

457.643 Structural Random Vibrations In-Class Material: Class 25

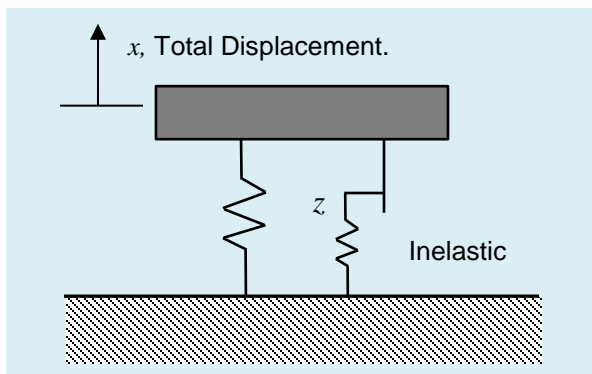
VI. Nonlinear Random Vibration Analysis

© (Differential equation based) hysteretic constitutive models in structural dynamics

“Hysteresis”

- Origin: ferromagnetic materials
- Memory-based multi-valued relation between an input signal & output (generally referring to “rate-independent” relationship only; ~~viscous materials~~)

Mechanical model for differential equation based hysteresis model



z : Auxiliary variable representing inelastic behavior (“internal variable” – Capecchi & de Felice 2001, ASCE JEM) ~ displacement of inelastic spring

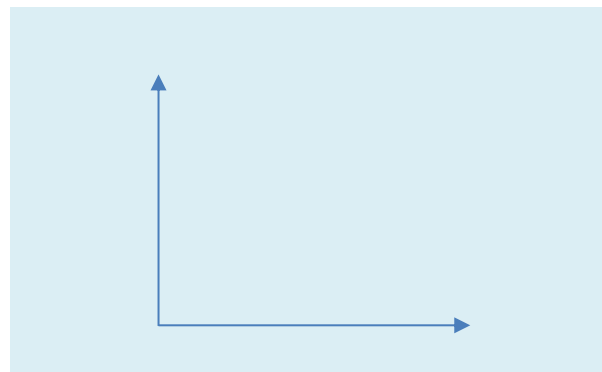
- ♦ $z = x$: no slide
- ♦ $z = 0$: slide

(nonlinearity determined by difference between z and x)

Resisting force:

$$f_s(x, z) = \alpha k_0 x + (1 - \alpha) k_0 z$$

- ♦ α : post-to-pre-yield stiffness ratio
 - ✓ $\alpha = 0$: perfectly plastic
 - ✓ $\alpha = 1$: linear elastic
- ♦ k_0 : initial stiffness



Evolution of z follows a nonlinear differential equation

$$\dot{z} = \dot{x} \cdot h(x, \dot{x}, z)$$

Meaning of the nonlinear function $h(\cdot)$?

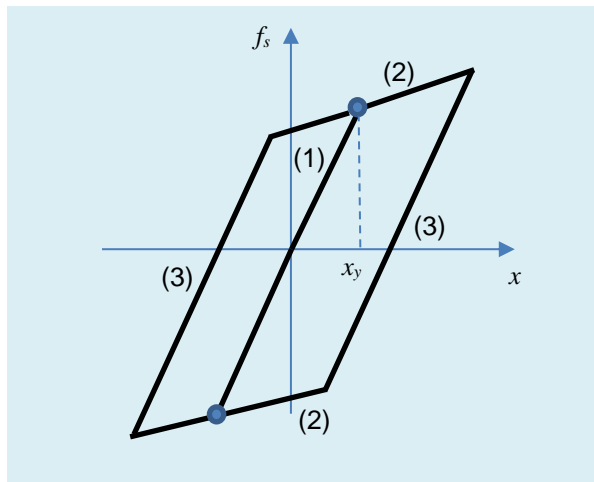
Song, J., and A. Der Kiureghian (2006). Generalized Bouc-Wen model for highly asymmetric hysteresis. *Journal of Engineering Mechanics*. ASCE, 132(6), 610-618.

$$\frac{dz}{dt} = \frac{dx}{dt} \cdot h(\cdot)$$

Therefore, $h(\cdot) (= \frac{dz}{dx})$ determines the slope of z with respect to x at a given time.

© **Bilinear model (Kaul & Penzien 1974 JEMD; Asano & Iwan 1984 EESD)**

Main idea: describe inelastic spring in the mechanical model by a Coulomb slider (i.e. no slide until it reaches the yield displacement)



(1) $-x_y < z < x_y$:

the Coulomb slider does not slide, i.e.

$$z = x \text{ and } \dot{z} = \dot{x}$$

$$f_s(x, z) = \alpha k_0 x + (1 - \alpha) k_0 x = k_0 x \text{ (linear)}$$

(2) $z > x_y, \dot{x} > 0$ or $z < -x_y, \dot{x} < 0$:

Coulomb slider slides (i.e. $\dot{z} = 0$)

(3) $z > x_y, \dot{x} < 0$ or $z < -x_y, \dot{x} > 0$:

Coulomb slider stops sliding $\dot{z} = \dot{x}$

Differential-equation model by Kaul & Penzien (1974):

$$\dot{z} = \dot{x} \cdot \{U(z + x_y) - U(z - x_y) + U(z - x_y) \cdot U(-\dot{x}) + U(-z - x_y) \cdot U(\dot{x})\}$$

where $U(\cdot)$ denotes the unit step function.

How to solve the nonlinear system differential equation, i.e.

E.O.M. with $f_s = \alpha k_0 x + (1 - \alpha) k_0 z$ plus $\dot{z} = \dot{x} \cdot h(x, \dot{x}, z)$

e.g. Runge-Kutta method (after transforming to state-space formulation $\dot{y} = g(y) + f$)

© **Bouc-Wen class model**

Bouc (1967) first proposed, and Wen (1976) later modified to the form

$$\dot{z} = \dot{x} \cdot [A - |z|^n \psi(x, \dot{x}, z)]$$

where

- A : scale of hysteresis loop
- n : smoothness of transition from pre-yielding to post-yielding
- $\psi(x, \dot{x}, z)$: “shape-control” function

Reviews are available in Song & ADK (2006, JEM), and Ismail et al. (2009, Archi. Comp. Meth. Engrg.)

1) Bouc (1967, 1971)

- ♦ $n = 1$
- ♦ $\psi(x, \dot{x}, z) = \gamma + \beta \text{sgn}(\dot{x}z)$

2) Wen (1976)

- ♦ n : generalized
- ♦ $\psi(x, \dot{x}, z) = \gamma + \beta \text{sgn}(\dot{x}z)$

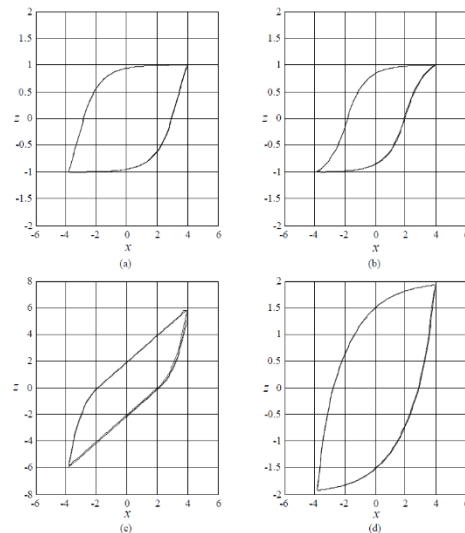


Figure 3.3 Hysteresis loops by Bouc-Wen model ($A=1, n=1$) (a) $\gamma=0.5, \beta=0.5$, (b) $\gamma=0.1, \beta=0.9$, (c) $\gamma=0.5, \beta=-0.5$ and (d) $\gamma=0.75, \beta=-0.25$

The parameters γ and β in the “shape-control” function determine the shapes of the hysteresis loops (Song and ADK 2006)

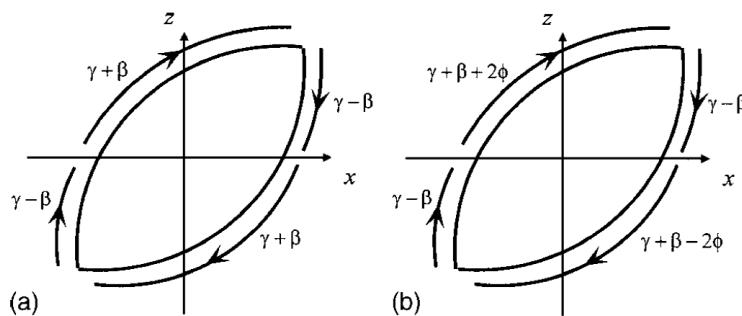


Fig. 2. Values of shape-control function for: (a) original Bouc–Wen model; and (b) model by Wang and Wen

- 3) Baber & Wen (1981): Considered the degradation effect by making the model parameters functions of ϵ , “the dissipated energy”
- 4) Baber & Noori (1984): Introduced additional parameters to describe “pinching” effect
- 5) Wang & Wen (1998): Described “asymmetric” shape by adding additional terms

$$\psi(x, \dot{x}, z) = \gamma + \beta \operatorname{sgn}(\dot{x}z) + \phi[\operatorname{sgn}(\dot{x}) + \operatorname{sgn}(z)]$$

➔ Added more DOFs (see the figure above)

- 6) Generalized Bouc-Wen (Song & ADK, 2006)

Generalize the “shape-control” function to describe highly asymmetric behavior

$$\psi(x, \dot{x}, z) = \beta_1 \operatorname{sgn}(\dot{x}z) + \beta_2 \operatorname{sgn}(x\dot{x}) + \beta_3 \operatorname{sgn}(xz) + \beta_4 \operatorname{sgn}(\dot{x}) + \beta_5 \operatorname{sgn}(z) + \beta_6 \operatorname{sgn}(x)$$

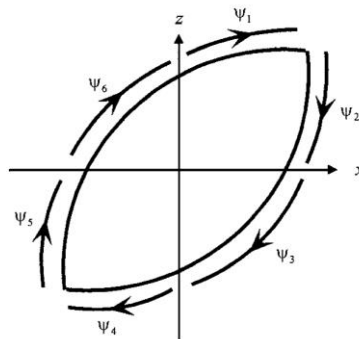


Fig. 3. Values of shape-control function for generalized Bouc–Wen model

Six phases can now have all different values, and the values are determined as

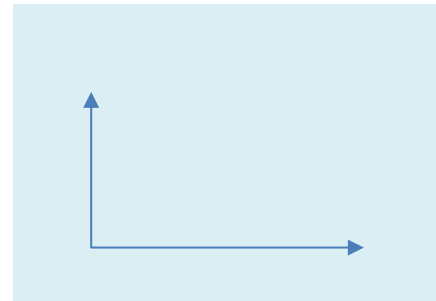
$$\begin{bmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \\ \psi_5 \\ \psi_6 \end{bmatrix} = \begin{bmatrix} 1 & 1 & 1 & 1 & 1 & 1 \\ -1 & -1 & 1 & -1 & 1 & 1 \\ 1 & -1 & -1 & -1 & -1 & 1 \\ 1 & 1 & 1 & -1 & -1 & -1 \\ -1 & -1 & 1 & 1 & -1 & -1 \\ 1 & -1 & -1 & 1 & 1 & -1 \end{bmatrix} \begin{bmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \\ \beta_4 \\ \beta_5 \\ \beta_6 \end{bmatrix}$$

The model parameters $\beta_i, i = 1, \dots, 6$ can be fitted by use of

$$\begin{bmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \\ \beta_4 \\ \beta_5 \\ \beta_6 \end{bmatrix} = \frac{1}{4} \begin{bmatrix} 1 & 0 & 1 & 1 & 0 & 1 \\ 0 & -1 & -1 & 0 & -1 & -1 \\ 1 & 1 & 0 & 1 & 1 & 0 \\ 1 & -1 & 0 & -1 & 1 & 0 \\ 0 & 1 & -1 & 0 & -1 & 1 \\ 1 & 0 & 1 & -1 & 0 & -1 \end{bmatrix} \begin{bmatrix} \psi_1 \\ \psi_2 \\ \psi_3 \\ \psi_4 \\ \psi_5 \\ \psi_6 \end{bmatrix}$$

※ Weakness of Bouc-Wen class models:

- ♦ can violate the requirement of classical plasticity theory (“Drucker’s postulate”; Bažant 1978); can create negative dissipative energy when “loading-unloading” occurs without load reversal
- ♦ But this problem is not critical if $E[f_s] \cong 0$ (Wen 1989, Hurtado & Barbat 1996)



※ Bouc-Wen class models are widely-used in structural dynamics and earthquake engineering because

- 1) Can describe a wide-class of phenomena (pinching, degradation, etc.)
- 2) Facilitates efficient time history analysis (no IF or THEN)
- 3) Facilitates efficient random vibration analysis

e.g. Nonlinear random vibration analysis for Bouc-Wen model by Equivalent Linearization Method (Wen 1980)

◎ **Nonlinear time-history analysis of structural system with Bouc-Wen class models**

$$\mathbf{M}\ddot{\mathbf{x}} + \mathbf{C}\dot{\mathbf{x}} + \mathbf{R}(\mathbf{x}, \dot{\mathbf{x}}, \mathbf{z}) = -\mathbf{M}\mathbf{1}\ddot{x}_g$$

where $\mathbf{R}(\mathbf{x}, \dot{\mathbf{x}}, \mathbf{z})$ uses $f_s = \alpha k_0 x + (1 - \alpha)k_0 z$ to describe the resistant force of each B-W element. The auxiliary variable follows the nonlinear differential equation $\dot{z} = \dot{x} \cdot h(x, \dot{x}, z)$.

Transformed to state-space formulation, i.e. $\mathbf{y} = \{x_1, \dot{x}_1, x_2, \dot{x}_2, \dots, z_1, \dots, z_m\}$

Example: Two connected equipment items in an electrical substation (Song, 2004)

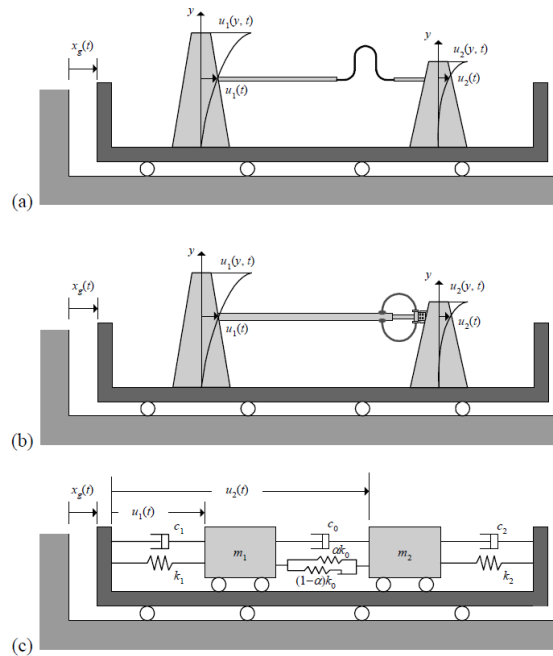


Figure 2.1 Mechanical models of equipment items connected by rigid bus connectors: (a) RB-FSC-connected system, (b) Bus-slider-connected system, and (c) idealized system with SDOF equipment models

$$\dot{\mathbf{y}} = \mathbf{g}(\mathbf{y}) + \mathbf{f}$$

where

$$\mathbf{y} = \{u_1, \dot{u}_1, u_2, \dot{u}_2, z\}^T$$

$$\mathbf{g}(\mathbf{y}) = \begin{Bmatrix} \dot{u}_1 \\ -\left(\frac{k_1 + \alpha k_0}{m_1}\right)u_1 - \left(\frac{c_1 + c_0}{m_1}\right)\dot{u}_1 + \frac{\alpha k_0}{m_1}u_2 + \frac{c_0}{m_1}\dot{u}_2 + \frac{(1-\alpha)k_0}{m_1}z \\ \dot{u}_2 \\ \frac{\alpha k_0}{m_2}u_1 + \frac{c_0}{m_2}\dot{u}_1 - \left(\frac{k_2 + \alpha k_0}{m_2}\right)u_2 - \left(\frac{c_2 + c_0}{m_2}\right)\dot{u}_2 - \frac{(1-\alpha)k_0}{m_2}z \\ \Delta \dot{\mathbf{u}} \cdot \mathbf{h}(\Delta \mathbf{u}, \Delta \dot{\mathbf{u}}, z) \end{Bmatrix}$$

$$\mathbf{f} = \left\{ 0 \quad -\frac{l_1}{m_1}\ddot{x}_g \quad 0 \quad -\frac{l_2}{m_2}\ddot{x}_g \quad 0 \right\}^T$$

Can solve the differential equation by a numerical method such as the fourth and fifth order Runge-Kutta-Fehlberg (RKF) algorithm.

© **Equivalent Linearization Method (ELM; a.k.a. stochastic linearization method)**

Among various methods such as Fokker-Planck equation, stochastic averaging, moment closure, perturbation (Lutes and Sarkani 2004), ELM is considered as a nonlinear random vibration approach with the highest potential for practical use (Pradlwarter & Schuëller 1991)

- Applicable to both stationary and nonstationary processes
- Applicable to a wide class of nonlinear behavior
- Can handle MDOF systems and FE models
- Takes significantly less efforts than Monte Carlo simulations (especially for low-probability events)

Consider an original nonlinear system: $\dot{\mathbf{y}} = \mathbf{g}(\mathbf{y}) + \mathbf{f}$:

One can find an “equivalent linear” system: $\dot{\mathbf{y}}_e = \mathbf{A} \cdot \mathbf{y}_e + \mathbf{f}$ such that the mean-square error (caused by linearization) $E[(\mathbf{g}(\mathbf{y}) - \mathbf{A}\mathbf{y})^T(\mathbf{g}(\mathbf{y}) - \mathbf{A}\mathbf{y})]$ is minimized.

Note: ELM based on the error definition above is considered “standard” ELM while the error measure $E[(\mathbf{g}(\mathbf{y}_e) - \mathbf{A}\mathbf{y}_e)^T(\mathbf{g}(\mathbf{y}_e) - \mathbf{A}\mathbf{y}_e)]$ is called “SPEC-alternative” ELM (Crandall 2001).

Crandall, S.H. (2001) Is stochastic equivalent linearization a subtly flawed procedure? *Probabilistic Engineering Mechanics*, 16:169-176

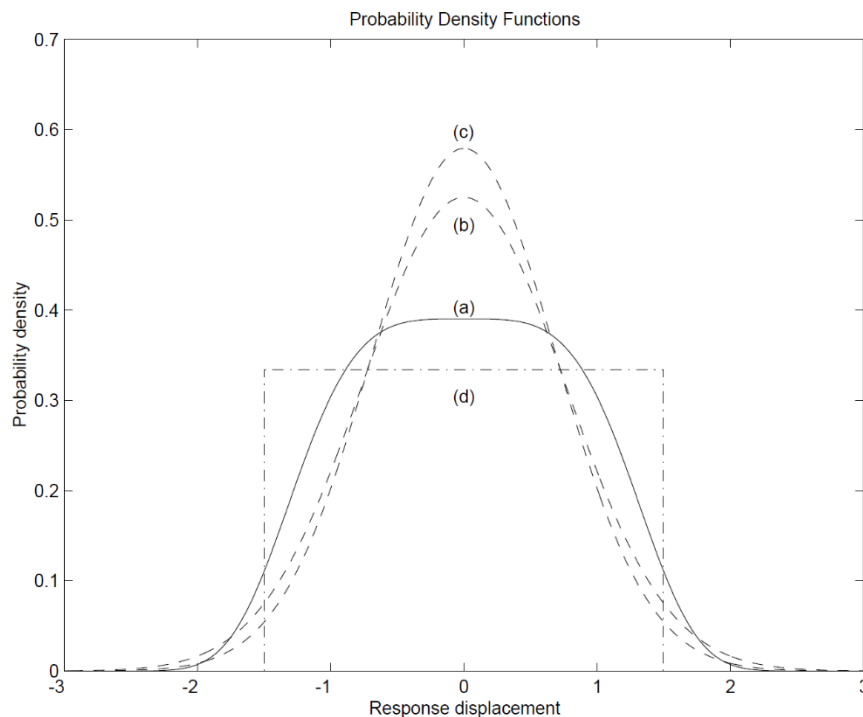


Fig. 2. Probability density functions for cubic oscillator.