

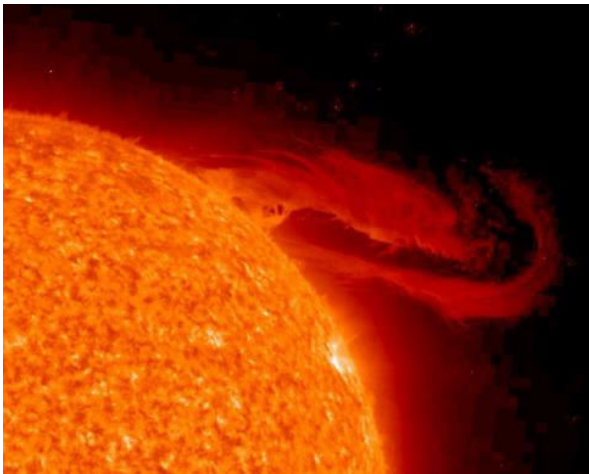
Introduction to Nuclear Fusion

Prof. Dr. Yong-Su Na

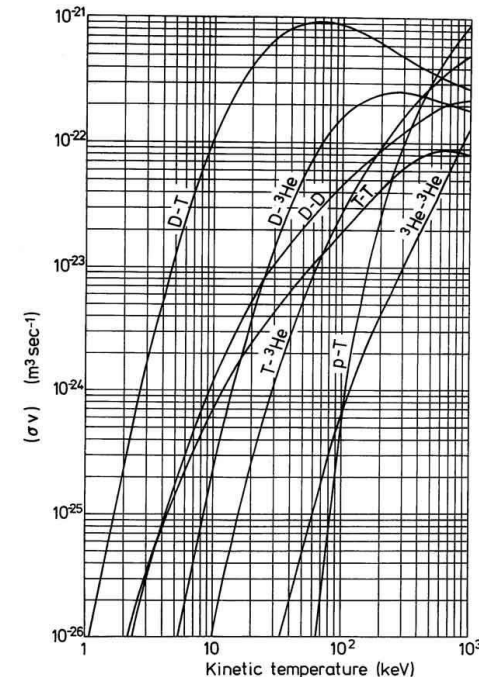
How to describe a plasma?

Description of a Plasma

- Thermonuclear fusion with high energy particles of 10-20 keV energy → plasma



- Three approaches to describe a plasma
 - Single particle approach
 - Kinetic theory
 - Fluid theory

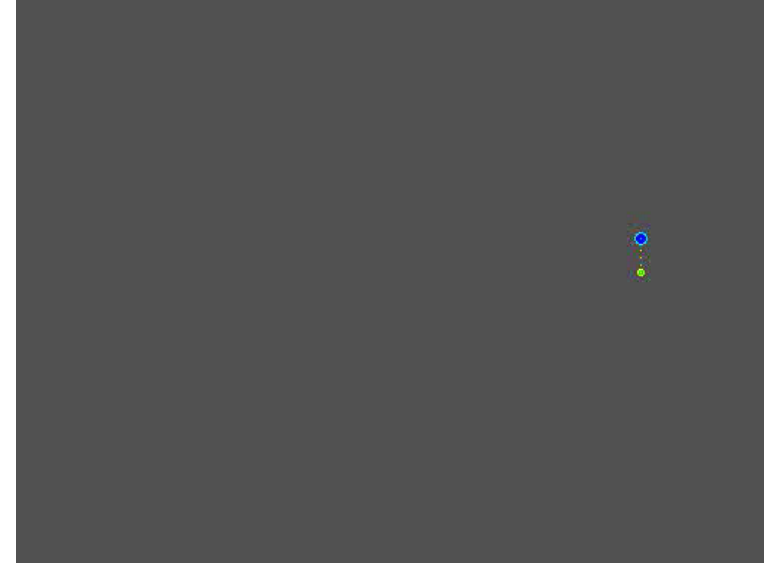
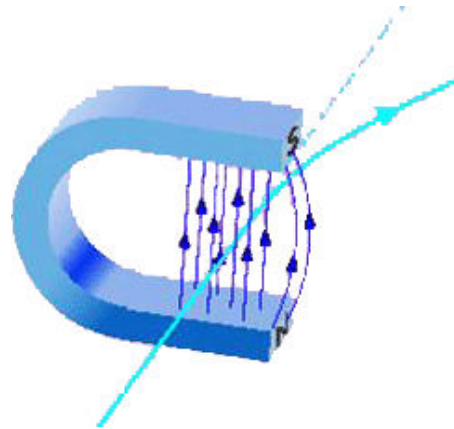


Individual Charge Trajectories

- **Equation of motion**

- Basic relation determining the motion of an individual charged particle of mass m and charge q in a combined electric ($\mathbf{E}$) and magnetic ($\mathbf{B}$) field
- Neglecting electromagnetic fields generated by movement of the charge itself

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B})$$



Individual Charge Trajectories

- **Homogeneous electric field**

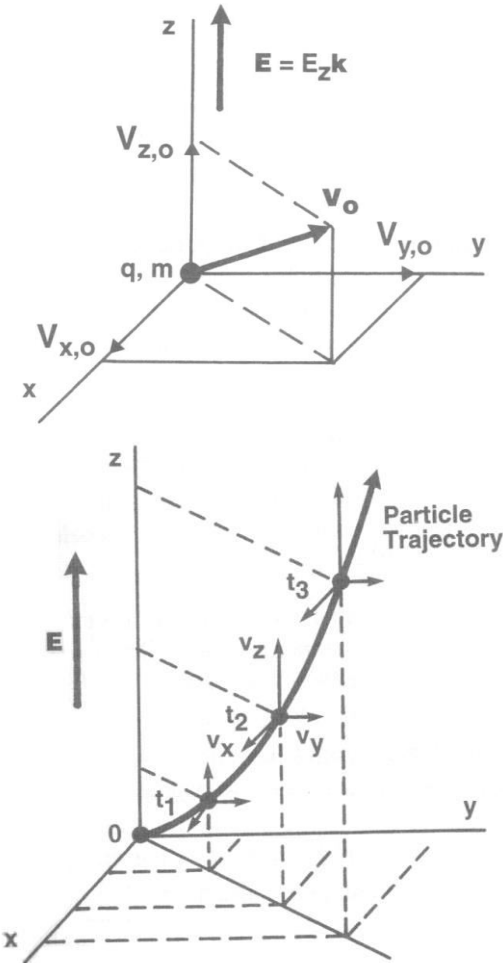
$$m \frac{d\mathbf{v}}{dt} = q\mathbf{E} = qE_z \mathbf{k}$$

$$v_x(0) = v_{x,0}, \quad v_y(0) = v_{y,0}, \quad v_z(0) = v_{z,0}$$

$$v_x(t) = v_{x,0}, \quad v_y(t) = v_{y,0}, \quad v_z(t) = v_{z,0} + \frac{q}{m} E_z t$$

$$x(t) = v_{x,0}t, \quad y(t) = v_{y,0}t,$$

$$z(t) = v_{z,0}t + \frac{1}{2} \left(\frac{q}{m} E_z \right) t^2$$



Individual Charge Trajectories

- **Homogeneous magnetic field**

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{v} \times \mathbf{B}) = q(\mathbf{v} \times B_z \mathbf{k})$$

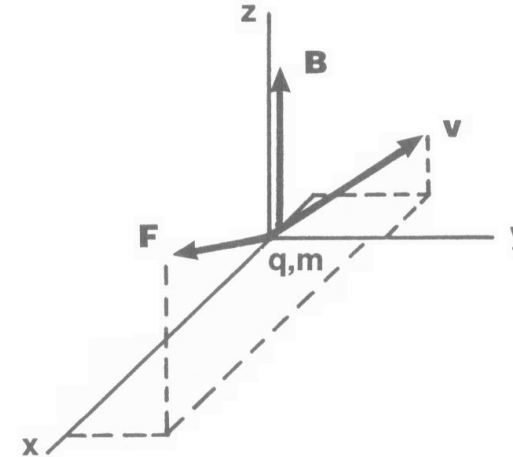
$$v_x(0) = v_{x,0}, \quad v_y(0) = v_{y,0}, \quad v_z(0) = v_{z,0} = v_{\parallel}$$

$$v_x(t) = v_{\perp} \cos(\omega_c t + \phi), \quad v_y(t) = \mp v_{\perp} \sin(\omega_c t + \phi), \quad v_z(t) = v_{z,0} = v_{\parallel}$$

$$v_{\perp} = \sqrt{v_x^2 + v_y^2}, \quad \tan(\phi) = \mp \frac{v_{y,0}}{v_{x,0}}, \quad \omega_c = \frac{|q|B_z}{m} \quad \text{Cyclotron frequency}$$

$$x(t) = x_0 + \frac{v_{\perp}}{\omega_c} \sin(\omega_c t + \phi), \quad y(t) = y_0 \pm \frac{v_{\perp}}{\omega_c} \cos(\omega_c t + \phi), \quad z(t) = z_0 + v_{\parallel} t$$

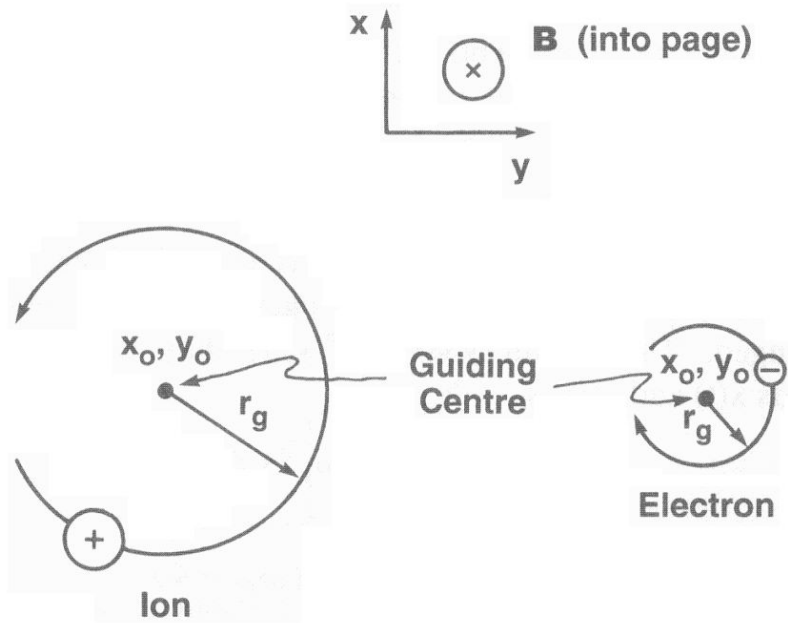
$$(x - x_0)^2 + (y - y_0)^2 = \left(\frac{v_{\perp}}{\omega_c} \right)^2 = \left(\frac{v_{\perp} m}{|q| B_z} \right)^2 = r_L^2 \quad \text{Larmor radius}$$



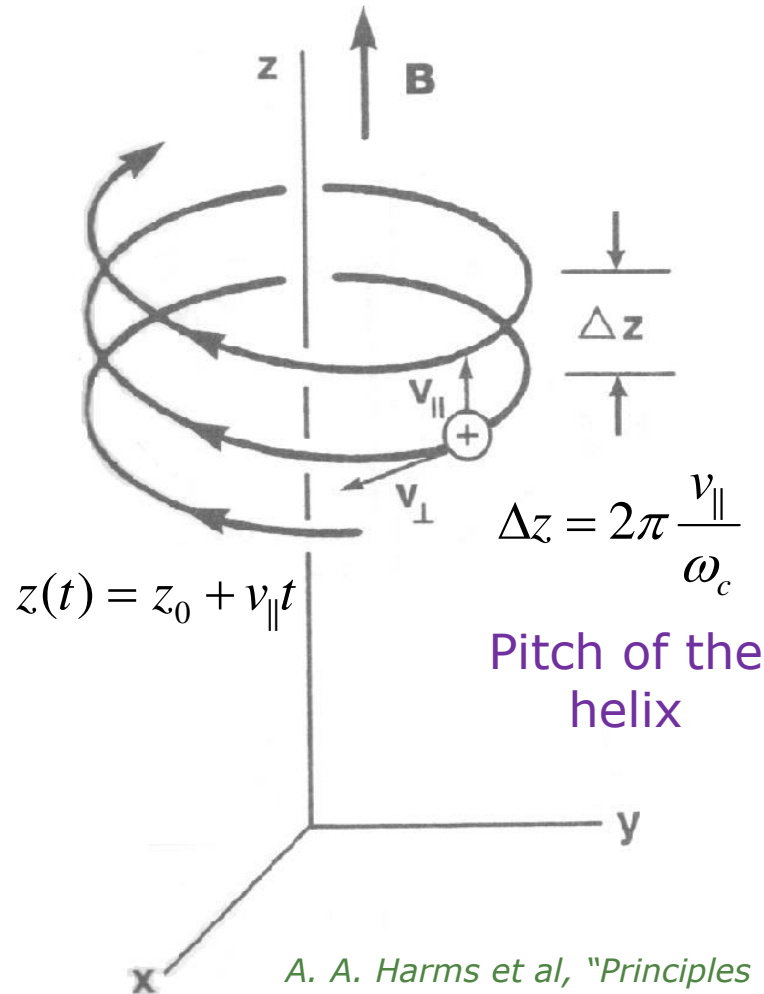
5 T, 10 keV
 $r_L = 0.05$ mm for e
 $r_L = 2.9$ mm for d
 $r_L = 3.5$ mm for t

Individual Charge Trajectories

- Homogeneous magnetic field



$$(x - x_0)^2 + (y - y_0)^2 = \left(\frac{v_{\perp} m}{|q| B_z} \right)^2 = r_L^2$$



Pitch of the helix

ion



Magnetic field

Individual Charge Trajectories

- Combined homogeneous electric and magnetic field

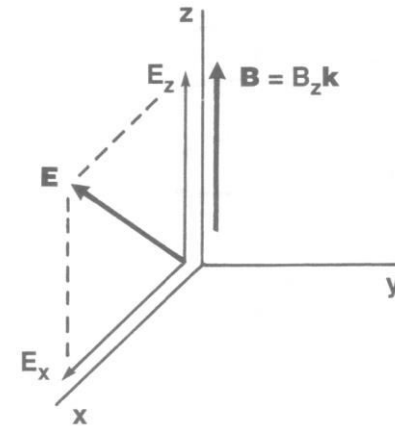
$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{E} + \mathbf{v} \times \mathbf{B}) = q[(E_x \mathbf{i} + E_z \mathbf{k}) + \mathbf{v} \times B_z \mathbf{k}]$$

$$v_x(0) = v_{x,0}, \quad v_y(0) = v_{y,0}, \quad v_z(0) = v_{z,0} = v_{\parallel}$$

$$v_x(t) = v_{\perp}^* \cos(\omega_c t + \phi^*), \quad v_y(t) = \mp v_{\perp}^* \sin(\omega_c t + \phi^*) - \frac{E_x}{B_z},$$

$$v_z(t) = v_{\parallel} + \left(\frac{qE_z}{m} \right) t$$

$$v_{\perp}^* = \sqrt{v_x^2 + \left(v_y + \frac{E_x}{B_z} \right)^2} = \frac{v_{\perp} \cos \phi}{\cos \phi^*}, \quad \tan(\phi^*) = \mp \frac{v_{y,0} + \frac{E_x}{B_z}}{v_{x,0}}$$



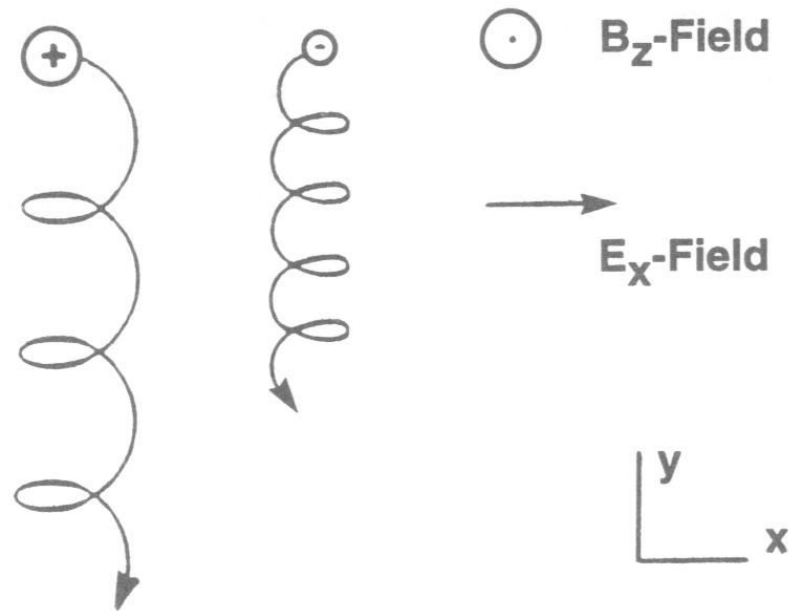
Individual Charge Trajectories

- Combined homogeneous electric and magnetic field

$$x(t) = x_0 + \frac{v_{\perp}^*}{\omega_c} \sin(\omega_c t + \phi^*), \quad y(t) = y_0 \pm \frac{v_{\perp}^*}{\omega_c} \cos(\omega_c t + \phi^*) - \frac{E_x}{B_z} t,$$

$$z(t) = z_0 + v_{\parallel} t + \frac{1}{2} \left(\frac{qE_z}{m} \right) t^2$$

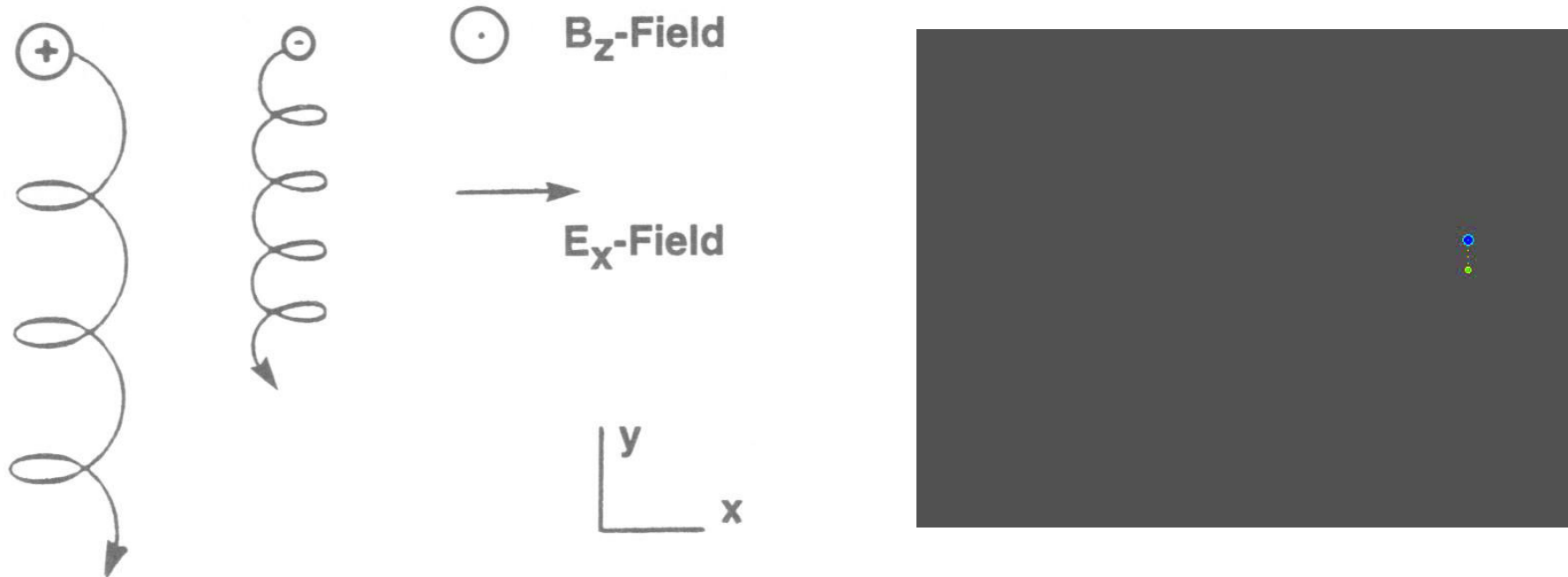
$$(x - x_0)^2 + \left(y - y_0 + \frac{E_x}{B_z} t \right)^2 = \left(\frac{v_0^*}{\omega_c} \right)^2$$



Individual Charge Trajectories

- Combined homogeneous electric and magnetic field

$$\mathbf{v} = \mathbf{v}_{gc} + \mathbf{v}_g \quad : \text{Guiding centre} + \text{Gyro motion}$$



Individual Charge Trajectories

- Combined homogeneous electric and magnetic field

$$\mathbf{v} = \mathbf{v}_{gc} + \mathbf{v}_g \quad : \text{Guiding centre} + \text{Gyro motion}$$

2.2

IN MAGNETIC FIELDS

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Because of this change of momentum, the *centre of curvature* of the projection on the xy -plane of the path is displaced a distance

$$\mathbf{D} = \mathbf{p}' - \mathbf{p} \quad (9)$$

or, because of (6)–(8),

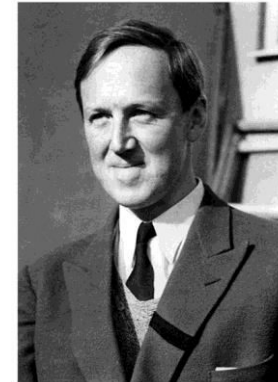
$$\mathbf{D} = -\frac{c}{eH^2}[\mathbf{H}\Delta\mathbf{p}_\perp], \quad (10)$$

$$\mathbf{D} = -\frac{c}{eH^2}\left[\mathbf{H} \int \mathbf{f}_\perp dt\right]. \quad (11)$$

This formula holds for a single collision, but of course also for a series of collisions. If $\mathbf{f}$ is a continuous force, $\mathbf{D}$ gives the displacement of that point where the centre of curvature would be if $\mathbf{f}$ vanished for a moment. This point shall be called the guiding centre. If $\mathbf{f}$ is continuous, *the guiding centre drifts with the velocity*

$$\mathbf{U}_\perp = \frac{d\mathbf{D}}{dt} = -\frac{c}{eH^2}[\mathbf{H}\mathbf{f}]. \quad (12)$$

H. Alfvén, Cosmical Electrodynamics, Oxford (1950)



Hannes Alfvén
(1908-1995)
"Nobel prize in
Physics (1970)"



Individual Charge Trajectories

- Combined homogeneous electric and magnetic field

$$\mathbf{v} = \mathbf{v}_{gc} + \mathbf{v}_g \quad : \text{Guiding centre} + \text{Gyro motion}$$

$$m \frac{d\mathbf{v}}{dt} = \mathbf{F} + q(\mathbf{v} \times \mathbf{B})$$

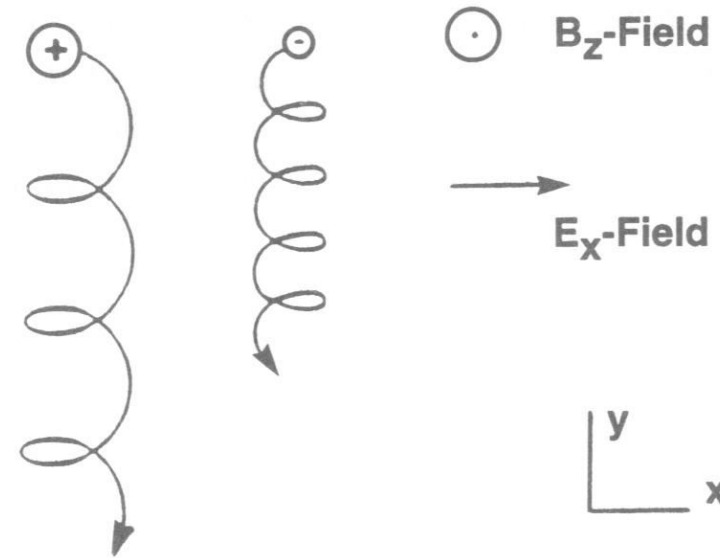
$$m \frac{d\mathbf{v}_{gc}}{dt} + m \frac{d\mathbf{v}_g}{dt} = \mathbf{F} + q(\mathbf{v}_{gc} \times \mathbf{B}) + q(\mathbf{v}_g \times \mathbf{B})$$

$$v_{gc,\parallel}(t) = v_{gc,\parallel}(0) + \frac{1}{m} \int F_{\parallel} dt$$

$$\bar{\mathbf{v}}_{gc,\perp} = \frac{\bar{\mathbf{F}} \times \mathbf{B}}{qB^2} = \mathbf{v}_{DF}$$

$$\mathbf{F} = \mathbf{E}q$$

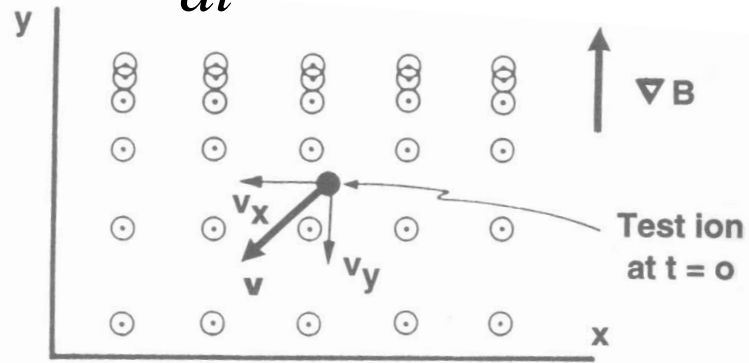
$$\mathbf{v}_{D,E} = \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$



Individual Charge Trajectories

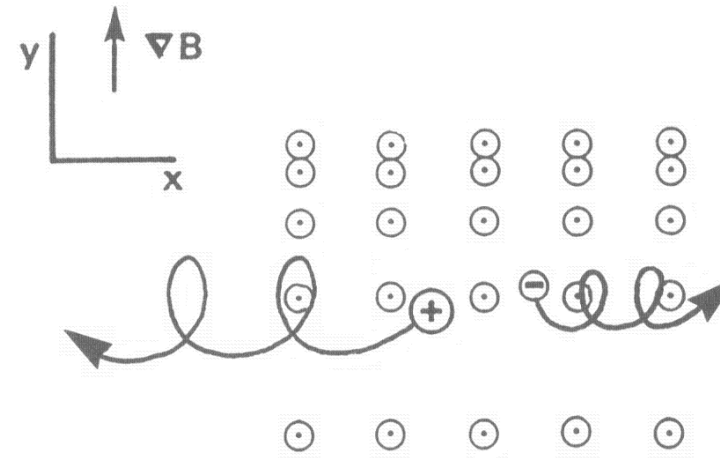
- Spatially varying magnetic field

$$m \frac{d\mathbf{v}}{dt} = q[\mathbf{v} \times \mathbf{B}(r)]$$



$$\begin{aligned} \bar{\mathbf{F}} &= q(\vec{v}_\perp \times \vec{B}) \approx qv_\perp \left(-\frac{r_L}{2} \right) \left(\frac{dB}{dy} \right) \mathbf{j} \\ &= -|q| \left(\frac{v_\perp^2}{2\omega_c} \right) \left(\frac{dB}{dy} \right) \mathbf{j} = -|q| \frac{v_\perp^2}{2\omega_c} \nabla B \\ &= -\frac{|q|}{2} v_\perp r_L \nabla B = -\frac{mv_\perp^2 / 2}{B} \nabla B \end{aligned}$$

$$\nabla B = \frac{\partial B_z}{\partial y} \mathbf{j}, \quad v_\parallel = 0$$

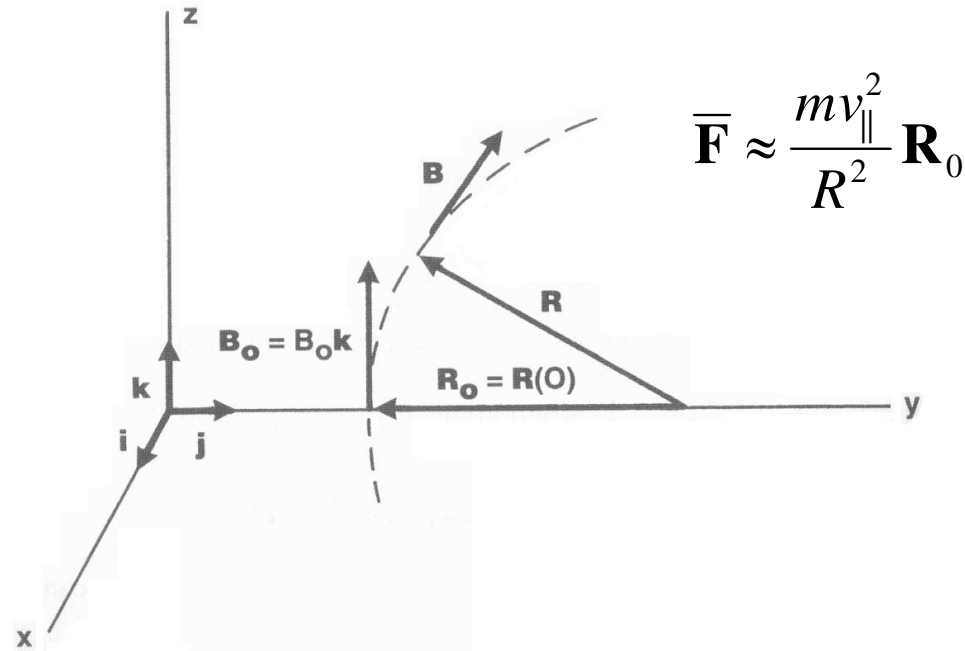


$$\begin{aligned} \mathbf{v}_{D,\nabla B} &= \frac{\bar{\mathbf{F}} \times \mathbf{B}}{qB^2} \\ &= \pm \frac{v_\perp^2}{2\omega_c} \frac{\mathbf{B} \times \nabla B}{B^2} = \pm \frac{1}{2} v_\perp r_L \frac{\mathbf{B} \times \nabla B}{B^2} \end{aligned}$$

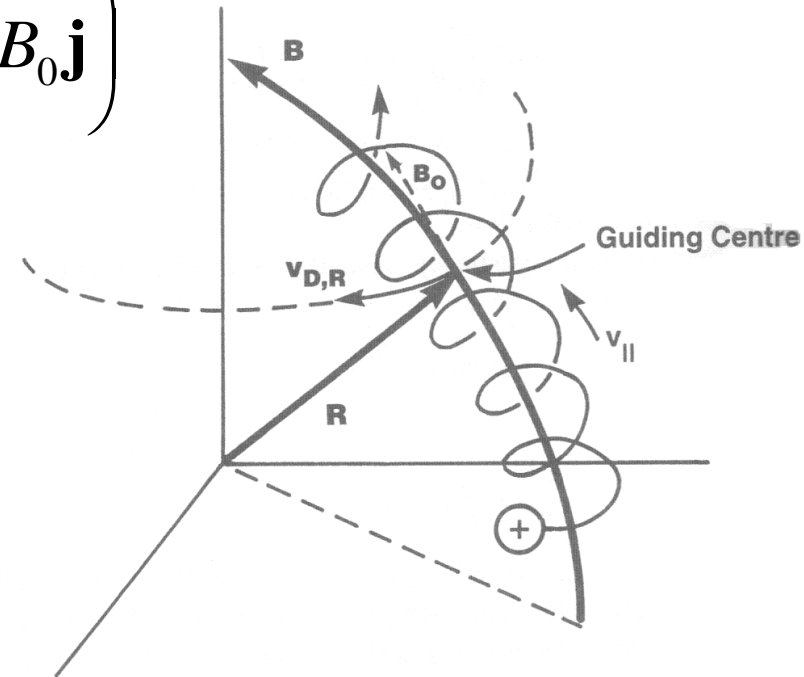
Individual Charge Trajectories

- **Curvature drift**

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{v} \times \mathbf{B}) = q(\mathbf{v} \times B_0 \mathbf{k}) + q \left(\mathbf{v} \times \frac{z}{R} B_0 \mathbf{j} \right)$$



$$\bar{\mathbf{F}} \approx \frac{mv_{\parallel}^2}{R^2} \mathbf{R}_0$$

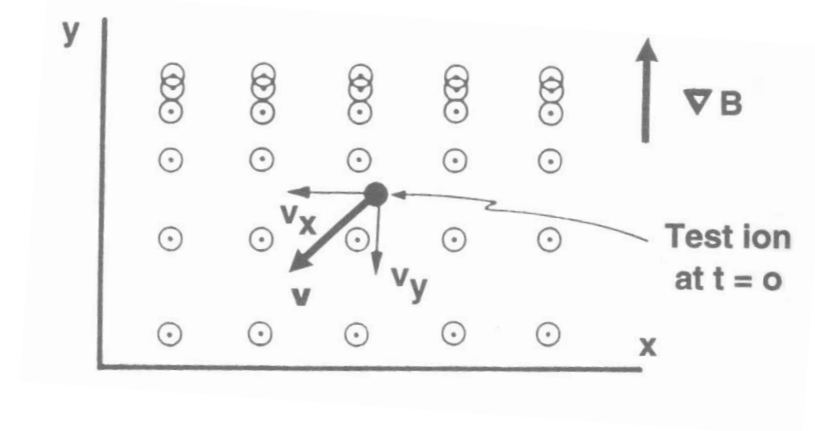
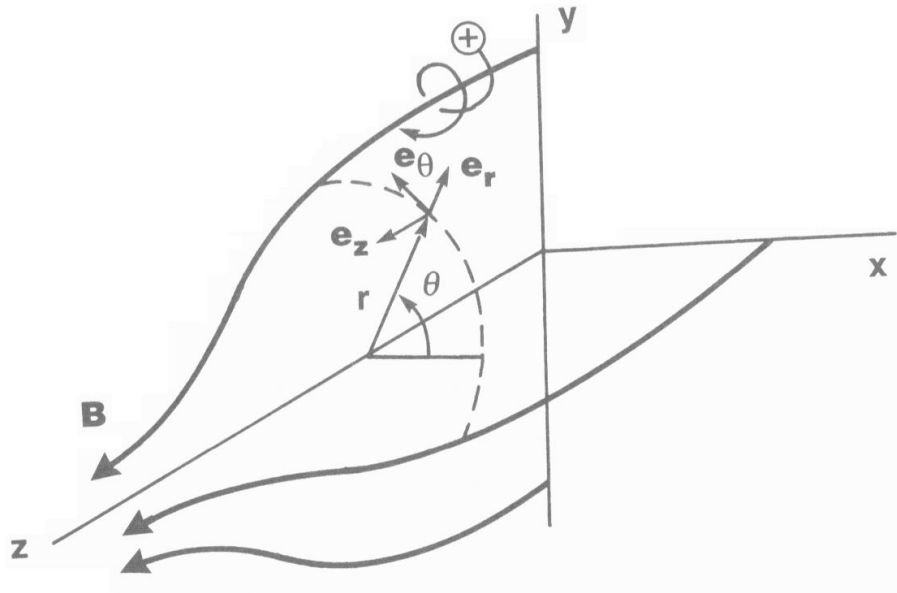


$$\mathbf{v}_{D,R} = \bar{\mathbf{v}}_{gc,x} = \frac{mv_{\parallel}^2}{qB_0^2} \frac{\mathbf{R}_0 \times \mathbf{B}_0}{R^2}$$

Individual Charge Trajectories

- Axial field variation**

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{v} \times B_z \mathbf{e}_z) + q(\mathbf{v} \times B_r \mathbf{e}_r)$$

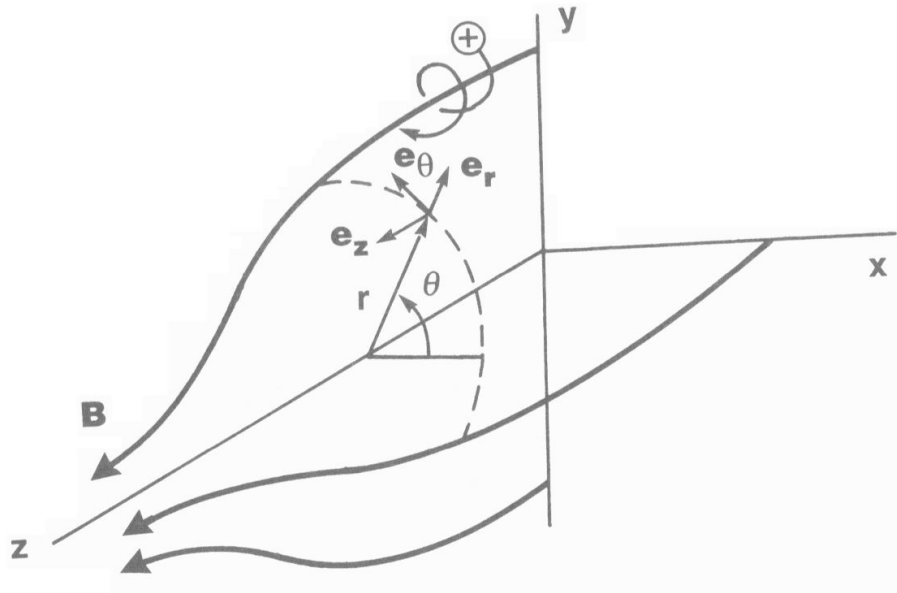


$$\begin{aligned} \bar{\mathbf{F}} &= q(\bar{\mathbf{v}}_{\perp} \times \vec{B}) \approx qv_{\perp} \left(-\frac{r_L}{2} \right) \left(\frac{dB}{dy} \right) \mathbf{j} \\ &= -|q| \left(\frac{v_{\perp}^2}{2\omega_c} \right) \left(\frac{dB}{dy} \right) \mathbf{j} = -|q| \frac{v_{\perp}^2}{2\omega_c} \nabla B \\ &= -\frac{|q|}{2} v_{\perp} r_L \nabla B = -\frac{mv_{\perp}^2 / 2}{B} \nabla B \end{aligned}$$

Individual Charge Trajectories

- **Axial field variation**

$$m \frac{d\mathbf{v}}{dt} = q(\mathbf{v} \times B_z \mathbf{e}_z) + q(\mathbf{v} \times B_r \mathbf{e}_r)$$



$$\nabla \cdot \mathbf{B} = 0 \rightarrow B_r \approx -\frac{r}{2} \frac{\partial B_z}{\partial z} \quad B_\theta = 0, \quad B_r(r=0) = 0$$

$$\mathbf{F}_\parallel = -\frac{1}{2} \frac{mv_\perp^2}{B} \nabla_\parallel B = -\mu \nabla_\parallel B$$

μ : magnetic moment of the gyrating particle

Individual Charge Trajectories

- **Invariant of motion**

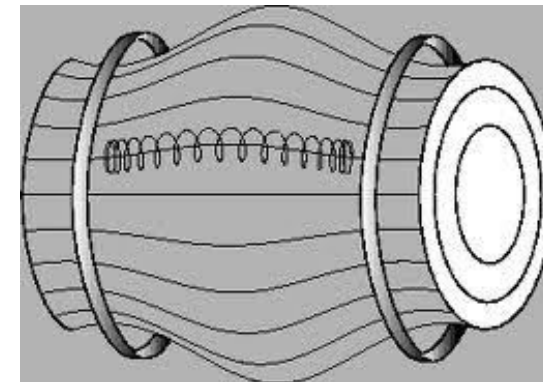
$$\frac{d}{dt} E_0 = \frac{d}{dt} \left(\frac{1}{2} m v_{\perp}^2 + \frac{1}{2} m v_{\parallel}^2 \right) = 0 \quad \mu = \frac{m v_{\perp}^2 / 2}{B} \quad m \frac{dv_{\parallel}}{dt} = -\frac{\mu}{v_{\parallel}} \frac{dB}{dt} \quad m v_{\parallel} \frac{dv_{\parallel}}{dt} = -\mu \frac{dB}{dt}$$

$$\mathbf{F}_{\parallel} = m \frac{d\mathbf{v}_{\parallel}}{dt} = -\mu \nabla_{\parallel} \mathbf{B} = -\mu \frac{\partial \mathbf{B}}{\partial s} = -\mu \frac{\partial \mathbf{B}}{\partial s} \cdot \frac{ds}{dt} \cdot \frac{dt}{ds} = -\mu \frac{\partial \mathbf{B}}{\partial s} \cdot \frac{ds}{dt} \cdot \frac{1}{v_{\parallel}} = -\frac{\mu}{v_{\parallel}} \frac{dB}{dt} \quad \rightarrow \quad \frac{d}{dt} \left(\frac{1}{2} m v_{\parallel}^2 \right) = -\mu \frac{dB}{dt}$$

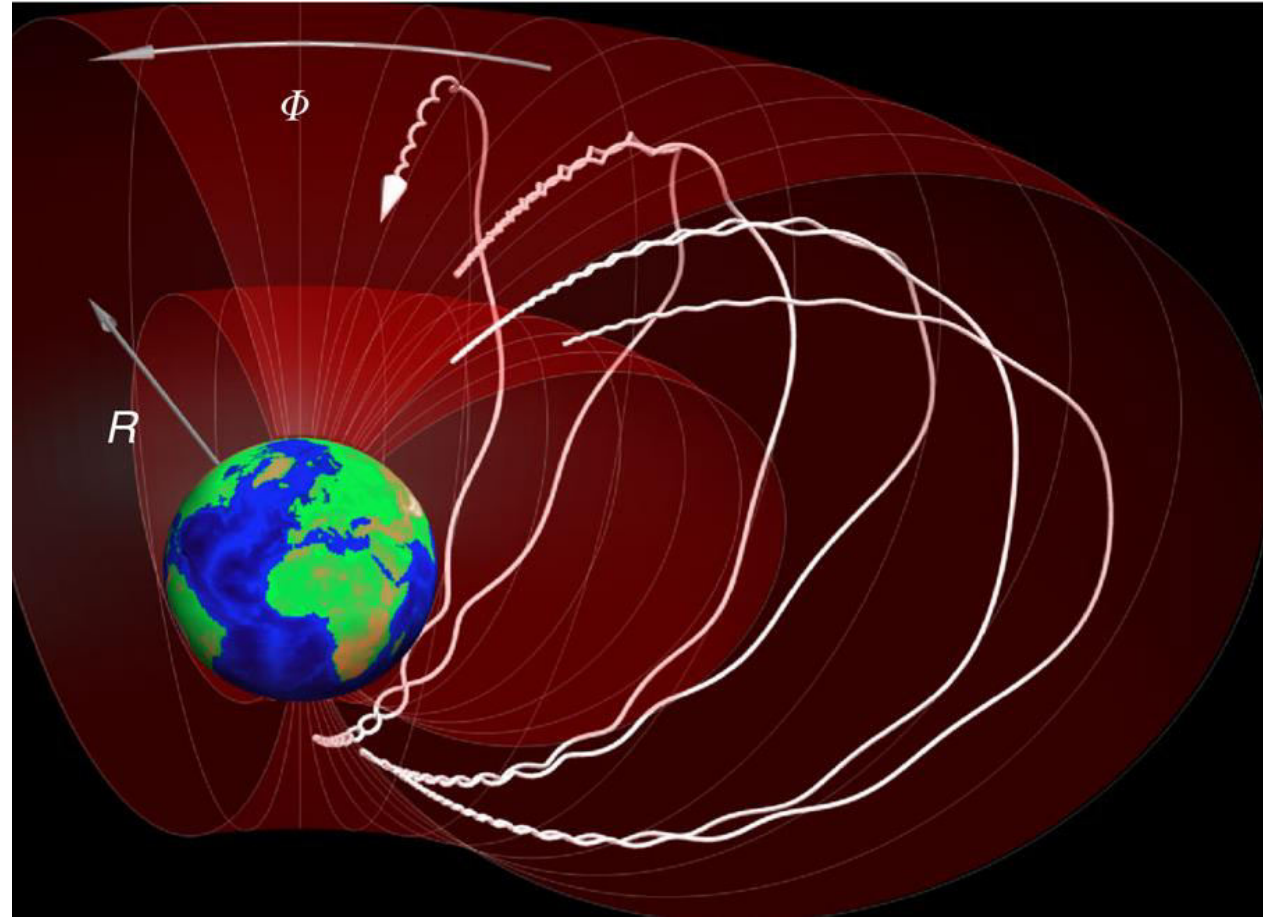
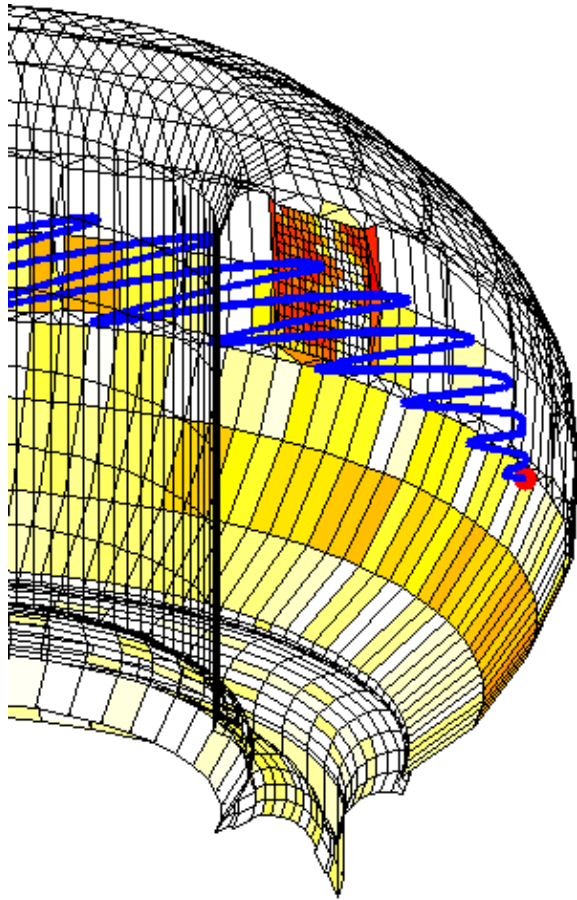
$$\rightarrow \frac{d}{dt} \left(\frac{1}{2} m v_{\perp}^2 + \frac{1}{2} m v_{\parallel}^2 \right) = \frac{d}{dt} (\mu B) + \left(-\mu \frac{dB}{dt} \right) = 0$$

$$\frac{d}{dt} (\mu B) = \mu \frac{dB}{dt} + B \frac{d\mu}{dt}$$

$$\rightarrow \frac{d\mu}{dt} = 0 : \text{adiabatic invariant}$$

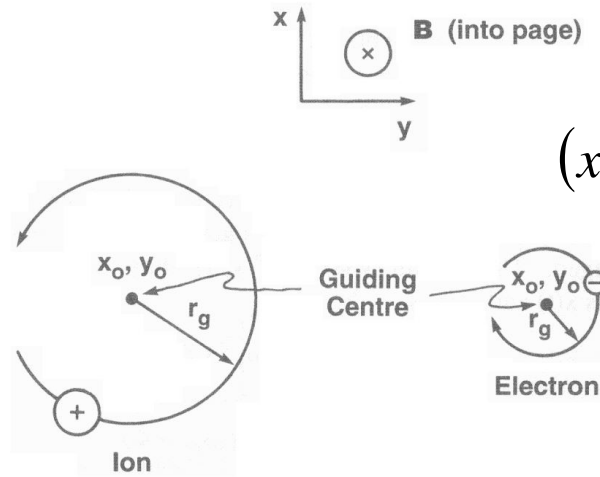


Individual Charge Trajectories

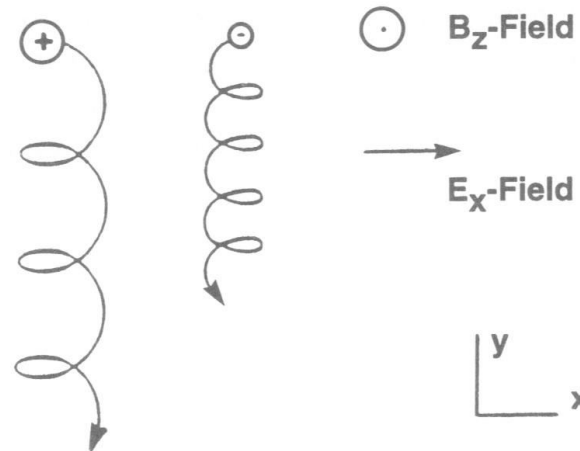
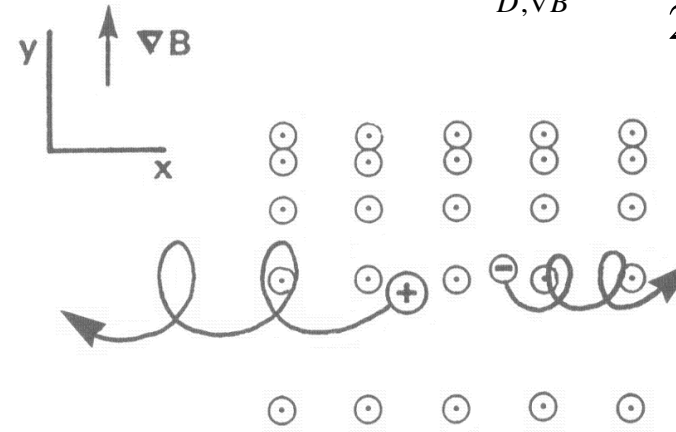


Individual Charge Trajectories

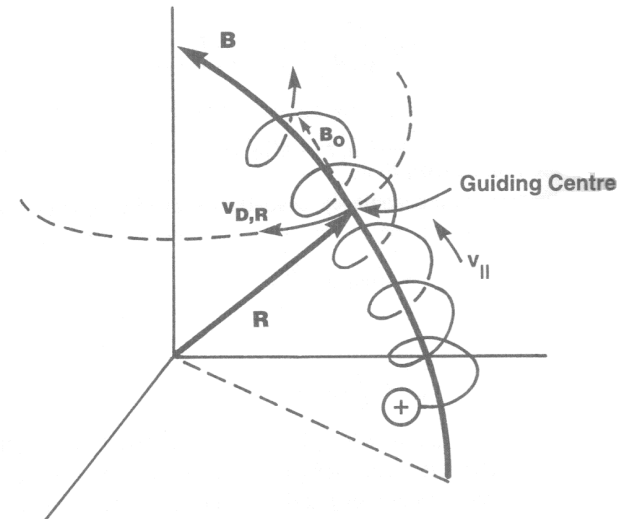
$$\mathbf{v}_{D,\nabla B} = \pm \frac{1}{2} v_{\perp} r_L \frac{\mathbf{B} \times \nabla B}{B^2}$$



$$(x - x_0)^2 + (y - y_0)^2 = r_L^2$$



$$\mathbf{v}_{D,E} = \frac{\mathbf{E} \times \mathbf{B}}{B^2}$$



$$\mathbf{v}_{D,R} = \frac{mv_{\parallel}^2}{qB_0^2} \frac{\mathbf{R}_0 \times \mathbf{B}_0}{R^2}$$