#### **Mechanics of Composite Materials**

CHAPTER 3. Ply Elasticity

SangJoon Shin School of Mechanical and Aerospace Engineering Seoul National University

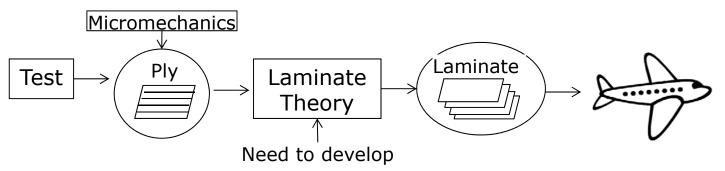


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#### Look at 3-D and 2-D anisotropic elasticity

- See Jones, Chap. 2 at Appendix A



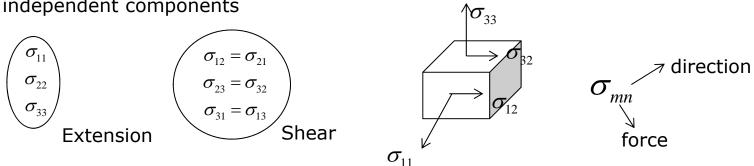
#### ✤ Make a brief review

- 1. Jones Book
- 2. Bisplinghoff, Mar and Pian, "Statics of Deformable Bodies" -> (tensor notation)
- 3. Herrmann, "Applied Anisotropic Elasticity"
- \* Notation Right hand coord. System,  $x_m$   $x_1$   $x_2$

#### \* Components of stress, $\sigma_{mn}$

"Stress tensor", 2 subscripts -> 2<sup>nd</sup> order

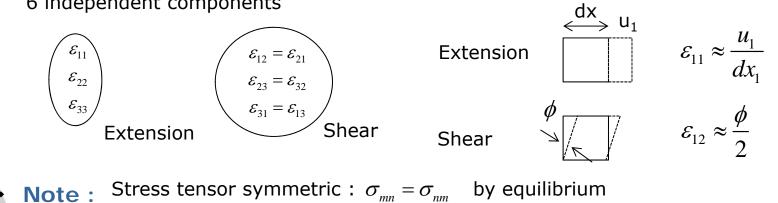
6 independent components



Components of strain,  $\mathcal{E}_{mn}$ \*\*

"Strain tensor", 2 subscripts -> 2<sup>nd</sup> order

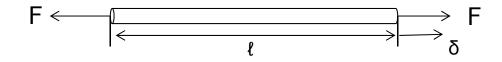
6 independent components



Strain tensor symmetric :  $\mathcal{E}_{mn} = \mathcal{E}_{nm}$  by Geometrical consideration

#### Stress – Strain Relations

Hooke's Law,  $F = k\delta$ : linear relation for rod



Can rewrite as  $\sigma = E\varepsilon$ 

where  $\sigma = \text{stress} = F/A$ 

$$\epsilon = \text{strain} = \delta/\ell$$

E = modulus of elasticity

Extending to 3-D stress, have "Generalized Hooke's Law"

 $\sigma_{mn} = E_{mnpq} \varepsilon_{pq}$ 

 $E_{mnpq}$  -> "Elasticity tensor" (3×3×3×3 = 81 components)

#### Recall Tensor Notation Rules

Latin subscripts (*m*, *n*, *p*, *q*, *r*, ...) -> 1, 2, 3

Greek subscripts (  $\alpha, \beta, \gamma$  , ... ) -> 1, 2

1. Subscripts that appear only once in a term are either 1, 2, or 3

$$\sigma_{i} = f(x_{i}) \quad \rightarrow \quad \left\{ \begin{array}{l} \sigma_{1} = f(x_{1}) \\ \sigma_{2} = f(x_{2}) \\ \sigma_{3} = f(x_{3}) \end{array} \right.$$

2. Subscripts repeated in a term are "dummy" subscripts -> Sum on them  $E_{ij}\varepsilon_j = E_{i1}\varepsilon_1 + E_{i2}\varepsilon_2 + E_{i3}\varepsilon_3 = \sum_{j=1}^3 E_{ij}\varepsilon_j$ 

3. No subscript can appear more than twice in a term

$$f_i C_{ij} D_i$$
 (**x**)

So, in general stress strain  $\sigma_{mn} = E_{mnpq} \varepsilon_{pq}$  9 equations are represented.

$$\sigma_{11} = E_{11pq} \varepsilon_{pq}$$

$$\sigma_{12} = E_{12pq} \varepsilon_{pq}$$

$$\sigma_{13} = E_{13pq} \varepsilon_{pq}$$

$$\sigma_{21} = E_{21pq} \varepsilon_{pq}$$

$$\vdots$$

$$\sigma_{31} = E_{31pq} \varepsilon_{pq}$$

$$\vdots$$

$$\sigma_{33} = E_{33pq} \varepsilon_{pq}$$

Look at 1<sup>st</sup> equations and sum over P

Overall symmetry (Energy considerations)

$$E_{mnpq} = E_{pqmn}$$

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#### So at most 21 independent constants

$E_{1111}$ $E_{1122}$	$E_{1212}$ $E_{1213}$	$E_{1112}$ $E_{2212}$ $E_{3312}$
$E_{2222}$ $E_{2233}$	$E_{1313}$ $E_{1323}$	$E_{1113}$ $E_{2213}$ $E_{3314}$
$E_{3333}$ $E_{3311}$	$E_{2323}$ $E_{2312}$	$E_{1123}$ $E_{2223}$ $E_{3323}$
Extension-Extension	Shear-Shear	Coupling Shear-Extension

Material with all 21 independent constants is <u>Anisotropic</u>

Have used Tensor Notation

To here mary books use a "contracted" notation (Jones, Tsai, etc) Also called "Engineering" Notation

3 major difference

① Subscript changes

Tensor	Contracted			Physical	
11	$\longrightarrow$	1	$\longrightarrow$	Extens.	in 1
22	$\longrightarrow$	2	$\longrightarrow$	Extens.	in 2
33	$\longrightarrow$	3	$\longrightarrow$	Extens.	in 3
23	$\longrightarrow$	4	$\longrightarrow$	Rotate	about 1
31	$\longrightarrow$	5	$\longrightarrow$	"	2
12	$\longrightarrow$	6	$\longrightarrow$	w	3

② Shear strain changes

Tensor shear strain is  $\frac{1}{2}$  of Engineering shear strain.

We change the notation from  $\varepsilon$  to  $\gamma$ 

Engineering		Tensor	Contracted
$\gamma_{12}$	=	$\mathcal{E}_{12} + \mathcal{E}_{21}$	${\cal E}_6$
$\gamma_{13}$	=	$\varepsilon_{13} + \varepsilon_{31}$	$\mathcal{E}_5$
$\gamma_{23}$	=	$\varepsilon_{23} + \varepsilon_{32}$	$\mathcal{E}_4$

(3) Elasticity constants represented by  $C_{ij}$  instead  $E_{mnpq}$  (Still 21 components)

Tensor Engineering

 $\begin{array}{cccc} E_{mnpq} & \rightarrow & C_{ij} \\ mn & \rightarrow & i \\ pq & \rightarrow & j \end{array}$ 

The "Generalized Hooke's Law" is

 $\sigma_{mn} = E_{mnpq} \varepsilon_{pq}$  (Tensor notation)  $\sigma_{i} = C_{ij} \varepsilon_{j}$  (Engineering notation)

Still use summation convention

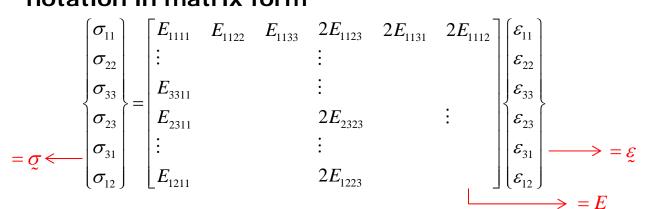
$$\sigma_i = \sum_j C_{ij} \varepsilon_j = C_{ij} \varepsilon_j$$
  
j = 1, 2, 3, ... 6

 $C_{ij} = C_{ji} \leftarrow$  Symmetry of the Elasticity constants still applies  $E_{mnpq} = E_{pqmn}$  other symm.s in  $E_{mnpq} = E_{nmpq}$ , etc. all automatically included in Engineering Notation.

Be careful of 2 in shear strain.

$$2 \mathcal{E}_{mn} = \mathcal{Y}_{mn}$$

Can see usefulness of Engineering notation by writing Tensor notation in matrix form

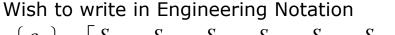


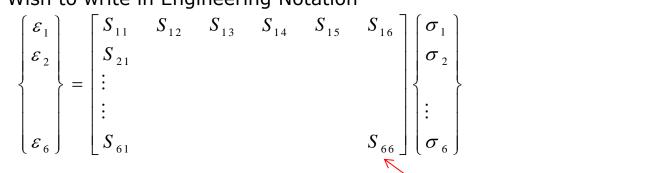
<u>E</u> matrix not symmetric, inconvenient; contracted(Engineering) more convenient

$$\sigma = C \quad \varepsilon \quad , \quad \sigma_i = C_{ij} \quad \varepsilon_j$$
Note  $\varepsilon_4 = 2 \quad \varepsilon_{23}$  etc.  
 $\sigma_{11} = \cdots \quad \cdots \quad 2E_{1123}\varepsilon_{23} + \cdots$   
 $\sigma_1 = \cdots \quad \cdots \quad C_{14}\varepsilon_4 + \cdots$   
C is symmetric,  $C_{12} = E_{1123}$ 

#### ✤ COMPLIANCE

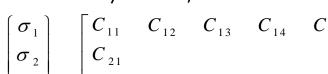
Just we have  $\sigma_{mn} = E_{mnpq} \varepsilon_{pq}$ Also have inverse  $\varepsilon_{mn} = S_{mnpq} \sigma_{pq}$  $S_{mnpq} \rightarrow Compliance tensor$  $\begin{array}{ccc} \sigma_{\tilde{c}} & = & E & \varepsilon_{\tilde{c}} \\ \varepsilon & = & E^{-1} \sigma \end{array} = S$ Same symmetries as E<sub>mppg</sub> Writing out Tensor Relations as matrices  $\begin{cases} \sigma_{11} \\ \sigma_{22} \\ \sigma_{12} \end{cases} = \begin{bmatrix} E_{1111} & E_{1122} & \cdots & 2E_{1112} \\ \vdots \\ E_{1211} & & 2E_{1212} \end{bmatrix} \begin{cases} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{12} \end{cases} \Rightarrow = \varepsilon$  $\begin{cases} \varepsilon_{11} \\ \varepsilon_{22} \\ \varepsilon_{12} \end{cases} = \begin{bmatrix} S_{1111} & S_{1122} & \cdots & 2S_{1112} \\ \vdots \\ S_{1211} & & 2S_{1212} \end{bmatrix} \begin{cases} \sigma_{11} \\ \sigma_{22} \\ \vdots \\ \sigma_{12} \end{cases} = \varepsilon$ 6x1 6x6 6x1





For Elasticity matrix, use

Symmetric : relate S<sub>ii</sub> to S<sub>mnpa</sub>



 $\begin{cases} \sigma_{1} \\ \sigma_{2} \\ \\ \sigma_{6} \end{cases} = \begin{bmatrix} C_{11} & C_{12} & C_{13} & C_{14} & C_{15} & C_{16} \\ C_{21} & & & & \\ \vdots \\ C_{61} & & & & \\ C_{61} & & & & \\ C_{ij}^{-1} & = S_{ij} & & \\ \end{bmatrix} \begin{cases} \varepsilon_{1} \\ \varepsilon_{2} \\ \\ \varepsilon_{6} \end{cases}$ 

(S.G) (\$/lb)

All symmetric matrices in Engineering Notation

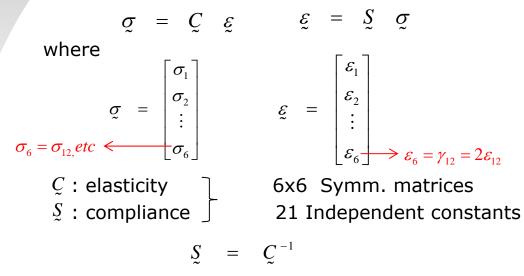
Problem Set #1 \*\*

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Weight Figure of Merit =  $\frac{\sigma_{ULT}}{(Specific Gravity)}$ 

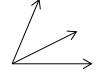
Cost Figure of Merit =  $\frac{\sigma_{ULT}}{\sigma_{ULT}}$ 

Stress-Strain relations (Engineering Notation)



#### ✤ Type of Materials

- Fully Anisotropic  $\rightarrow$  21 constants



Vary along non-orthogonal axis. Different stiffness along each direction (Same crystals)

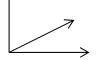
- Monoclinic Material  $\rightarrow$  13 constants



1 axis  $\perp$  other two Different stiffness along each direction (Some crystals, some composites)

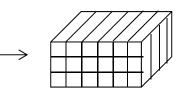


- Orthotropic Material  $\rightarrow$  9 constants

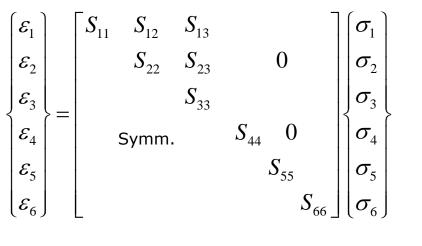


3 axis  $\perp$  to each other Different stiffness along each direction

Important Practical Case



Crystals. Composites



But no shear

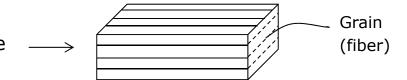
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3 perpendicular axis  $x_1$  is stiffer than  $x_2 = x_3$ same stiffness any direction in  $x_2$ ,  $x_3$  plane

Wood or composite \_



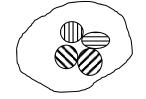
Like orthotropic, but additionally

5 constants

#### 3-14

Isotropic Material  $\rightarrow$  2 constants

Same properties in all directions Most metals, Resin



Many crystals randomly oriented Polycrystalline material

Same as Transversely Isotropic, but additionally

$$S_{22} = S_{11}$$
  

$$S_{13} = S_{12}$$
  

$$S_{55} = S_{44} = 2(S_{11} - S_{12})$$
  

$$S_{66} = S_{66} = 2(S_{11} - S_{12})$$

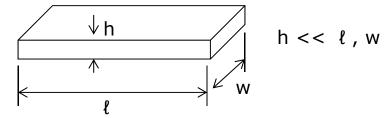
Only 2 constants  $S_{11}$  and  $S_{12}$ 

 $\rm S_{mn}{'}s$  traditionally expressed in terms of Modulus of Elasticity E and Poisson's Ratio  $~{\cal V}$ 

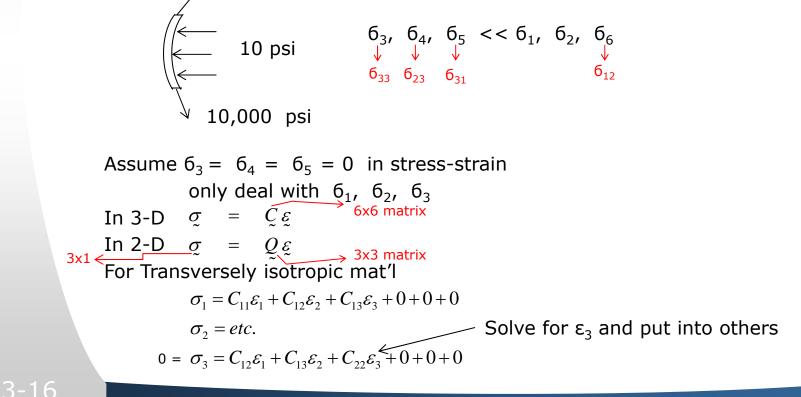
$$S_{11} = \frac{1}{E}, \quad S_{12} = -\frac{v}{E}$$
  
with these  
$$\varepsilon_{1} = \frac{1}{E} [\sigma_{1} - v\sigma_{2} - v\sigma_{3}]$$
  
:  
$$\varepsilon_{4} = \frac{2(1+v)}{E} \sigma_{4}$$

2-Dim. Plane stress approximations

Many structures are thin (plate)



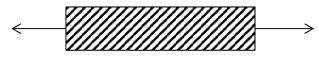
Also, not heavily loaded through thickness



$$\sigma_{6} = C_{66}\varepsilon_{6}$$

$$\Box = \begin{pmatrix} \sigma_{1} = Q_{11}\varepsilon_{1} + Q_{12}\varepsilon_{2} \\ \sigma_{2} = Q_{21}\varepsilon_{1} + Q_{22}\varepsilon_{2} \\ \sigma_{6} = Q_{66}\varepsilon_{6} \end{pmatrix}$$
Transversely isotropic

In general, for fibers not along the axis looks



$$\sigma_{1} = Q_{11}\varepsilon_{1} + Q_{12}\varepsilon_{2} + Q_{16}\varepsilon_{6}$$
  

$$\sigma_{2} = \cdots$$

$$\sigma_{6} = Q_{61}\varepsilon_{1} + Q_{62}\varepsilon_{2} + Q_{66}\varepsilon_{6}$$
[Q] 3x3, matrix symm.

Homework Prob.  $\rightarrow$  Relation between 3-D and 2-D

#### Properties of Single Ply

Ply → flat → plane stress 2-D stress-strain Eqns are  $\sigma = Q \in \mathcal{E}$ 

 $x_1, T$ 

 $\begin{array}{c} L : \text{ longitudinal} \\ \hline \end{array} \\ x_2, L \\ T : \text{ Transverse} \end{array}$ 

 $\stackrel{\checkmark}{\rightarrow}$  3x3 matrix (matrix of 6 constants)

On this set of axes - orthotropic

	$\sigma_1$		$Q_{11}$	$Q_{12}$	0 ]	$\left( \mathcal{E}_{1} \right)$	
<	$\sigma_{_2}$	=	$Q_{12}$	$Q_{\scriptscriptstyle 22}$	$\begin{bmatrix} 0 \\ 0 \\ Q_{66} \end{bmatrix}$	$\left\{ \mathcal{E}_{2} \right\}$	-
	$\left[\sigma_{_{6}} ight]$		0	0	$Q_{66}$	$\left[ \mathcal{E}_{6} \right]$	

4 constants

Also cloth (0/90) weave works this way, but some funny products may not. From Strength of Materials, we are familiar with Engineering Constants  $E_L = E_T - v_{LT} - G_{LT}$ Those are obtained from experimental tests. Formal definitions from

$$\begin{aligned} \varepsilon_{1} &= \frac{1}{E_{L}} (\sigma_{1} - v_{LT} \sigma_{2}) \\ \varepsilon_{2} &= \frac{1}{E_{T}} (\sigma_{2} - v_{TL} \sigma_{1}) \\ \varepsilon_{6} &= \frac{1}{G_{LT}} \sigma_{6} \end{aligned} \qquad \begin{cases} \sigma_{1} \\ \sigma_{2} \\ \sigma_{6} \end{cases} = \left| \frac{\varepsilon_{L}}{\sigma_{1}} - v_{LT} v_{TL} - v_{LT} v_{T}$$

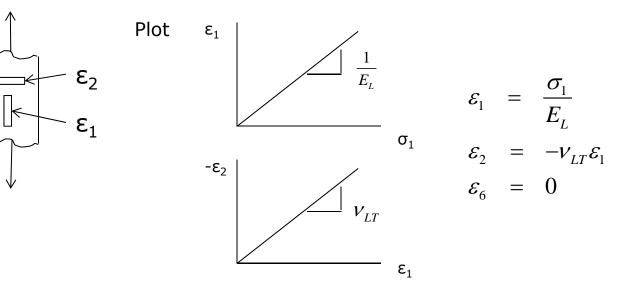
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Questions

- a) How to find Engineering Constants ?
- b) How to relate them to Elastic Constants Q<sub>ii</sub>
- a) Tests for Engineering constants
  - ① Longitudinal Tests

Apply known P (dead weight, calibrated machine) long, narrow specimen Know  $\sigma_1 = P/A$  (except near ends, reinforce there)  $\sigma_2 = 0$ ,  $\sigma_6 = 0$ Measure  $\sigma_1 = \sigma_2$  (c. 2) with strain space

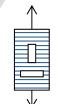
Measure  $\varepsilon_1$ ,  $\varepsilon_2$  ( $\varepsilon_6$  ?) with strain gages



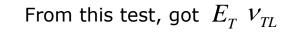
From this test  $\rightarrow$  2 constants

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#### 2 Transverse Tension



Same deal, apply known P

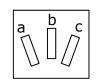


 $\mathcal{E}_6 = e_c - e_a$ 

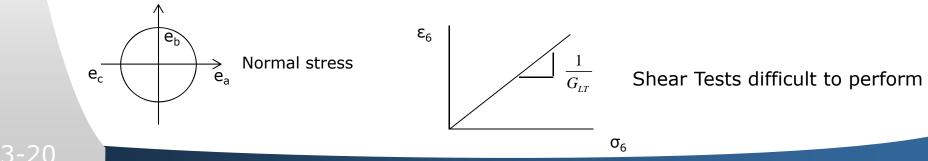
3 Shear Tests



Apply known shear  $\sigma_6$  (not too easy)

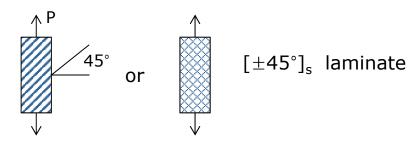


Measure shear strain with a rosette

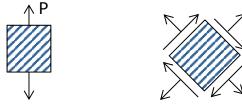


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#### Easier to test 45° Ply in tension



This gives a mixed state of stress in axis system of the material



But can untangle to get G

So, we have

$$\varepsilon_{1} = \frac{1}{E_{L}} \sigma_{1} - \frac{v_{TL}}{E_{T}} \sigma_{2}$$

$$\varepsilon_{2} = -\frac{v_{LT}}{E_{L}} \sigma_{1} + \frac{1}{E_{T}} \sigma_{2}$$

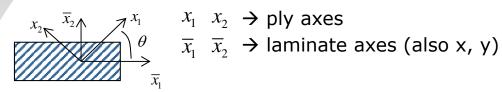
$$\varepsilon_{6} = \frac{1}{G_{LT}} \sigma_{6}$$

1

Because of symmetry

$$\frac{v_{TL}}{E_T} = \frac{v_{LT}}{E_L} \qquad \begin{array}{c} v_{LT} : \text{Major Poisson's Ratio} & \sim 0.3 \\ v_{TL} : E_T/E_L \cdot v_{LT} & \sim 0.02 \end{array}$$

#### Rotation of Plies



Ply at angle  $\theta$  from lamina axis  $\overline{x}_1$  $(+\theta \rightarrow x_1 \text{ going towards } x_2)$ In  $x_1$ ,  $x_2$  (Ply axes)  $\rightarrow$  2-D orthotropic Material  $\begin{cases} \sigma_1 \\ \sigma_2 \\ \sigma_6 \end{cases} = \begin{bmatrix} Q_{11} & Q_{12} & 0 \\ Q_{12} & Q_{22} & 0 \\ 0 & 0 & Q_{66} \end{bmatrix} \begin{cases} \varepsilon_1 \\ \varepsilon_2 \\ \varepsilon_6 \end{cases} \quad \text{or} \quad \sigma = Q \varepsilon$ 

Q's from  $E_I = E_T = V_{IT} = G_{IT}$ To find stress-strain in  $\overline{x}_1$ ,  $\overline{x}_2$  (laminate axes) First relate stresses in 2 axis systems.

 $\sigma_{mn} = l_{m\overline{p}} \ l_{n\overline{q}} \quad \overline{\sigma}_{pq} \leftarrow Standard transform Law$   $\swarrow \qquad Stress in \ \overline{x}_1 \ \overline{x}_2$ Stress tensor Direction cosine in  $x_1, x_2$  = cos(angle  $x_m$  and  $\overline{x}_p$ )

Table of cosines

	$x_1$	<i>x</i> <sub>2</sub>
$\overline{x}_1$	$\cos \theta$	$\cos(90+\theta) = -\sin\theta$
$\overline{X}_2$	$\cos(90 - \theta) = \sin \theta$	$\cos \theta$

$$\sigma_{11} = l_{1\overline{1}} l_{1\overline{1}} \overline{\sigma}_{11} + l_{1\overline{2}} l_{1\overline{2}} \overline{\sigma}_{22} + l_{1\overline{1}} l_{1\overline{2}} \overline{\sigma}_{12} + l_{1\overline{2}} l_{1\overline{1}} \overline{\sigma}_{21}$$
  
$$= \cos^2 \theta \overline{\sigma}_{11} + \sin^2 \theta \overline{\sigma}_{22} + \cos \theta \sin \theta \overline{\sigma}_{12} + \cos \theta \sin \theta \overline{\sigma}_{21}$$
  
$$\sigma_{22} = etc.$$
  
$$\sigma_{12} = etc.$$

So obtain

$$\begin{cases} \sigma_1 \\ \sigma_2 \\ \sigma_6 \end{cases} = \begin{bmatrix} c^2 & s^2 & 2cs \\ s^2 & c^2 & -2cs \\ -cs & cs & (c^2 - s^2) \end{bmatrix} \begin{bmatrix} \overline{\sigma}_1 \\ \overline{\sigma}_2 \\ \overline{\sigma}_6 \end{bmatrix}$$
Ply
$$T_{\sigma}$$
Laminate
or
 $\sigma = T_{\sigma}\overline{\sigma}$ 
where  $c = \cos\theta$   $s = \sin\theta$ 

Also for strain

$$\begin{vmatrix} \varepsilon_{1} \\ \varepsilon_{2} \\ \frac{1}{2} \\ \varepsilon_{6} \end{vmatrix} = T_{\sigma} \begin{cases} \overline{\varepsilon}_{1} \\ \overline{\varepsilon}_{2} \\ \frac{1}{2} \\ \overline{\varepsilon}_{6} \end{cases}$$
 Recall 
$$\varepsilon_{12} = \frac{1}{2}\gamma_{12} = \frac{1}{2}\varepsilon_{6}$$

Absorb the  $\frac{1}{2}$  into  $\frac{T}{\sigma}$  gives

$$\begin{cases} \varepsilon_1 \\ \varepsilon_2 \\ \varepsilon_6 \end{cases} = \begin{bmatrix} c^2 & s^2 & cs \\ s^2 & c^2 & -cs \\ -2cs & 2cs & (c^2 - s^2) \end{bmatrix} \begin{cases} \overline{\varepsilon}_1 \\ \overline{\varepsilon}_2 \\ \overline{\varepsilon}_6 \end{cases}$$
Ply T<sub>\varepsilon</sub> Laminate

or  $\underline{\varepsilon} = \underline{T}_{\underline{\varepsilon}} \overline{\underline{\varepsilon}}$ 

Placing into Ply axes stress strain

$$\sigma = Q \varepsilon$$

$$T_{\sigma} \overline{\sigma} = Q T_{\varepsilon} \overline{\varepsilon} \qquad \overline{Q}$$

$$\overline{\sigma} = \overline{T_{\sigma}^{-1} Q T_{\varepsilon}} \overline{\varepsilon}$$

$$o r \quad \overline{\sigma} = \overline{Q} \overline{\varepsilon} \qquad \text{Stress-strain Relation}$$

$$\text{In laminate } \overline{x_1} \quad \overline{x_2} \text{ axes}$$

Now Note Inverses

$$T_{\sigma}^{-1} = T_{\sigma}(-\theta) = T_{\varepsilon}^{T}$$
$$T_{\varepsilon}^{-1} = T_{\varepsilon}(-\theta) = T_{\sigma}^{T}$$

So rotated Q matrix is

$$\overline{Q} = T_{\varepsilon}^{T} Q T_{\varepsilon}$$

Q fully populated now

Also in Jones Notation,

Laminated Axes  $\overline{x}_1, \overline{x}_2 \rightarrow x, y$ Laminated stress  $\overline{\sigma}_i \rightarrow \sigma_x, \sigma_y, \tau_{xy}$ Laminated strain  $\overline{\varepsilon}_i \rightarrow \varepsilon_x, \varepsilon_y, \gamma_{xy}$ 

Final Laminated stress – strain Eqns

$$\begin{cases} \overline{\sigma}_{x} \\ \overline{\sigma}_{y} \\ \overline{\tau}_{xy} \end{cases} = \begin{bmatrix} \overline{Q}_{11} & \overline{Q}_{12} & \overline{Q}_{13} \\ \overline{Q}_{12} & \overline{Q}_{22} & \overline{Q}_{26} \\ \overline{Q}_{16} & \overline{Q}_{26} & \overline{Q}_{66} \end{bmatrix} \begin{cases} \overline{\varepsilon}_{x} \\ \overline{\varepsilon}_{y} \\ \overline{\gamma}_{xy} \end{cases}$$

Multiplies, out matrices

$$\begin{aligned} \overline{Q}_{11} &= Q_{11}\cos^4\theta + 2(Q_{12} + 2Q_{66})\sin^2\theta\cos^2\theta + Q_{22}\sin^4\theta \\ \overline{Q}_{12} &= (Q_{11} + Q_{22} - 4Q_{66})\sin^2\theta\cos^2\theta + Q_{12}(\sin^4\theta + \cos^4\theta) \\ \overline{Q}_{22} &= etc. \\ \overline{Q}_{66} &= etc. \end{aligned}$$

Similarly can transform compliances.

from  $\varepsilon = S\sigma$ obtain  $\overline{\varepsilon} = \overline{S}\overline{\sigma}$ where  $\overline{S} = T_{\sigma}^{T}ST_{\sigma}$ 

Alternate ways of rotating

 $\overline{E}_{mnpq} = l_{m\overline{r}} l_{n\overline{s}} l_{p\overline{t}} l_{q\overline{u}} E_{rstu}$ Also, can mathematically reduce  $\overline{Q}_{ij}$  by  $\cos^2 \theta = \frac{1 + \cos 2\theta}{2}$ , etc. can then express

> $\overline{Q}_{ij} = A_{ij} + B_{ij} \cos 2\theta + C_{ij} \cos 4\theta + D \sin 2\theta + E \sin 4\theta$ A, B, C, D, E  $\rightarrow$  depend only on 4 invariants