relation: $|V_o(j\omega)| = |T(j\omega)| |V_i(j\omega)|$

Filter Types

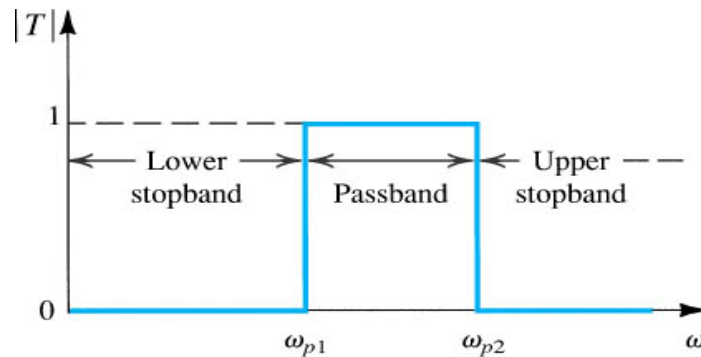
- Frequency-selection function $\left\{ \begin{array}{l} \text{passing} \rightarrow \text{passband: } |T| = 1 \\ \text{stopping} \rightarrow \text{stopband: } |T| = 0 \end{array} \right.$
- Brick-wall response



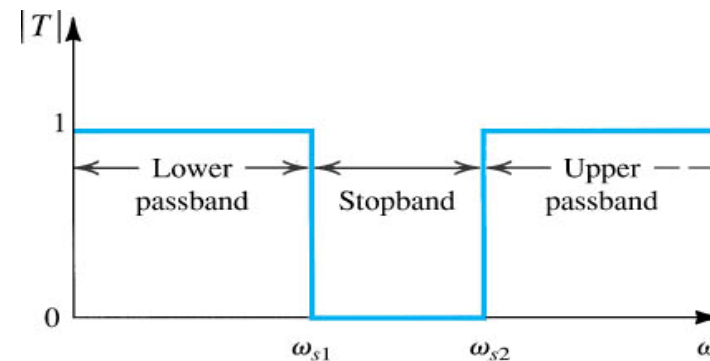
(a) Low-pass (LP)



(b) High-pass (HP)

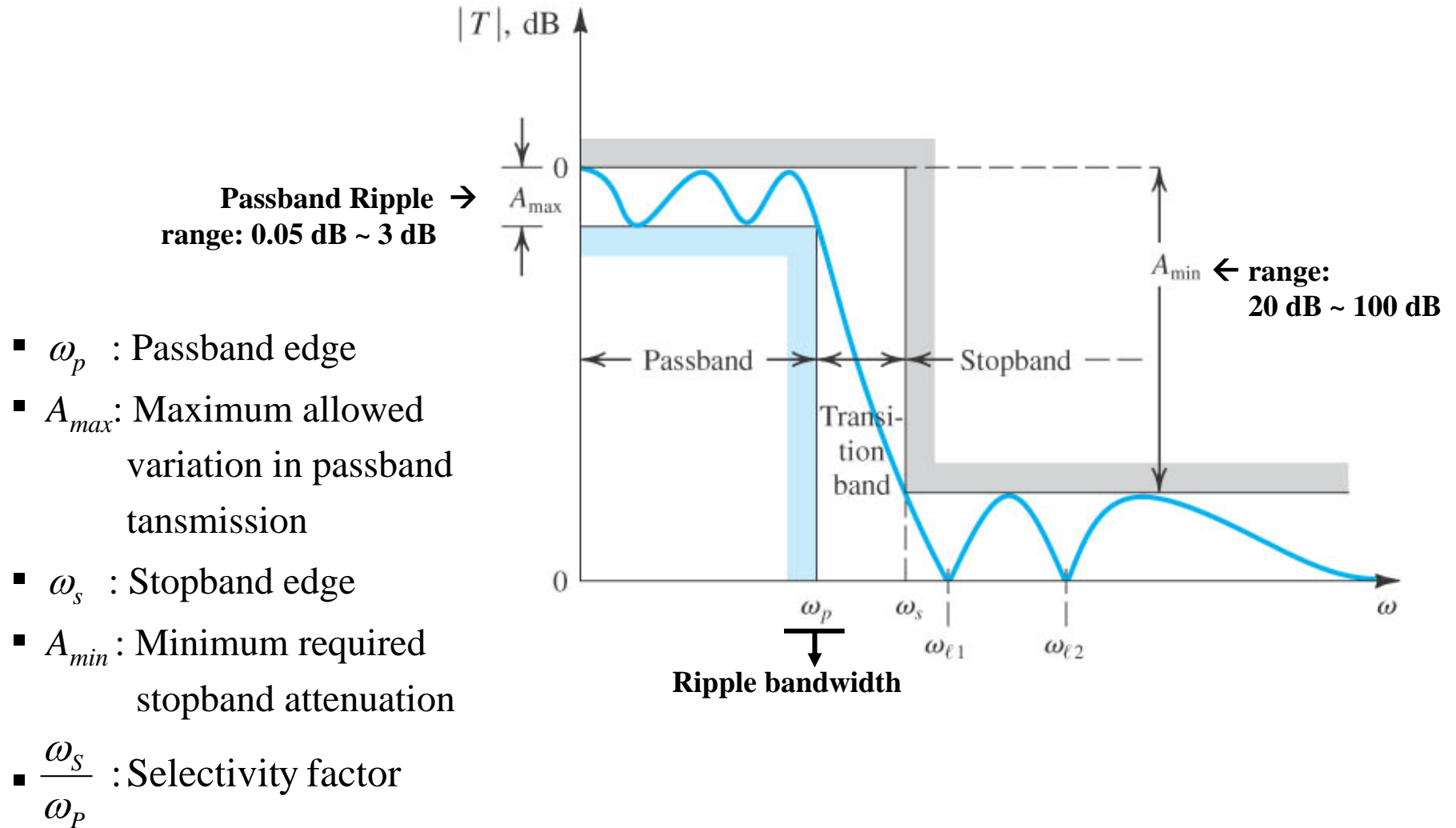


(c) Bandpass (BP)



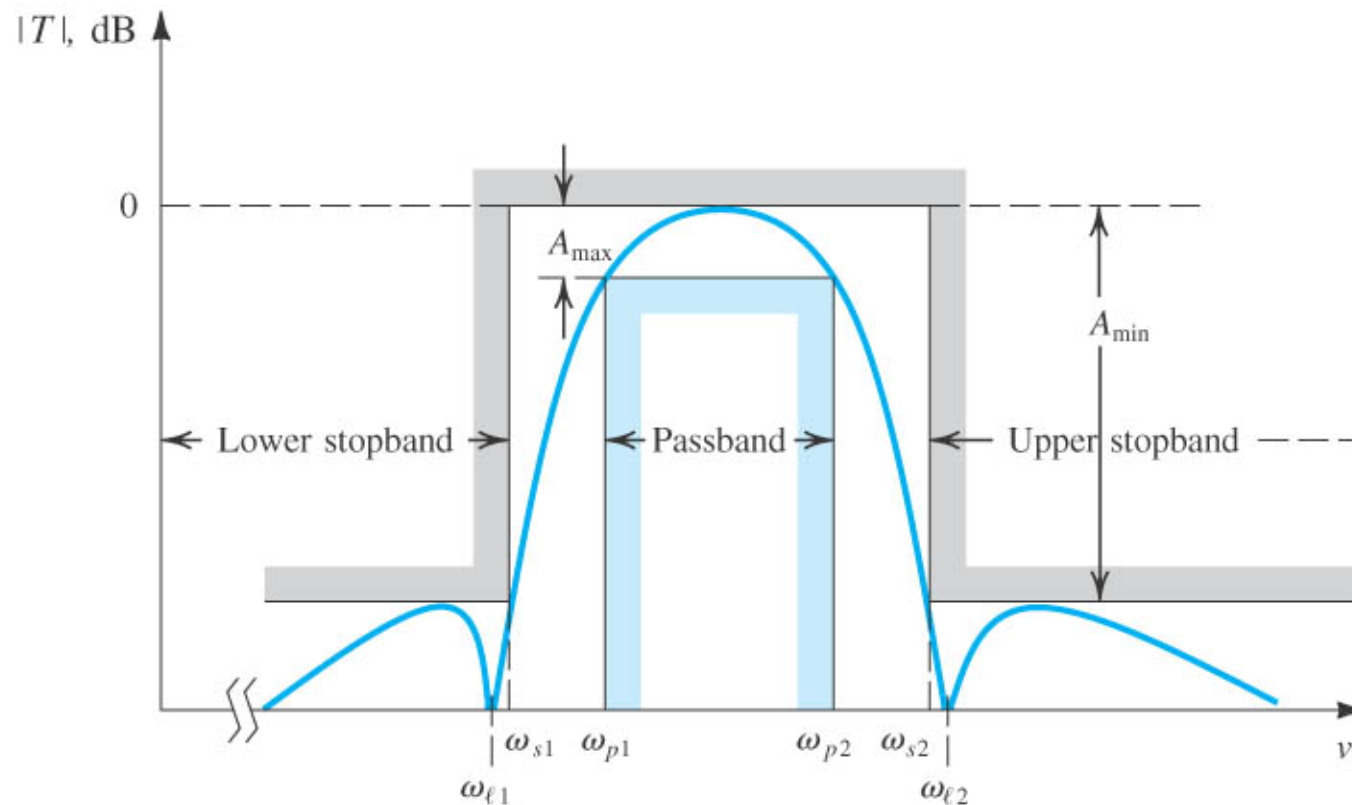
(d) Bandstop (BS)

Filter Specification



Filter Specification

- Filter approximation
 - The process of obtaining a transfer function that meets given specifications
 - Performed using computer programs(Snelgrove, 1982;Ouslis and Sedra, 1995), filter design table(Zverev, 1967) or closed-form expressions(Section 12.3)



The Filter Transfer Function

□ Filter Transfer Function $T(s)$

$$\blacksquare T(s) = \frac{a_M s^M + a_{M-1} s^{M-1} + \cdots + a_0}{s^N + b_{N-1} s^{N-1} + \cdots + b_0}$$

– N : Filter order

– If $N \geq M$, stable

– a_0, \dots, a_M & b_0, \dots, b_{N-1} : real numbers

→ Poles need be in conjugate pairs or negative real numbers

$$\blacksquare T(s) = \frac{a_M (s - z_1)(s - z_2) \cdots (s - z_M)}{(s - p_1)(s - p_2) \cdots (s - p_N)}$$

– z_1, \dots, z_M = transfer - function zeros = transmission zeros

– p_1, \dots, p_N = transfer - function poles = natural modes

→ All poles must lie in left half plane.

The Filter Transfer Function (zeros)

- Since in the stopband the transmission is zero or small
 → the zeros are usually, placed on the $j\omega$ axis at stopband frequencies

1. zeros at $s = +j\omega_{l1}$ & $+j\omega_{l2}$

also at $s = -j\omega_{l1}$ & $-j\omega_{l2}$

— Numerator polynomial

$$(s + j\omega_{l1})(s - j\omega_{l1})(s + j\omega_{l2})(s - j\omega_{l2})$$

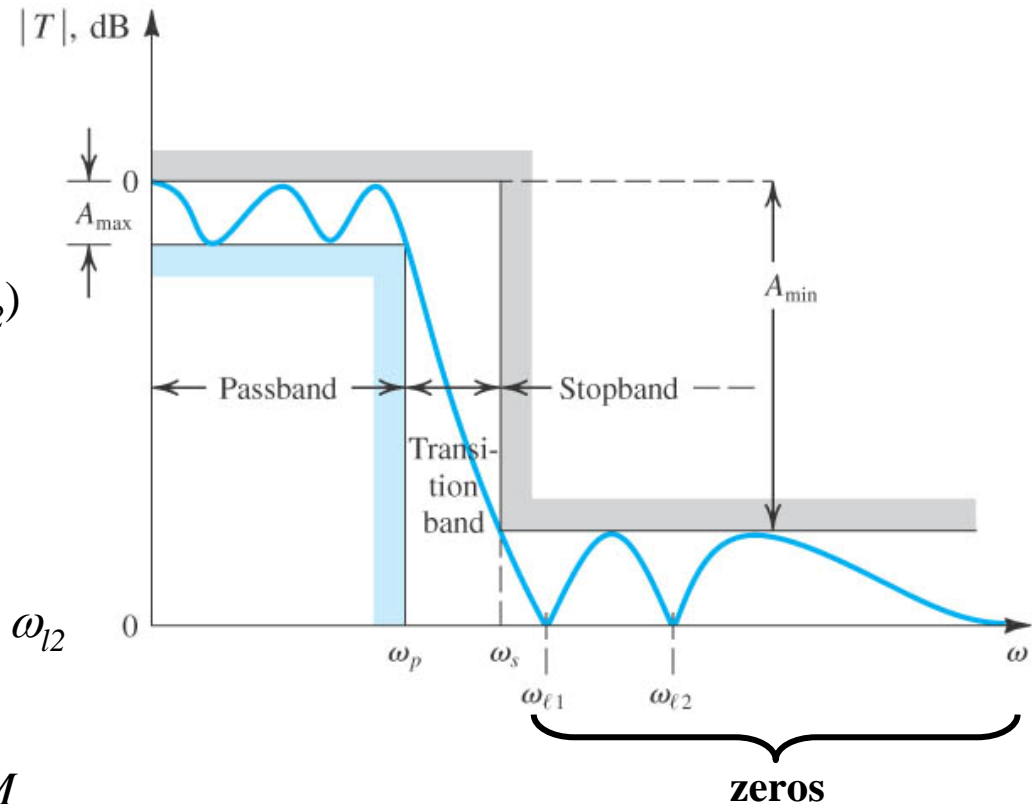
$$= (s^2 + \omega_{l1}^2)(s^2 + \omega_{l2}^2)$$

→ for $s = j\omega$,

$$(s^2 + \omega_{l1}^2)(s^2 + \omega_{l2}^2)$$

$$= (-\omega^2 + \omega_{l1}^2)(-\omega^2 + \omega_{l2}^2)$$

which is zero at $\omega = \omega_{l1}$ and $\omega = \omega_{l2}$

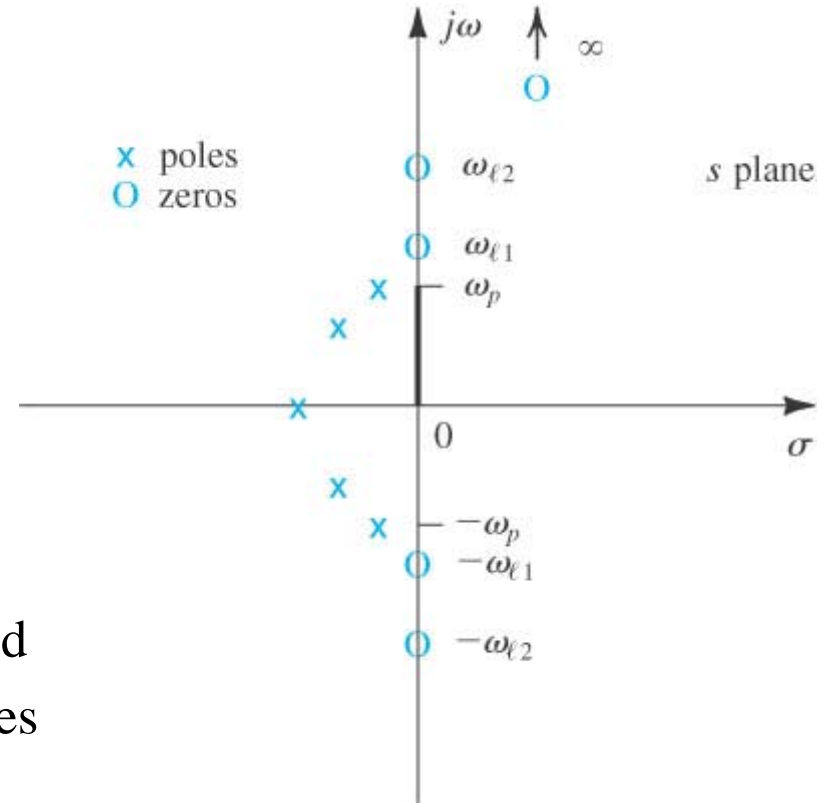
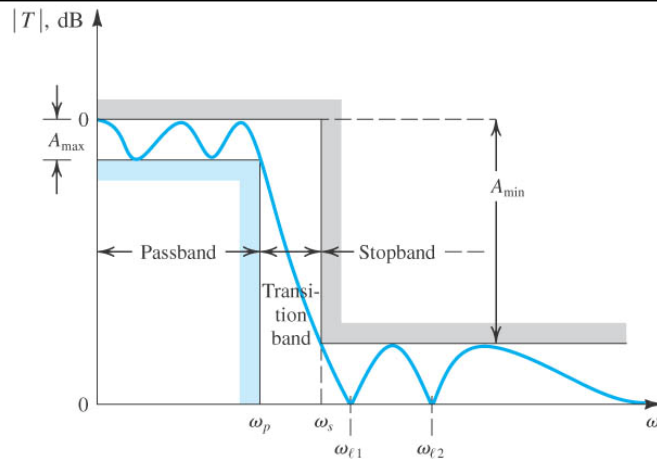


2. zeros at $s = \infty$

the numbers of zeros at $s = \infty = N - M$

∴ as $s \rightarrow \infty$, $T(s) \rightarrow \frac{a_M}{s^{N-M}} \rightarrow$ Denominator make zeros at infinity

The Filter Transfer Function (ex. 1 : LPF)

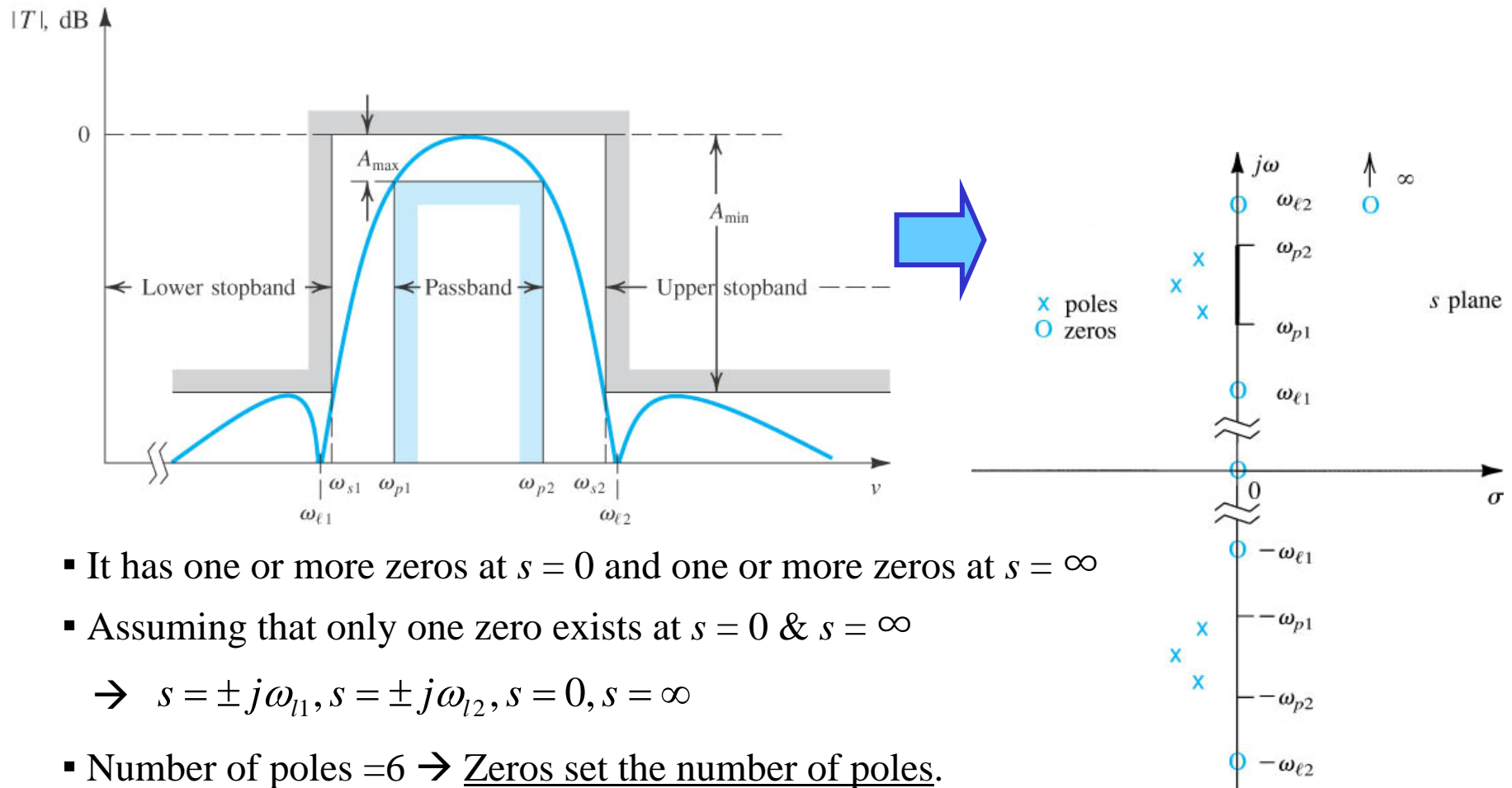


- Number of **poles** = 5
- Two pairs of complex-conjugate poles and real-axis pole
 - all the poles lie in the vicinity of passband
 - high transmission at pass band frequencies

▪ **Zeros** : $s = \pm j\omega_{\ell 1}$ & $\pm j\omega_{\ell 2}$ & ∞

$$T(s) = \frac{a_4(s^2 + \omega_{\ell 1}^2)(s^2 + \omega_{\ell 2}^2)}{s^5 + b_4s^4 + b_3s^3 + b_2s^2 + b_1s + b_0}$$

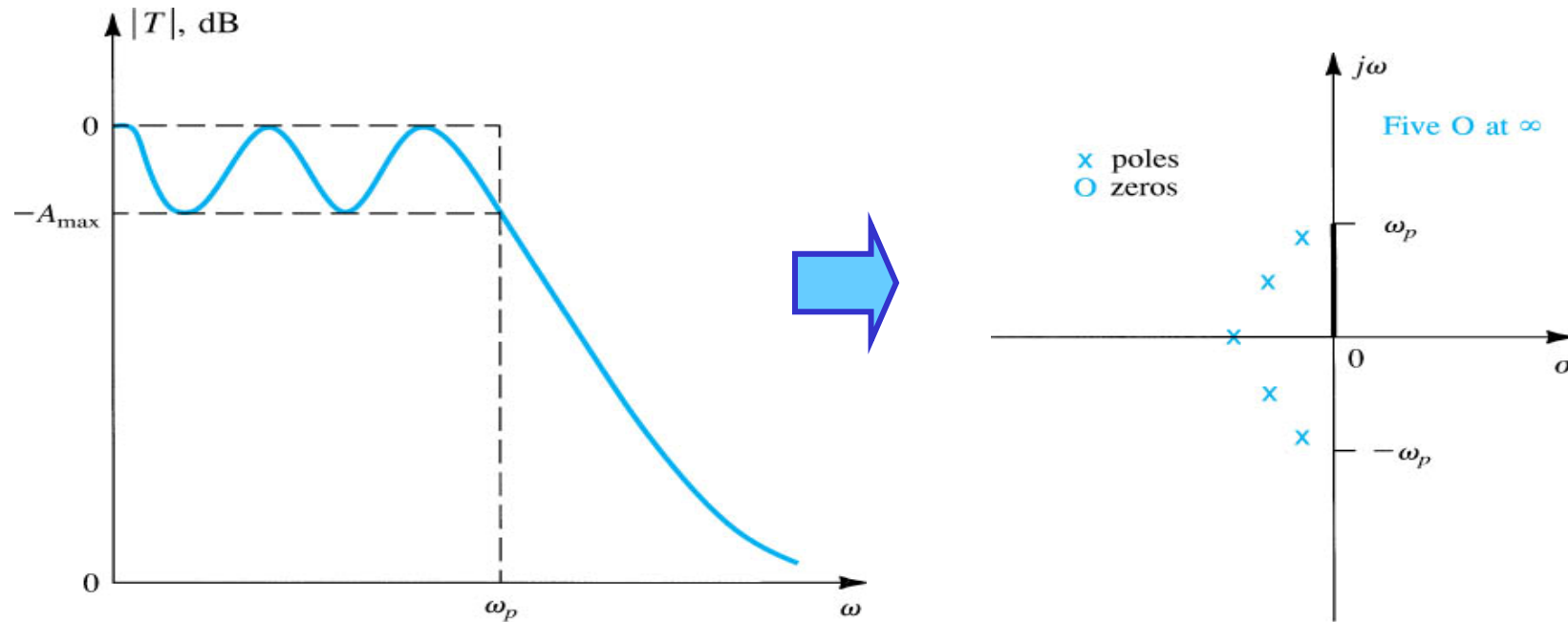
The Filter Transfer Function (ex. 2 : BPF)



- It has one or more zeros at $s = 0$ and one or more zeros at $s = \infty$
- Assuming that only one zero exists at $s = 0$ & $s = \infty$
 $\rightarrow s = \pm j\omega_{l1}, s = \pm j\omega_{l2}, s = 0, s = \infty$
- Number of poles = 6 \rightarrow Zeros set the number of poles.

$$T(s) = \frac{a_5 s (s^2 + \omega_{l1}^2) (s^2 + \omega_{l2}^2)}{s^6 + b_5 s^5 + \dots + b_0}$$

The Filter Transfer Function (ex. 3 : all-pole LPF)



- It is possible that all zeros are at $s = \infty$

$$T(s) = \frac{a_0}{s^N + b_{N-1}s^{N-1} + \dots + b_0} \rightarrow \text{all-pole filter}$$

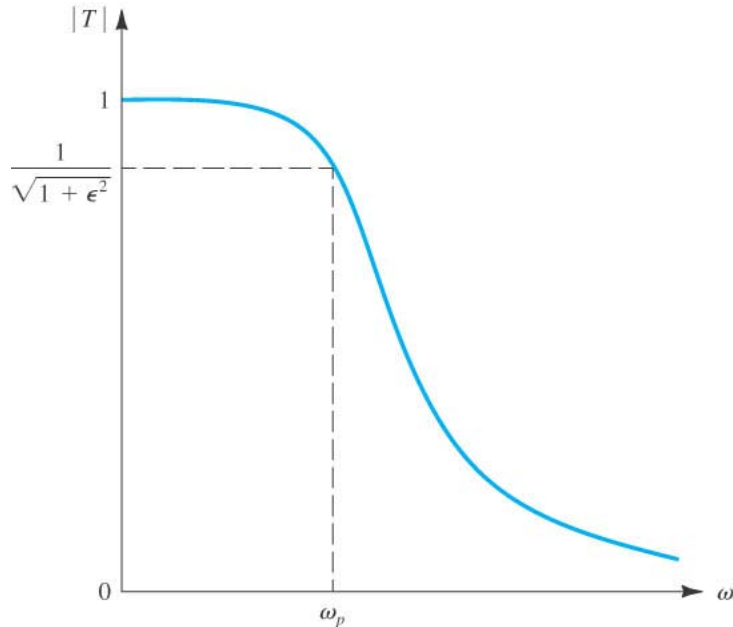
- The more selective the required filter response is, the higher its order must be, and the closer its natural modes are to the $j\omega$ axis.

Butterworth and Chebyshev Filters

- In this section, we present two functions that are frequently used in **approximating the transmission characteristics** of low-pass filters.
: Closed-form expressions

Butterworth Filters : Filter Shape and $|T(j\omega)|$

□ The Butterworth Filter



- Monotonically decreasing transmission
- All the transmission zero at $\omega = \infty \rightarrow$ **all-pole**
- The magnitude function for an N^{th} -order Butterworth filter with a passband edge ω_p is

$$|T(j\omega)| = \frac{1}{\sqrt{1 + \varepsilon^2 \left(\frac{\omega}{\omega_p}\right)^{2N}}}$$

\rightarrow at $\omega = \omega_p$,

$$|T(j\omega_p)| = \frac{1}{\sqrt{1 + \varepsilon^2}}$$

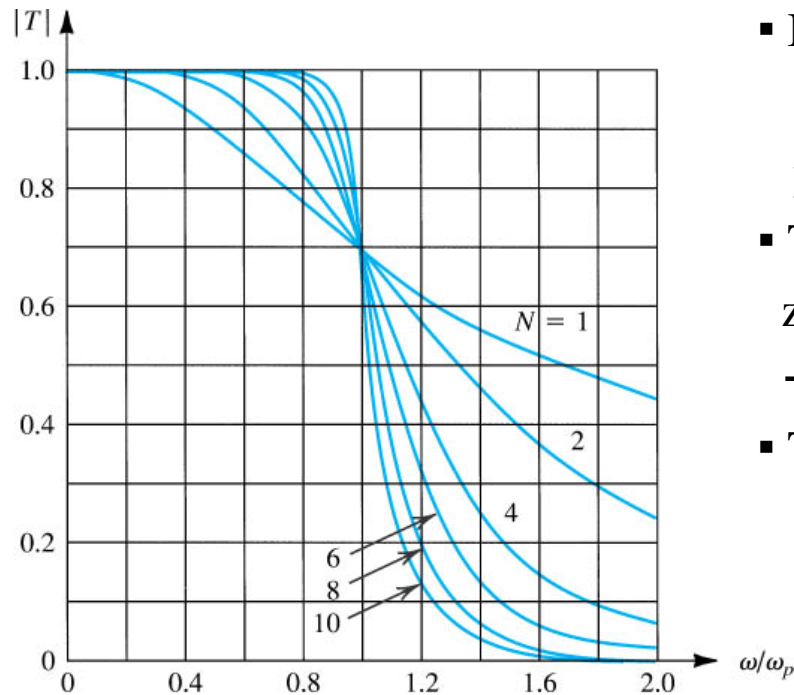
- Thus, the parameter ε determines the maximum variation in passband transmission,

$$A_{\max} = 20 \log \sqrt{1 + \varepsilon^2}$$

- Conversely, given A_{\max} ,

$$\varepsilon = \sqrt{10^{A_{\max}/10} - 1}$$

The Butterworth Filter : N effects



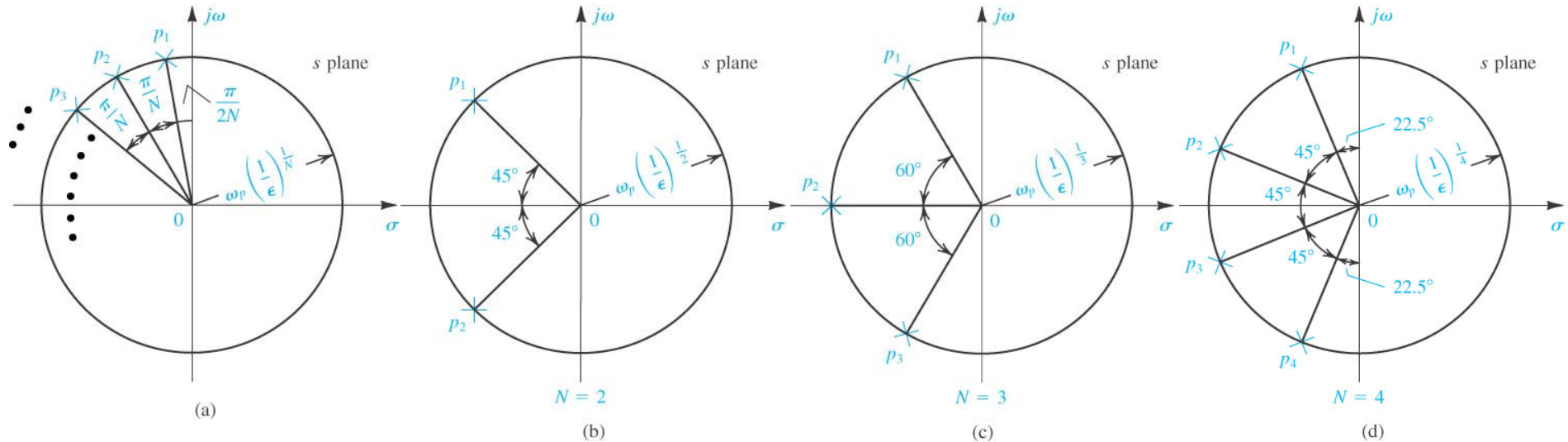
- In the Butterworth response the maximum deviation in passband transmission occurs at the passband edge, ω_p , only
- The first $2N-1$ derivatives of $|T|$ relative to ω are zero at $\omega = 0$
 - very flat near $\omega = 0$ (**maximally flat** response)
- The degree of passband flatness increases as the order N is increased
 - as the order N is increased the filter response approaches the ideal **brick-wall**

- At the edge of the stopband, $\omega = \omega_s$, attenuation is

$$A(\omega_s) = -20 \log \left[\frac{1}{\sqrt{1 + \varepsilon^2 (\omega_s / \omega_p)^{2N}}} \right] = 10 \log \left[1 + \varepsilon^2 (\omega_s / \omega_p)^{2N} \right]$$

- The required filter order = the lowest integer value of N that yields $A(\omega_s) \geq A_{\min}$

The Butterworth Filter : Poles



- The natural modes of an N^{th} -order Butterworth filter can be determined from the graphical construction above.
- Natural modes lies on a circle of radius $\omega_p(1/\epsilon)^{1/N}$
 → same frequency of $\omega_0 = \omega_p(1/\epsilon)^{1/N}$
- Space by equal angles of π/N , with the first mode at an angle $\pi/2N$ from the $+j\omega$ axis.
- Transfer function is

$$T(s) = \frac{K\omega_0^N}{(s - p_1)(s - p_2)\cdots(s - p_N)} \quad \rightarrow K \text{ is a constant dc gain of the filter}$$

The Butterworth Filter

- How to find a Butterworth transfer function

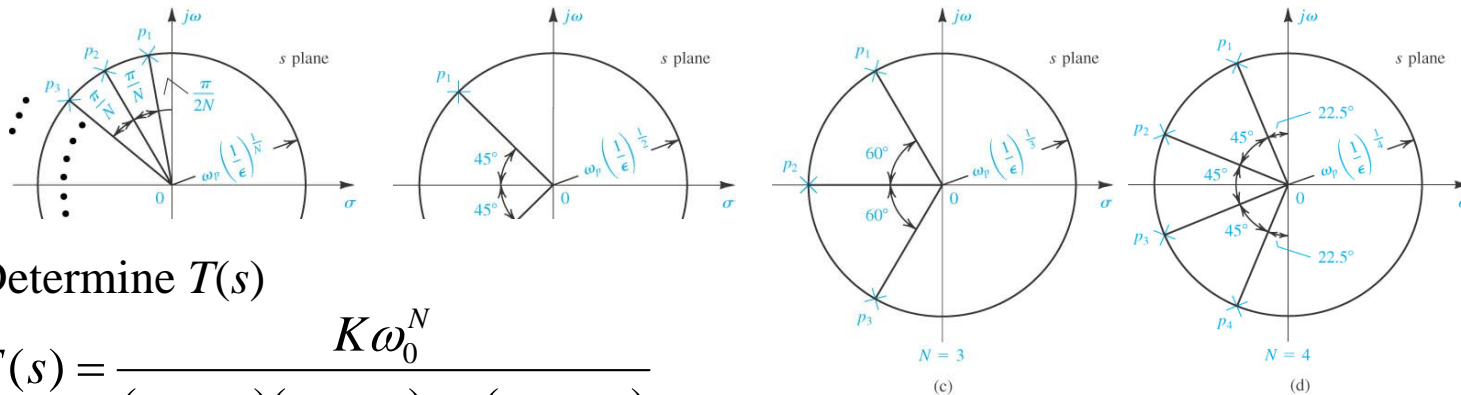
1. Determine ε .

$$\varepsilon = \sqrt{10^{A_{\max}/10} - 1}$$

2. Determine the required filter order as the lowest integer value of N that results in $A(\omega_S) \geq A_{\min}$.

$$A(\omega_S) = -20 \log \left[1 / \sqrt{1 + \varepsilon^2 (\omega_S / \omega_P)^{2N}} \right] = 10 \log \left[1 + \varepsilon^2 (\omega_S / \omega_P)^{2N} \right]$$

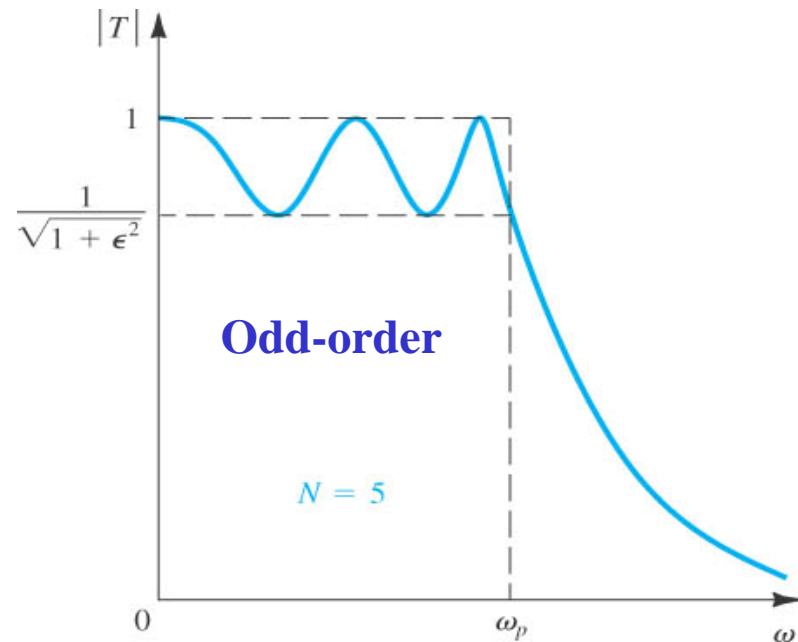
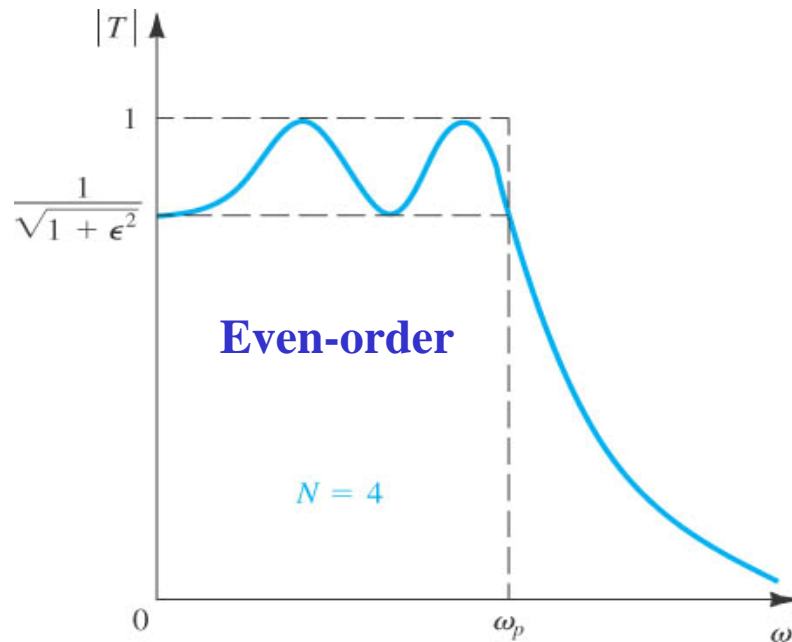
3. Determine the N natural modes.



4. Determine $T(s)$

$$T(s) = \frac{K\omega_0^N}{(s - p_1)(s - p_2)\cdots(s - p_N)}$$

The Chebyshev Filter



- **Equi-ripple response** (A_{\max} = the peak ripple) in the passband and a monotonically decreasing transmission in the stopband.
- The odd-order filter, $|T(0)|=1$
The even-order filter exhibits its maximum magnitude deviation at $\omega = 0$.
- Total number of passband maxima and minima equals the order of the filter, N .
- All the zeros are at $\omega = \infty$. → **all-pole filter**

The Chebyshev Filter : Equation

- The magnitude of the transfer function with a passband edge ω_p is

$$|T(j\omega)| = \frac{1}{\sqrt{1 + \varepsilon^2 \cos^2[N \cos^{-1}(\omega/\omega_p)]}} \quad \text{for } \omega \leq \omega_p$$

$$|T(j\omega)| = \frac{1}{\sqrt{1 + \varepsilon^2 \cosh^2[N \cosh^{-1}(\omega/\omega_p)]}} \quad \text{for } \omega \geq \omega_p$$

$$|T(j\omega_p)| = \frac{1}{\sqrt{1 + \varepsilon^2}} \quad \text{for } \omega = \omega_p$$

- Maximum passband ripple :

$$A_{\max} = 10 \log(1 + \varepsilon^2)$$

conversely,

$$\varepsilon = \sqrt{10^{A_{\max}/10} - 1}$$

- The attenuation at the stopband edge ($\omega = \omega_s$) is

$$A(\omega_s) = 10 \log[1 + \varepsilon^2 \cosh^2(N \cosh^{-1}(\omega_s/\omega_p))]$$

→ Required order N calculation by

finding the lowest integer value of N that yields $A(\omega_s) \geq A_{\min}$.

- Increasing the order N of the Chebyshev filter causes its magnitude function to approach the ideal brick-wall low-pass response.

The Chebyshev Filter

- The poles are

$$p_k = -\omega_p \sin\left(\frac{2k-1}{N} \frac{\pi}{2}\right) \sinh\left(\frac{1}{N} \sinh^{-1} \frac{1}{\varepsilon}\right) + j\omega_p \cos\left(\frac{2k-1}{N} \frac{\pi}{2}\right) \cosh\left(\frac{1}{N} \sinh^{-1} \frac{1}{\varepsilon}\right) \quad k = 1, 2, \dots, N$$

- The transfer function is

$$T(s) = \frac{K \omega_p^N}{\varepsilon 2^{N-1} (s - p_1)(s - p_2) \cdots (s - p_N)}$$

- How to find the transfer function

1. Determine ε
2. Determine the order required, $A(\omega_s)$
3. Determine the poles, p_k
4. Determine the transfer function, $T(s)$

First-Order and Second-Order Filter Functions

- Nth-order response is very hard to visualize → Simple filter transfer functions
 - first and second order
- Cascade design.
 - possible for the design of active filters (utilizing op amps and RC circuits).
 - OP-amp output : low impedance
- High-order transfer function $T(s)$ can be factored into the product of first-order and second-order functions.

First-Order Filter Function

□ First-Order Filter Function

$$T(s) = \frac{a_1 s + a_0}{s + \omega_0}$$

→ bilinear transfer function

— natural mode at $s = -\omega_0$

— transmission zero at $s = -\frac{a_0}{a_1}$

— high-frequency gain = a_1

— The numerator coefficients, a_0 and a_1 , determine the type of filter

First-Order Filter Function

Filter Type and $T(s)$	s-Plane Singularities	Bode Plot for $ T $	Passive Realization	Op Amp-RC Realization
(a) Low pass (LP) $T(s) = \frac{a_0}{s + \omega_0}$	<div style="border: 1px solid blue; padding: 2px; display: inline-block; color: blue; font-weight: bold;">Poles can only be at real axis</div> 		$CR = \frac{1}{\omega_0}$ DC gain = 1	$CR_2 = \frac{1}{\omega_0}$ DC gain = $-\frac{R_2}{R_1}$
(b) High pass (HP) $T(s) = \frac{a_1 s}{s + \omega_0}$			$CR = \frac{1}{\omega_0}$ High-frequency gain = 1	$CR_1 = \frac{1}{\omega_0}$ High-frequency gain = $-\frac{R_2}{R_1}$
(c) General $T(s) = \frac{a_1 s + a_0}{s + \omega_0}$			$(C_1 + C_2)(R_1 // R_2) = \frac{1}{\omega_0}$ $C_1 R_1 = \frac{a_1}{a_0}$ DC gain = $-\frac{R_2}{R_1 + R_2}$ HF gain = $-\frac{C_1}{C_1 + C_2}$	$C_2 R_2 = \frac{1}{\omega_0}$ $C_1 R_1 = \frac{a_1}{a_0}$ DC gain = $-\frac{R_2}{R_1}$ HF gain = $-\frac{C_1}{C_2}$

First-Order Filter Function

$T(s)$	Singularities	$ T $ and ϕ	Passive Realization	Op Amp-RC Realization
<p>All pass (AP)</p> $T(s) = -a_1 \frac{s - \omega_0}{s + \omega_0}$ <p>$a_1 > 0$</p>			<p>$CR = 1/\omega_0$ Flat gain (a_1) = 0.5</p>	<p>$CR = 1/\omega_0$ Flat gain (a_1) = 1</p>

- Although the transmission is constant, its phase shows frequency selectivity
- All-pass filters are used as phase shifters and in systems that require phase shaping

Second-Order Filter Function

□ Second-Order-Filters

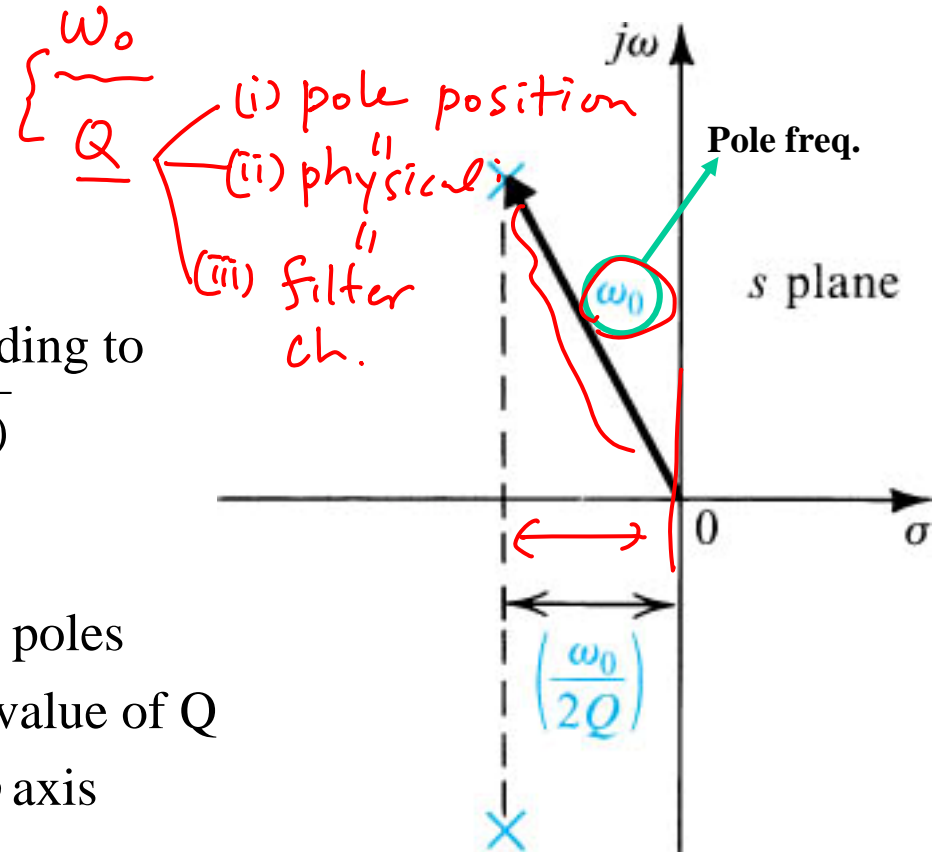
$$T(s) = \frac{a_2 s^2 + a_1 s + a_0}{s^2 + (\omega_0 / Q) s + \omega_0^2}$$

- Where ω_0 and Q determine the natural modes (poles) according to

$$p_1, p_2 = -\frac{\omega_0}{2Q} \pm j\omega_0 \sqrt{1 - (1/4Q^2)}$$

— ω_0 = pole frequency

- Q determines the distance of the poles from the $j\omega$ axis: the higher the value of Q the closer the poles are to the $j\omega$ axis
→ more selective
- $Q < 0$ → poles are in the RHP → oscillations
- Q = pole quality factor = pole Q



Filter Type and $T(s)$	s-Plane Singularities	$ T $
<p>(a) Low pass (LP)</p> $T(s) = \frac{a_0}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>DC gain = $\frac{a_0}{\omega_0^2}$</p>	<p>s-Plane Singularities</p>	<p>T</p> <p>Butterworth</p>
<p>(b) High pass (HP)</p> $T(s) = \frac{a_2 s^2}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>High-frequency gain = a_2</p>	<p>s-Plane Singularities</p>	<p>T</p>
<p>(c) Bandpass (BP)</p> $T(s) = \frac{a_1 s}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>Center-frequency gain = $\frac{a_1 Q}{\omega_0}$</p>	<p>s-Plane Singularities</p>	<p>T</p>

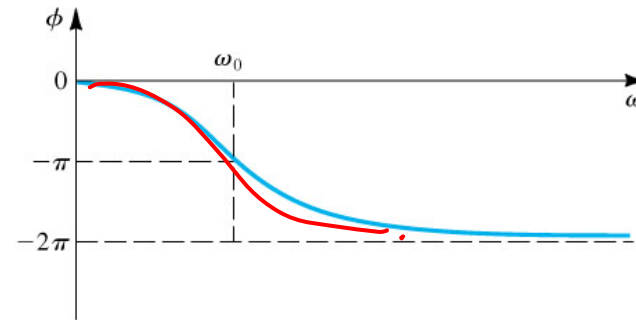
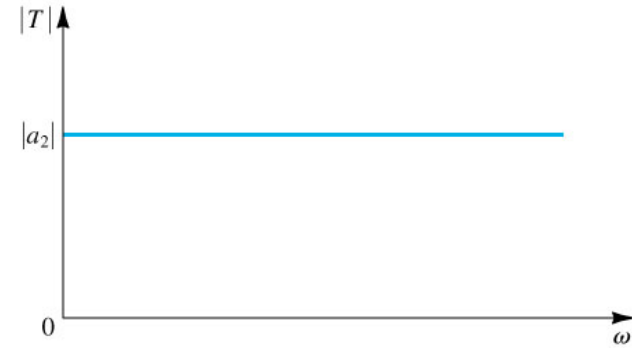
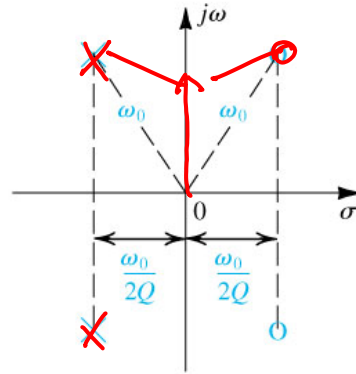
Filter Type and $T(s)$	s-Plane Singularities	$ T $
<p>(d) Notch</p> $T(s) = a_2 \frac{s^2 + \omega_0^2}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>DC gain = a_2 High-frequency gain = a_2</p>		
<p>(e) Low-pass notch (LPN)</p> $T(s) = a_2 \frac{s^2 + \omega_n^2}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>DC gain = $a_2 \frac{\omega_n^2}{\omega_0^2}$ High-frequency gain = a_2</p> <p>$\omega_n \geq \omega_0$</p>		$\omega_{\max} = \omega_0 \sqrt{\frac{\left(\frac{\omega_n^2}{\omega_0^2}\right)\left(1 - \frac{1}{2Q^2}\right) - 1}{\frac{\omega_n^2}{\omega_0^2} + \frac{1}{2Q^2} - 1}}$
<p>(f) High-pass notch (HPN)</p> $T(s) = a_2 \frac{s^2 + \omega_n^2}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$ <p>DC gain = $a_2 \frac{\omega_n^2}{\omega_0^2}$ High-frequency gain = a_2</p> <p>$\omega_n \leq \omega_0$</p>		$T_{\max} = \frac{ a_2 \omega_n^2 - \omega_{\max}^2 }{\sqrt{(\omega_0^2 - \omega_{\max}^2)^2 + \left(\frac{\omega_0}{Q}\right)^2 \omega_{\max}^2}}$

Second-Order Filter Function

(g) All pass (AP)

$$T(s) = a_2 \frac{s^2 - s\frac{\omega_0}{Q} + \omega_0^2}{s^2 + s\frac{\omega_0}{Q} + \omega_0^2}$$

Flat gain = a_2



Second-Order Filter Function

- LP case: The peak occurs only for $Q > \frac{1}{\sqrt{2}}$

$$Q = \frac{1}{\sqrt{2}} \rightarrow \text{Butterworth, or maximally flat}$$

- HP case: Transmission zeros at $s=0$

Dual to LP

- BP case: Transmission zeros at $s=0$ and $s=\infty$

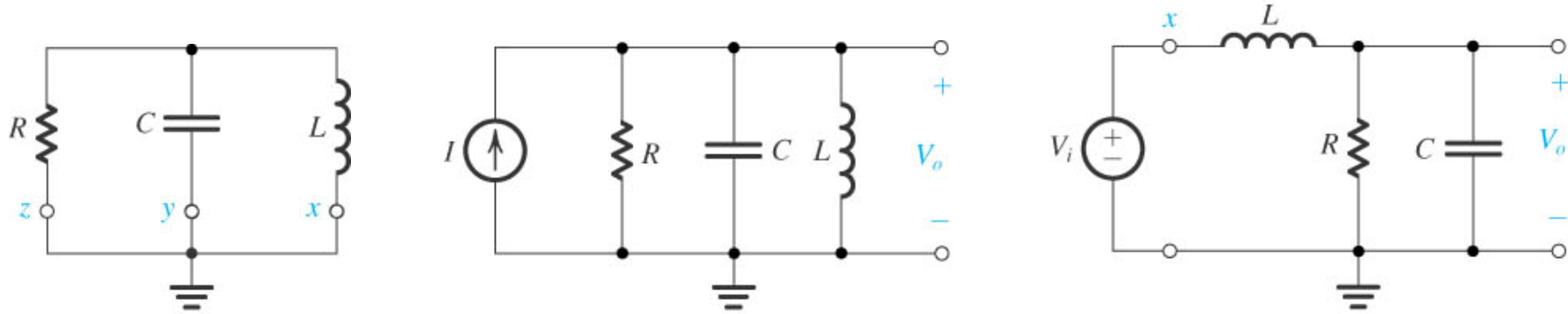
Magnitude response peaks at $\omega = \omega_o = \text{center frequency}$

$$3\text{dB: } \omega_1, \omega_2 = \omega_o \sqrt{1 + \frac{1}{4Q^2}} \pm \frac{\omega_o}{2Q}$$

$$\text{BW} = \omega_2 - \omega_1 = \frac{\omega_o}{Q} \quad \text{:as } Q \uparrow \text{ BW } \downarrow \text{ more selective}$$

The Second-order LCR Resonator : Natural Modes

□ The Resonator Natural Modes



- The natural modes can be determined by applying an excitation that does not change the natural structure of the circuit

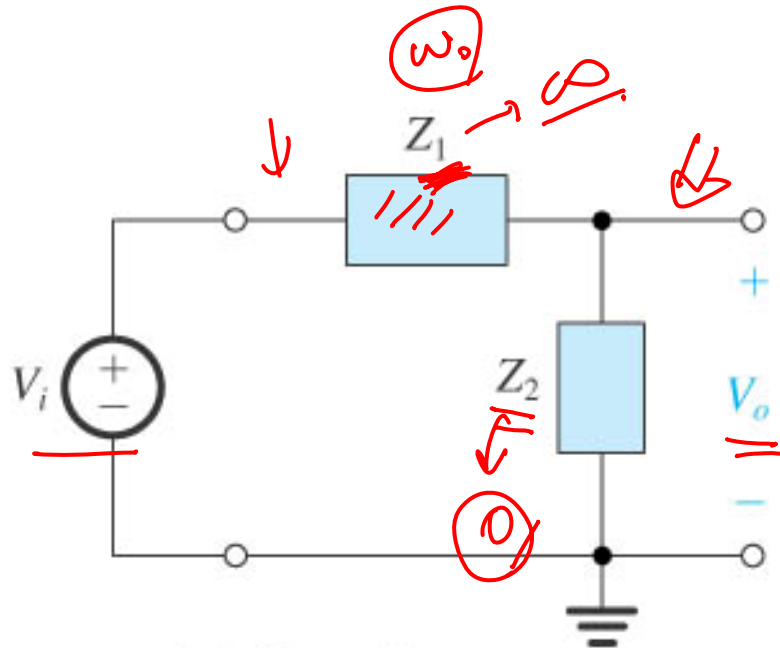
$$\downarrow \frac{V_o}{I} = \frac{1}{Y} = \frac{1}{\frac{1}{sL} + sC + \frac{1}{R}} = \frac{\frac{s}{C}}{s^2 + s\left(\frac{1}{CR}\right) + \frac{1}{LC}} \quad \leftarrow = \quad \frac{as}{s^2 + \left(\frac{\omega_0}{Q}\right)s + \omega_0^2} \quad \left| \begin{array}{l} \omega_0 \\ Q \end{array} \right.$$

$$\omega_0^2 = \frac{1}{LC} \quad ; \quad \frac{\omega_0}{Q} = \frac{1}{CR} \quad ; \quad \omega_0 = \frac{1}{\sqrt{LC}} \quad ; \quad Q = \omega_0 CR = \omega_0 \frac{L}{R} \quad ||$$

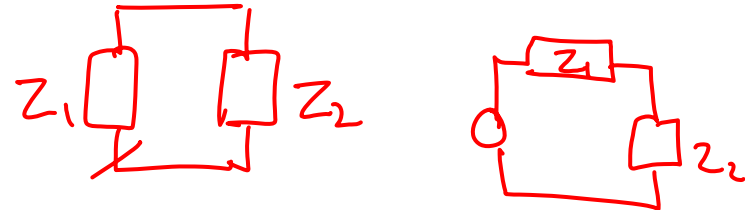
$\omega_0 L = \frac{1}{\omega_0 C}$

The Second-order LCR Resonator : Adding Zeros

□ Realization of Transmission Zeros

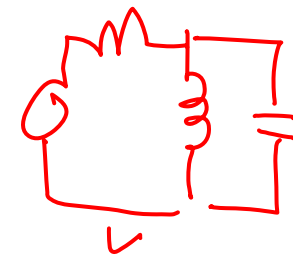
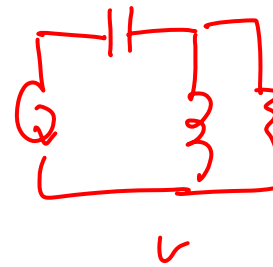
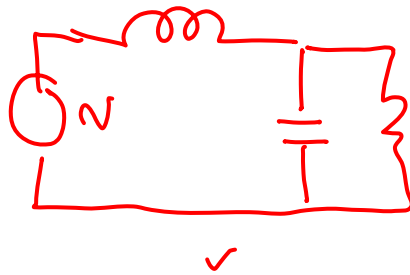
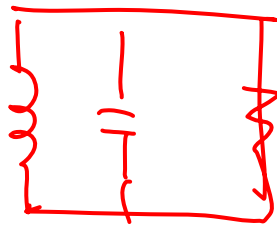
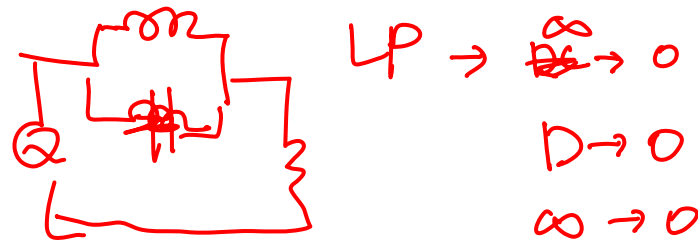


(a) General structure



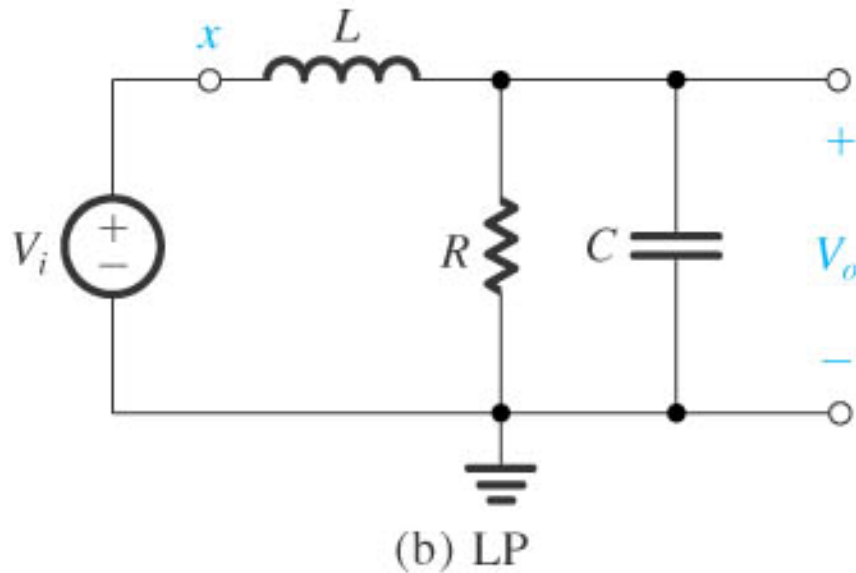
$$T(s) = \frac{V_o(s)}{V_i(s)} = \frac{Z_2(s)}{Z_1(s) + Z_2(s)}$$

zeros: $Z_2(s) = 0$ while $Z_1(s) \neq 0$
 $Z_1(s) = \infty$ while $Z_2(s) \neq \infty$



The Second-order LCR Resonator

□ Realization of the Low-Pass Function



■ zeros $sL = \infty$

$$\frac{1}{(sC + \frac{1}{R})} = 0$$

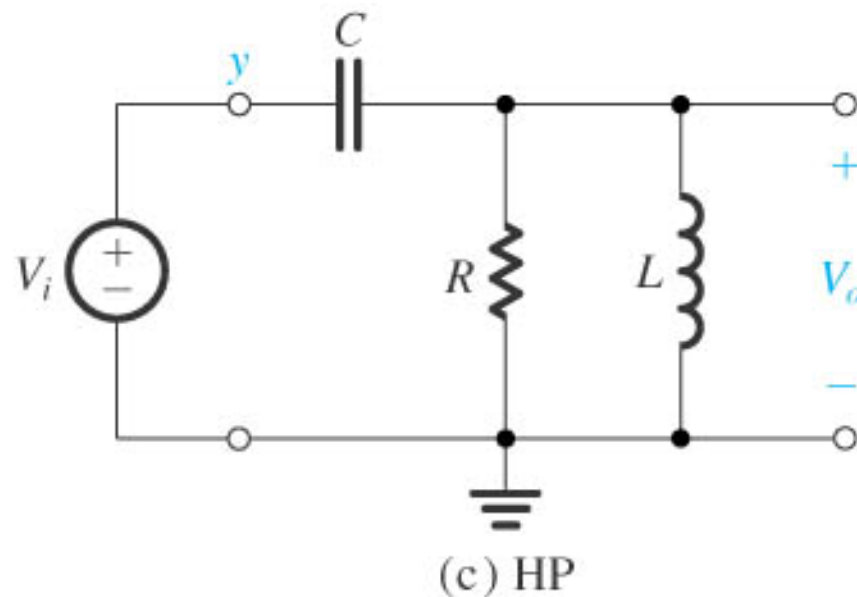
\Rightarrow two zeros at $s = \infty$

$$T(s) \equiv \frac{V_o}{V_i} = \frac{Z_2}{Z_1 + Z_2} = \frac{Y_1}{Y_1 + Y_2} = \frac{\frac{1}{sL}}{\frac{1}{sL} + sC + \frac{1}{R}}$$

$$= \frac{1/LC}{s^2 + s(1/CR) + 1/LC}$$

The Second-order LCR Resonator

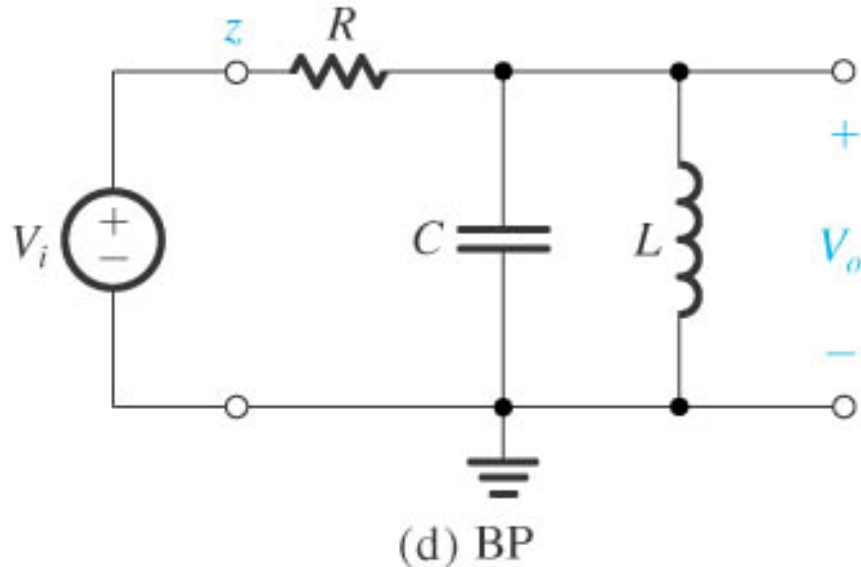
□ Realization of the High-Pass Function



- zeros $\begin{cases} s = 0 : \text{Capacitor} \\ s = 0 : \text{Inductor} \end{cases}$
- $$T(s) \equiv \frac{V_o}{V_i} = \frac{a_2 s^2}{s^2 + s\left(\frac{\omega_o}{Q}\right) + \omega_o^2}$$
- $$a_2 = 1$$

The Second-order LCR Resonator

□ Realization of the Band-Pass Function



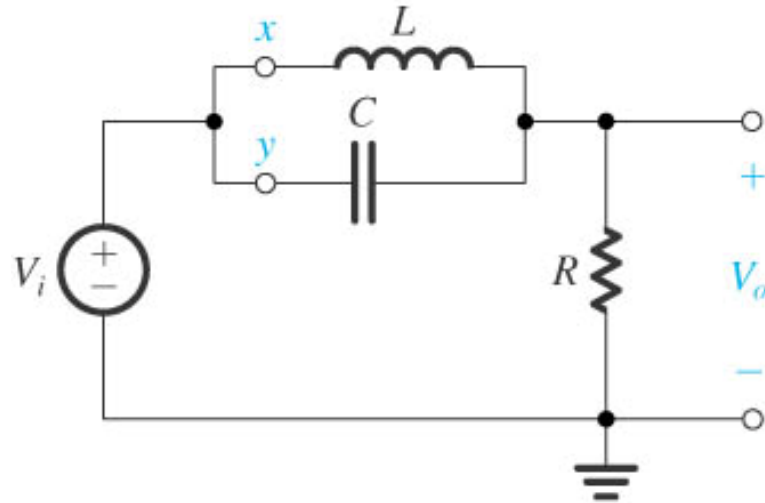
- zeros $\begin{cases} s = 0 : \text{Inductor} \\ s = \infty : \text{Capacitor} \end{cases}$

at ω_0 ,
 LC-tuned circuit exhibits an infinite impedance \rightarrow no current flows
 the center freq. gain is unity

$$T(s) = \frac{Y_R}{Y_R + Y_L + Y_C} = \frac{\frac{1}{R}}{\frac{1}{R} + \frac{1}{sL} + sC}$$

$$= \frac{s\left(\frac{1}{CR}\right)}{s^2 + s\left(\frac{1}{CR}\right) + \left(\frac{1}{LC}\right)}$$

Realization of the Notch Functions



(e) Notch at ω_0

- The impedance of the LC circuit becomes infinite at $\omega_0 = 1/\sqrt{LC}$
 \rightarrow zero transmission

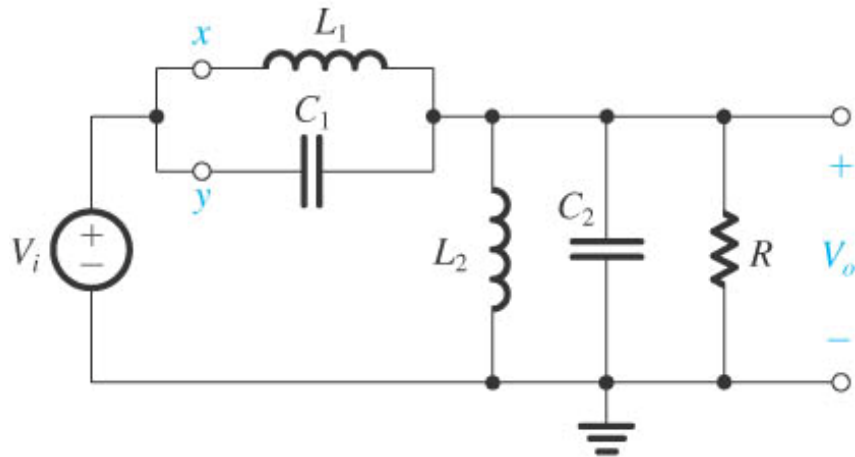
- The resistor does not introduce zeros.

$$T(s) = a_2 \frac{s^2 + \omega_0^2}{s^2 + s(\omega_0/Q) + \omega_0^2}$$

- To obtain arbitrary ω_n

LPN

$$L_1 C_1 = \frac{1}{\omega_n^2}$$



(f) General notch

- $L_1 C_1$ tank will introduce a pair of zeros at $\pm j\omega_n$, provided the $L_2 C_2$ tank is not resonant at ω_n .

- The natural modes have not been altered,
 $C_1 + C_2 = C$; $L_1 \parallel L_2 = L$

- It is obtained from the original LCR resonator by lifting part of L and part of C off ground

Realization of the All-Pass Function

- The all-pass transfer function

$$T(s) = \frac{s^2 - s(\omega_o/Q) + \omega_o^2}{s^2 + s(\omega_o/Q) + \omega_o^2}$$

$$= 1 - \boxed{\frac{s2(\omega_o/Q)}{s^2 + s(\omega_o/Q) + \omega_o^2}} \rightarrow \text{bandpass function with a center-frequency gain of 2}$$

↓

- All pass realization with a flat gain of 0.5

$$T(s) = 0.5 - \frac{s(\omega_o/Q)}{s^2 + s(\omega_o/Q) + \omega_o^2}$$

