Chapter 7

Applications of the Second Law

Advanced Thermodynamics Min Soo Kim



For reversible process, $\delta q_r = du + P \, dv$

- 1. Adiabatic process : $\delta q_r = 0, ds = 0, s = constant$
- 2. Isothermal process : $s_2 s_1 = \int_1^2 \frac{\delta q_r}{r} = \frac{q_r}{r}$



3. Isothermal (and isobaric) change of phase : $s_2 - s_1 = \frac{l}{r}$

4. Isochoric process :
$$s_2 - s_1 = \int_1^2 c_v \frac{dT}{T} = c_v \ln \frac{T_2}{T_1}$$

$$\frac{\delta q_r}{T} = \frac{dh}{T} - \frac{v}{T}dP = ds$$

$$s_2 - s_1 = \int_1^2 c_p \frac{dT}{T} = c_p \ln \frac{T_2}{T_1}$$



The total quantity of heat transferred in a reversible process from state 1 to state 2 is given by

$$q_r = \int_1^2 T \, ds$$

The T-s diagram is simple rectangle for a Carnot cycle. The area under the curve is

$$\oint Tds = \sum q_r = w$$

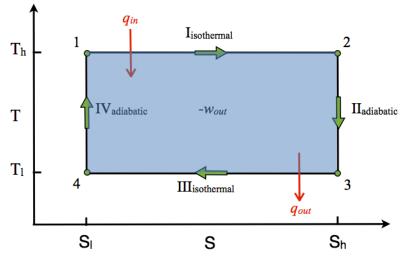


Figure7.1 T-s diagram for a Carnot cycle [1]

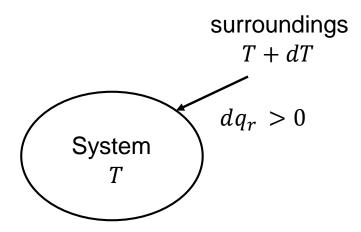


The heat flow out of the surroundings at every point is equal in magnitude and opposite in sign to the heat flow into the system.

$$dq_{in} = dq_{out} = dq_r$$

For a reversible process,

 $dT \ll 1$





7.3 Entropy Change of the Surroundings (Reversible)

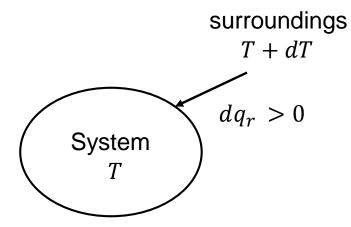
So,

$$\left(\frac{\delta q_r}{T+dT}\right)_{surroundings} \approx \left(\frac{\delta q_r}{T}\right)_{surroundings} = (\delta s)_{surroundings}$$

and from

 $ds_{system} + ds_{surroundings} = ds_{universe}$,

$$|ds|_{surroundings} = -|ds|_{system} \quad \& \quad |ds|_{universe} = 0$$



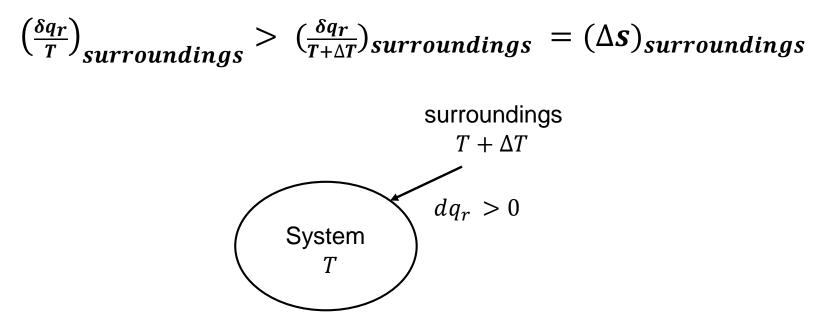


7.3* Entropy Change of the Surroundings (Irreversible)

However for an irreversible case,

 $\Delta T > 0$

and



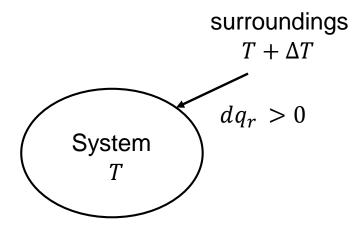


7.3* Entropy Change of the Surroundings (Irreversible)

So,

 $\Delta s_{system} + \Delta s_{surroundings} = \Delta s_{universe} > 0$

(Entropy generation!)





7.4 Entropy change for an Ideal Gas

With $du = c_v dT$, we have

$$\frac{\vec{d}q_r}{T} = \frac{c_v dT}{T} + \frac{P}{T} dv = ds$$

For a reversible process, For an ideal gas, P/T = R/v, so

$$ds = c_v \frac{dT}{T} + R \frac{dv}{v}$$

Integrating, we have

$$s_2 - s_1 = c_v \ln \left(\frac{T_2}{T_1}\right) + R \ln \left(\frac{v_2}{v_1}\right)$$



From the combined first and second laws, Tds = du + P dv

$$Tds = c_{v} dT + T \left(\frac{\partial P}{\partial T}\right)_{v} dv \qquad (s = s(T, v))$$

$$Tds = c_p dT - T \left(\frac{\partial v}{\partial T}\right)_p dP \qquad (s = s(T, P))$$

$$Tds = c_p(\frac{\partial T}{\partial \nu})_p \, d\nu + c_\nu \, (\frac{\partial T}{\partial P})_\nu \, dP \ (s = s(\nu, P))$$



Let T and P be the independent variables .

The enthalpy is $h \equiv u + P v$ thus,

$$Tds = dh - vdP = \left(\frac{\partial h}{\partial T}\right)_P dT + T \left(\frac{\partial h}{\partial P}\right)_T dP - vdP$$
$$ds = \frac{1}{T} \left(\frac{\partial h}{\partial T}\right)_P dT + \frac{1}{T} \left[\left(\frac{\partial h}{\partial P}\right)_T - v\right] dP$$



7.5 The Tds Equations

With s = s(T, P), we have

$$ds = (\frac{\partial s}{\partial T})_P dT + (\frac{\partial s}{\partial P})_T dP$$

Since T and P are independent, it follows that

$$(\frac{\partial s}{\partial T})_P = \frac{1}{T} (\frac{\partial h}{\partial T})_P \qquad and \qquad (\frac{\partial s}{\partial T})_T = \frac{1}{T} [(\frac{\partial h}{\partial P})_T - v]$$



The differential ds is exact. Therefore,

$$\begin{bmatrix}\frac{\partial}{\partial P}(\frac{\partial s}{\partial T})_P\end{bmatrix}_T = \frac{\partial^2 s}{\partial P \partial T} = \frac{\partial^2 s}{\partial T \partial P} = \begin{bmatrix}\frac{\partial}{\partial P}(\frac{\partial s}{\partial T})_T\end{bmatrix}_P$$

Substituting last two Equations from previous slide, we get

$$\frac{1}{T}\frac{\partial^2 h}{\partial P \partial T} = \frac{1}{T} \left[\frac{\partial^2 h}{\partial T \partial P} - \left(\frac{\partial \nu}{\partial T} \right)_P \right] - \frac{1}{T^2} \left[\left(\frac{\partial h}{\partial P} \right)_T - \nu \right]$$

$$\leftrightarrow \ (\frac{\partial h}{\partial P})_T = -T \ (\frac{\partial v}{\partial T})_P + v$$



7.5 The Tds Equations

For a reversible process, $(\partial h/\partial T)_P = c_P$

$$Tds = c_P dT - T \left(\frac{\partial v}{\partial P}\right)_P dP$$

Finally, since $(\partial v / \partial T)_P = v\beta$, we have

$$Tds = c_P dT - T \, \nu \beta dP$$

