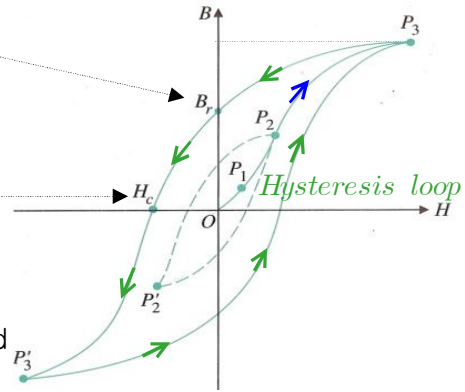


4) Remnant flux density B_r :

The residual magnetic flux density that does not go to zero after the applied magnetic field is removed.

Coercive field intensity H_c :

The magnetic field intensity applied in the opposite direction to make the magnetic flux density of a magnetized medium vanish.



5) In a uniform magnetic field, $\mathbf{F}_m = I \oint d\mathbf{l} \times \mathbf{B} = \mathbf{0}$

$$\mathbf{T} = \mathbf{m} \times \mathbf{B} = \pi a^2 \hat{\mathbf{l}} \times \mathbf{B}$$

6) From $\nabla \cdot \mathbf{B} = 0$, vector magnetic potential: $\mathbf{B} = \nabla \times \mathbf{A}$ ①

① in Faraday's law $\nabla \times \mathbf{E} = -\partial \mathbf{B} / \partial t$:

$$\nabla \times \left(\mathbf{E} + \frac{\partial \mathbf{A}}{\partial t} \right) = \mathbf{0} \Rightarrow \mathbf{E} + \frac{\partial \mathbf{A}}{\partial t} = -\nabla V$$

$$\Rightarrow \mathbf{E} = -\nabla V - \frac{\partial \mathbf{A}}{\partial t} \equiv \mathbf{E}_V + \mathbf{E}_A \quad (\text{V/m}) \quad ②$$

② in Gauss's law $\nabla \cdot \mathbf{D} = \rho_v$:

$$\nabla^2 V = -\frac{\rho_v}{\epsilon} - \frac{\partial}{\partial t} (\nabla \cdot \mathbf{A}) \quad ③$$

①, ② in Ampere's law $\nabla \times \mathbf{H} = \mathbf{J} + \frac{\partial \mathbf{D}}{\partial t}$ using $\nabla \times (\nabla \times \mathbf{A}) = \nabla (\nabla \cdot \mathbf{A}) - \nabla^2 \mathbf{A}$:

$$\nabla^2 \mathbf{A} - \mu \epsilon \frac{\partial^2 \mathbf{A}}{\partial t^2} = -\mu \mathbf{J} + \nabla (\nabla \cdot \mathbf{A} + \mu \epsilon \frac{\partial V}{\partial t}) \quad ④$$

Using Lorentz condition, $\nabla \cdot \mathbf{A} + \mu \epsilon \frac{\partial V}{\partial t} = 0$ by Lorentz gauge transformations, ③ and ④ are simplified as

$$\left(\nabla^2 - \mu \epsilon \frac{\partial^2}{\partial t^2} \right) \begin{Bmatrix} V \\ \mathbf{A} \end{Bmatrix} = \begin{Bmatrix} -\rho_v / \epsilon \\ -\mu \mathbf{J} \end{Bmatrix} : \text{Wave equation}$$

3. 1) Axially symmetric ($\partial / \partial \phi = 0$) and No edge effect ($\partial / \partial z = 0$)

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \mu_o I \Rightarrow B(r) 2\pi r = \mu_o I \Rightarrow \mathbf{B}(r) = \hat{\phi} \frac{\mu_o I}{2\pi r}$$

$$2) \nabla^2 = \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2}{\partial \phi^2} + \frac{\partial^2}{\partial z^2} = \frac{1}{r} \frac{d}{dr} \left(r \frac{\partial}{\partial r} \right)$$

$\nabla^2 \mathbf{A} = -\mu_o \mathbf{J}$: vector Poisson's equation

BVP in a current-free ($\mathbf{J} = \mathbf{0}$) region ($a < r < b$) $\Rightarrow \nabla^2 \mathbf{A} = \mathbf{0}$

$$\frac{d}{dr} \left(r \frac{dA_z}{dr} \right) = 0 \quad ①$$

BCs:

$$A_z(r)|_{r=b} = 0 \quad (2)$$

$$\oint_C \mathbf{B} \cdot d\mathbf{l} = \mu_o I \Rightarrow \oint_C (\nabla \times \mathbf{A}) \cdot d\mathbf{l} = \mu_o I \Rightarrow \oint_C \left(-\frac{\partial A_z}{\partial r} \hat{\phi} \right) \cdot (\hat{\phi} r d\phi) = \mu_o I$$

$$\Rightarrow - \int_0^{2\pi} \left(\frac{\partial A_z}{\partial r} \right) (r d\phi) = \mu_o I \quad (3)$$

Integrating (1) twice, $r \frac{dA_z}{dr} = C_1 \Rightarrow dA_z = C_1 \frac{dr}{r} \Rightarrow A_z(r) = C_1 \ln r + C_2 \quad (4)$

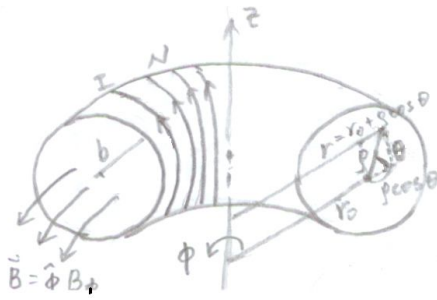
(2) in (4): $C_2 = C_1 \ln b \quad (5), \quad (5) \text{ in } (4): A_z(r) = C_1 \ln \frac{r}{b} \quad (6)$

(6) in (3): $-\int_0^{2\pi} \left(\frac{C_1}{r} \right) (r d\phi) = \mu_o I \Rightarrow -2\pi C_1 = \mu_o I \Rightarrow C_1 = -\mu_o I / 2\pi \quad (7)$

(7) in (6): $A_z(r) = \hat{z} \frac{\mu_o I}{2\pi} \ln \frac{b}{r}$

$$\therefore \mathbf{B}(r) = \nabla \times \mathbf{A} = -\hat{\phi} \frac{dA_z}{dr} = \hat{\phi} \frac{\mu_o I}{2\pi} \left(\frac{1}{r} \right), \quad (a < r < b)$$

4. 1) Ampere's circuital law: $\oint_C \mathbf{B} \cdot d\mathbf{l} = \mu_o NI, \quad (b-a) < r < (b+a)$



$$\Rightarrow \mathbf{B}(r) = \hat{\phi} \frac{\mu_o NI}{2\pi} \frac{1}{r}, \quad (b-a) < r < (b+a)$$

2) $\Phi = \int \mathbf{B} \cdot d\mathbf{s} = \int_0^b \int_0^{2\pi} B_\phi \rho d\theta d\rho$

$$= \frac{\mu_o NI}{2\pi} \left[\int_0^b \left(\int_0^{2\pi} \frac{d\theta}{r_o + \rho \cos \theta} \right) \rho d\rho \right]$$

$$= \frac{\mu_o NI}{2\pi} \int_0^b \frac{2\pi}{\sqrt{r_o^2 + \rho^2}} \rho d\rho$$

$$= -\mu_o NI \sqrt{r_o^2 - \rho^2} \Big|_0^b = \mu_o NI (r_o - \sqrt{r_o^2 - b^2})$$

3) $\therefore L = \frac{N\Phi}{I} = \mu_o N^2 (r_o - \sqrt{r_o^2 - b^2})$

5. 1) Using Ampere's circuital law, $\mathbf{B}_1(y) = \hat{x} \frac{\mu_o I}{2\pi y}$

2) motional emf $= \oint (\mathbf{u} \times \mathbf{B}_1) \cdot d\mathbf{l}$

$$= \frac{\mu_o I_o u_o}{2\pi} \left[\int_{\downarrow} \left(\hat{y} \times \frac{\hat{x}}{d} \right) \cdot (-\hat{z} dl) + \int_{\uparrow} \left(\hat{y} \times \frac{\hat{x}}{d+w} \right) \cdot (\hat{z} dl) \right]$$

$$= \frac{\mu_o I_o u_o h}{2\pi} \left(\frac{1}{d} - \frac{1}{d+w} \right) = \frac{\mu_o I_o u_o h w}{2\pi d(d+w)}$$

3) $i_2 = -\frac{emf}{R} = -\frac{\mu_o I_o u_o h w}{2\pi d(d+w)R}$ flowing in the counter-clock-wise direction

to increase Φ in loop to oppose the decrease in Φ due to loop movement.