Chapter 17. Diffusivity and the mechanisms of mass transport

- Fick's law of binary diffusion (molecular mass transport)
- Temperature and pressure dependence of diffusivities
- Theory of diffusion in gases at low density, in binary liquids, in colloids suspensions, and polymers
- Mass and molar transport by convection
- Summary of mass and molar fluxes
- The Maxwell-Stefan equations for multicomponent diffusion in gases at low density

Behaviour of polymeric liquids



At steady state

$$\frac{w_{Ay}}{A} = \rho \mathfrak{D}_{AB} \frac{\omega_{A0} - 0}{Y}$$

Fick's law of

binary diffusion

$$j_{Ay} = -\rho \mathfrak{D}_{AB} \frac{d\omega_A}{dy}$$

Fick's law of binary diffusion

$$j_{Ay} = -\rho \mathcal{D}_{AB} \frac{d\omega_A}{dy}$$

 j_{Ay} , the molecular mass flux ω_A , the mass fraction \mathfrak{D}_{AB} the diffusivity ρ is the density

• Mass average velocity for a binary mixture

$$v_y = \omega_A v_{Ay} + \omega_B v_{By}$$

• Mass flux

$$j_{Ay} = \rho \omega_A (v_{Ay} - v_y)$$

Dimensionless numbers

mass diffusivity DAB.

thermal diffusivity $\alpha = k/\rho \hat{C}_{p}$

momentum diffusivity $\nu = \mu / \rho$

The Prandtl number: $\Pr = \frac{\nu}{\alpha} = \frac{\hat{C}_p \mu}{\nu}$ The Schmidt number:² Sc = $\frac{\nu}{\mathcal{D}_{AB}} = \frac{\mu}{\rho \mathcal{D}_{AB}}$ k The Lewis number:² $Le = \frac{\alpha}{2}$

$$\frac{d}{\mathcal{D}_{AB}} = \frac{1}{\rho \hat{C}_p \mathcal{D}_{AB}}$$

17.2 Temperature and pressure dependence of Diffusivities

• For gas mixtures at low pressure (kinetic theory)

$$\frac{p\mathfrak{D}_{AB}}{(p_{cA}p_{cB})^{1/3}(T_{cA}T_{cB})^{5/12}(1/M_A + 1/M_B)^{1/2}} = a \left(\frac{T}{\sqrt{T_{cA}T_{cB}}}\right)^b$$

- Diffusivities
 - are inversely proportional to the pressure
 - increases with increasing temperature
 - almost independent of composition

17.3 Theory of diffusion in gases at low density

• Self diffusivity

$$\mathfrak{D}_{AA^{*}} = \frac{2}{3} \frac{\sqrt{\kappa T / \pi m_{A}}}{\pi d_{A}^{2}} \frac{1}{n} = \frac{2}{3\pi} \frac{\sqrt{\pi m_{A} \kappa T}}{\pi d_{A}^{2}} \frac{1}{\rho}$$

• For binary mixtures

$$\mathcal{D}_{AB} = \frac{2}{3} \sqrt{\frac{\kappa T}{\pi}} \sqrt{\frac{1}{2} \left(\frac{1}{m_A} + \frac{1}{m_B}\right)} \frac{1}{\pi (\frac{1}{2} (d_A + d_B))^2} \frac{1}{n}$$

Chapman-Enskog kinetic theory

$$c\mathcal{D}_{AB} = \frac{3}{16} \sqrt{\frac{2RT}{\pi} \left(\frac{1}{M_A} + \frac{1}{M_B}\right)} \frac{1}{\tilde{N}\sigma_{AB}^2 \Omega_{\mathcal{D},AB}} \qquad \text{collision integral}$$

17.7 Mass and molar transport by convection

- Mass and molar concentrations
 - Mass average and molar average velocities
 - Molecular mass and molar fluxes
 - Convective mass and molar fluxes

Mass and molar concentrations

 $\rho_{\alpha} = \text{mass concentration of species } \alpha$ $\rho = \sum_{\alpha=1}^{N} \rho_{\alpha} = \text{mass density of solution}$ $\omega_{\alpha} = \rho_{\alpha} / \rho = \text{mass fraction of species } \alpha$

 c_{α} = molar concentration of species α $c = \sum_{\alpha=1}^{N} c_{\alpha}$ = molar density of solution $x_{\alpha} = c_{\alpha}/c$ = mole fraction of species α

Mass and molar concentrations

"Transport Phenomena" 2nd ed., R.B. Bird, W.E. Stewart, E.N. Lightfoot



Differential relations:

$$\nabla x_a = -\frac{M^2}{M_a} \sum_{\substack{\gamma=1\\\gamma\neq a}}^N \left[\frac{1}{M} + \omega_a \left(\frac{1}{M_\gamma} - \frac{1}{M_\alpha} \right) \right] \nabla \omega_\gamma \tag{P}^a$$

$$\nabla \omega_a = -\frac{M_a}{M^2} \sum_{\substack{\gamma=1\\\gamma\neq\alpha}}^N [M + x_a (M_\gamma - M_\alpha)] \nabla x_\gamma \tag{Q}^a$$

*Equations (P) and (Q), simplified for binary ystems, are

$$\nabla x_A = \frac{\frac{1}{M_A M_B} \nabla \omega_A}{\left(\frac{\omega_A}{M_A} + \frac{\omega_B}{M_B}\right)^2} \qquad (P') \qquad \qquad \nabla \omega_A = \frac{M_A M_B \nabla x_A}{\left(x_A M_A + x_B M_B\right)^2} \qquad (Q')$$

Mass average and molar averages velocity

• Mass average velocity

$$\mathbf{v} = \frac{\sum_{\alpha=1}^{N} \rho_{\alpha} \mathbf{v}_{\alpha}}{\sum_{\alpha=1}^{N} \rho_{\alpha}} = \frac{\sum_{\alpha=1}^{N} \rho_{\alpha} \mathbf{v}_{\alpha}}{\rho} = \sum_{\alpha=1}^{N} \omega_{\alpha} \mathbf{v}_{\alpha}$$

• Molar average velocity

$$\mathbf{v}^* = \frac{\sum\limits_{\alpha=1}^N c_\alpha \mathbf{v}_\alpha}{\sum\limits_{\alpha=1}^N c_\alpha} = \frac{\sum\limits_{\alpha=1}^N c_\alpha \mathbf{v}_\alpha}{\sum\limits_{\alpha=1}^N c_\alpha} = \sum\limits_{\alpha=1}^N x_\alpha \mathbf{v}_\alpha$$

17.8 Summary of mass and molar fluxes

Equivalent forms of Fick's law of binary diffusion

Table 17.8-2 Equivalent Forms of Fick's (First) Law of Binary Diffusion

Flux	Gradient	Form of Fick's Law
j 4	$\nabla \omega_A$	$\mathbf{j}_{A}=- ho\mathfrak{D}_{AB} abla \omega_{A}$
\mathbf{J}^*_{Λ}	∇x_A	$\mathbf{J}^*_A = -c \mathfrak{D}_{AB} \nabla x_A$
n _A	$\nabla \omega_A$	$\mathbf{n}_{A} = \boldsymbol{\omega}_{A}(\mathbf{n}_{A} + \mathbf{n}_{B}) - \boldsymbol{\rho} \boldsymbol{\mathfrak{D}}_{AB} \nabla \boldsymbol{\omega}_{A} = \boldsymbol{\rho}_{A} \mathbf{v} - \boldsymbol{\rho} \boldsymbol{\mathfrak{D}}_{AB} \nabla \boldsymbol{\omega}_{A}$
\mathbb{N}_{Λ}	∇x_A	$\mathbb{N}_{A} = x_{A}(\mathbb{N}_{A} + \mathbb{N}_{B}) - c \mathcal{D}_{AB} \nabla x_{A} = c_{A} \mathbf{v}^{*} - c \mathcal{D}_{AB} \nabla x_{A}$

"Transport Phenomena" 2nd ed., R.B. Bird, W.E. Stewart, E.N. Lightfoot

Molecular mass and molar fluxes

"Transport Phenomena" 2nd ed., R.B. Bird, W.E. Stewart, E.N. Lightfoot

Table 17.0-1 Trotation 101 Mass and Motal Truxes	Table 17.8-1	Notation	for	Mass	and	Molar	Fluxes
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Quantity	With respect to stationary axes		With respect to mass average velocity v	With respect to molar average velocity v*		
Velocity of species α (cm/s)	ν _α	(A)	$\mathbf{v}_{lpha} = \mathbf{v}$	(B)	$\mathbf{v}_{lpha} = \mathbf{v}^{*}$	(C)
Mass flux of species α (g/cm ² s)	$\mathbf{n}_{\alpha} = ho_{lpha} \mathbf{v}_{lpha}$	(D)	$\mathbf{j}_{\alpha} = \boldsymbol{\rho}_{\alpha}(\mathbf{v}_{\alpha} - \mathbf{v})$	(E)	$\mathbf{j}_{\alpha}^{*} = \rho_{\alpha}(\mathbf{v}_{\alpha} - \mathbf{v}^{*})$	(F)
Molar flux of species α (g-moles/cm ² s)	$\mathbb{N}_{\alpha} = c_{\alpha} \mathbb{V}_{\alpha}$	(G)	$\mathbf{J}_{\alpha} = c_{\alpha}(\mathbf{v}_{\alpha} - \mathbf{v})$	(H)	$\mathbf{J}^*_{\alpha} = c_{\alpha}(\mathbf{v}_{\alpha} - \mathbf{v}^*)$	(I)
Sums of mass fluxes	$\sum_{\alpha=1}^{N} \mathbf{n}_{\alpha} = \boldsymbol{\rho} \mathbf{v}$	(J)	$\sum_{\alpha=1}^{N} \mathbf{j}_{\alpha} = 0$	(K)	$\sum_{\alpha=1}^{N} \mathbf{j}_{\alpha}^{*} = \rho(\mathbf{v} - \mathbf{v}^{*})$	(L)
Sums of molar fluxes	$\sum_{\alpha=1}^{N} \mathbb{N}_{\alpha} = c \mathbf{v}^*$	(M)	$\sum_{\alpha=1}^{N} \mathbf{J}_{\alpha} = c(\mathbf{v}^* - \mathbf{v})$	(N)	$\sum_{\alpha=1}^{N} \mathbf{J}_{\alpha}^{*} = 0$	(O)
Relations between mass and molar fluxes	$\mathbf{n}_{\alpha} = M_{\alpha} \mathbb{N}_{\alpha}$	(P)	$\mathbf{j}_{\alpha} = M_{\alpha} \mathbf{J}_{\alpha}$	(Q)	$\mathbf{j}_{\alpha}^{*} = M_{\alpha} \mathbf{J}_{\alpha}^{*}$	(R)
Interrelations among mass fluxes	$\mathbf{n}_{\alpha} = \mathbf{j}_{\alpha} + \boldsymbol{\rho}_{\alpha} \mathbf{v}$	(S)	$\mathbf{j}_{\alpha} = \mathbf{n}_{\alpha} - \boldsymbol{\omega}_{\alpha} \sum_{\beta=1}^{N} \mathbf{n}_{\beta}$	(T)	$\mathbf{j}_{\alpha}^{*} = \mathbf{n}_{a} - x_{\alpha} \sum_{\beta=1}^{N} \frac{M_{\alpha}}{M_{\beta}} \mathbf{n}_{\beta}$	(U)
Interrelations among molar fluxes	$\mathbb{N}_{\alpha} = \mathbb{J}_{\alpha}^{*} + c_{\alpha} \mathbf{v}^{*}$	(V)	$\mathbf{J}_{\alpha} = \mathbf{N}_{\alpha} - \boldsymbol{\omega}_{\alpha} \sum_{\beta=1}^{N} \frac{M_{\beta}}{M_{\alpha}} \mathbf{N}_{\beta}$	(W)	$\mathbf{J}_{\alpha}^{*} = \mathbf{N}_{\alpha} - x_{\alpha} \sum_{\beta=1}^{N} \mathbf{N}_{\beta}$	(X)

Maxwell-Stefan equations for multicomponent diffusion in gases at low pressure

Maxwell-Stefan equations

$$\nabla x_{\alpha} = -\sum_{\beta=1}^{N} \frac{x_{\alpha} x_{\beta}}{\mathfrak{D}_{\alpha\beta}} (\mathbf{v}_{\alpha} - \mathbf{v}_{\beta}) = -\sum_{\beta=1}^{N} \frac{1}{c \mathfrak{D}_{\alpha\beta}} (x_{\beta} \mathbb{N}_{\alpha} - x_{\alpha} \mathbb{N}_{\beta}) \qquad \alpha = 1, 2, 3, \dots, N$$

- Other cases in multicomponent diffusion
 - Reverse diffusion, a species move against its own gradient
 - Osmotic diffusion, a species diffuses though its concentration gradient is zero
 - Diffusion barrier, a species does not diffuse through its concentration gradient is nonzero