

# Electromagnetics:

## Wave Characteristics on Finite Transmission Lines

(9-4)

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# Lines with Resistive Termination (1)

Recall:

$$V(z') = \frac{I_L}{2} [(Z_L + Z_0)e^{\gamma z'} + (Z_L - Z_0)e^{-\gamma z'}]$$

$$I(z') = \frac{I_L}{2Z_0} [(Z_L + Z_0)e^{\gamma z'} - (Z_L - Z_0)e^{-\gamma z'}]$$

$$\rightarrow V(z') = \frac{I_L}{2} (Z_L + Z_0) e^{\gamma z'} \left[ 1 + \frac{Z_L - Z_0}{Z_L + Z_0} e^{-2\gamma z'} \right] \quad \leftarrow \Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = |\Gamma| e^{j\theta_\Gamma}$$

$$= \frac{I_L}{2} (Z_L + Z_0) e^{\gamma z'} \left[ 1 + \Gamma e^{-2\gamma z'} \right]$$

*Voltage reflection coefficient*

$$\rightarrow I(z') = \frac{I_L}{2Z_0} (Z_L + Z_0) e^{\gamma z'} \left[ 1 - \Gamma e^{-2\gamma z'} \right]$$

For a lossless transmission line:

$$V(z') = \frac{I_L}{2} (Z_L + R_0) e^{j\beta z'} \left[ 1 + \Gamma e^{-j2\beta z'} \right] = \frac{I_L}{2} (Z_L + R_0) e^{j\beta z'} \left[ 1 + |\Gamma| e^{j(\theta_\Gamma - 2\beta z')} \right]$$

$$I(z') = \frac{I_L}{2Z_0} (Z_L + R_0) e^{j\beta z'} \left[ 1 - \Gamma e^{-j2\beta z'} \right] = \frac{I_L}{2Z_0} (Z_L + R_0) e^{j\beta z'} \left[ 1 - |\Gamma| e^{j(\theta_\Gamma - 2\beta z')} \right]$$

# Lines with Resistive Termination (2)

*Also Recall:*

$$V(z') = I_L (Z_L \cosh \gamma z' + Z_0 \sinh \gamma z')$$

$$I(z') = \frac{I_L}{Z_0} (Z_L \sinh \gamma z' + Z_0 \cosh \gamma z')$$

For a lossless transmission line:

$$\rightarrow V(z') = V_L \cos \beta z' + j I_L R_0 \sin \beta z'$$

$$\rightarrow I(z') = I_L \cos \beta z' + j \frac{V_L}{R_0} \sin \beta z'$$

For purely resistive termination:

$$\rightarrow Z_L = R_L$$

$$\rightarrow V_L = I_L R_L$$

$$|V(z')| = V_L \sqrt{\cos^2 \beta z' + (R_0 / R_L)^2 \sin^2 \beta z'}$$

$$|I(z')| = I_L \sqrt{\cos^2 \beta z' + (R_L / R_0)^2 \sin^2 \beta z'}$$

← *Standing waves*

# Lines with Resistive Termination (3)

Recall for a lossless transmission line:

$$V(z') = \frac{I_L}{2} (Z_L + R_0) e^{j\beta z'} \left[ 1 + |\Gamma| e^{j(\theta_\Gamma - 2\beta z')} \right]$$

$$I(z') = \frac{I_L}{2Z_0} (Z_L + R_0) e^{j\beta z'} \left[ 1 - |\Gamma| e^{j(\theta_\Gamma - 2\beta z')} \right]$$

$$\leftarrow \Gamma = \frac{Z_L - R_0}{Z_L + R_0} = |\Gamma| e^{j\theta_\Gamma}$$

Standing-wave ratio (SWR):  $S = \frac{|V_{\max}|}{|V_{\min}|} = \frac{|I_{\max}|}{|I_{\min}|} = \frac{1 + |\Gamma|}{1 - |\Gamma|} \rightarrow |\Gamma| = \frac{S - 1}{S + 1}$

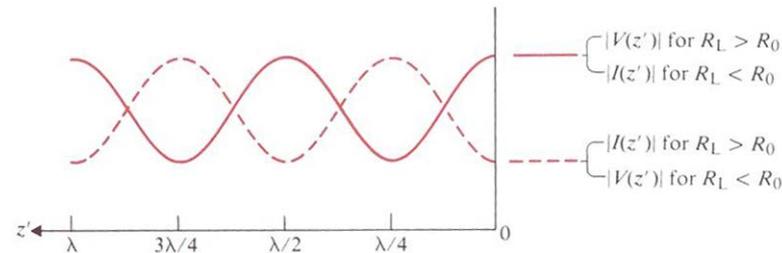
$\Gamma = 0 \rightarrow S = 1 \leftarrow Z_L = Z_0 \rightarrow$  *Matched load*  
 $\Gamma = -1 \rightarrow S = \infty \leftarrow Z_L = 0 \rightarrow$  *Short circuit*  
 $\Gamma = +1 \rightarrow S = \infty \leftarrow Z_L = \infty \rightarrow$  *Open circuit*

$\theta_\Gamma - 2\beta z'_M = -2n\pi \rightarrow V_{\max} \ \& \ I_{\min}$   
 $\theta_\Gamma - 2\beta z'_m = -(2n + 1)\pi \rightarrow V_{\min} \ \& \ I_{\max}$

For resistive terminations:  $\rightarrow \Gamma = \frac{R_L - R_0}{R_L + R_0}$

$R_L > R_0 \rightarrow V_{\max}$  *at the terminating resistance*

$R_L < R_0 \rightarrow V_{\min}$  *at the terminating resistance*



# Lines with Resistive Termination (4)

For resistive terminations:

$$|V(z')| = V_L \sqrt{\cos^2 \beta z' + (R_0 / R_L)^2 \sin^2 \beta z'} \quad \rightarrow \Gamma = \frac{R_L - R_0}{R_L + R_0}$$

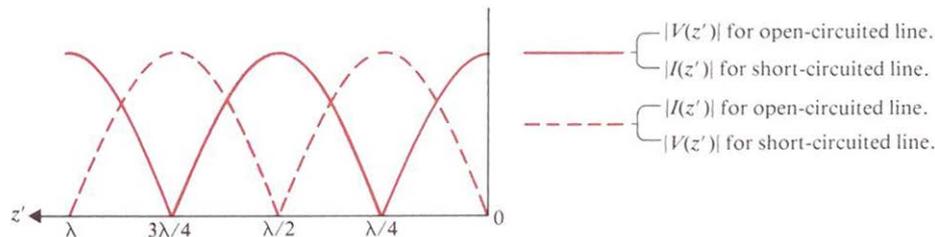
$$|I(z')| = I_L \sqrt{\cos^2 \beta z' + (R_L / R_0)^2 \sin^2 \beta z'}$$

$$R_L = \infty \quad \rightarrow |V(z')| = V_L |\cos \beta z'| \quad \rightarrow \Gamma = 1 \text{ \& } S \rightarrow \infty$$

$$\rightarrow |I(z')| = \frac{V_L}{R_0} |\sin \beta z'|$$

$$R_L = 0 \quad \rightarrow |V(z')| = I_L R_0 |\sin \beta z'| \quad \rightarrow \Gamma = -1 \text{ \& } S \rightarrow \infty$$

$$\rightarrow |I(z')| = I_L |\cos \beta z'| \quad \rightarrow \text{All the minima go to zero.}$$



# Lines with Arbitrary Termination

What if the terminating impedance is not a pure resistance?

$$\rightarrow Z_L = R_L + jX_L \leftarrow TL(l_m) \text{ terminated by } R_m$$

$$\leftarrow R_m < R_0$$

Recall:

$$Z_i = Z_0 \frac{Z_L + Z_0 \tanh \gamma l}{Z_0 + Z_L \tanh \gamma l}$$

$$\rightarrow R_i + jX_i = R_0 \frac{R_m + jR_0 \tan \beta l_m}{R_0 + jR_m \tan \beta l_m}$$

"Equivalent"

Experimental procedure to determine  $Z_L$ :

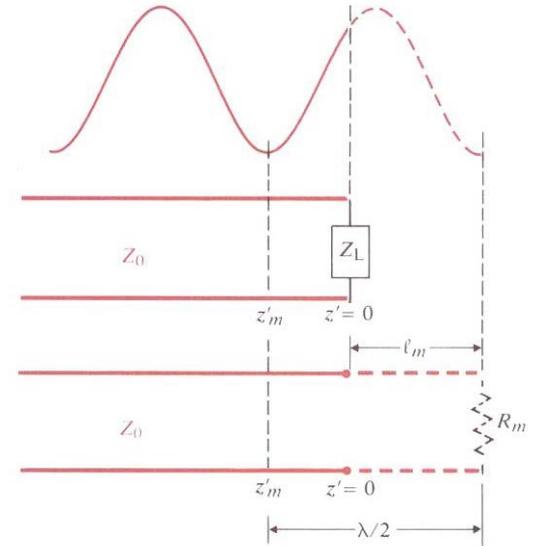
1. Find  $|\Gamma|$  from  $S$ :  $|\Gamma| = \frac{S-1}{S+1}$

2. Find  $\theta_\Gamma$  from  $z'_m$ :  $\theta_\Gamma - 2\beta z'_m = -\pi$

3. Find  $Z_L$ :  $Z_L = R_L + jX_L = R_0 \frac{1 + |\Gamma| e^{j\theta_\Gamma}}{1 - |\Gamma| e^{j\theta_\Gamma}}$

Recall:

$$\Gamma = \frac{Z_L - Z_0}{Z_L + Z_0} = |\Gamma| e^{j\theta_\Gamma}$$



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

Also recall:

$$V(z') = V_L \cos \beta z' + jI_L R_0 \sin \beta z'$$

$$= V_L \cos \beta z' + jV_L \frac{R_0}{R_L} \sin \beta z'$$

$$z'_m + l_m = \lambda / 2$$

$$\leftarrow R_L = R_m$$

$$S = R_0 / R_m$$

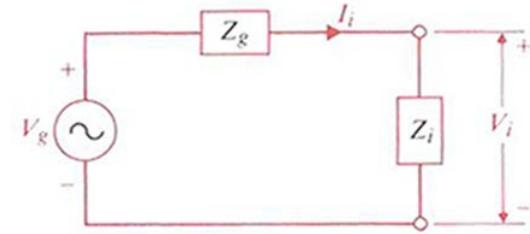
# Transmission-Line Circuits (1)

Voltage generator  $V_g$  with an internal impedance  $Z_g$ :

$$V_i = V_g - I_i Z_g$$

$$V_i = \frac{I_L}{2} (Z_L + Z_0) e^{\gamma l} [1 + \Gamma e^{-2\gamma l}]$$

$$I_i = \frac{I_L}{2Z_0} (Z_L + Z_0) e^{\gamma l} [1 - \Gamma e^{-2\gamma l}]$$



D. K. Cheng, Field and Wave Electromagnetics, 2nd ed., Addison-Wesley, 1989.

$$\rightarrow \frac{I_L}{2} (Z_L + Z_0) e^{\gamma l} [1 + \Gamma e^{-2\gamma l}] = V_g - \frac{I_L}{2Z_0} (Z_L + Z_0) e^{\gamma l} [1 - \Gamma e^{-2\gamma l}] Z_g$$

$$\rightarrow \frac{I_L}{2} (Z_L + Z_0) e^{\gamma l} = \frac{Z_0 V_g}{Z_0 + Z_g} \frac{1}{[1 - \Gamma_g \Gamma e^{-2\gamma l}]} \quad \leftarrow \Gamma_g = \frac{Z_g - Z_0}{Z_g + Z_0}$$

*Voltage reflection coefficient at the generator end*

$$\leftarrow z' = l - z$$

$$\rightarrow V(z') = \frac{I_L}{2} (Z_L + Z_0) e^{\gamma z'} [1 + \Gamma e^{-2\gamma z'}] = \frac{Z_0 V_g}{Z_0 + Z_g} e^{-\gamma z} \left( \frac{1 + \Gamma e^{-2\gamma z'}}{1 - \Gamma_g \Gamma e^{-2\gamma l}} \right)$$

$$\rightarrow I(z') = \frac{I_L}{2Z_0} (Z_L + Z_0) e^{\gamma z'} [1 - \Gamma e^{-2\gamma z'}] = \frac{V_g}{Z_0 + Z_g} e^{-\gamma z} \left( \frac{1 - \Gamma e^{-2\gamma z'}}{1 - \Gamma_g \Gamma e^{-2\gamma l}} \right)$$

# Transmission-Line Circuits (2)

*Recall:*  $V(z') = \frac{Z_0 V_g}{Z_0 + Z_g} e^{-\gamma z} \left( \frac{1 + \Gamma e^{-2\gamma l'}}{1 - \Gamma_g \Gamma e^{-2\gamma l}} \right)$

*What if:*  $\Gamma = 0$  or  $\Gamma_g = 0$ ?

$$= \frac{Z_0 V_g}{Z_0 + Z_g} e^{-\gamma z} (1 + \Gamma e^{-2\gamma l'}) (1 + \Gamma_g \Gamma e^{-2\gamma l} + \Gamma_g^2 \Gamma^2 e^{-4\gamma l} + \dots)$$

$$= \frac{Z_0 V_g}{Z_0 + Z_g} [e^{-\gamma z} + (\Gamma e^{-\gamma l'}) e^{-\gamma z'} + \Gamma_g (\Gamma e^{-2\gamma l}) e^{-\gamma z} + \dots]$$

$$= V_1^+ + V_1^- + V_2^+ + V_2^- + \dots$$

$$\leftarrow V_1^+ = \frac{V_g Z_0}{Z_0 + Z_g} e^{-\gamma z} = V_M e^{-\gamma z}$$

$$\leftarrow V_1^- = \Gamma (V_M e^{-\gamma l'}) e^{-\gamma z'}$$

$$\leftarrow V_2^+ = \Gamma_g (\Gamma V_M e^{-2\gamma l}) e^{-\gamma z}$$

⋮

