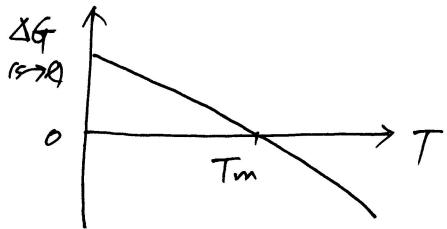


1. (a)



$$\left(\frac{\partial \Delta G}{\partial T}\right) = -\Delta S < 0, \text{ at } T=T_m, \Delta G=0.$$

$$\left(\frac{\partial \Delta G}{\partial T^2}\right) = -\frac{\Delta H}{T^2} < 0, \quad T_m = 0^\circ C \text{ at } 1 \text{ atm.}$$

at  $T < T_m \rightarrow$  solid ~~stable~~ is stable

at  $T > T_m \rightarrow$  liquid is stable.

(b) condensed phase  $\rightarrow$  vapor phase.  $\Delta V = V_v - V_c \approx V_v$ .

from Clapeyron equation,  ~~$\left(\frac{\partial P}{\partial T}\right)_V = \frac{\Delta H}{T \Delta V}$~~

$$\frac{dP}{dT} = \frac{\Delta H}{TV_v} = \frac{\Delta H}{RT^2} P.$$

rearranging gives,

$$\frac{dP}{P} = \frac{\Delta H}{RT^2} dT. \quad \therefore d\ln P = \frac{\Delta H}{RT^2} dT.$$

If  $\Delta H$  is independent of temperature,  $c_{p(v)} = c_{p(c)}$ , integration gives

$$\ln P = -\frac{\Delta H}{RT} + \text{constant.},$$

↳ explains an exponential increase in the saturation vapor pressure exerted by a condensed phase with increasing temperature.

(c)

$$(i) \ln p(\text{atm}) = -\frac{15180}{T} - 0.155 \ln T + 19.25$$

$$(ii) \ln p(\text{atm}) = -\frac{15250}{T} - 1.255 \ln T + 21.19$$

$$\text{for (i), } \frac{d\ln p}{dT} = \frac{15180}{T^2} - \frac{0.155}{T} = \frac{\Delta H_{(i)}}{RT^2}$$

$$\therefore \Delta H_{(i)} = R(15180 - 0.155T).$$

$$\text{for (ii), } \frac{d\ln p}{dT} = \frac{15250}{T^2} - \frac{1.255}{T} = \frac{\Delta H_{(ii)}}{RT^2}$$

$$\therefore \Delta H_{(ii)} = R(15250 - 1.255T).$$

since  $\Delta H_{\text{sub}} > \Delta H_{\text{evap}} \Leftrightarrow \Delta H_{(i)} > \Delta H_{(ii)}$

$\therefore (i)$  is for solid

2.(a). when the solute B obeys Henry's law.  $\rightarrow a_B = k_B X_B$ .

$$\rightarrow \ln a_B = \ln k_B + \ln X_B$$

$$d\ln a_B = d\ln X_B$$

by using Gibbs-Duhem equation,  $X_A d\ln a_A + X_B d\ln a_B = 0$ .

$$d\ln a_A = -\frac{X_B}{X_A} d\ln a_B = -\frac{X_B}{X_A} d\ln X_B.$$

$$= -\frac{X_B}{X_A} \frac{dX_B}{X_B} = -\frac{dX_B}{X_A}.$$

$$\text{and, } -\frac{dX_B}{X_A} = \frac{dX_A}{X_A} = d\ln X_A \quad \leftarrow dX_A + dX_B = 0.$$

$$\therefore d\ln a_A = d\ln X_A$$

$$\ln a_A = \ln X_A + \text{constant}.$$

$$a_A = \text{constant} \times X_A.$$

$$= X_A. \quad (X_A = 1, a_A = 1. \quad \therefore \text{constant} = 1.)$$

(b)

$$\begin{aligned} \Delta G'^M &= (n_A \bar{G}_A + n_B \bar{G}_B) - (n_A G_A^\circ + n_B G_B^\circ) \\ &= n_A (\bar{G}_A - G_A^\circ) + n_B (\bar{G}_B - G_B^\circ) \\ &= n_A \Delta \bar{G}_A^M + n_B \Delta \bar{G}_B^M \\ &= RT (n_A \ln a_A + n_B \ln a_B). \end{aligned}$$

when ideal,  $a_A = X_A$ ,  $a_B = X_B$ , 1 mole of solution,

$$\therefore \Delta G^{M,\text{id}} = RT (X_A \ln X_A + X_B \ln X_B). //$$

from the Gibbs-Helmholtz equation,

$$\left[ \frac{\partial(\Delta \bar{G}_i^M/T)}{\partial T} \right]_{P,\text{comp}} = -\frac{\Delta \bar{H}_i^M}{T^2}, \quad \Delta \bar{G}_i^{\text{M,id}} = RT \ln X_i$$

$$\therefore \frac{d(R \ln X_i)}{dT} = -\frac{\Delta \bar{H}_i^M}{T^2}, \quad \text{and. as } X_i \text{ is independent of } T.$$

$$\therefore \Delta \bar{H}_i^{\text{M,id}} = 0. \quad \therefore \Delta H^{\text{M,id}} = 0. \quad (\Delta H^M = n_A \Delta \bar{H}_A^M + n_B \Delta \bar{H}_B^M).$$

$$\left(\frac{\partial \Delta G^M}{\partial T}\right)_{P, \text{conf}} = -\Delta S^M$$

$$\Delta S^{M,id} = -R(X_A \ln X_A + X_B \ln X_B),$$

(c). mixing of  $N_A$  particles of A with  $N_B$  particles of B,

$$\Delta S'_{\text{conf}} = k \ln \frac{(N_A + N_B)!}{N_A! N_B!}$$

$$= k[\ln(N_A + N_B)! - \ln N_A! - \ln N_B!]$$

Application of Stirling's theorem gives,

$$\Delta S'_{\text{conf}} = -k \left[ N_A \ln \left( \frac{N_A}{N_A + N_B} \right) + N_B \ln \left( \frac{N_B}{N_A + N_B} \right) \right].$$

$$\text{Now, } \frac{N_A}{N_A + N_B} = \frac{n_A}{n_A + n_B} = X_A, \quad \frac{N_B}{N_A + N_B} = X_B.$$

$$\Delta S'_{\text{conf}} = -k N_0 (n_A \ln X_A + n_B \ln X_B)$$

$$\therefore \Delta S_{\text{conf}} = -R(X_A \ln X_A + X_B \ln X_B).$$

which is identical with  $\Delta S^{M,id}$ .

(d) ideal solution :  $\Delta \bar{H}_x^M = 0, \Delta \bar{S}_x^M = -R \ln X_x$ .

regular solution :  $\Delta \bar{H}_x^M \neq 0, \Delta \bar{S}_x^M = \Delta \bar{S}_x^{M,id} = -R \ln X_x$ .

3.(a). The critical value of  $\alpha$  is seen to be that which makes  $\frac{\partial^2 \Delta G^M}{\partial X_B^2}$ , and  $\frac{\partial^3 \Delta G^M}{\partial X_B^3}$  simultaneously equal to zero at that composition at which immiscibility becomes imminent.

$$\Delta G^M = RT(X_A \ln X_A + X_B \ln X_B) + RT\alpha X_A X_B$$

$$\frac{\partial \Delta G^M}{\partial X_B} = RT \left[ \ln \frac{X_B}{X_A} + \alpha (X_A - X_B) \right].$$

$$\frac{\partial^2 \Delta G^M}{\partial X_B^2} = RT \left( \frac{1}{X_A} + \frac{1}{X_B} - 2\alpha \right) - \textcircled{1}$$

$$\frac{\partial^3 \Delta G^M}{\partial X_B^3} = RT \left( \frac{1}{X_A^2} - \frac{1}{X_B^2} \right) - \textcircled{2}$$

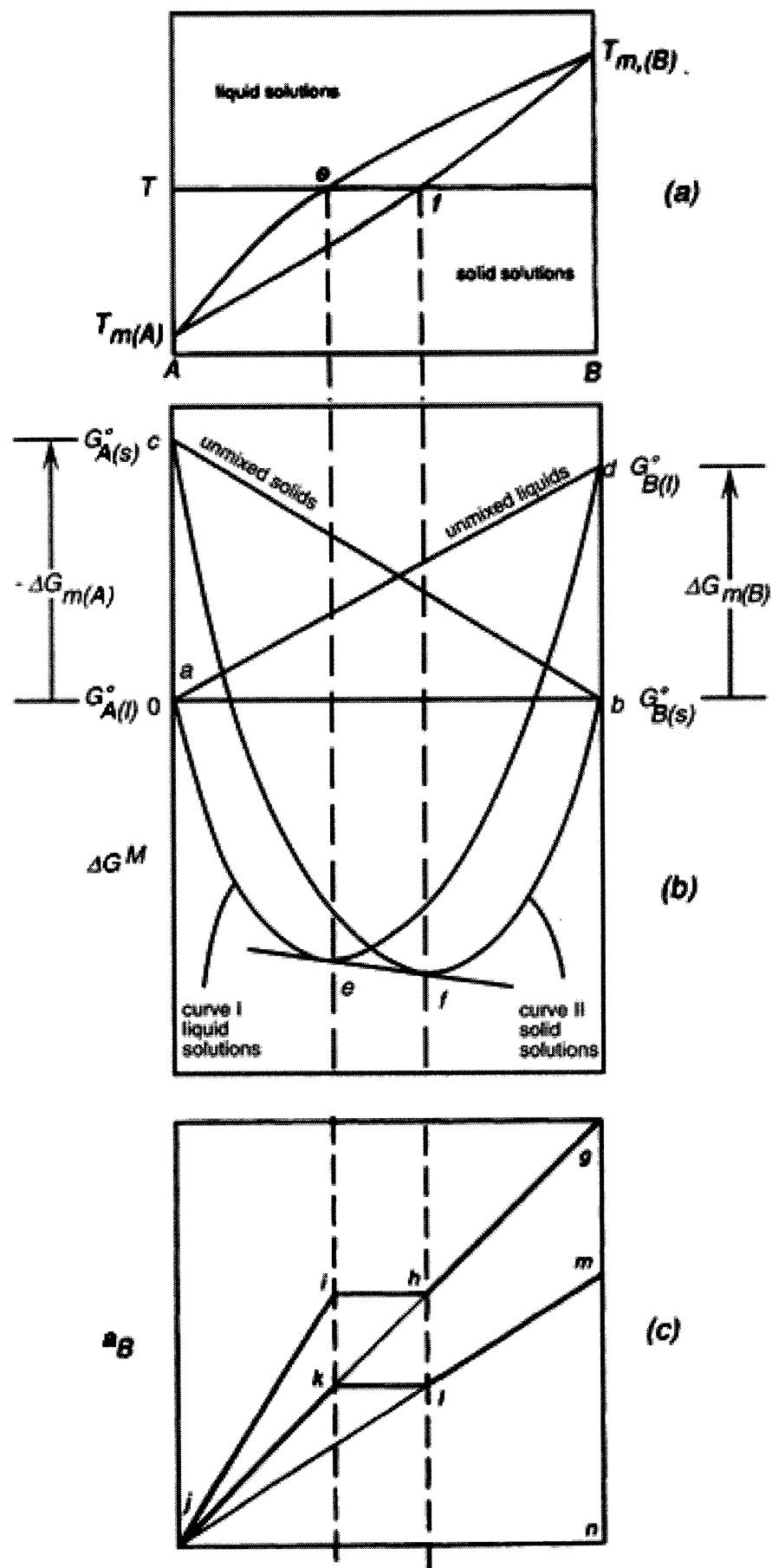
$\textcircled{2}$  equals to zero when  $X_A = X_B = 0.5$

$\textcircled{1}$  equals to zero at  $X_A = X_B = 0.5$  when  $\alpha = 2$ .

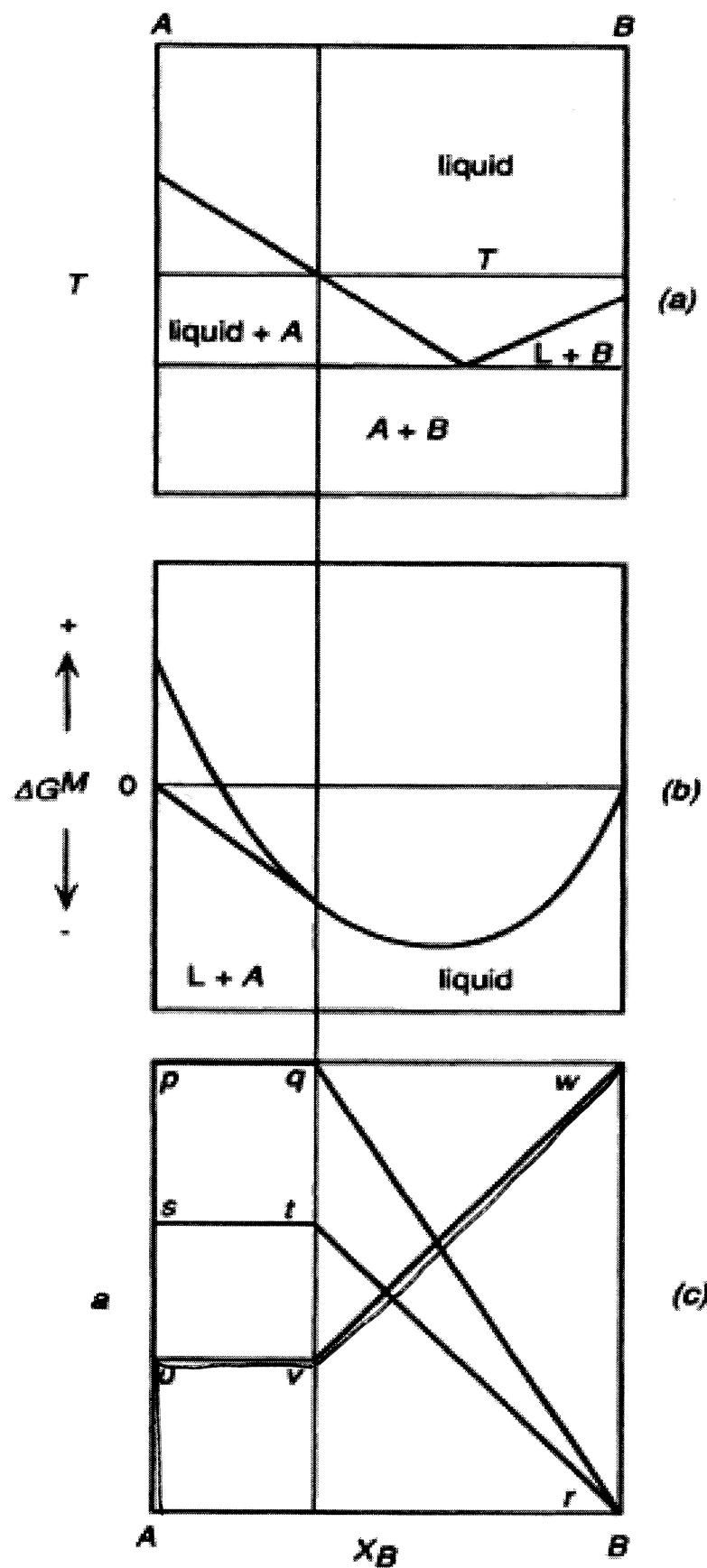
$\alpha = 2$  is thus the critical value.

Since  $\mathcal{R} = RT\alpha$ , The critical temperature  $T_{cr} = \frac{\mathcal{R}}{2R}$ .

3.(b)



3.(c)



4. (a) Gibbs - Helmholtz equation.

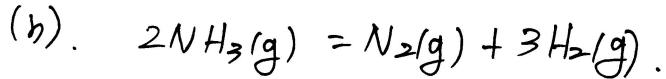
$$\left[ \frac{\partial(\Delta G^\circ/T)}{\partial T} \right]_P = -\frac{\Delta H^\circ}{T^2}$$

As  $\Delta G^\circ = -RT \ln K_p$ , then

$$\frac{\partial \ln K_p}{\partial T} = \frac{\Delta H^\circ}{RT^2} \quad \text{or} \quad \frac{\partial \ln K_p}{\partial \left(\frac{1}{T}\right)} = -\frac{\Delta H^\circ}{R}.$$

If  $\Delta H^\circ > 0$ ,  $K_p$  increases with increasing temperature.

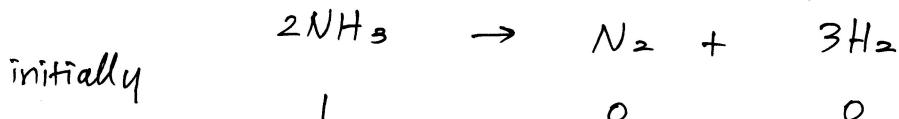
If  $\Delta H^\circ < 0$ ,  $K_p$  decreases with increasing temperature.



constant total pressure of 1 atm, at 400°C (673 K).

$$\Delta G^\circ = 81030 - 25.8 T \ln T - 31.1 T \text{ J.}$$

$$\Delta G^\circ_{673K} = -47390 \text{ J}, \quad K_p \cdot 673K = 4748.$$



$$n_T = 1-2x+x+3x = 1+2x \text{ moles.}$$

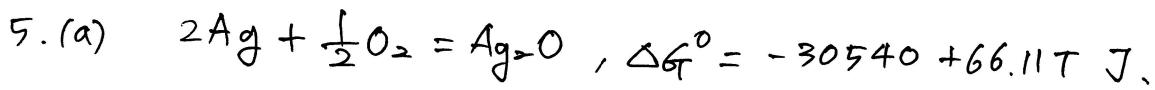
$$P_{H_2} = \frac{3x}{1+2x}, \quad P_{N_2} = \frac{x}{1+2x}, \quad P_{NH_3} = \frac{1-2x}{1+2x}.$$

$$\begin{aligned} K_p \cdot 673K &= \frac{P_{H_2}^3 P_{N_2}}{P_{NH_3}^2} = \frac{27x^4}{(1+2x)^2(1-2x)^2} \\ &= \frac{27x^4}{[1-(2x)^2]^2} \quad \text{or} \quad K_p^{1/2} = \frac{5.196x^2}{(1-4x^2)} \end{aligned}$$

$$\therefore x = 0.4954.$$

$$P_{H_2} = \frac{3x}{1+2x} = 0.7465 \text{ atm}$$

$$P_{N_2} = \frac{x}{1+2x} = 0.2488 \text{ atm}, \quad P_{NH_3} = \frac{1-2x}{1+2x} = 0.0047 \text{ atm}.$$



$$\Delta G^\circ = -RT\ln K, \text{ where } K = 1/p_{\text{O}_2}^{1/2}.$$

(i) pure oxygen at 1 atm.

$$\Delta G^\circ = -30540 + 66.11T = -RT\ln 1 = 0.$$

$$\therefore T = 462 \text{ K.}$$

(ii) in air

$$\Delta G^\circ = -30540 + 66.11T = -8.3144 \times T \times \ln \left( \frac{1}{(0.21)^{1/2}} \right)$$

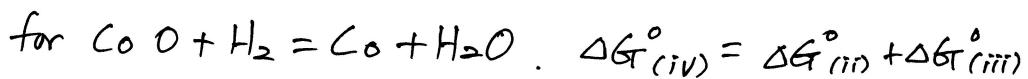
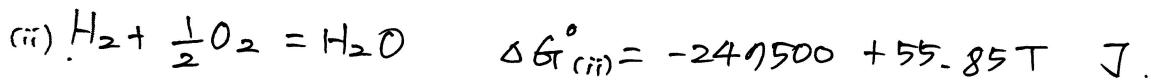
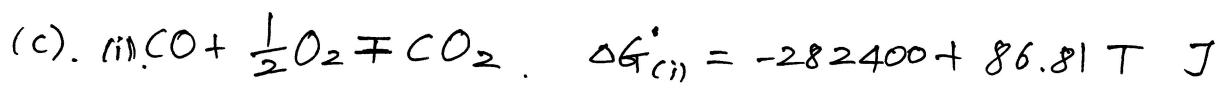
$$= \frac{1}{2} \times 8.3144 \times T \times \ln 0.21$$

$$\therefore T = 421 \text{ K.}$$

(b) From the Ellingham diagram, ~~line~~ line for  $\text{CO} + \frac{1}{2}\text{O}_2 = \text{CO}_2$  (i)  
and line for  $\text{H}_2 + \frac{1}{2}\text{O}_2 = \text{H}_2\text{O}$  (ii) intersects each other at 85°C (1123 K).

$T > 1123 \text{ K.}, \Delta G^\circ_{(\text{ii})} < \Delta G^\circ_{(\text{i})} \rightarrow \text{H}_2$  is the more efficient reducing agent.

$T < 1123 \text{ K.}, \Delta G^\circ_{(\text{ii})} < \Delta G^\circ_{(\text{i})} \rightarrow \text{CO}$  is the more efficient reducing agent.



$$= -13600 - 16 T \text{ J.}$$

$$\text{at } 1673K, \quad \Delta G_{(iv)}^\circ = -40368 \text{ J.}$$

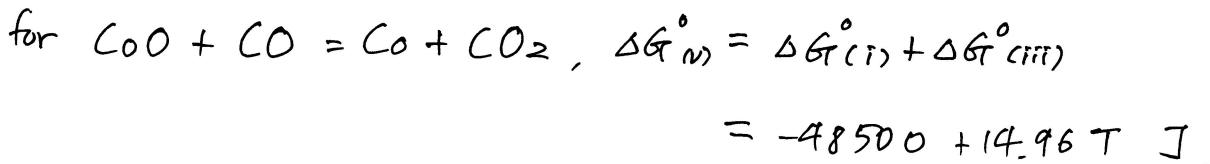
$$K_{(iv), 1673K} = \exp\left(\frac{-40368}{8.3144 \times 1673}\right) = 18.2 = \left(\frac{P_{\text{H}_2\text{O}}}{P_{\text{H}_2}}\right)_{\text{eq}}.$$

$\frac{18.2}{19.2} = 0.95$ . and thus 1 mole of  $\text{H}_2$  is required to reduce 0.95 moles of  $\text{CoO}$ .

$$\text{at } 873K, \quad \Delta G_{(iv)}^\circ = -21568 \text{ J.}$$

$$K_{(iv), 873K} = 44.6.$$

$$\frac{44.6}{45.6} = 0.98. \quad 0.98 \text{ moles of } \text{CoO} \text{ is reduced.}$$



$$\text{at } 1673K, \quad \Delta G_{(v)}^\circ = -23470 \text{ J.}$$

$$K_{(v), 1673K} = 5.40.$$

$$\frac{5.4}{6.4} = 0.84. \quad 0.84 \text{ moles of } \text{CoO} \text{ is reduced.}$$

$$\text{at } 873K, \quad \Delta G_{(v)}^\circ = -35440 \text{ J.}$$

$$K_{(v), 873K} = 132.$$

$$\frac{132}{133} = 0.99. \quad 0.99 \text{ moles of } \text{CoO} \text{ is reduced.}$$

$$6.(a). \quad aA + bB + cC + dD.$$

$$\Delta G = c\bar{G}_c + d\bar{G}_D - a\bar{G}_A - b\bar{G}_B. \quad - \textcircled{1}$$

$$\Delta G^\circ = cG^\circ_c + dG^\circ_D - aG^\circ_A - bG^\circ_B \quad - \textcircled{2}$$

$$\textcircled{1} - \textcircled{2}$$

$$\Delta G - \Delta G^\circ = c(\bar{G}_c - G^\circ_c) + d(\bar{G}_D - G^\circ_D) - a(\bar{G}_A - G^\circ_A) - b(\bar{G}_B - G^\circ_B)$$

$$\text{since, } \bar{G}_X = G_X^\circ + RT \ln a_X,$$

$$\begin{aligned}\Delta G - \Delta G^\circ &= c(RT \ln a_c) + d(RT \ln a_D) - a(RT \ln a_A) - b(RT \ln a_B) \\ &= RT \ln \left( \frac{a_c^c a_D^d}{a_A^a a_B^b} \right) = RT \ln Q.\end{aligned}$$

Reaction equilibrium is established when  $\Delta G = 0$ .

$$\therefore \Delta G^\circ = -RT \ln Q^\circ.$$

$$\text{and } \Delta G^\circ = -RT \ln K.$$

$$\therefore Q^\circ = K.$$



$$\ln K_p = -\frac{\Delta G^\circ}{RT} = \ln \left( \frac{P_{\text{CO}}}{P_{\text{CO}_2}} \times \alpha_{\text{FeO}} \right)$$

at. 1600°C (1873 K)

$$\ln \left( \frac{P_{\text{CO}}}{P_{\text{CO}_2}} \times \alpha_{\text{FeO}} \right) = -\frac{\Delta G^\circ}{RT} = -\frac{-28370}{8.3144 \times 1873} \neq 6.18$$

$$\alpha_{\text{FeO}} = \frac{P_{\text{CO}}}{P_{\text{CO}_2}} \times \exp \left( \frac{-28370}{8.3144 \times 1873} \right) = 6.18.$$

$$\therefore \alpha_{\text{FeO}} = 6.18 \times 10^{-5}$$

$$(c). \quad F = \underbrace{P(C+1)}_{\Theta} - \underbrace{(P-1)(C+2)}_{\Theta} \\ = C+2+P$$

①.  $P$  phases is determined by specification of its  $T, P$ , and  $(C-1)$  compositions.

$\therefore P(C+1)$  variables are fixed, when the state of the entire system is specified.

② The conditions that the entire system be at complete equilibrium are

$T_\alpha = T_\beta = T_\gamma = \dots (P-1)$  equalities of temperature

$P_\alpha = P_\beta = P_\gamma = \dots (P-1)$  equalities of pressure

$\alpha_{i(\alpha)} = \alpha_{i(\beta)} = \alpha_{i(\gamma)} = \dots (P-1)$  equalities of the activity of the species  $i$ .

and so on for each  $C$  chemical species.

Thus the total number of equilibrium conditions is

$$(P-1)(C+2).$$