
Chapter 21

Elastomers

Structure

Thermodynamics

Statistical approach

Mechanical behavior

Elastomer

- rubber (← eraser), ゴム (← gum), elastomer [彈性體]
- Requirements for rubber
 - (1) stretch to $> 100\%$ and (2) retract to original dimension instantly
 - flexible chain $T_g < \text{room temp}$
 - no crystalline phase EPR vs PE or PP
 - (lightly) crosslinked how lightly? segmental motion
 - chemically \sim vulcanization [加黃] by sulfur or peroxide
 - physically \sim TPE

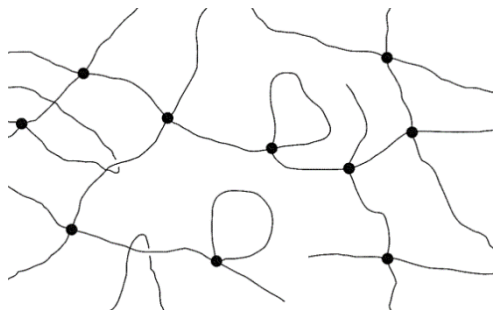
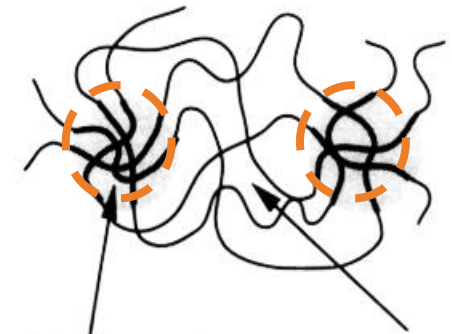
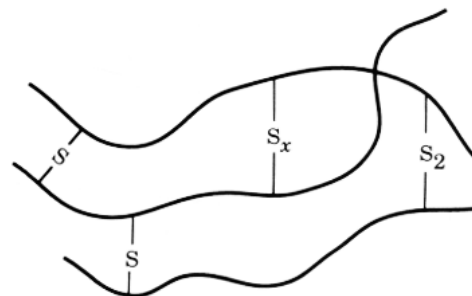


Fig 21.1 p513

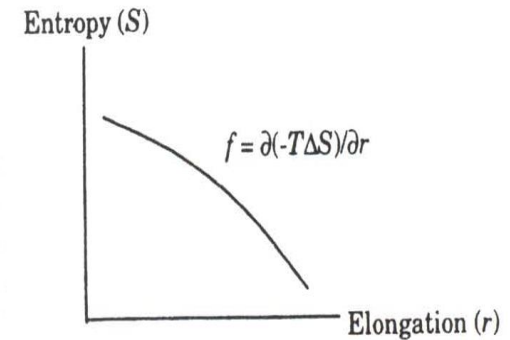
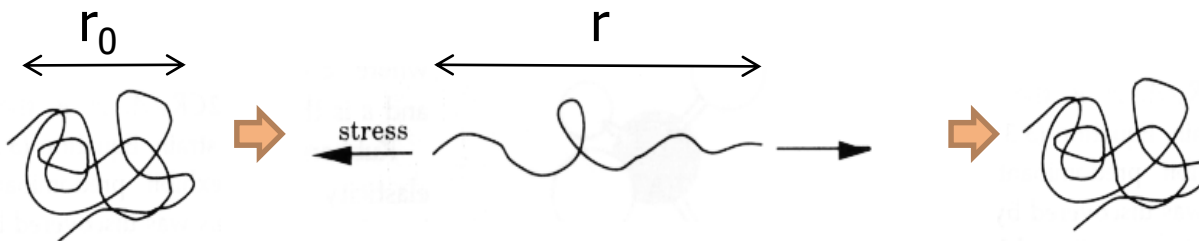


Hard domains with stiff segments

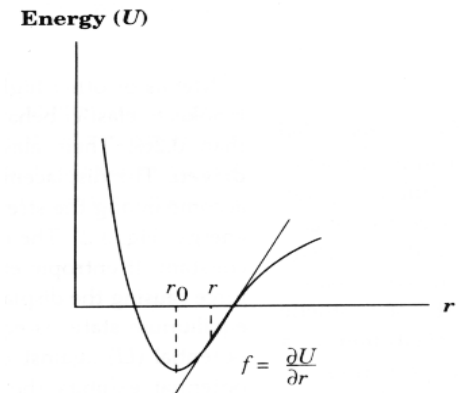
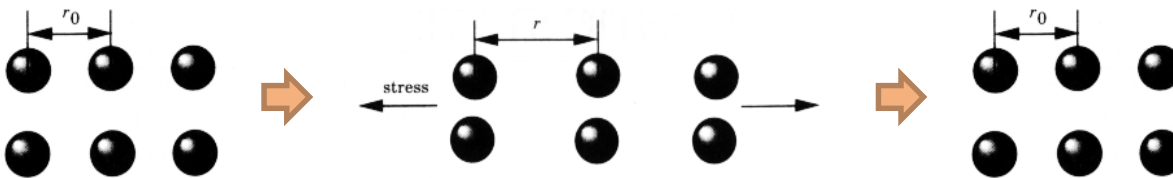
Soft domain with flexible segments

Rubber is an entropy spring.

- thermoelastic effect
 - Rubbers contract when heated, and
 - give out heat [get warm] when stretched.
- rubber spring \sim entropy-driven elasticity



- metal? energy-driven elasticity



(Classical) thermodynamics

- internal energy U

$$dU = \delta Q - \delta W \xrightarrow[\delta W = -fdl]{\delta Q = TdS} fdl = dU - TdS$$

- Helmoltz free energy \leftarrow constant volume $\leftarrow v = 0.5$ for rubbers

$$A = U - TS \rightarrow dA = dU - TdS \rightarrow fdl = dA \text{ (at constant } T\text{)}$$

- retractive force f

$$f = \left(\frac{\partial A}{\partial l} \right)_T = \left(\frac{\partial U}{\partial l} \right)_T - T \left(\frac{\partial S}{\partial l} \right)_T$$

energetic [f_e]
0 or small

entropic [f_s]
dominant

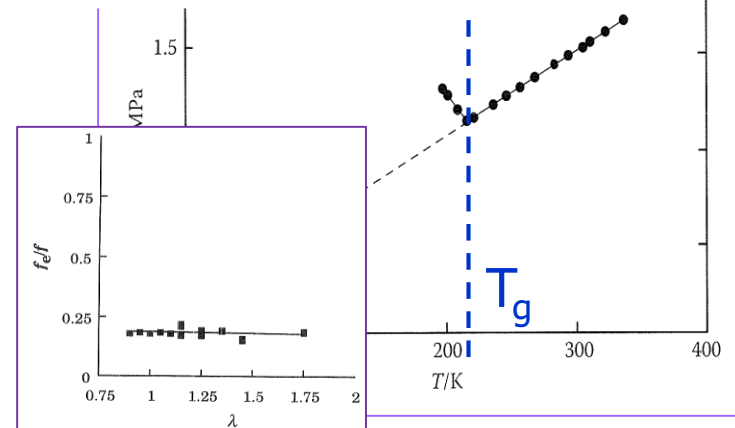
eqn (21.9-14)

$$f = \left(\frac{\partial U}{\partial l} \right)_T + T \left(\frac{\partial f}{\partial T} \right)_l$$

Fig 21.3

thermodynamic 'eqn of state' for rubber elasticity

$dU \approx 0 \rightarrow dW = dQ$
adiabatic stretching \rightarrow get warm



Statistical theory

- ΔS with deformation of one chain from (x,y,z) to (x',y',z')

$$W(x,y,z) = \left(\frac{\beta}{\pi^{1/2}} \right)^3 \exp(-\beta^2 r^2)$$
$$\beta = (3/(2nl^2))^{1/2}$$

$$S = k \ln \Omega$$

$$S = c - k\beta^2 r^2$$

$$S = c - k\beta^2 (x^2 + y^2 + z^2)$$

$$x' = \lambda_1 x, \quad y' = \lambda_2 y, \quad z' = \lambda_3 z$$

$$S' = c - k\beta^2 (\lambda_1^2 x^2 + \lambda_2^2 y^2 + \lambda_3^2 z^2)$$

$$\Delta S_i = S' - S = -k\beta^2 [(\lambda_1^2 - 1)x^2 + (\lambda_2^2 - 1)y^2 + (\lambda_3^2 - 1)z^2]$$

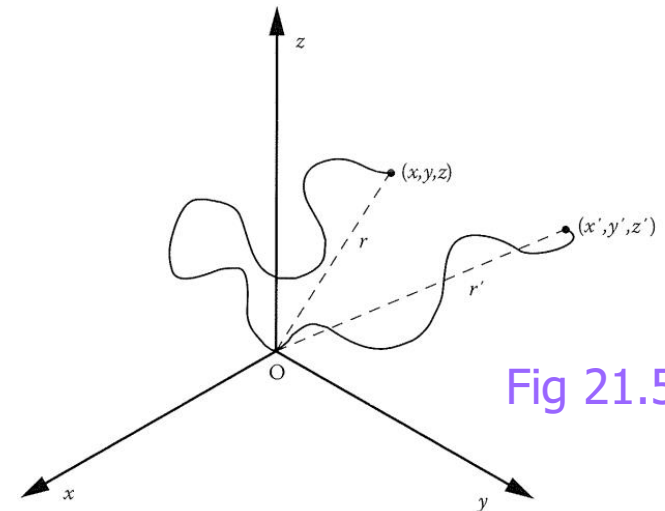
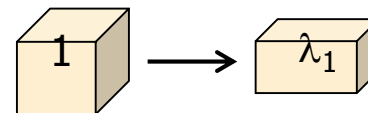


Fig 21.5



$\lambda = \text{extension ratio} = L/L_0$

□ ΔS of N chains (per unit volume)

$$dN = NW(x, y, z)dx dy dz$$

$$dN = N \left(\frac{\beta}{\pi^{1/2}} \right)^3 \exp[-\beta^2(x^2 + y^2 + z^2)] dx dy dz$$

$$\Delta S = \int \Delta S_i dN = -\frac{1}{2} Nk(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 - 3)$$

□ work of deformation $w = \frac{1}{2} G(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 - 3)$

$$dU = dQ - dW \xrightarrow{dU = 0} dw = (-)T\Delta S$$

$$\rho = \frac{NM_c}{N_A} \rightarrow N = \frac{\rho N_A}{M_c} = \frac{\rho R}{M_c k}$$

$$G = \frac{\rho RT}{M_c}$$

Fig 21.1

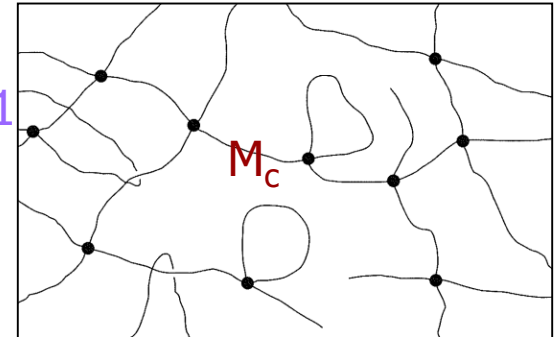
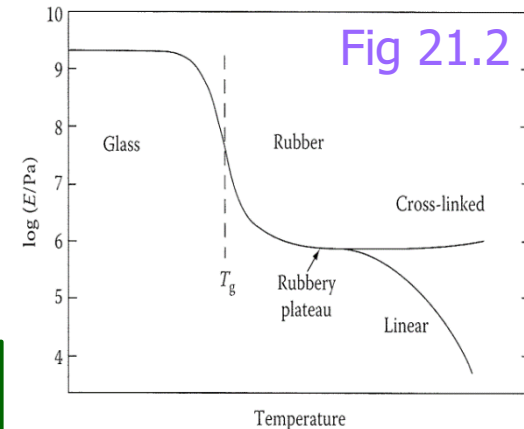


Fig 21.2



□ $T \uparrow \rightarrow$ modulus $\uparrow \sim$ another characteristic of elastomer

□ for linear polymers, $M_e = \frac{\rho RT}{G_N^o}$

Stress-strain behavior

$$w = \frac{1}{2} G(\lambda_1^2 + \lambda_2^2 + \lambda_3^2 - 3) = \frac{1}{2} G \left(\lambda^2 + \frac{2}{\lambda} - 3 \right)$$

↑

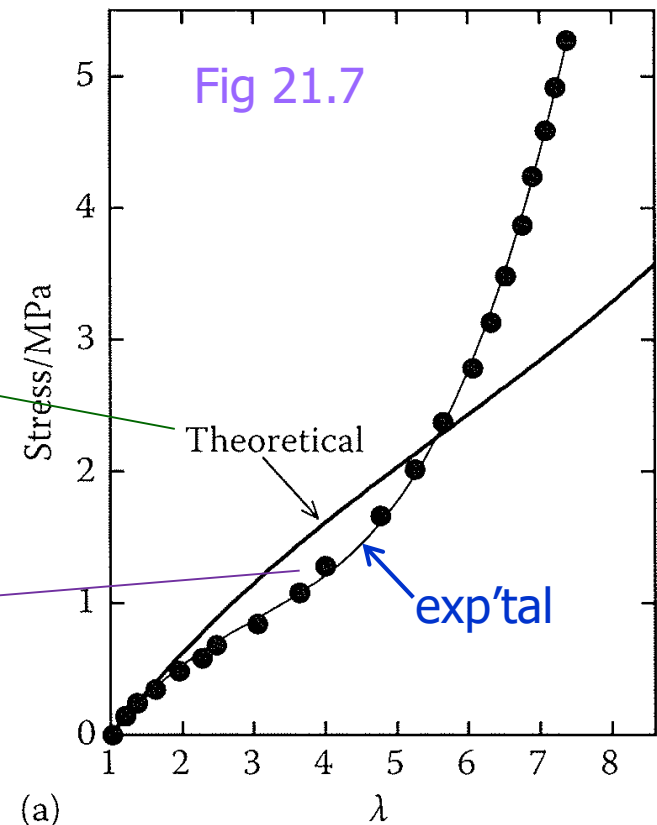
$$\lambda_1 = \lambda \text{ and } \lambda_1 \lambda_2 \lambda_3 = 1$$

$$\sigma_n = \frac{dw}{d\lambda} = G \left(\lambda - \frac{1}{\lambda^2} \right)$$

statistical mechanical equation for rubber elasticity

comparison with exp't

- at low λ , $\sigma(\text{theo}) > \sigma(\text{exp't})$
 - affine deformation vs phantom network
- at high λ , $\sigma(\text{theo}) < \sigma(\text{exp't})$
 - Gaussian to non-Gaussian chain segments
 - stressed, extended, fewer conformations
 - strain-induced crystallization



Chapter Extra 1-1

Polymer Rheology

Shear and elongational viscosity

Normal stress difference

Rheometry

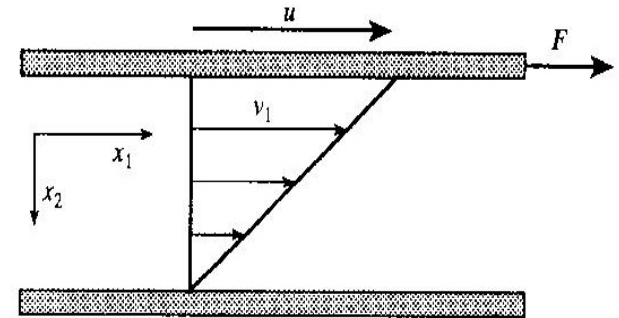
Rheology

□ rheology [流變學]

- study of flow and deformation of fluids [liquids]
- constitutive [stress-strain] relation of fluids

□ shear flow

- shear stress $\tau = F/A$ [$\text{N/m}^2 = \text{Pa}$]
- shear strain $\gamma = dx_1/dx_2$
- shear (strain) rate $d\gamma/dt$ [s^{-1}]
 - $dv_1 = dx_1/dt$
 - $d\gamma/dt = dv_1/dx_2 \sim$ velocity gradient



□ constitutive eqn: **Newton's law:** $\tau = \eta_{(s)} (d\gamma/dt)$

□ (shear) viscosity $\eta_{(s)}$

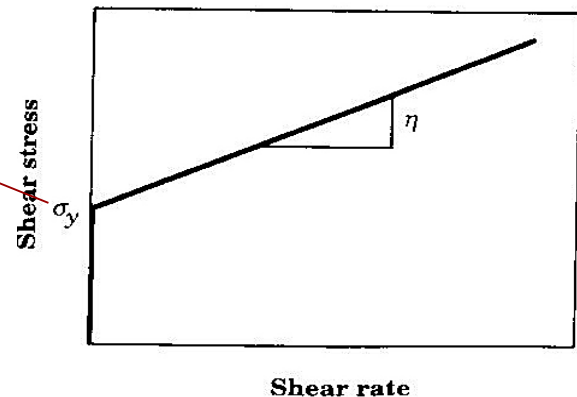
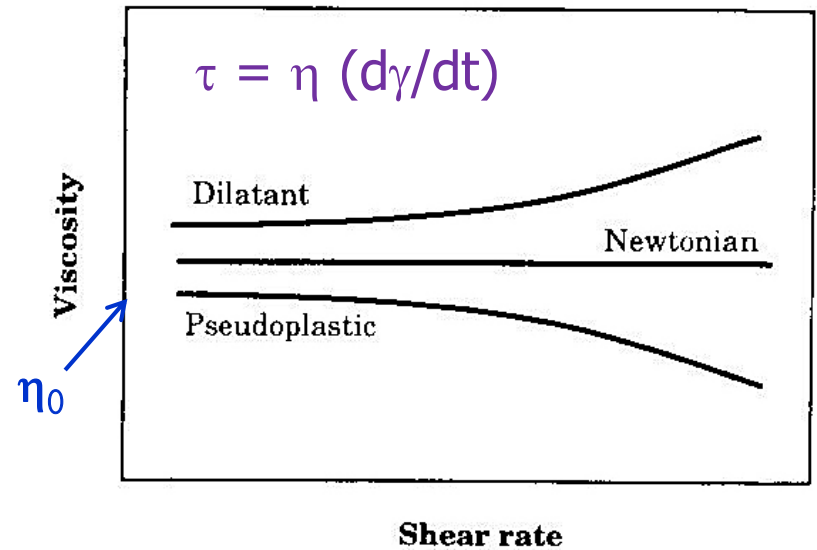
- η of water $\sim 10^{-3} \text{ Pa s} = 1 \text{ cP}$
- η of polymer melt $\sim 10^3 \text{ Pa s}$

$$1 \text{ Pa s} = 10 \text{ P(oise)}$$

Shear viscosity

- ❑ Temperature dependence of η
 - ❑ WLF eqn \sim for $T_g - 50 < T < T_g + 100$ K
 - ❑ Arrhenius-type relation \sim for $T > T_g + 100$ K
 - $\eta = A \exp[-B/T]$
 - $T \uparrow \rightarrow \eta \downarrow$
 - not for large ΔT (> 20 K)
- ❑ Pressure dependence of η
 - ❑ $\eta = A \exp[B/P]$
 - ❑ $p \uparrow \rightarrow \eta \uparrow$
 - ❑ inter-particle distance
- increasing P and decreasing T the same effect
 - 100 atm is equiv to 30–50 K

- Shear rate dependence of η
 - Newtonian \sim constant viscosity
 - many solutions and melts
 - non-Newtonian
 - dilatant = shear-thickening
 - suspensions
 - pseudoplastic = shear-thinning
 - polymer melts
 - chains aligned to shear direction
 - zero-shear(-rate) viscosity η_0
 - Bingham plastic \sim yielding
 - slurries, margarine



□ molecular weight dependence of η

□ $\eta = K M_w^{1.0}$ for $M < M_c$

$\eta = K M_w^{3.4}$ for $M > M_c$

□ $M_c = 2 - 3 M_e$

□ $M_e \sim$ entanglement MM

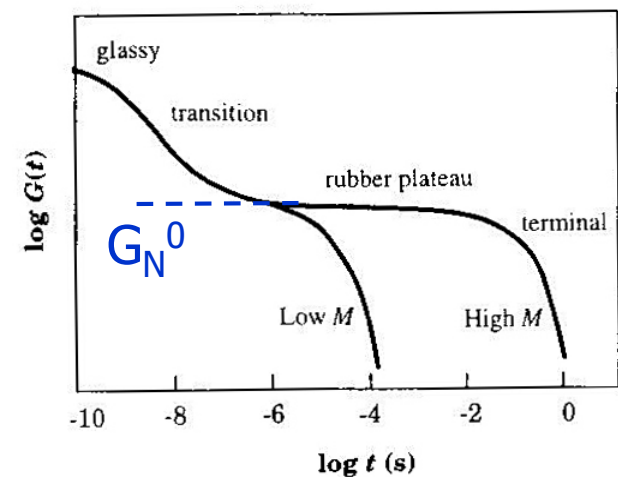
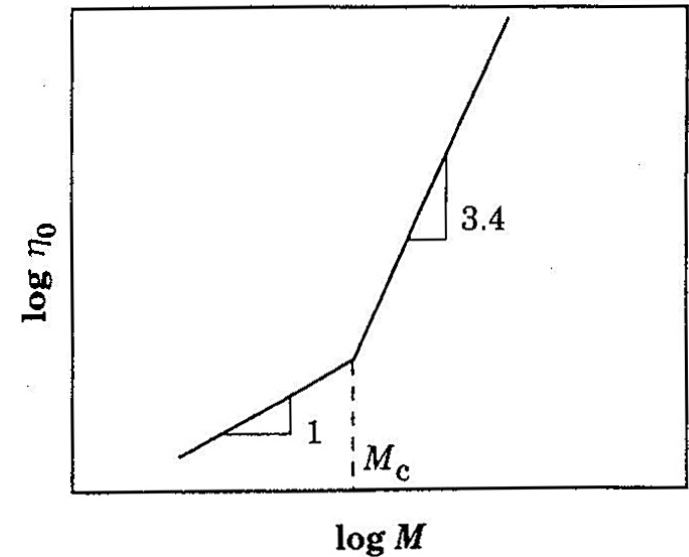
■ $M_e = \rho RT / G_N^0$

■ M_e depends on structure of chain

■ chain stiffness and interactions

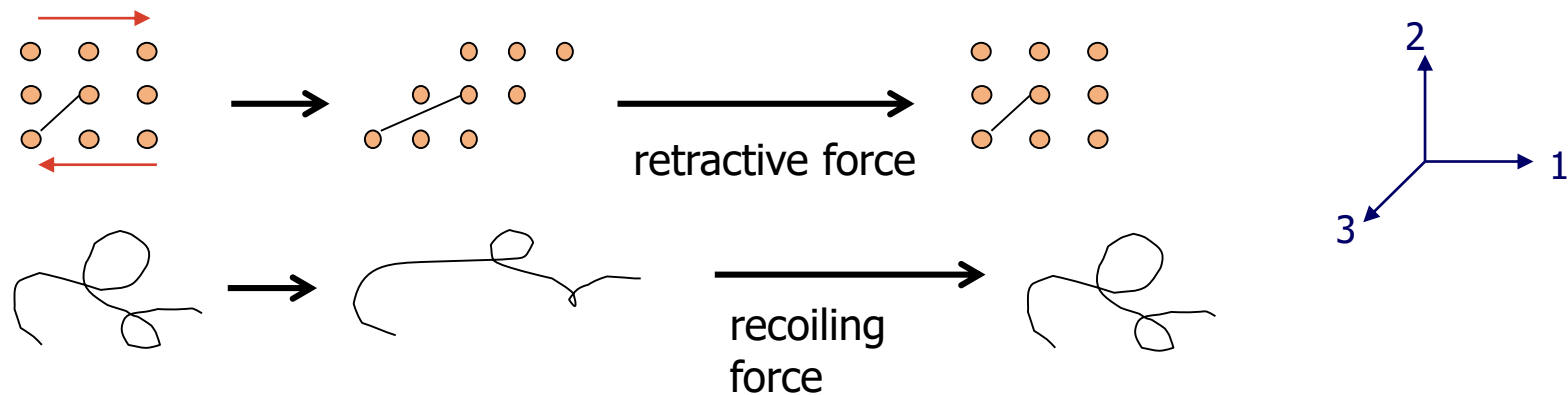
■ PE \sim 1200, PS \sim 20000, PC \sim 2500

■ $M_e(\text{step polymers}) < M_e(\text{chain polymers})$



Normal stress difference

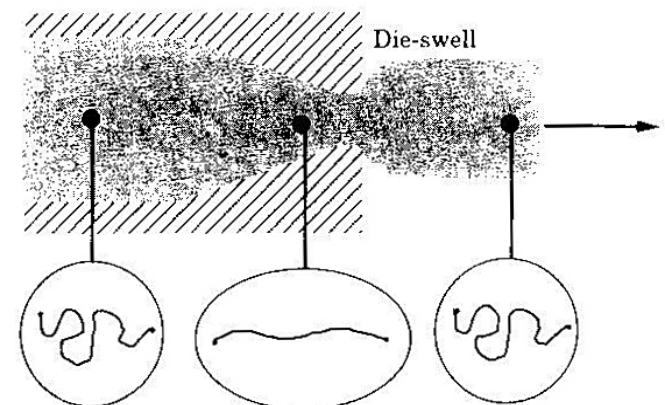
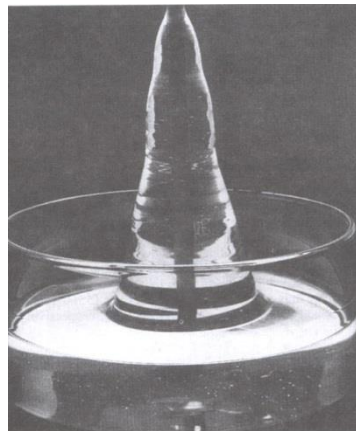
- normal stress **caused by** shear flow



- $\sigma_1 - \sigma_2 = N_1 > 0 \sim$ 1st normal stress difference
- $\sigma_2 - \sigma_3 = N_2 \approx 0 \sim$ 2nd normal stress difference

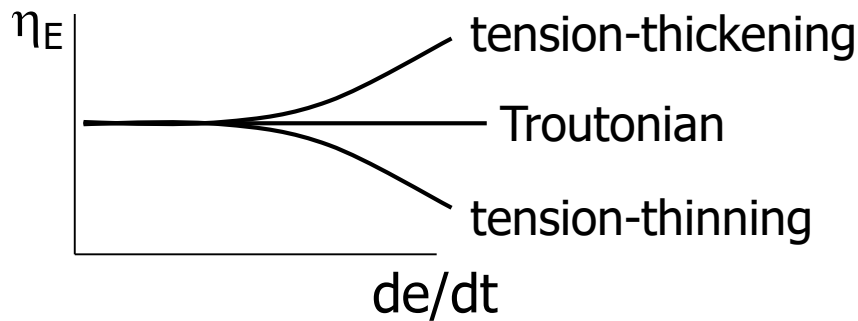
- results of NSD

- Weissenberg effect
 - rod-climbing
- die swell

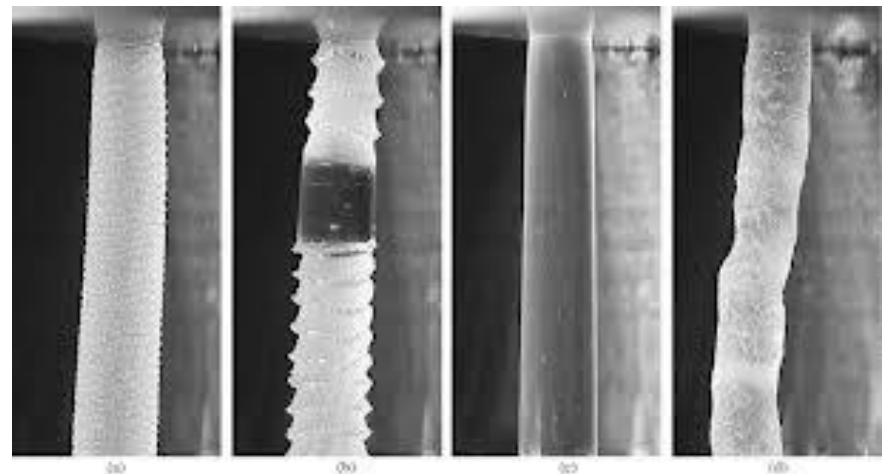
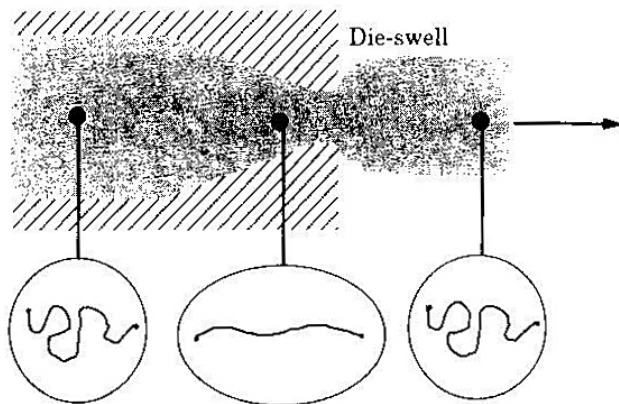


Elongational viscosity

- $\sigma = \eta_E (de/dt)$
 - $\eta_E \sim$ elongational [tensile] viscosity



- η_E determines melt strength [stability of melt]



Rheometry

□ Shear flow

□ rotational rheometers ~ drag

- cylinder, Brookfield
- parallel-plate rheometer
- cone-and-plate rheometer

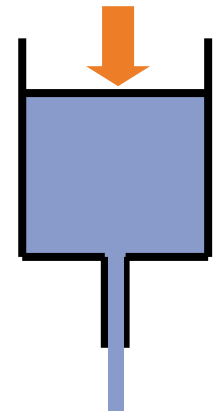
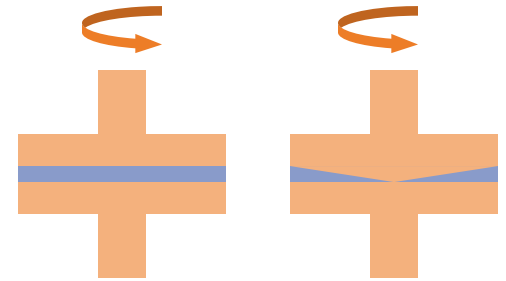
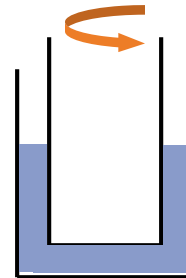
□ capillary or slit rheometer ~ pressure drop

➤ dynamic rheometry

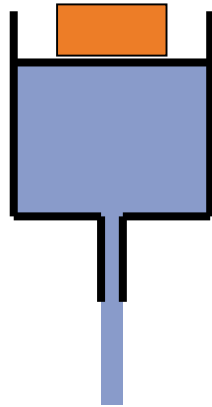
- $\eta' = G''/\omega \sim$ dynamic viscosity
- $\eta'' = G'/\omega$
- $\eta^* = \eta' - i\eta'' \sim$ complex viscosity

□ elongational flow

- difficult to perform expt
- possible only for very small strain rate



- ❑ melt index (MI) or melt flow index (MFI)
 - ❑ melt indexer ~ a simple capillary viscometer



- ❑ $M(F)I = \text{g of resin}/10 \text{ min}$
 - at specified weight and temperature
 - high MI = low η = low MM of a polymer

