### Session 6-2 Equation of motion







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Apply Newton's 2nd law of motion

$$\sum \vec{F} = m\vec{a}$$

 $\Delta F_x = \Delta m a_x$ 

- External forces = surface force + body force
  - Surface force:
  - ~ normal force + tangential force
  - Body forces:
  - ~ due to gravitational or electromagnetic fields, no contact
  - ~ act at the centroid of the element  $\rightarrow$  centroidal force





(A)

Consider only gravitational force

$$\vec{g} = \vec{i}g_x + \vec{j}g_y + \vec{k}g_z$$

LHS of (A):



STRESS-STRAIN RELATIONS

 $\sigma_x + \frac{\partial \sigma_x}{\partial x} \Delta x$ 

#### Divide (B) by volume of element

$$\frac{\Delta F_x}{\Delta x \Delta y \Delta z} = \rho g_x + \frac{\partial \sigma_x}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z}$$
(C)

RHS of (A):

$$\frac{\Delta ma_x}{\Delta x \Delta y \Delta z} = \rho a_x$$

(D)





### Combine (C) and (D)

$$\rho g_{x} + \frac{\partial \sigma_{x}}{\partial x} + \frac{\partial \tau_{yx}}{\partial y} + \frac{\partial \tau_{zx}}{\partial z} = \rho a_{x}$$

$$\rho g_{y} + \frac{\partial \tau_{xy}}{\partial x} + \frac{\partial \sigma_{y}}{\partial y} + \frac{\partial \tau_{zy}}{\partial z} = \rho a_{y}$$

$$\rho g_{z} + \frac{\partial \tau_{xz}}{\partial x} + \frac{\partial \tau_{yz}}{\partial y} + \frac{\partial \sigma_{z}}{\partial z} = \rho a_{z}$$

(6.21)





#### 6.4.1 Navier-Stokes equations

- Eq (6.21) ~ general equation of motion containing 9 unknowns
- For Newtonian fluids (with single viscosity coeff.), use stress-strain relation given in (5.29) and (5.30) to reduce the number of unknowns
   → Navier-Stokes equations
- Eq. (5.29):

$$\sigma_x = \underline{-p} + 2\mu \frac{\partial u}{\partial x} - \left(\frac{2}{3}\right)\mu(\nabla \cdot \vec{q})$$

pressure normal stress due to fluid deformation and viscosity





$$\sigma_{y} = -p + 2\mu \frac{\partial v}{\partial y} - \left(\frac{2}{3}\right) \mu \left(\nabla \cdot \vec{q}\right)$$

$$\sigma_{z} = -p + 2\mu \frac{\partial w}{\partial z} - \left(\frac{2}{3}\right) \mu \left(\nabla \cdot \vec{q}\right)$$
(6.22)
(6.23)

Eq. (5.30):

$$\tau_{yx} = \tau_{xy} = \mu \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right)$$
$$\tau_{yz} = \tau_{zy} = \mu \left( \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \right)$$
$$\tau_{zx} = \tau_{xz} = \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right)$$





Substitute Eqs. (5.29) & (5.30) into (6.21)

$$\rho g_{x} - \frac{\partial p}{\partial x} + \frac{\partial}{\partial x} \left[ 2\mu \frac{\partial u}{\partial x} - \frac{2}{3}\mu \left( \nabla \cdot \vec{q} \right) \right] + \frac{\partial}{\partial y} \left[ \mu \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right) \right] + \frac{\partial}{\partial z} \left[ \mu \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right] = \rho a_{x}$$

Assume <u>constant viscosity</u> (neglect effect of pressure and temperature on viscosity variation)

$$\rho g_{x} - \frac{\partial p}{\partial x} + \mu \frac{\partial}{\partial x} \left[ 2 \frac{\partial u}{\partial x} - \frac{2}{3} (\nabla \cdot \vec{q}) \right] + \mu \frac{\partial}{\partial y} \left[ \left( \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} \right) \right] + \mu \frac{\partial}{\partial z} \left[ \left( \frac{\partial u}{\partial z} + \frac{\partial w}{\partial x} \right) \right] = \rho a_{x}$$
$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}$$





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normal stress + shear stress





$$\rho g_{x} - \frac{\partial p}{\partial x} + \mu \left[ \frac{\partial^{2} u}{\partial x^{2}} + \frac{\partial^{2} u}{\partial y^{2}} + \frac{\partial^{2} u}{\partial z^{2}} \right] + \frac{1}{3} \mu \frac{\partial}{\partial x} (\nabla \cdot \vec{q}) = \rho a_{x}$$

$$\rho g_{y} - \frac{\partial p}{\partial y} + \mu \left[ \frac{\partial^{2} v}{\partial x^{2}} + \frac{\partial^{2} v}{\partial y^{2}} + \frac{\partial^{2} v}{\partial z^{2}} \right] + \frac{1}{3} \mu \frac{\partial}{\partial y} (\nabla \cdot \vec{q}) = \rho a_{y}$$

$$\rho g_{z} - \frac{\partial p}{\partial z} + \mu \left[ \frac{\partial^{2} w}{\partial x^{2}} + \frac{\partial^{2} w}{\partial y^{2}} + \frac{\partial^{2} w}{\partial z^{2}} \right] + \frac{1}{3} \mu \frac{\partial}{\partial z} (\nabla \cdot \vec{q}) = \rho a_{z}$$
(6.24)

→ <u>Navier-Stokes equation for compressible fluids</u> with constant viscosity





Vector form

$$\rho \vec{g} - \nabla p + \mu \nabla^2 \vec{q} + \frac{\mu}{3} \nabla (\nabla \cdot \vec{q}) = \rho \frac{\partial \vec{q}}{\partial t} + \rho (\vec{q} \cdot \nabla) \vec{q}$$
  
where  $\vec{a} = \frac{d \vec{q}}{dt} = \frac{\partial \vec{q}}{\partial t} + (\vec{q} \cdot \nabla) \vec{q}$  --- Eq. (2.5)

1) For inviscid (ideal) fluid flow,  $(\mu = 0) \rightarrow$  viscous forces are neglected.

$$\rho \vec{g} - \nabla p = \rho \frac{\partial \vec{q}}{\partial t} + \rho (\vec{q} \cdot \nabla) \vec{q}$$

 $\rightarrow$  Euler equations for ideal fluid





2) For incompressible fluids,  $\nabla \cdot \vec{q} = 0$  (Continuity Eq.)

$$\rho \vec{g} - \nabla p + \mu \nabla^2 \vec{q} = \rho \frac{\partial \vec{q}}{\partial t} + \rho (\vec{q} \cdot \nabla) \vec{q}$$

(6.25)

#### Define acceleration due to gravity as

$$g_{x} = -g \frac{\partial h}{\partial x}$$

$$g_{y} = -g \frac{\partial h}{\partial y}$$

$$g_{z} = -g \frac{\partial h}{\partial z}$$

$$\vec{g} = -g \nabla h$$





where h = vertical direction measured positive upward For Cartesian axes oriented so that h and z coincide

$$g_x = g_y = 0$$
,  $\frac{\partial h}{\partial z} = 1$  (6.26)  
 $g_z = -g$  (6.27)

- $\rightarrow$  minus sign indicates that acceleration due to gravity is in the negative *h* direction
- Then, <u>N-S equation for incompressible fluids</u> and isothermal flows are





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- Eq. (6.28): <u>4 unknowns</u> *u, v, w, p*
- → We need one more equation to obtain a solution when the boundary conditions are specified.
- $\rightarrow$  Eq. of continuity for incompressible fluid

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0$$

Boundary conditions

1) kinematic BC: velocity normal to any rigid boundary (wall) equal the

boundary velocity (velocity = 0 for stationary boundary)

2) physical BC: <u>no slip condition</u> (continuum stick to a rigid boundary)

 $\rightarrow$  tangential velocity <u>relative to the wall</u> vanish at the wall surface



- General solutions for Navier-Stocks equations are <u>not available</u> because of the <u>nonlinear, 2nd-order nature</u> of the partial differential equations.
- $\rightarrow$  Only particular solutions may be obtained by simplifications.
- → Numerical solutions are usually sought.
- Navier-Stocks equations in cylindrical coordinates for constant density and viscosity
  - *r*-component:

$$\rho \left( \frac{\partial v_r}{\partial t} + v_r \frac{\partial v_r}{\partial r} + \frac{v_{\theta}}{r} \frac{\partial v_r}{\partial \theta} - \frac{v_{\theta}^2}{r} + v_z \frac{\partial v_r}{\partial z} \right) \\
= \rho g_r - \frac{\partial p}{\partial r} + \mu \left[ \frac{\partial}{\partial r} \left\{ \frac{1}{r} \frac{\partial}{\partial r} (rv_r) \right\} + \frac{1}{r^2} \frac{\partial^2 v_r}{\partial \theta^2} - \frac{2}{r^2} \frac{\partial v_{\theta}}{\partial \theta} + \frac{\partial^2 v_r}{\partial z^2} \right]$$



#### $\theta$ - component:

$$\begin{split} \rho \bigg( \frac{\partial v_{\theta}}{\partial t} + v_r \frac{\partial v_{\theta}}{\partial r} + \frac{v_{\theta}}{r} \frac{\partial v_{\theta}}{\partial \theta} - \frac{v_r v_{\theta}}{r} + v_z \frac{\partial v_{\theta}}{\partial z} \bigg) \\ = \rho g_{\theta} - \frac{1}{r} \frac{\partial p}{\partial \theta} + \mu \bigg[ \frac{\partial}{\partial r} \bigg\{ \frac{1}{r} \frac{\partial}{\partial r} (r v_{\theta}) \bigg\} + \frac{1}{r^2} \frac{\partial^2 v_{\theta}}{\partial \theta^2} + \frac{2}{r^2} \frac{\partial v_r}{\partial \theta} + \frac{\partial^2 v_{\theta}}{\partial z^2} \bigg] \end{split}$$

*z* - component:

$$\rho \left( \frac{\partial v_z}{\partial t} + v_r \frac{\partial v_z}{\partial r} + \frac{v_\theta}{r} \frac{\partial v_z}{\partial \theta} + v_z \frac{\partial v_z}{\partial z} \right) \\
= \rho g_z - \frac{\partial p}{\partial z} + \mu \left[ \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial v_z}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 v_z}{\partial \theta^2} + \frac{\partial^2 v_z}{\partial z^2} \right]$$
(6.29)



Continuity eq. for incompressible fluid

$$\frac{1}{r}\frac{\partial}{\partial r}(rv_r) + \frac{1}{r}\frac{\partial}{\partial \theta}(v_\theta) + \frac{\partial}{\partial z}(v_z) = 0$$
(6.30)

Normal & shear stresses for constant density and viscosity

$$\sigma_{r} = -p + 2\mu \frac{\partial v_{r}}{\partial r}$$

$$\sigma_{\theta} = -p + 2\mu \left(\frac{1}{r} \frac{\partial v_{\theta}}{\partial \theta} + \frac{v_{r}}{r}\right)$$

$$\sigma_{z} = -p + 2\mu \frac{\partial v_{z}}{\partial z}$$





$$\tau_{r\theta} = \mu \left[ r \frac{\partial}{\partial r} \left( \frac{v_{\theta}}{r} \right) + \frac{1}{r} \frac{\partial v_{r}}{\partial \theta} \right]$$

$$\tau_{\theta z} = \mu \left[ \frac{\partial v_{\theta}}{\partial z} + \frac{1}{r} \frac{\partial v_{z}}{\partial \theta} \right]$$

$$\tau_{zr} = \mu \left[ \frac{\partial v_r}{\partial z} + \frac{\partial v_z}{\partial r} \right]$$



