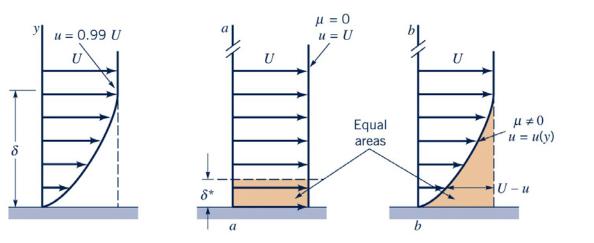


Turbulent Boundary-Layer Flows







Contents

- 9.1 Introduction
- 9.2 Structure of a turbulent boundary layer
- 9.3 Mean-flow characteristics for turbulent boundary layer

Objectives

- Study wall turbulence
- Derive equations of velocity distribution and friction coefficient for both smooth and rough walls





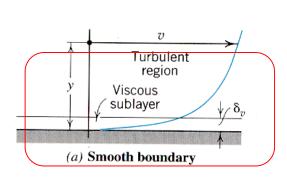
- Relations describing the mean-flow characteristics in boundary layer
- It is useful to predict <u>velocity magnitude</u> and <u>relation between velocity and</u> <u>wall shear</u> or pressure gradient forces.
- It is desirable that these relations should <u>not require knowledge of the</u> <u>turbulence details</u>.
- For laminar boundary layer, it is possible to obtain a solution by integrating the equations of motion (<u>Prandtl's 2D boundary layers equation</u>, Eq. (10.7))
- However, the turbulent boundary layer is composed of zones of different types of flow, and effective viscosity varies from wall out through the layer.
- Thus, theoretical solution is not practical for the general non-uniform boundary layer \rightarrow use semiempirical procedure

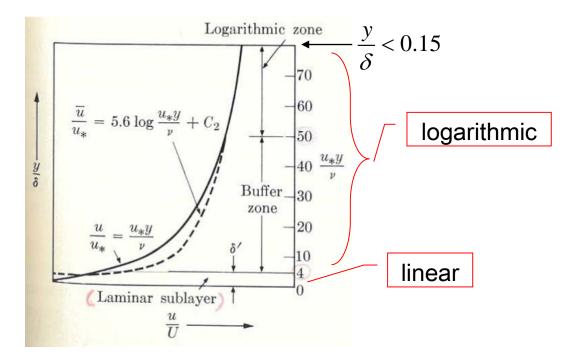




9.3.1 Universal velocity and friction laws: smooth walls

I. Inner law (wall law)

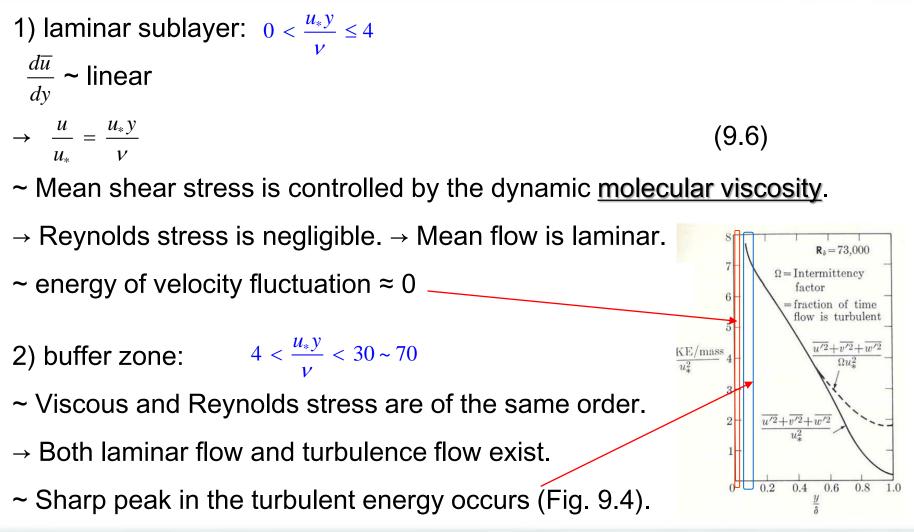




Velocity profile on smooth walls











- 3) turbulent zone inner region: $\frac{u_*y}{v} > 30 \sim 70$, and $\frac{y}{\delta} < 0.15$
- ~ fully turbulent flow
- ~ inner law zone/wall law
- ~ Intensity of turbulence decreases.
- ~ velocity equation: logarithmic function

 $\frac{\overline{u}}{u_*} = 5.6 \log \frac{u_* y}{v} + C_2$

Outer law (velocity defect law)

$$\frac{\delta}{\frac{U-\overline{u}}{u_*}} = -8.6 \log\left(\frac{y}{\delta}\right)$$
$$\frac{\overline{u}}{u_*} = \frac{U}{u_*} + 8.6 \log\left(\frac{y}{\delta}\right)$$

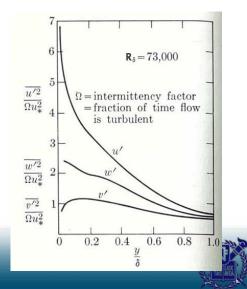
 $\frac{y}{-} > 0.15$

4) turbulent zone-<u>outer region</u>: $0.15\delta < y < 0.4\delta$

~ outer law, velocity-defect law

5) intermittent zone: $0.4\delta < y < 1.2\delta$

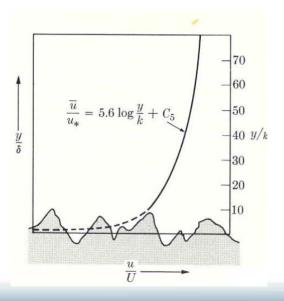
~ Flow is intermittently turbulent and non-turbulent.



- 6) non-turbulent zone
- ~ external flow zone
- ~ potential flow

[Cf] Rough wall:

 \rightarrow Laminar sublayer is destroyed by the <u>roughness elements.</u>







 $f(\overline{u}, u, \overline{v}) = 0$

- Law of wall (Inner law) $y < 0.15\delta$
- close to smooth boundaries (molecular viscosity dominant)
- Law of wall assumes that the relation between <u>wall shear stress and</u> <u>velocity at distance y</u> from the wall depends only on fluid density and viscosity
- Dimensional analysis yields

$$f(u, u_*, y, \rho, \mu) = 0$$

$$\frac{\overline{u}}{u_*} = f\left(\frac{u_* y}{v}\right)$$
(9.3)





i) Laminar sublayer

- mean velocity, $\overline{u} \equiv u$

- velocity gradient,
$$\frac{\partial u}{\partial y} \sim \text{constant} \equiv \frac{u}{y}$$

- shear stress,
$$\tau \approx \tau_0 = \mu \frac{\partial u}{\partial y} \bigg|_{y=0} \equiv \mu \frac{u}{y}$$
 (9.5)

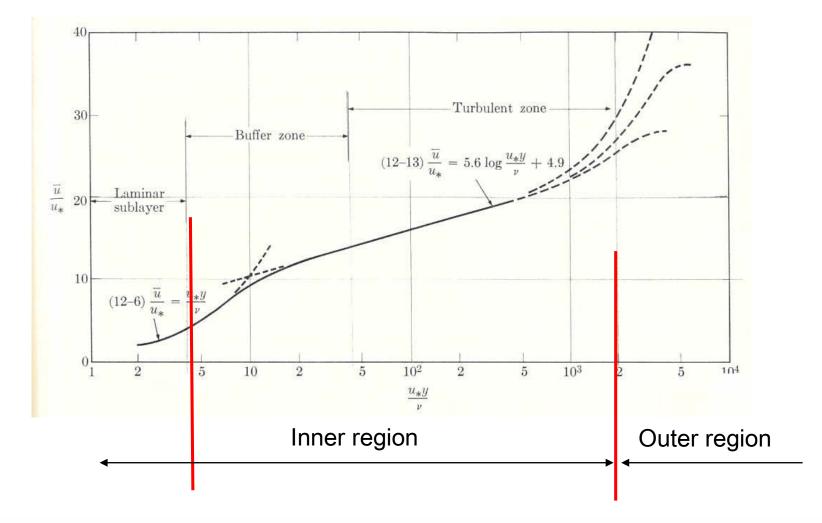
V

- shear velocity,
$$u_* = \sqrt{\frac{\tau}{\rho}} \rightarrow u_*^2 = \frac{\tau}{\rho} = \frac{\mu}{\rho} \frac{u}{y}$$

$$\frac{u}{u_*} = \frac{u_* y}{v} \tag{9.6}$$











Thickness of laminar sublayer

 $\frac{u_*y}{-} = 4$

- define thickness of laminar sublayer (δ') as the value of y which makes

$$\delta' = \frac{4\nu}{u_*} = \frac{4\nu}{\sqrt{\frac{\tau_0}{\rho}}} = \frac{4\nu}{\left(\frac{c_f \rho U^2 / 2}{\rho}\right)^{\frac{1}{2}}} = \frac{4\nu}{U\sqrt{c_f / 2}}$$
(9.7)

Where $\tau_0 = c_f \rho \frac{u^2}{2}$; $c_f = local \ shear \ stress \ coeff$.

- c_f decreases slowly with increasing Reynolds number, $\text{Re}_x = \frac{Ux}{v}$ Thus, δ' increases with distance along the surface.

Re] For the laminar sublayer in pipe flows

$$\delta' = 11.6 \frac{v}{u_*}$$



ii) Turbulent region

- Start with equation of 2D turbulent boundary layer, Eq. (8.8a)

$$\rho\left(\frac{\partial \overline{u}}{\partial t} + \overline{u}\frac{\partial \overline{u}}{\partial x} + \overline{v}\frac{\partial \overline{u}}{\partial y}\right) = -\frac{\partial \overline{p}}{\partial x} - \rho\frac{\partial \overline{u'^2}}{\partial x} + \mu\frac{\partial^2 \overline{u}}{\partial y^2} - \rho\frac{\partial \overline{u'v'}}{\partial y}$$

- In the turbulent zone of near wall region, the mean shear stress remains nearly equal to the wall shear; $\tau \approx \tau_0$

$$\rho\left(\frac{\partial \overline{u}}{\partial t} + \overline{u}\frac{\partial \overline{u}}{\partial x} + \overline{v}\frac{\partial \overline{u}}{\partial y}\right) + \frac{\partial \overline{p}}{\partial x} + \rho\frac{\partial \overline{u'^2}}{\partial x} = \frac{\partial}{\partial y}\left(\mu\frac{\partial \overline{u}}{\partial y} - \rho\overline{u'v'}\right) = \frac{\partial\tau}{\partial y} \approx \frac{\partial\tau_0}{\partial y} = 0$$

- Hence a solution is obtained if we have a relation for au_0

- We can use Prandtl's mixing length theory for near wall, $l = \kappa y$

$$\tau_0 \approx \tau = \rho l^2 \left| \frac{d\overline{u}}{dy} \right| \frac{d\overline{u}}{dy} = \rho \kappa^2 y^2 \left(\frac{d\overline{u}}{dy} \right)^2$$
(1)





Integrate (1)

$$\overline{u} = \frac{u_*}{\kappa} \ln y + C_1 \qquad u_* = \sqrt{\frac{\tau_0}{\rho}}$$

$$\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln y + C_1 \qquad (9.10)$$

Substitute BC [$\overline{u} = 0$ at y = y'] into (9.10)

$$0 = \frac{1}{\kappa} \ln y' + C_1$$

$$\therefore C_1 = -\frac{1}{\kappa} \ln y'$$
(5)

Assume $y' \propto \frac{v(m^2/s)}{u_*(m/s)} \rightarrow y' = C \frac{v}{u_*}$

Then (5) becomes

$$C_1 = -\frac{1}{\kappa} \ln y' = -\frac{1}{\kappa} \ln \left(C \frac{\nu}{u_*} \right) = C_2 - \frac{1}{\kappa} \ln \frac{\nu}{u_*}$$





(9.11)

Substitute (9.11) into (9.10)

$$\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln y + C_2 - \frac{1}{\kappa} \ln \frac{v}{u_*} = \frac{1}{\kappa} \ln \left(\frac{u_* y}{v}\right) + C_2$$

Changing the logarithm to base 10 yields

$$\frac{\overline{u}}{u_*} = \frac{2.3}{\kappa} \log_{10} \left(\frac{u_* y}{v} \right) + C_2$$
(9.12)

Empirical values of κ and C_2 for inner region of the boundary layer

$$\kappa = 0.41; C_2 = 4.9$$

$$\frac{\overline{u}}{u_*} = 5.6 \, \log\left(\frac{u_*y}{v}\right) + 4.9, \ 30 \sim 70 < \frac{u_*y}{v}, \ and \ \frac{y}{\delta} < 0.15$$
(9.13)

→ Prandtl's velocity distribution law; inner law; wall law





[Re] Prandtl's turbulent boundary layer equation

X-eq.
$$\rho\left(\frac{\partial \overline{u}}{\partial t} + \overline{u}\frac{\partial \overline{u}}{\partial x} + \overline{v}\frac{\partial \overline{u}}{\partial y}\right) = -\frac{\partial \overline{p}}{\partial x} - \rho\frac{\partial \overline{u'^2}}{\partial x} + \mu\frac{\partial^2 \overline{u}}{\partial y^2} - \rho\frac{\partial \overline{u'v'}}{\partial y}$$
 (8.8a)

y-eq.
$$0 = -\frac{\partial}{\partial y} (\overline{p} + \rho \overline{v'^2})$$
(8.8b)

 $\frac{\partial \overline{u}}{\partial x} + \frac{\partial \overline{v}}{\partial y} = 0$ Continuity eq.: (8.8c)

Solution to these equations: $\overline{u}(x, y, t), \overline{v}(x, y, t), \overline{p}(x, y, t)$

$$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_*y}{v}\right) + 4.9 \quad \rightarrow \quad \overline{u} = \quad \overline{u}(y)$$





II. Outer law; Velocity-defect law

- In the outer reaches of the turbulent boundary layer <u>for both smooth</u> <u>and rough walls</u>, Reynolds stresses dominate the viscous stresses to produce the velocity profile.
- It was observed that the <u>velocity defect (reduction) at y-values</u> was dependent on the magnitude of the <u>wall shear stress</u>

$$\frac{U-\overline{u}}{u_*} = g\left(\frac{y}{\delta}\right) \tag{9.14}$$

- A logarithmic relation for the function g can be obtained by assuming that Eq. (9.12) will give $\overline{u} = U$ at $y = \delta$

$$\frac{U}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_2^{-1}$$
(2)





Subtract (9.12) from (2)

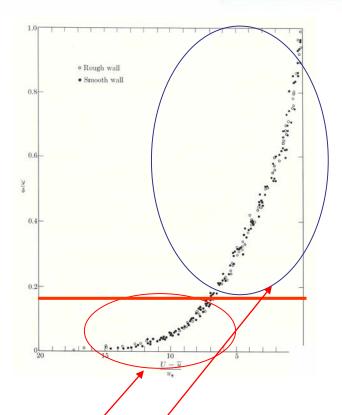
$$\frac{U}{u_{*}} = \frac{2.3}{\kappa} \log\left(\frac{u_{*}\delta}{v}\right) + C_{2}'$$

$$- \left| \frac{\overline{u}}{u_{*}} = \frac{2.3}{\kappa} \log\left(\frac{u_{*}y}{v}\right) + C_{2}\right|$$

$$\frac{U - \overline{u}}{u_{*}} = \frac{2.3}{\kappa} \left\{ \log\left(\frac{u_{*}\delta}{v}\right) - \log\left(\frac{u_{*}y}{v}\right) \right\} + C_{2}' - C_{2}$$

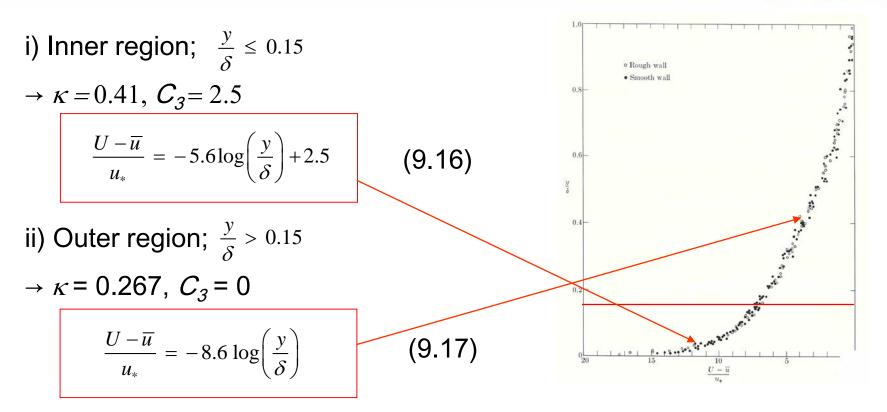
$$= \frac{2.3}{\kappa} \log\left(\frac{u_{*}\delta}{v}\frac{v}{u_{*}y}\right) + C_{3} = -\frac{2.3}{\kappa} \log\left(\frac{y}{\delta}\right) + C_{3}$$

$$\frac{U - \overline{u}}{u_{*}} = -\frac{2.3}{\kappa} \log\left(\frac{y}{\delta}\right) + C_{3}$$
(9.15)



- A single log-equation does not fit the data over the entire boundary layer.
- Instead <u>one equation will fit an inner region overlapping with Eq. (9.12)</u>, while <u>second equation will approximate the outer region</u>.





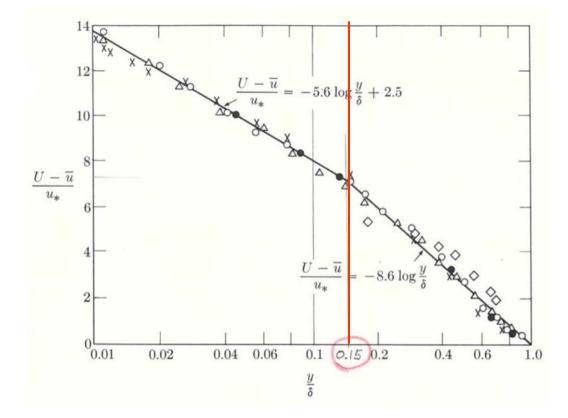
 \rightarrow Eqs. (9.16) & (9.17) apply to both smooth and rough surfaces.

→ Eq. (9.16) = Eq. (9.13)
$$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_*y}{v}\right) + 4.9, \ 30 \sim 70 < \frac{u_*y}{v}, \ and \ \frac{y}{\delta} < 0.15$$





- The velocity-defect law is applicable for both smooth and rough walls







Velocity profile equations

	Wall law	Velocity-defect law	
	Smooth wall	Smooth wall	Rough wall
Outer region $\frac{y}{\delta} > 0.15$	_	$\frac{U - \overline{u}}{u_*} = -8.6 \log\left(\frac{y}{\delta}\right)$	
Inner region $30 \sim 70 < \frac{u_* y}{v}$ $\frac{y}{\delta} < 0.15$	$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_* y}{v}\right) + 4.9$	$\frac{U - \overline{u}}{u_*} = -5.6 \log\left(\frac{y}{\delta}\right) + 2.5$	
Laminar sublayer	$\frac{u}{u_*} = \frac{u_* y}{v}$	_	





III. Surface-resistance formulas

(1) Local shear-stress coefficient on smooth walls

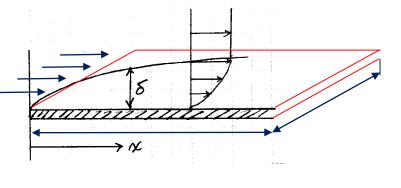
velocity profile ↔ shear-stress equations

$$u_* = \sqrt{\tau / \rho} = \sqrt{\frac{c_f}{2}} U \rightarrow \frac{U}{u_*} = \sqrt{\frac{2}{c_f}}$$

where $c_f = local shear stress coefficient$

$$[\mathsf{Re}] \quad \tau_0 = \frac{1}{2}\rho c_f U^2 \rightarrow \frac{\tau_0}{\rho} = \frac{c_f}{2} U^2 \qquad D_f = \tau_0 A = \frac{1}{2}A\rho c_f U^2 D_f = \tau_0 A = \frac{1}{2}A\rho c_f U^2 D = D_f + D_p (9.18)$$







(i) Assume that logarithmic law will give the relation

Substituting $\overline{u} = U at y = \delta$ into Eq. (9.12) yields

$$\frac{U}{u_*} = \frac{2.3}{\kappa} \log\left(\frac{u_*\delta}{v}\right) + C_4 \tag{A}$$

Substitute (9.18) into (A)

$$\sqrt{\frac{2}{c_f}} = \frac{2.3}{\kappa} \log\left(\frac{U\delta}{\nu}\sqrt{\frac{c_f}{2}}\right) + C_4 = \frac{2.3}{\kappa} \log\left(\operatorname{Re}_{\delta}\sqrt{\frac{c_f}{2}}\right) + C_4$$
(9.19)

- c_f is a function of Reynolds number for smooth walls.
- c_f is not given explicitly.





(ii) For explicit expression, use displacement thickness and momentum thickness θ instead of δ

Clauser:
$$\frac{1}{\sqrt{c_f}} = 3.96 \log \operatorname{Re}_{\delta^*} + 3.04$$
 (9.20)

Squire and Young:
$$\frac{1}{\sqrt{c_f}} = 4.17 \log \text{Re}_{\theta} + 2.54$$
 (9.21)

where
$$\operatorname{Re}_{\delta^*} = \frac{U\delta^*}{v}$$
; $\operatorname{Re}_{\theta} = \frac{U\theta}{v}$

$$\operatorname{Re}_{\delta^*}, \operatorname{Re}_{\theta} = f(\operatorname{Re}_x)$$
 (9.22)





iii) Karman's relation

~ assume <u>turbulence boundary layer</u> all the way from the leading edge (i.e., no preceding stretch of laminar boundary layer)

(9.23)

$$\frac{1}{\sqrt{c_f}} = 4.15\log(\text{Re}_x c_f) + 1.7$$

 Karman's equation is useful for ships and aircraft wings where the laminar boundary layer is insignificant. Vo Vo Vo Vo Vo Vo Vo Vo Viscous sublayer Laminar boundary layer Viscous sublayer



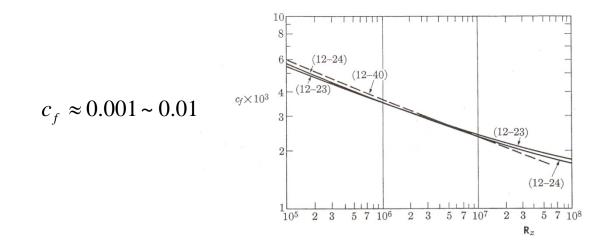


iv) Explicit equation by Schultz-Grunow (1940)

$$c_f = \frac{0.370}{(\log \text{Re}_x)^{2.58}}$$

(9.24)

Comparison of (9.23) and (9.24)





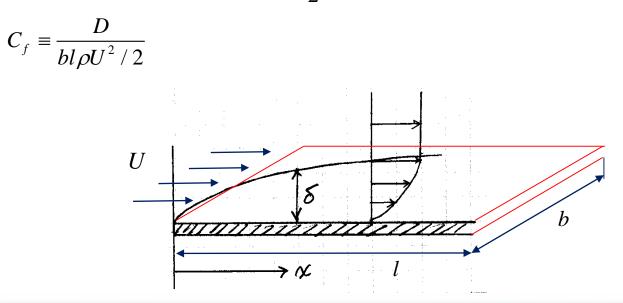




(2) Average shear-stress coefficient on smooth walls

Consider average shear-stress coefficient over a distance / along a flat plate of a width *b*

Total frictional drag $(D_f) = \tau \times bl = \frac{1}{2}C_f \rho U^2 bl$







i) Schoenherr (1932)

$$\frac{1}{\sqrt{C_f}} = 4.13\log(\operatorname{Re}_l C_f)$$

- (9.26)
- assume turbulence boundary layer all the way from the leading edge
- Similar to von Karman's equation

where $\operatorname{Re}_{l} = \frac{Ul}{V}$

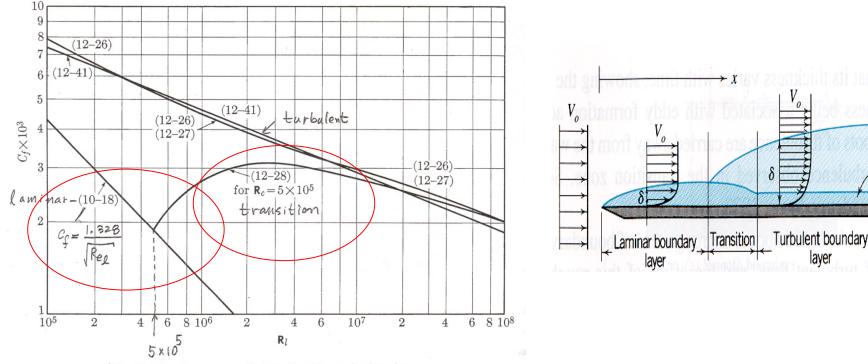
ii) Schultz-Grunow

$$C_{f} = \frac{0.427}{(\log \operatorname{Re}_{l} - 0.407)^{2.64}} \quad , \ 10^{6} < \operatorname{Re}_{l} < 10^{9}$$
 (9.27)





Comparison of (9.26) and (9.27)









Viscous

sublayer

- 3) Transition formula
- Boundary layer developing on a smooth flat plate
- ~ At upstream end (leading edge), laminar boundary layer develops because viscous effects dominate due to inhibiting effect of the smooth wall on the development of turbulence.
- ~ Thus, when there is a significant stretch of laminar boundary layer preceding the <u>turbulent layer</u>, total friction is the laminar portion up to x_{crit} plus the turbulent portion from x_{crit} to *l*.
- ~ Therefore, <u>average shear-stress coefficient is lower than the prediction by Eqs.</u> (9.26) or (9.27).
- → Use transition formula

$$C_{f} = \frac{0.427}{(\log \operatorname{Re}_{l} - 0.407)^{2.64}} - \frac{A}{\operatorname{Re}_{l}}$$





(9.28)

where A/Re_l = correction term = $f(\text{Re}_{crit})$, $\text{Re}_{crit} = \frac{Ux_{crit}}{v}$ \rightarrow A = 1,060~3,340 (Table 9.2, p. 240)

- Since A is a function of Re_{crit} , the curve in Fig. 10 is given for $\text{Re}_{crit} = 5 \times 10^5$
- For flow along smooth walls

 $300,000 < \text{Re}_{crit} < 600,000$

- For laminar flow:

$$C_f = \frac{1.328}{\text{Re}_l^{1/2}}$$





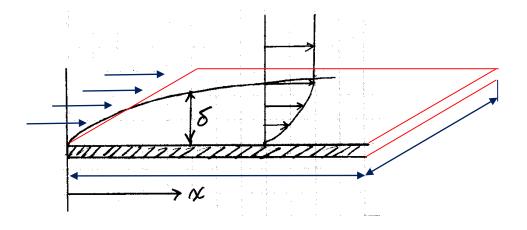
SURFACE RESISTANCE FORMULAS FOR BOUNDARY LAYERS WITH $d\overline{p}/dx = 0$ Smooth walls Rough walls LOCAL SHEAR Universal equations (12-46) $\frac{1}{\sqrt{c_{\ell}}} = 3.96 \log \frac{\delta}{k} + C_8$ $1/\sqrt{c_f} = 3.96 \log R_{\delta^*} + 3.04$ Clauser (12-20) $C_8 = f$ (size, shape, and distribution of roughness) $1/\sqrt{c_f} = 4.17 \log \mathbf{R}_{\theta} + 2.54$ Squire and Young (12–21) $1/\sqrt{c_f} = 4.15 \log (\mathbf{R}_x c_f) + 1.7$ von Kármán (12-23) $c_f = \frac{0.370}{(\log \mathbf{R}_r)^{2.58}}$ Schultz-Grunow (12-24) $c_f = \frac{0.0466}{\mathbf{R}_{\star}^{1/4}} =$ 0.059Power law (12-40)R1/5 AVERAGE SHEAR Universal equations $1/\sqrt{C_f} = 4.13 \log (\mathbf{R}_l C_f)$ Schoenherr (12-26) $C_f = \frac{0.427}{(\log \mathbf{R}_l - 0.407)^{2.64}}$ Schultz-Grunow (12-27) $C_f = \frac{0.074}{\mathbf{R}_f^{1/5}}$ Power law (12-41) Transition formula $C_f = \frac{0.427}{(\log \mathbf{R}_l - 0.407)^{2.64}} - \frac{A}{\mathbf{R}_l}$ Schultz-Grunow-Prandtl (12-28) $A = f(\mathbf{R}_{crit})$ as given in Table 12-2





[Ex. 9.1] Turbulent boundary-layer velocity and thickness An <u>aircraft</u> flies at 25,000 ft with a speed of 410 mph (600 ft/s). Compute the following items for the boundary layer at a distance 10 ft from

the leading edge of the wing of the craft.

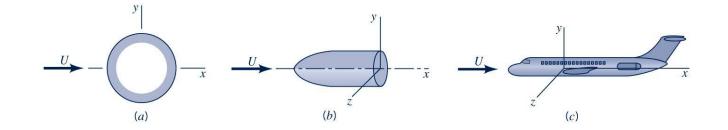






- ① Frictional drag = surface resistance = skin drag
- ② Pressure drag = form drag
- ~ depends largely on shape or form of the body

For airfoil, hydrofoil, and slim ships: surface drag > form drag For bluff objects like spheres, bridge piers: surface drag < form drag





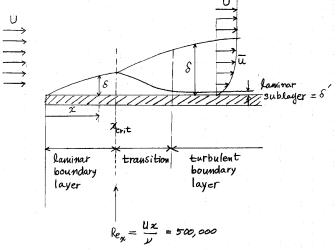


(a) Thickness δ' (laminar sublayer; $\delta' = \frac{4\nu}{u_*}$) at x = 10 ftAir at El. 25,000 ft: $\nu = 3 \times 10^{-4} ft^2 / s$

$$\rho = 1.07 \times 10^{-3} \ slug \ / \ ft^3$$

Find
$$\operatorname{Re}_{crit} = 5 \times 10^5 = \frac{600(x)}{3 \times 10^{-4}}$$

$$\therefore x_{crit} = 0.25 ft \sim \text{negligible comparaed to } l = 10 ft$$



Therefore, assume that <u>turbulent boundary layer</u> develops <u>all the way</u> <u>from the leading edge</u>.

$$\operatorname{Re}_{x} = \frac{Ux}{v} = \frac{600(10)}{(3 \times 10^{-4})} = 2 \times 10^{7}$$





Use Schultz-Grunow eq., (9.24) to compute c_f

$$c_{f} = \frac{0.370}{(\log \operatorname{Re}_{x})^{2.58}} = \frac{0.370}{\left\{\log\left(2 \times 10^{7}\right)\right\}^{2.58}} = \frac{0.370}{(7.30)^{2.58}} = 0.0022$$

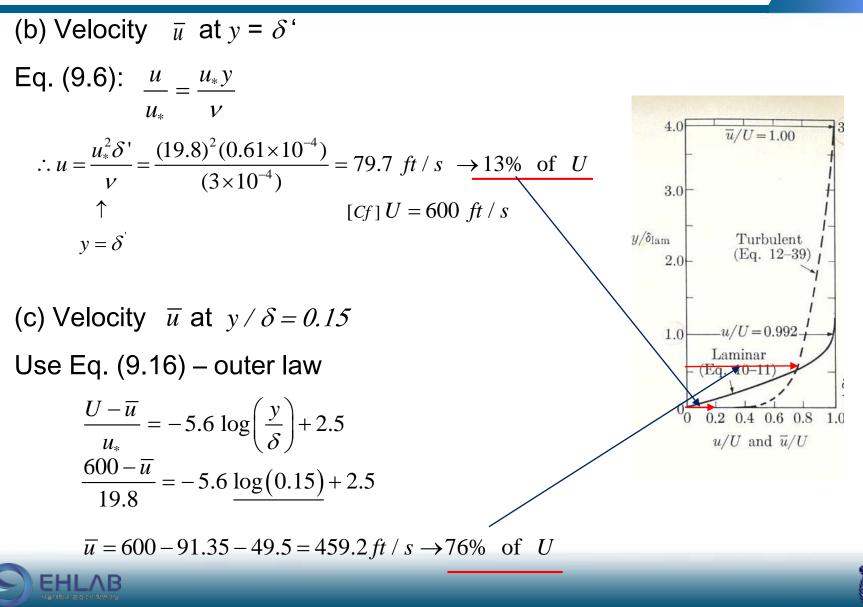
$$\tau_{0} = \frac{\rho}{2} c_{f} U^{2} = \frac{1}{2} (1.07 \times 10^{-3}) (0.0022) (600)^{2} = 0.422 lb / ft^{2}$$

$$u_{*} = \sqrt{\frac{\tau_{0}}{\rho}} = \sqrt{\frac{0.422}{1.07 \times 10^{-3}}} = 19.8 \, ft / s$$

$$\delta' = \frac{4\nu}{u_{*}} = \frac{4(3 \times 10^{-4})}{19.8} = 0.61 \times 10^{-4} \, ft = 7.3 \times 10^{-4} \, in$$







$$\begin{bmatrix} Cf \end{bmatrix} \quad \overline{u} = U + 5.6 \ u_* \ \log\left(\frac{y}{\delta}\right) - 2.5u_*$$

(d) Distance y at $y / \delta = 0.15$ and thickness δ

Use Eq. (9.13) – inner law

$$\frac{\overline{u}}{u_*} = 5.6 \log\left(\frac{u_* y}{v}\right) + 4.9$$

$$At \frac{y}{\delta} = 0.15: \frac{459}{19.8} = 5.6 \log\left(\frac{19.8y}{3 \times 10^{-4}}\right) + 4.9$$
$$\log\left(\frac{19.8y}{3 \times 10^{-4}}\right) = 3.26; \quad \frac{19.8y}{3 \times 10^{-4}} = 1839$$
$$y = 0.028' = 0.33in \approx 0.8cm$$

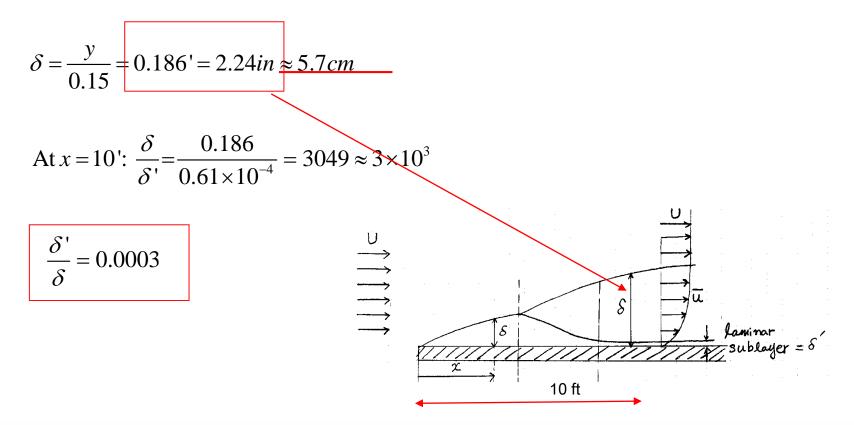




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(B)

Substitute (B) into







[Ex. 9.2] Surface resistance on a smooth boundary given as Ex. 9.1 (a) Displacement thickness δ^*

$$\delta^{*} = \int_{0}^{h} \left(1 - \frac{u}{U} \right) dy \qquad (7.9)$$

$$\frac{\delta^{*}}{\delta} = \int_{0}^{h/\delta} \left(1 - \frac{\overline{u}}{U} \right) d\left(\frac{y}{\delta} \right), \quad h/\delta \ge 1 \qquad (A)$$

Neglect laminar sublayer and approximate buffer zone with Eq. (9.16)

(i)
$$y/\delta < 0.15$$
, $\frac{U-\overline{u}}{u_*} = -5.6\log\left(\frac{y}{\delta}\right) + 2.5 \leftarrow (9.16)$
Divide (9/16) by U
 $\therefore 1 - \frac{\overline{u}}{U} = -5.6\frac{u_*}{U}\log\frac{y}{\delta} + 2.5\frac{u_*}{U} = \left(-2.43\ln\frac{y}{\delta} + 2.5\right)\frac{u_*}{U}$ (B)





(*ii*)
$$y/\delta > 0.15$$
, $\frac{U-\overline{u}}{u_*} = -8.6\log\left(\frac{y}{\delta}\right) \quad \leftarrow (9.17)$
$$\therefore 1 - \frac{\overline{u}}{U} = -8.6\frac{u_*}{U}\log\left(\frac{y}{\delta}\right) = -3.74\frac{u_*}{U}\ln\left(\frac{y}{\delta}\right)$$

Substituting (B) and (C) into (A) yields

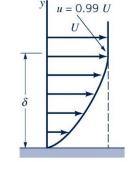


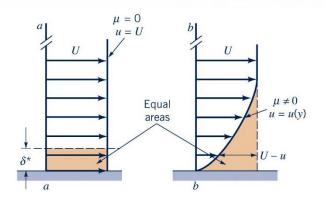


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(C)

$$\delta^* = 0.1184\delta = 0.1184 \ (0.186) = 0.022 \ ft$$
$$\frac{\delta^*}{\delta} = 0.1184 \ \to 11.8\%$$





(b) Local surface-resistance coeff. c_f Use Eq. (9.20) by Clauser

$$\frac{1}{\sqrt{c_f}} = 3.96 \log \operatorname{Re}_{\delta^*} + 3.04 \quad \leftarrow \quad \operatorname{Re}_{\delta^*} = \frac{600 \times 0.022}{3 \times 10^{-4}} = 44,000$$
$$= 3.96 \log (44,000) + 3.04$$

$$\therefore c_f = 2.18 \times 10^{-3} = 0.00218$$

[Cf] $c_f = 0.0022$ by Schultz-Grunow Eq.



