
Chapter 16

Amorphous State

Structure in amorphous state

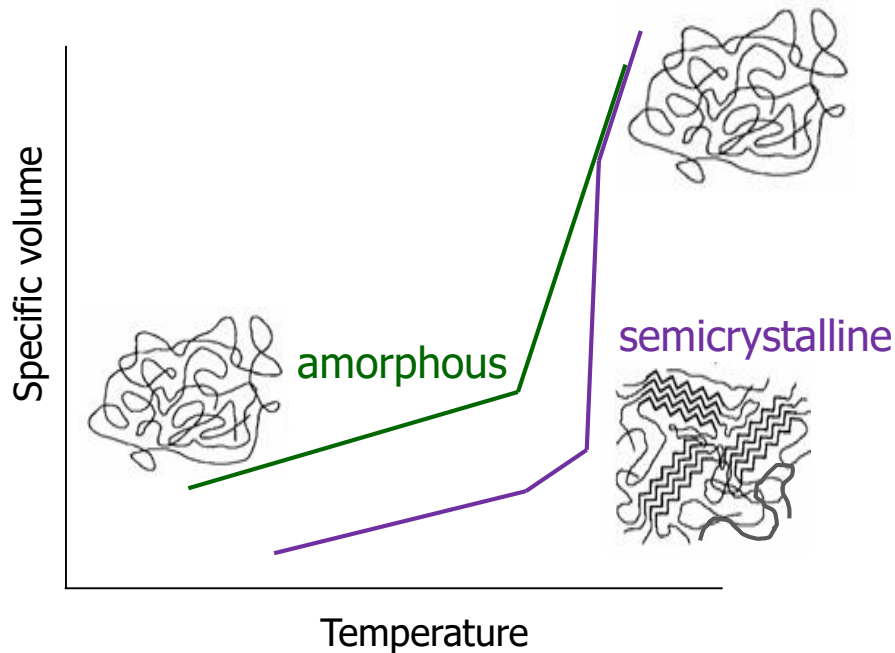
Glass transition

Macromolecular dynamics

Solid-state structure of polymers

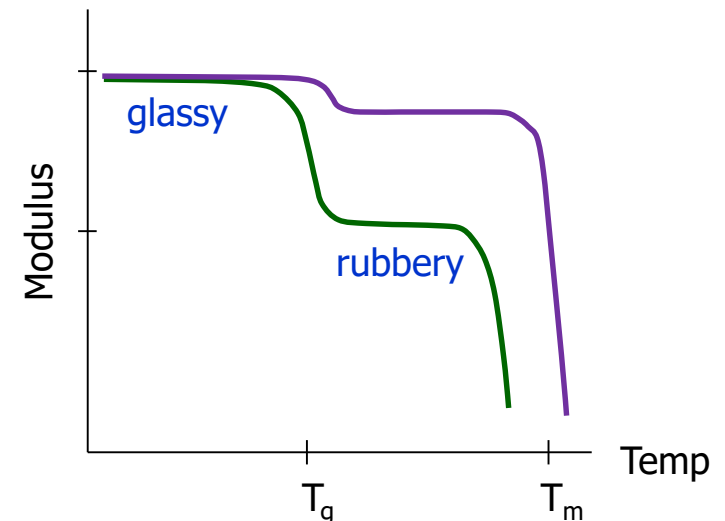
Ch 16 sl 2

- Polymers upon **cooling** from melt



- **amorphous** or **semicrystalline**
- depending on
 - regularity of chain
 - rate of formation [cooling rate]

upon **heating** from solid



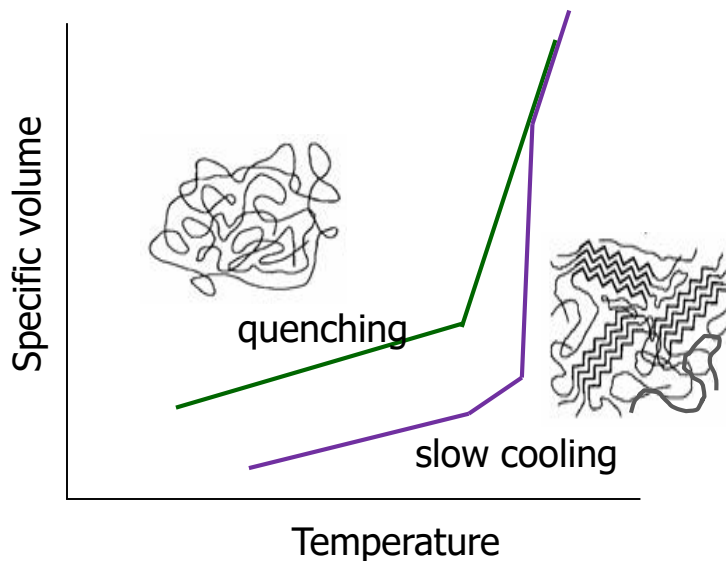
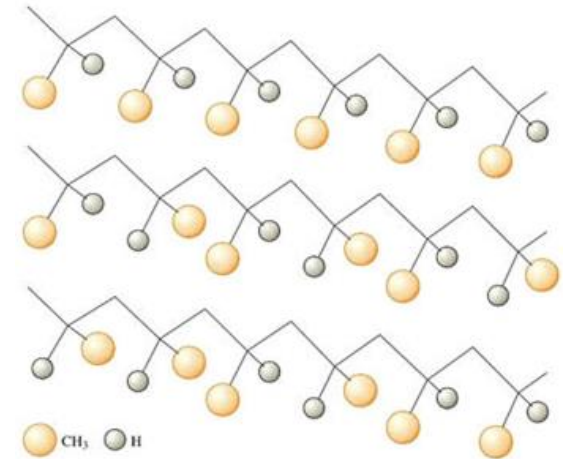
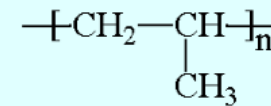
glass(-rubber) transition - only for amorphous polymers or regions
melting (of crystal)



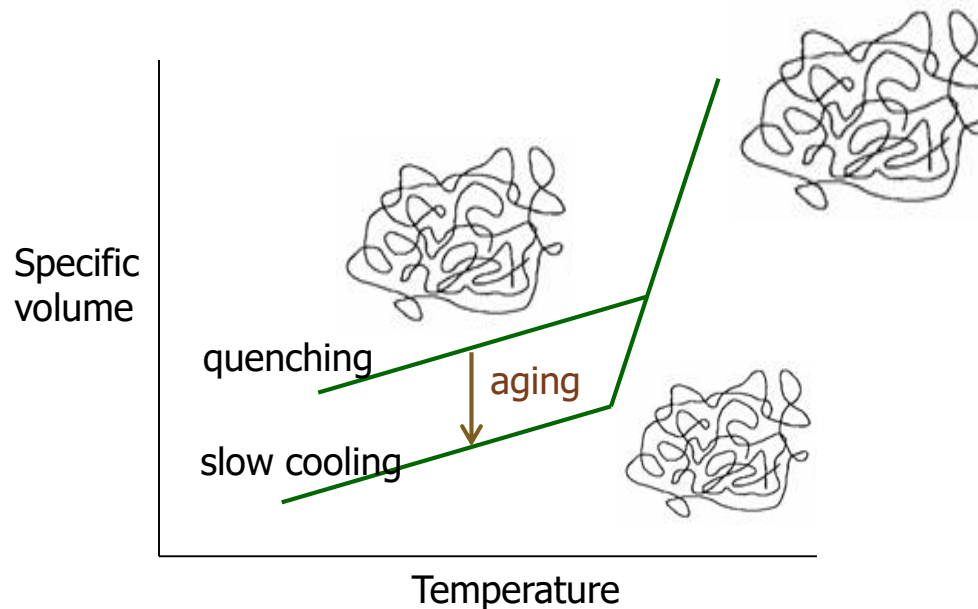
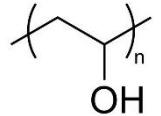
Chain regularity and cooling rate

Ch 16 sl 3

- regular chain → crystallizable
 - stereoregular [isotactic or syndiotactic] vinyl polymers
 - iPP, sPS, ---
 - main-chain polymers (w/ no pendant group)
 - PE, polyester, nylon, ---
 - **semicrystalline** when slow-cooled;
amorphous upon quenching



- ❑ irregular chain → not crystallizable
 - ❑ atactic vinyl polymers
 - ❑ (massively) branched polymers
 - ❑ **amorphous** upon quenching;
aged amorphous upon slow-cooling
 - ❑ **semicrystalline** when side group is (very) small (like PVA)
or crystallizable itself (>C8 or mesogenic)



Structure in amorphous polymers

Ch 16 sl 5

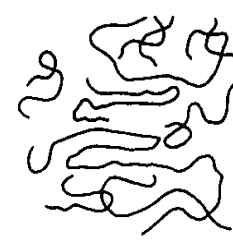
- ❑ morphology of amorphous state?
 - = order in disordered state?
 - = Is it completely disordered?
 - ❑ arguing **partial ordering**
 - electron diffraction, WAXS, (SEM ~ artifact)
 - ❑ arguing **no ordering** over 1 nm
 - birefringence, WAXS, Raman, NMR, etc
 - **natural parallelization** ~ 'spaghetti'
 - **Order should be local, incomplete, and in a very short range, if any.**
- ❑ **chain dimension** in amorphous state
 - ❑ can be determined by SANS
 - ❑ same to that determined in theta condition [**unperturbed state**]



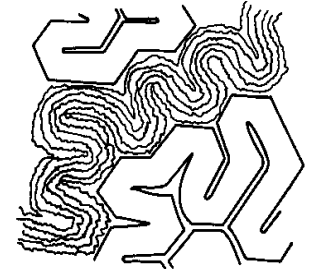
(a) P. J. Flory



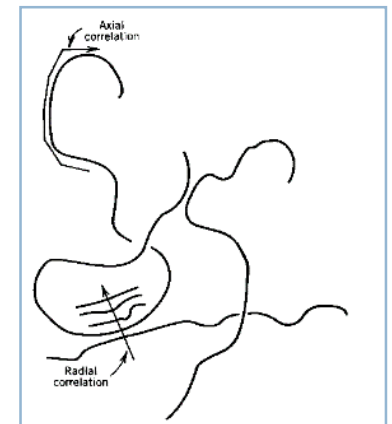
(b) V. P. Privalko and Y. S. Lipatov



(c) G. S. Y. Yeh



(d) W. Pechhold et al.



Glass transition 1: Thermodynamics Ch 16 sl 6

□ phase transition (by Ehrenfest)

■ 1st-order transition

- 1st derivatives of free energy G is discontinuous

$$\left(\frac{\partial G}{\partial P}\right)_T = V \quad \left(\frac{\partial(G/T)}{\partial(1/T)}\right)_P = H \quad \left(\frac{\partial G}{\partial T}\right)_P = -S$$

- step change in **V, H, S**
- melting, boiling, polymorphism, ---

■ 2nd-order transition

- 2nd derivative of free energy is discontinuous

- step change in **α, κ, C_p**

$$\left(\frac{\partial}{\partial T} \left(\frac{\partial G}{\partial P}\right)_T\right)_P = \left(\frac{\partial V}{\partial T}\right)_P = \alpha V \quad \left(\frac{\partial^2 G}{\partial P^2}\right)_T = \left(\frac{\partial V}{\partial P}\right)_T = -\kappa V$$

$$-\left(\frac{\partial^2 G}{\partial T^2}\right)_P = \left(\frac{\partial S}{\partial T}\right)_P = \frac{C_p}{T} \quad \left(\frac{\partial H}{\partial T}\right)_P = \left(\frac{\partial Q}{\partial T}\right)_P = C_p$$

- glass-rubber, normal-superconducting, ---

Fig 16.3

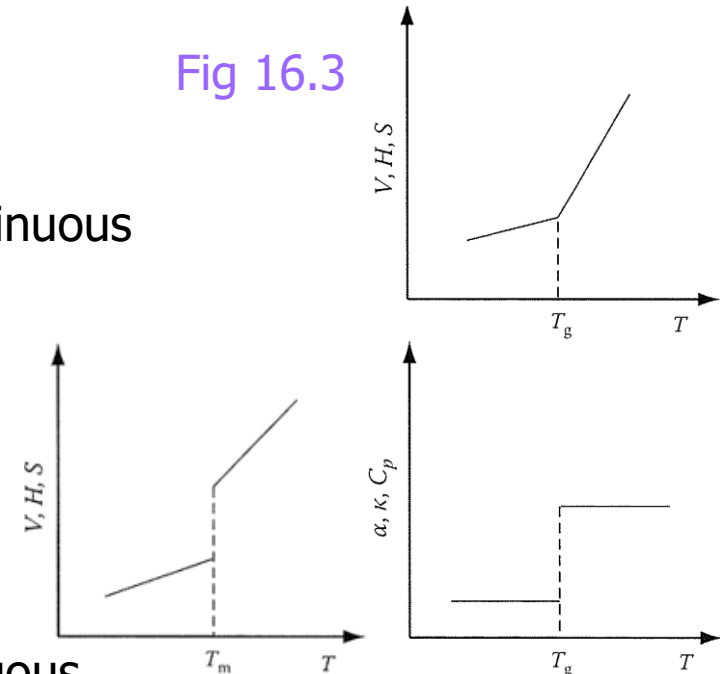
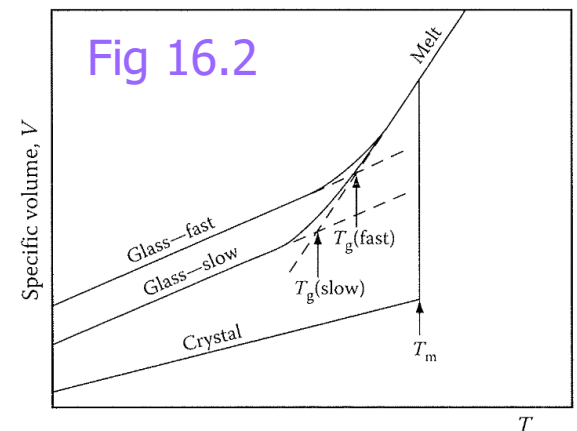


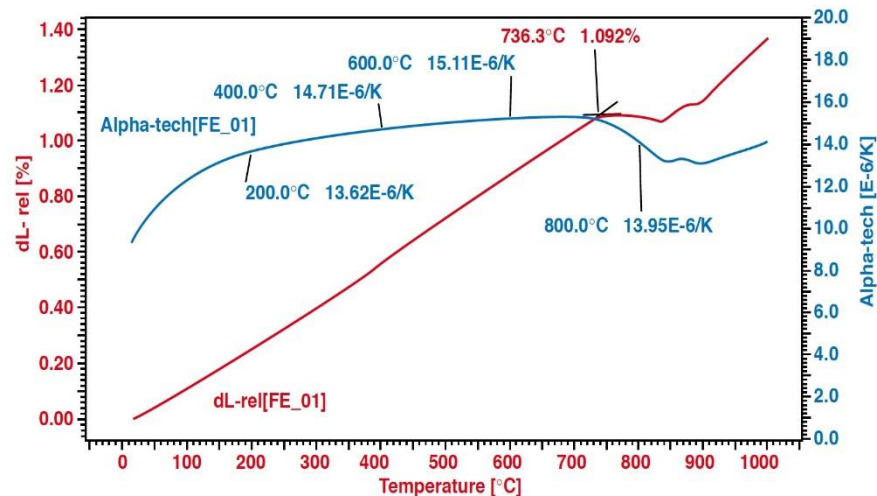
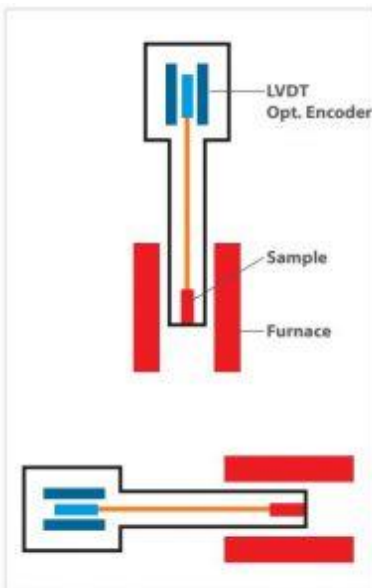
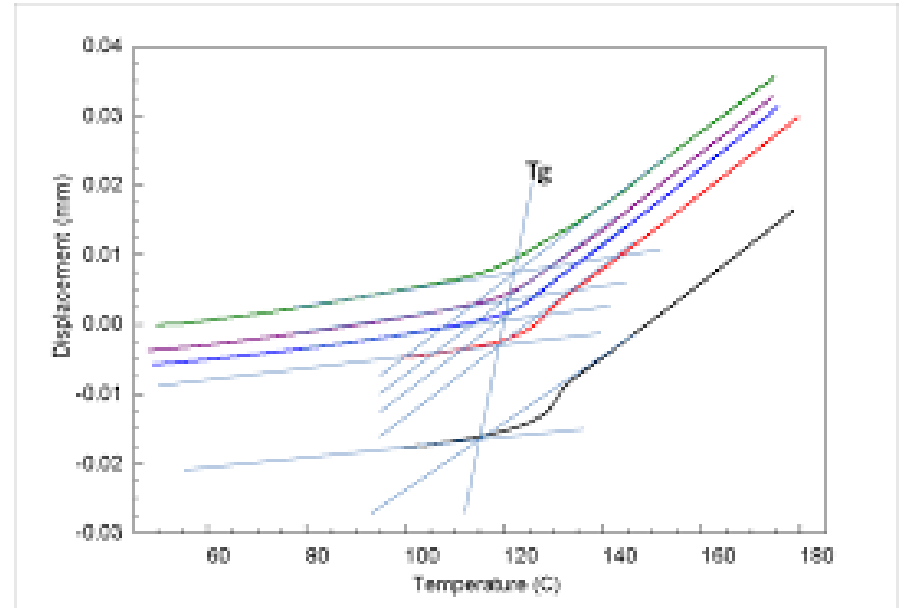
Fig 16.2



Determination of T_g

□ dilatometry

- change in volume (or α)
 - $\Delta V \leftarrow \Delta L$ if isotropic
- dilatometer or thermomechanical analyzer [TMA]



□ thermal analysis

- change in heat (or C_p)
- differential scanning calorimeter [DSC]
 - power-compensation type $\sim \Delta E$
 - heat-flux type $\sim \Delta T$
 - differential thermal analyzer [DTA], actually

Fig 17.33

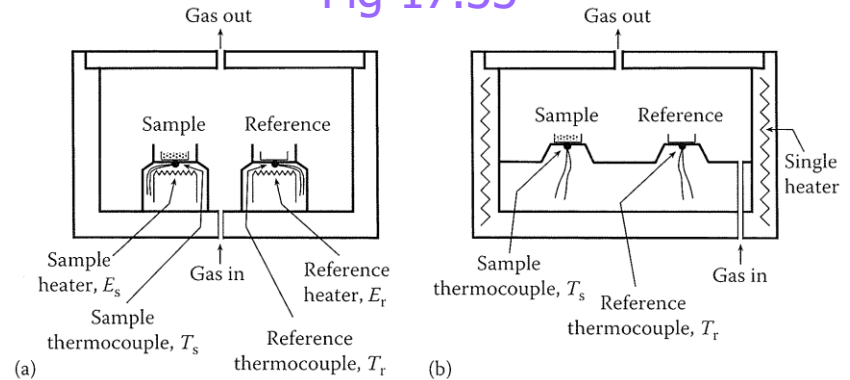
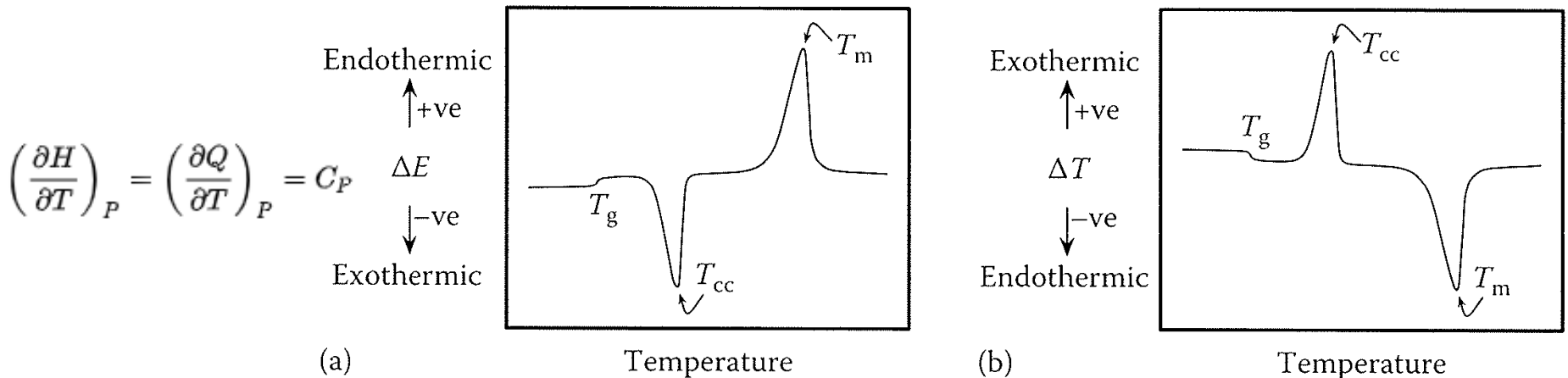


Fig 17.34



mechanical analysis

- mechanical damping $[\delta]$
- dynamic mechanical (thermal) analyzer [DM(T)A]

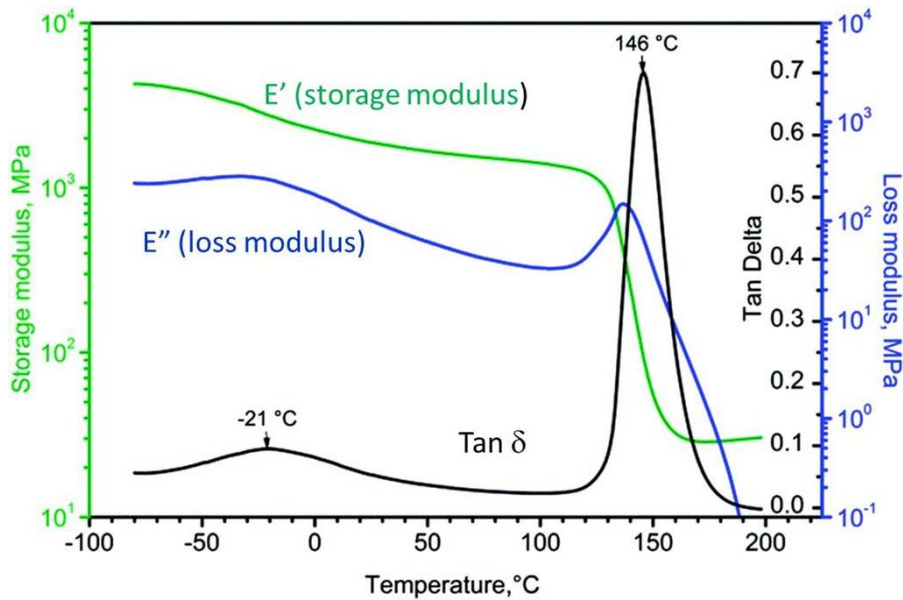
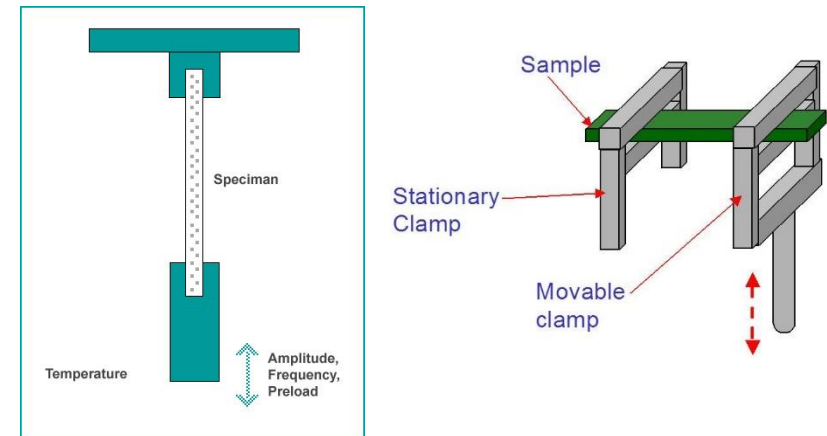
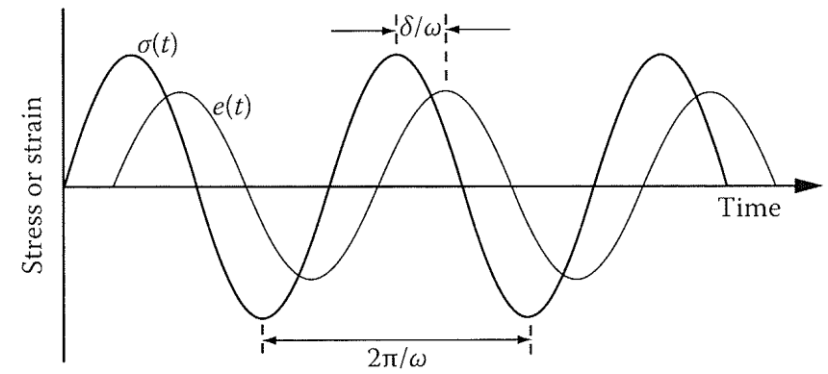
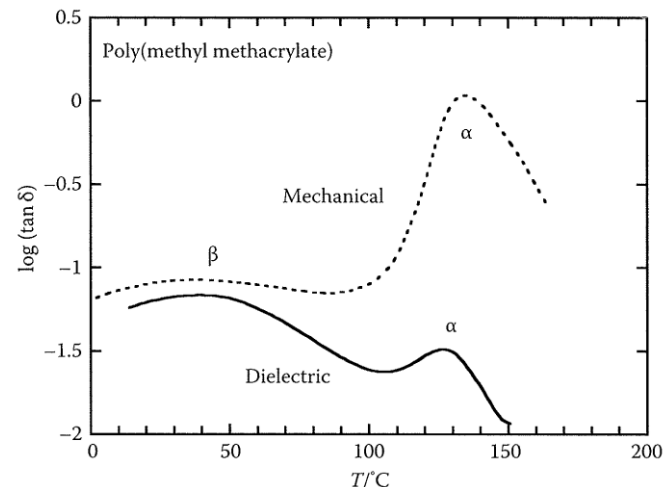
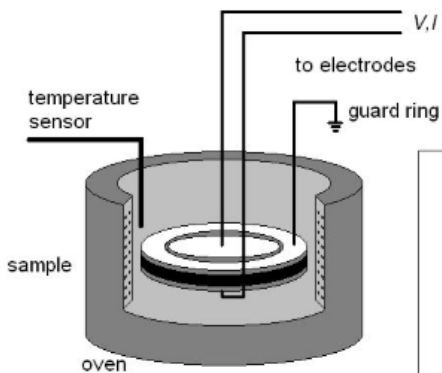


Fig 20.6



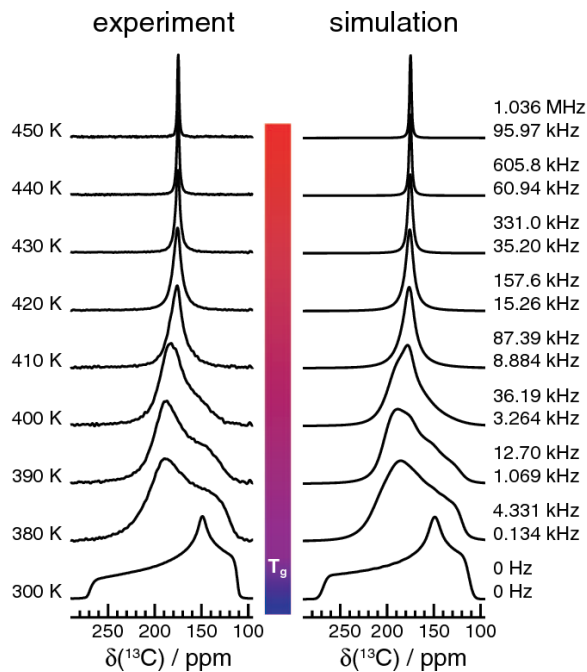
dielectric analysis

- ▣ electrical damping $[\delta]$
- ▣ dielectric thermal analyzer [DE(T)A]



NMR

- ▣ relaxation time $[\tau]$



Glass transition 2: Kinetic theory

Ch 16 sl 11

- T_g is rate-dependent.
 - cooling rate \sim dilatometry
 - $dT/dt \sim$ DSC
 - frequency \sim DMA
- Glass transition is a **pseudo-2nd-order** phase transition.

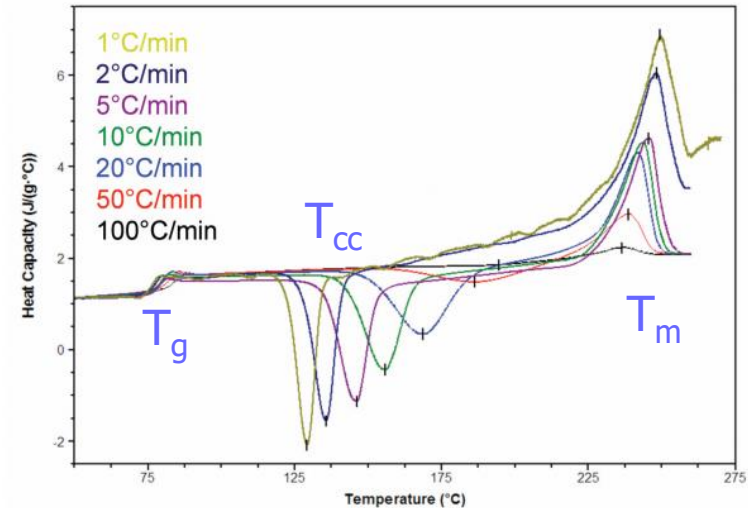
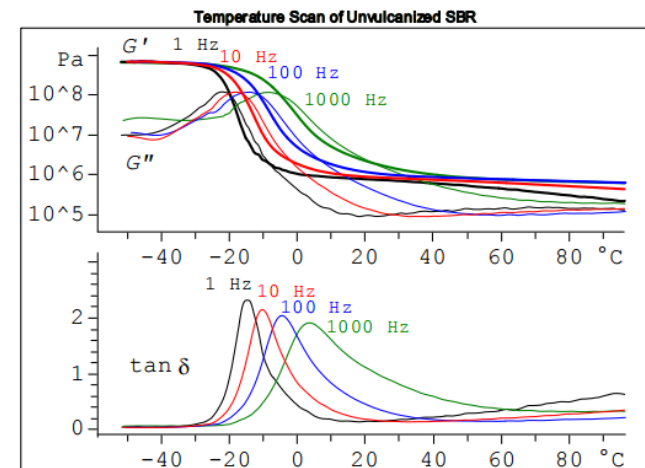


Figure 7: DSC thermal curves of PET at various scan rates, showing the glass transition, crystallisation and melting events.



□ Glass transition when $t = \tau$.

□ Deborah number $De = \tau/t$

- τ = material [relaxation] time; t = experimental [observation] time
- $De > 1$ for (elastic) solid; $De < 1$ for (viscous) liquid

□ Every material is viscoelastic.

- Deborah said “The mountains flowed before the Lord”

□ T_g when $De = 1$

□ exp't rate $\uparrow \rightarrow t \downarrow \rightarrow De \uparrow \rightarrow T_g \uparrow$

➤ glass? glassy?

➤ glassy vs rubbery

➤ glassy vs semicrystalline

➤ glassy polymer vs (semi)crystalline polymer

Glass transition 3: Free volume theory Ch 16 sl 13

- ❑ free volume
 - ❑ empty volume? expansion volume? fluctuation (sweep) volume?
 - ❑ Free volume is not a real physical volume.
 - Free volume is volume for molecular motion.
- ❑ Free volume expands at above T_g .
 - ❑ free volume $V_f = V_f^* + (T - T_g) \left(\frac{\partial V}{\partial T} \right)$
 - ❑ free vol fraction $f_V = f_g + (T - T_g) \alpha_f$
- ❑ Glassy state is an **iso-free-volume state**.
 - ❑ f_g is universal (2.5%?). Chapt 20

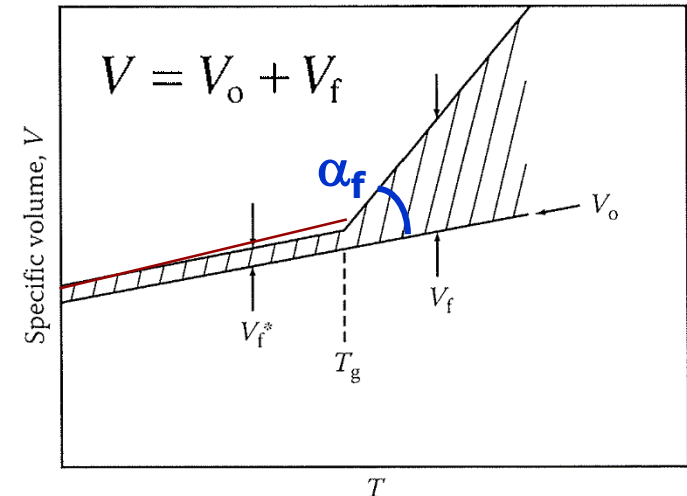
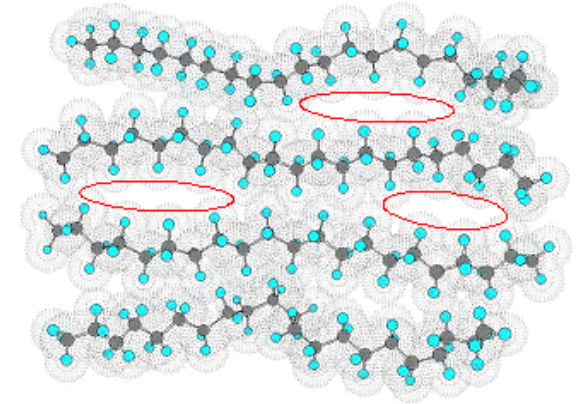


Fig 16.4

Factors controlling T_g

- (repeat unit) chemical structure
 - ▣ chain stiffness (**intra**molecular steric hindrance)
 - aromatic > aliphatic
 - substituent, branching > linear
 - # and size $\uparrow \rightarrow T_g \uparrow$

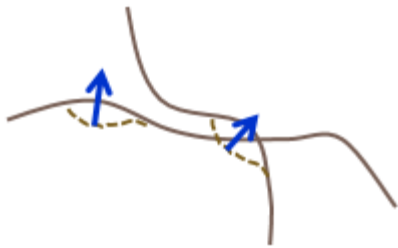


Table 16.1

p389
 polar groups restrict rotation
 more than non-polar subs?
Not really. inter rather than intra

TABLE 16.1
Approximate Values of Glass Transition Temperature, T_g , for Various Polymers

Repeat Unit	Polymer Name	$T_g/^\circ\text{C}$
$-\text{CH}_2-\text{CH}_2-$	Polyethylene	-130 to -10
$-\text{CH}_2-\text{CH}_2-\text{O}-$	poly(ethylene oxide)	-67
$-\text{H}_2\text{C}-\text{C}_6\text{H}_4-\text{CH}_2-$	poly(<i>p</i> -xylylene)	80
Vinyl polymers	Side group (X)	
$-\text{CH}_2-\text{CHX}-$	$-\text{CH}_3$	-23
	$-\text{CH}_2-\text{CH}_3$	-24
	$-\text{CH}_2-\text{CH}_2-\text{CH}_3$	-40
	$-\text{CH}_2-\text{CH}(\text{CH}_3)_2$	30
	$-\text{C}_6\text{H}_5$	100
	$-\text{Cl}$	81
	$-\text{OH}$	85
	$-\text{CN}$	97

□ chemical structure (cont'd)

□ intermolecular interactions [2ndary bondings]

■ London dispersion forces [VdW forces]

- subs size $\uparrow \rightarrow$ distance $\uparrow \rightarrow$ forces $\downarrow \rightarrow T_g \downarrow$
- compete with stiffening effect: PE < PP > PB > C3 -- < C8 < C9

■ dipole-dipole interactions

- subs dipole moment $\uparrow \rightarrow T_g \uparrow$: PE < PVC

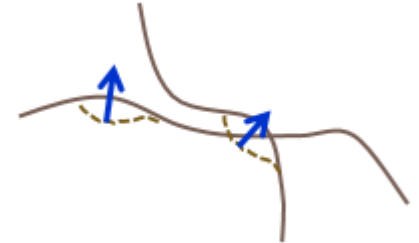
■ H-bonding

- polyamides, polyurethanes

□ intra- and intermolecular factors

$$■ T_g = \frac{2\delta^2}{mR} + C_1$$

- $\delta \sim$ solubility parameter
- $m \sim$ mobility of single chain



copolymers and blends

block or graft copolymers and immiscible blends

2 phase → 2 T_g

random copolymers and miscible blends

1 phase → 1 T_g

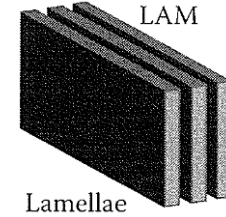
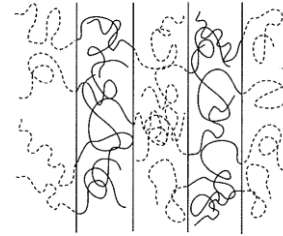
rule of mixture ~ upper bound

$$T_g^{cop} = T_g^A w_A + T_g^B w_B$$

lower bound ~ better fit to exp't

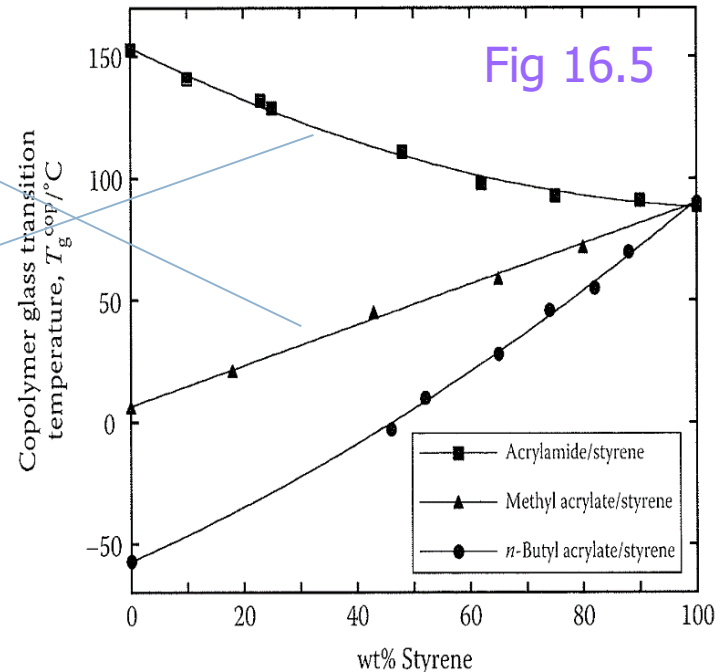
$$\frac{1}{T_g^{cop}} = \frac{w_A}{T_g^A} + \frac{w_B}{T_g^B}$$

from free volume? [p390](#) not necessary.



A-A-A-A-A-A-A-A-A-A-B-B-B-B-B-B-B-B-B-B

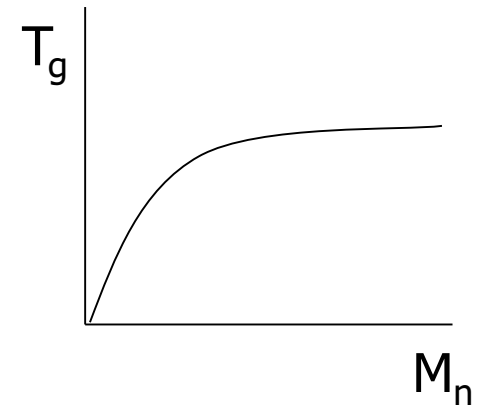
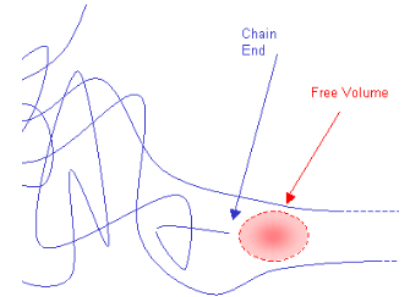
~B-B-B-A-B-B-A-B-A-A~



□ molar mass

$$□ T_g^\infty = T_{g0} + \frac{K}{M}$$

- MM ↑ → # of chain ends ↓ → free vol ↓ → T_g ↑
- $M = M_n$
- asymptotic
- from free volume, $K = 2\rho N_A \theta / \alpha_f$
 - not that quantitative ~ ignore



□ branching

- small # of branches lowers T_g
 - introduce free volume
- larger # of branches → higher T_g
 - stiffen chain like substituents

$$T_{g0} = T_g^\infty - \frac{\rho y N_A \theta}{\alpha_f \bar{M}_n}$$

y chain ends, $y-2$ branches
ignore also

□ crosslinking

- crosslinking density $\uparrow \rightarrow$ free vol $\downarrow \rightarrow T_g \uparrow$

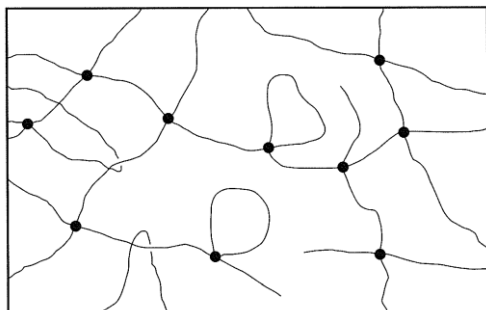


Fig 21.1

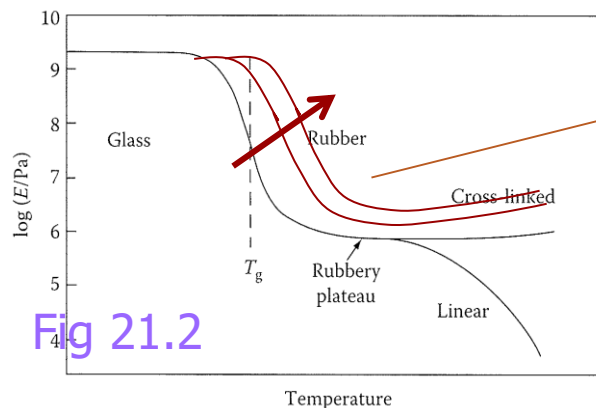


Fig 21.2

effect very much exaggerated

□ film thickness

- more mobility in the surface area
- effect only in (very) thin films

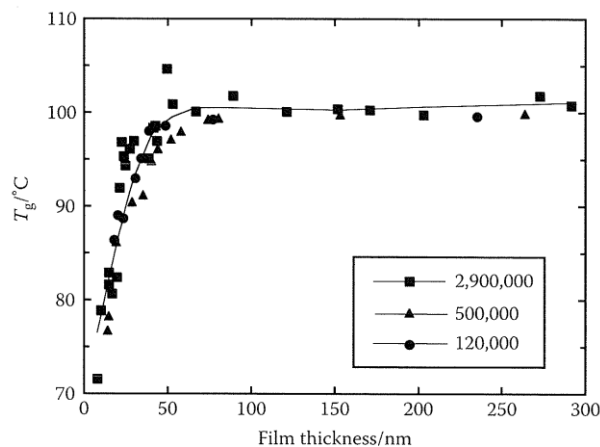
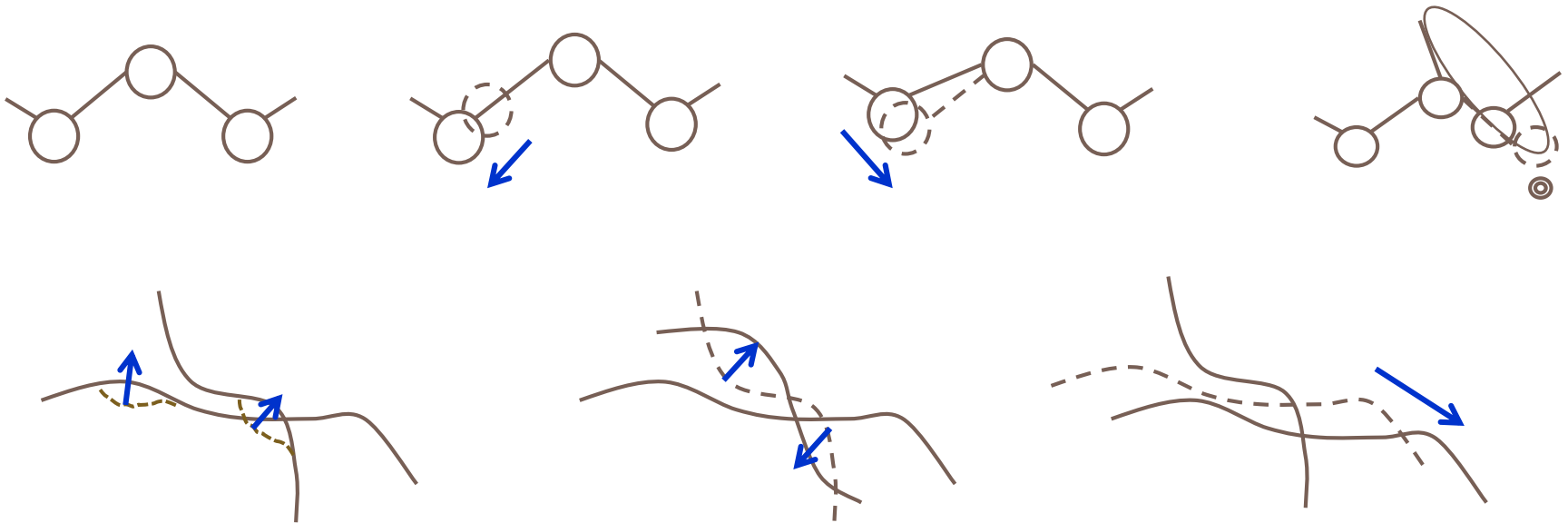


Fig 16.6

Macromolecular dynamics

Ch 16 sl 19

- motions of polymers
 - deformation of bond length and angle \sim local \sim *elastic*
 - change in conformation \sim segmental \sim *viscoelastic*
 - translation \sim whole chain \sim *viscous*



Models for dynamics

- Rouse(-Bueche-Zimm) theory
 - 'bead-and-spring'
 - models single chain in **dilute solution**
 - describes mechanical behavior like creep
 - describes melt viscosity for **low MM (before entanglement)** only $\rightarrow \eta \propto M$

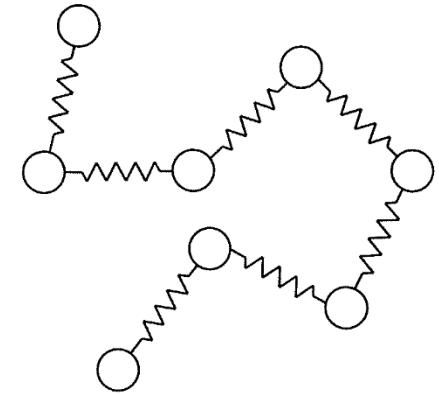
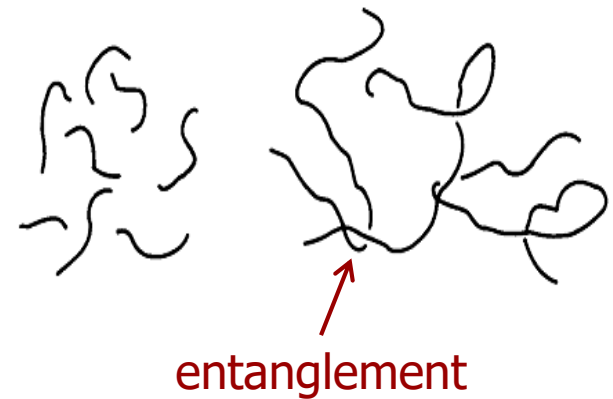
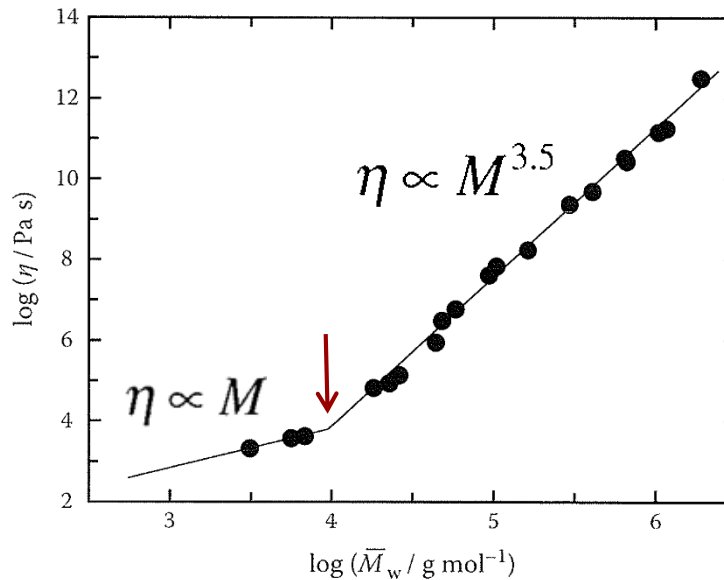


Fig 16.7



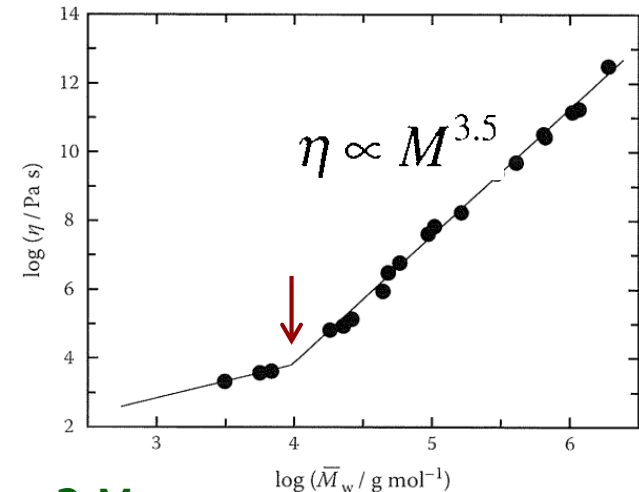
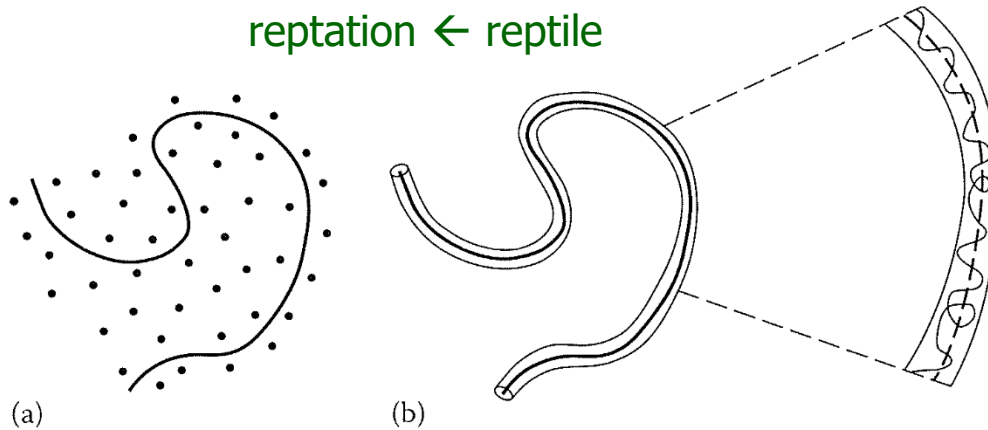
□ reptation theory (by de Gennes)

Nobel prize 1991

- considers **entanglement** [= interchain interaction]
 - for concentrated solution, bulk
 - 'chain in a tube'
 - time for motion of chain in tube $\propto M_w$
 - time for motion of tube $\propto L^2 \propto M_w^2$
- $\rightarrow \eta \propto M_w^3$



reptation ← reptile



$M_{cr} \approx 2 M_e$
 M_e from G_N^0

$M_{cr(itical)}$