

## D. Instabilities of tokamaks

1) **Ideal MHD modes** [ $n=0$ ,  $t = O(10^{-6})$ ] : *Linear pert. theory, Energy principle*

### a. Ideal kink modes

Driven by plasma currents  $j$  ( $\nabla j$ )

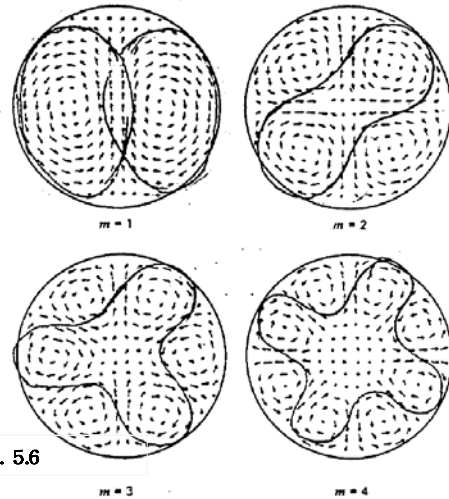
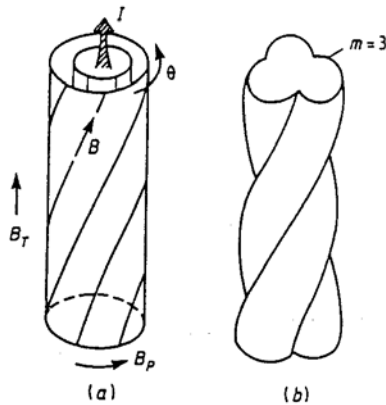
→ perturbation  $\propto e^{i(m\theta - nt)}$  ⇒ (m, n) mode

#### i) External modes

= Free surface modes (m=0 sausage, m=1 helical kink, m=2, 3, ... surface kinks)

⇒ (m, n) modes fall on plasma surface  $r=a$  (vac. region),  
i.e., mode rational surface  $q(a) = m/n$

Fastest and most dangerous



Gross<sup>3)</sup> Fig. 5.6

Stabilizing effects and stability conditions :

- ① Conducting wall stabilization for low n modes
- ② Strong  $B_\phi$
- ③  $q(a) > m/n$  for (m, n) modes w/o conducting wall

Kruskal-Shafranov limit for m = n = 1 mode:

$$q(a) > 1 \rightarrow I_\phi \text{ limitation} \quad (36)$$

$$(e.g.) I_\phi < \frac{2\pi B_\phi^0}{q(a)_{\text{H.o.}}} \left(\frac{a}{R_o}\right)^2 R \quad (37)$$

$$\Rightarrow R_o/a=3, R_o=11 \text{ m}, B_\phi^0 = 5 \text{ T} \rightarrow I_{\text{max}} \approx 12 \text{ MA}$$

- ④ Centrally peaked  $j_\phi(r)$  for  $m \geq 2$

$$\Rightarrow \frac{q(a)}{q(0)} \geq \nu + 1 \approx 2 \sim 3 \text{ for } j_\phi = j_o(1 - r^2/a^2)^\nu \quad (38)$$

For  $\nu=0$  (uniform), always unstable against any (m, n) mode

- ⑤ Magnetic shear ( $S \propto \frac{i'}{i} = \frac{q'}{q}$ )

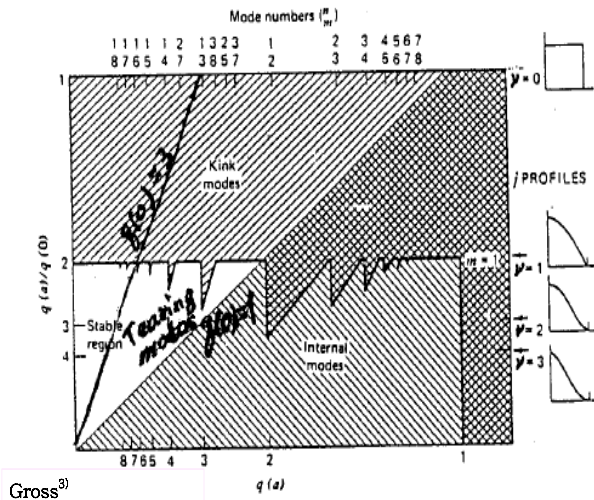


FIGURE 5.7. Stability diagram for ideal kink and internal modes of a large-aspect-ratio circular cross-section tokamak. The current distribution is  $j_z(r) = j_z(0)(1 - r/a)^2$ . Complete stability against kink modes may be obtained for any  $q(a) > 1$  by sufficient peaking of the current profile (larger value of  $\nu$ ). The maximum stable current for this model is obtained with a parabolic current distribution with  $q(0) = 1$  and  $q(a) = 2$ . (Data from J.A. Wesson [17].)

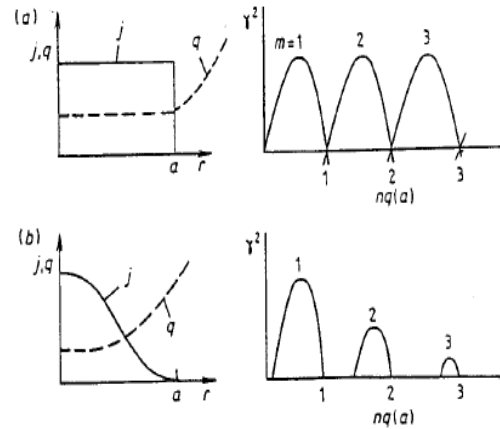


Figure 5.2 Growth rates of ideal MHD kink modes for different current density profiles: (a) homogeneous current density distribution; (b) natural current profile with decreased current density near the walls.

## ii) Internal modes

Fixed boundary modes

∃ Singular surface inside plasma

Localized near rational surface  $r = r_s$  where  $q(r_s) = m/n$

Stability condition for  $m=1, n=1$  mode :  $q(0) > 1$  (39)

( $m \geq 2, \forall n$  modes  $\rightarrow$  stable)

## b. Ideal interchange modes

Driven by  $\nabla p$

Internal modes : localized near mode rational surface  $r_s$

where  $q(r_s) = m/n$

No threat to confinement unless  $q(0) \ll 1$

Stabilization :

- ① Min-B
- ② Magnetic shear
- ③ Mercier necessary condition

$$\left(\frac{q'}{q}\right)^2 + \frac{4}{r} \frac{P'}{B_\phi^2/2\mu_0} (1 - q^2) > 0 \quad (40)$$

$$\Rightarrow q(r) > 1 \text{ for all } r$$

$$\Rightarrow q(0) > 1 \quad (41)$$

shear  
stabilization

(cf) Suydam criterion for cylindrical plasma

$$\left(\frac{q'}{q}\right)^2 + \frac{4}{r} \frac{P'}{B_\phi^2/2\mu_0} > 0 \quad (42)$$

- ④ Elongated outward triangular cross section

### c. Ideal ballooning modes

Driven by  $\nabla p$  at bad-curvature surface region  
 Localized high- $n$  interchange mode at outbound edge of circular  
 high- $\beta$  tokamak or at the tips of an elongated plasma  
 Most dangerous and limiting MHD instability

Stabilization :

- $\exists$  critical  $\beta$  above which  
 $\exists$  ball. modes at any  $q$  values.

$$\beta < \beta_{\max} \approx \frac{\varepsilon}{q^2} \quad (43)$$

- Large  $R$  and  $\nabla p (\propto 1/a)$  easily produce ball. modes  $\rightarrow$  LAR
- Upper limit on  $n$

$$\frac{nk(T_e + T_i)}{B^2/2\mu_0} < \frac{a}{q^2 R_o}$$

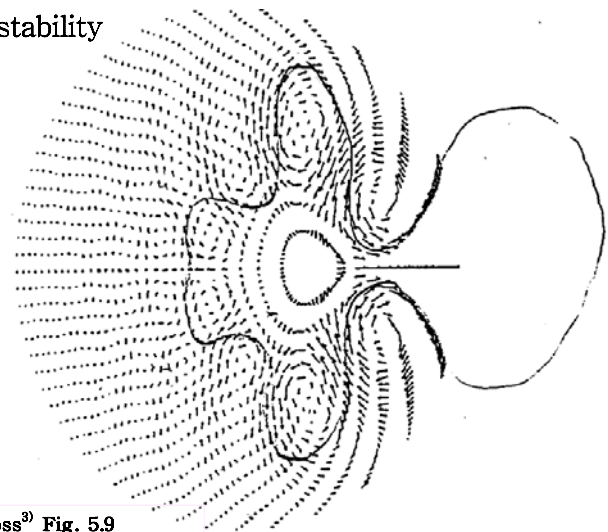
$$\Rightarrow n < n_{\max} \approx \frac{a}{6R_o} \frac{1}{k(T_e + T_i)} \frac{B^2}{2\mu_0} \quad (44)$$

(e.g.)  $R_o/a = 3, B = 9 T, T_e = T_i = 10 keV$   
 $\rightarrow n_{\max} \approx 5 \times 10^{20} m^{-3}$

- Upper limit on  $\beta$

$\beta_{\max} \uparrow$  as  $\varepsilon \uparrow$ , Elongation  $\uparrow$  (noncircular),

Conducting wall, Low  $n \downarrow$  modes (W/ conducting wall)



Gross<sup>9)</sup> Fig. 5.9

5.9. Flow pattern for a tokamak ballooning mode with  $n = 3$  at  $(\beta^*) = 3\%$ . Beta star

### Notes)

Limit on  $\beta$  due to ideal MHD instabilities :

$$\begin{aligned} \beta &\equiv \frac{p}{B^2/2\mu_0} \approx \beta_p \left( \frac{B_\theta}{B_\phi} \right)^2 = \beta_p \left( \frac{a/R_o}{q(a)} \right)^2 \\ &= \underbrace{\left( \frac{a}{R_o} \beta_p \right)}_{\textcircled{1}} \underbrace{\left( \frac{a}{R_o} \right)}_{\textcircled{2}} \underbrace{\left( \frac{q(o)}{q(a)} \right)^2}_{\textcircled{3}} \underbrace{\left( \frac{1}{q(o)} \right)^2}_{\textcircled{4}} \end{aligned} \quad (45)$$

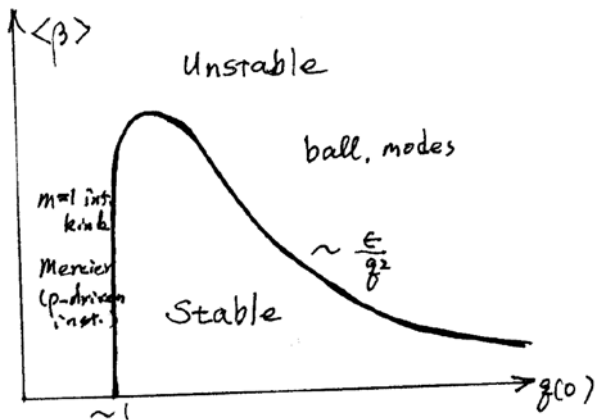
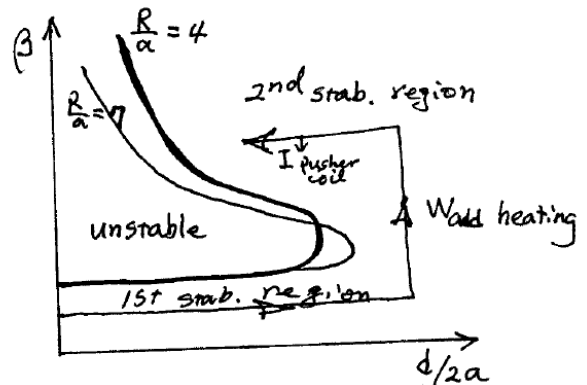
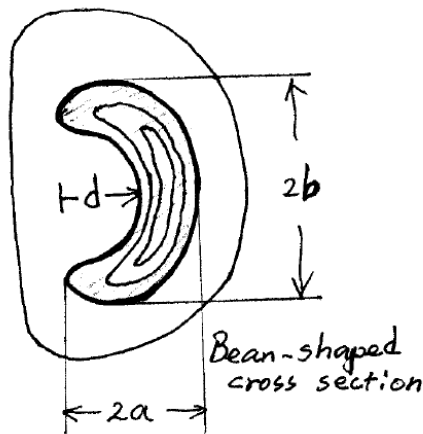
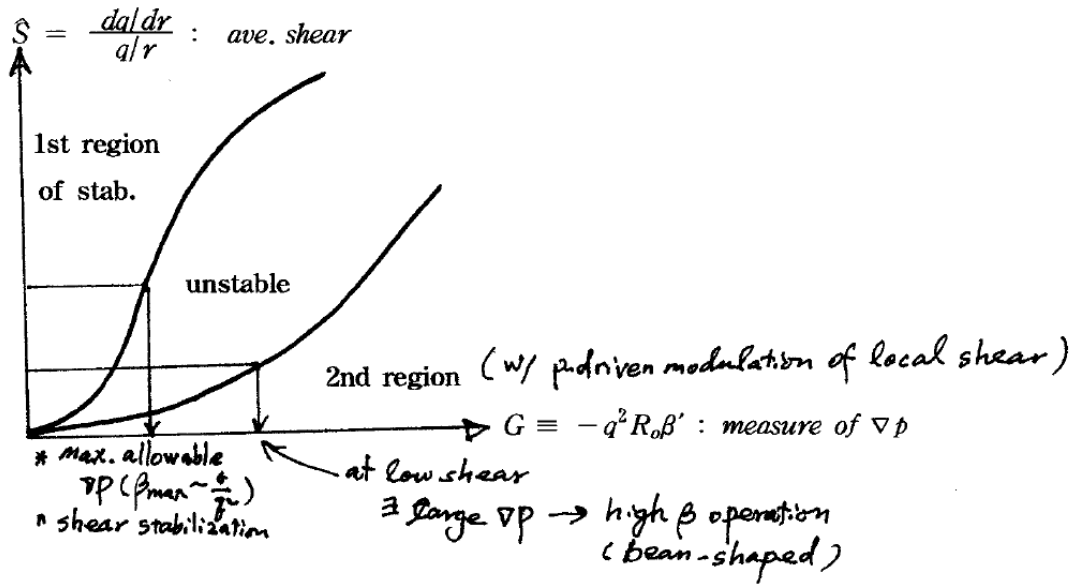
①  $\leq 1$  : ballooning modes limit

②  $\leq 1/3$  : space limit (geometry, shielding, maintenance, heating, etc)

③  $\leq 0.2$  : surface kinks

④  $\leq 1$  : internal modes

(cf) Energy principle  $\Rightarrow$



Troyon limit for { Mercier criterion  
ballooning  
 $m=1$  internal kink  
 $n=1$  external ball. kink

$$\beta_{crit} (\%) = \beta_N \frac{I(MA)}{a(m)B(T)} \quad (46)$$

where normalized  $\beta_N \approx 2.8 \sim 5$

► Troyon factor  $C_T$

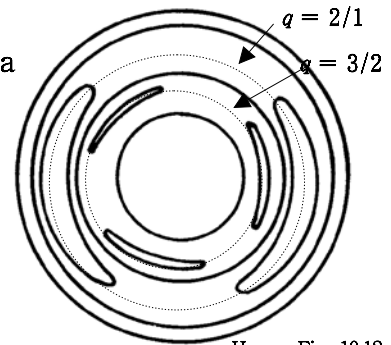
## 2) Resistive MHD modes [ $n \neq 0$ , $t = O(10^{-4} \sim 10^{-2})$ ]

### a. Tearing modes = Resistive internal kink modes ( $m \geq 2$ )

Driven by perturbed  $B$  induced by current layer ( $\nabla j$ ) in plasma

Magnetic island formation (local pinch)

Mode rational surface  $r_s$  [ $q(r_s) = m/n$ ] falls in plasma



Harms Fig. 10.12

Saturation at some fraction of plasma width

( $\sim$  a few tenth of plasma radius  $a$ )

Growth rate  $\gamma \propto n^{3/5}$

More tolerable & lower than ideal modes

Unstable region  $\downarrow$  as sharpness of current profile  $v \uparrow$

$m \uparrow$

closeness of wall to plasma

$[q(a)/q(0)] \uparrow$  ; shear

Stability condition :

$$q(0) \geq 3$$

(47)

### b. $m=1$ tearing mode

Always occurs late in the sequence of unstable modes whenever  $q(0) < 1$

for  $m = n = 1$

Independent of current profile

$$\gamma \propto n^{1/3}$$

### c. Nonlinear tearing modes (Disruptive instability)

Nonlinear evolutions of resistive modes:

{	Sawtooth oscillation	( $m = n = 1$ )
	Mirnov osc. : fish-bone	( $n = 1, m = 1, 2, 3$ )
	Major disruption	( $m = 2, n = 1$ : trigger)

#### i) Internal disruption

(1,1) mode disruption

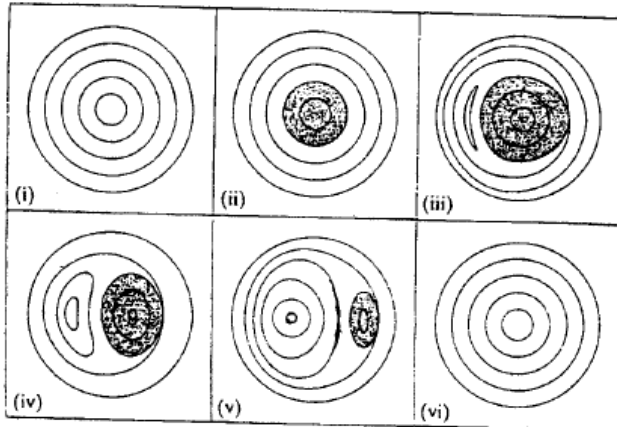
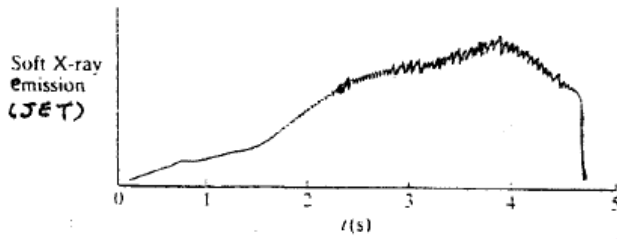
Observation of Sawtooth oscillation (soft X-ray, 2 ~ 20 keV)

Mirnov oscillation ( $n=1, m=1,2,3 \dots$ ): Fish-bone mode

Enhanced energy transport in plasma center

Wessen<sup>13)</sup> Fig.7.1.2

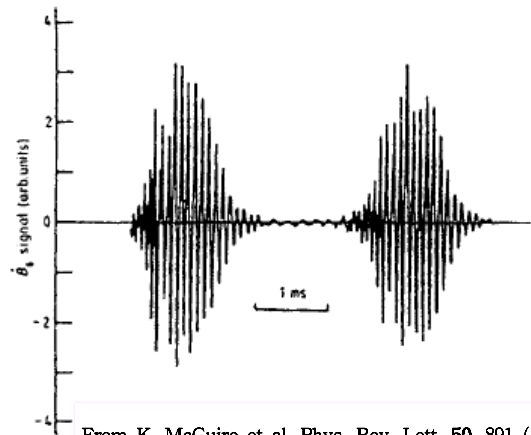
### Sawtooth Oscillation



Wessen<sup>13)</sup> Fig.7.6.4

Possible development of the magnetic field structure during the rapid fall stage of the sawtooth instability. The  $m = 1$  instability displaces the  $q < 1$  region (shown shaded) and restores  $q_0$  to a value  $> 1$ .

### Fish-bone oscillation



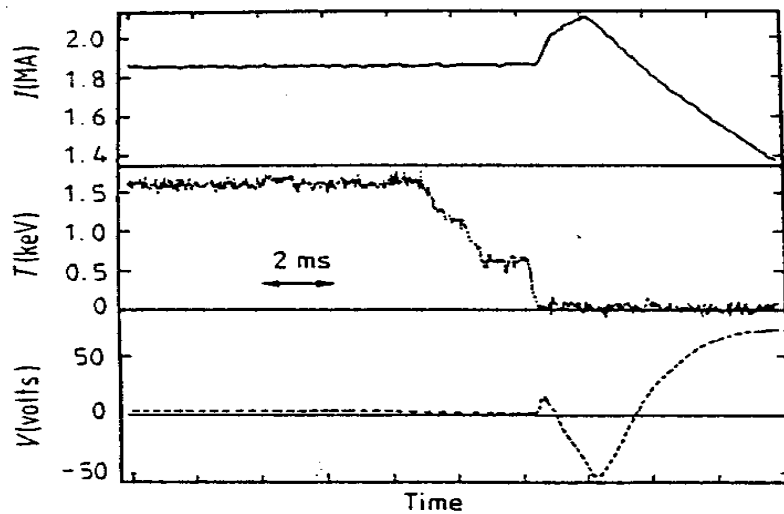
From K. McGuire et al. Phys. Rev. Lett. 50, 891 (1983)

Oscillations of the poloidal-field time derivative as an indication of fish-bone instability

### ii) Major disruption

Nonlinear evolution of  $m \geq 2$  tearing

(2, 1) mode triggering external disruption when  $q(0) < 1$  &  $q(a) < 2.5$



From J.A. Wessen et al. Nucl. Fusion 30, 1011 (1990)

Termination of confinement:

- Most dangerous and unpredictable
- Large negative voltage spike in transformer
- Damage to wall of PFCs
- Rapid cooling of plasma
- Hard X-ray burst
- Plasma column shift

Control:

- Keep  $I_\phi$  and n within stable limits
- Feedback control by helical coils
- Additional heating to raise  $T_e$
- Close fitting conducting wall

#### d. Resistive interchange modes

Weaker stabilizing effect of shear

$$\gamma \propto \eta^{1/3}$$

Stable for  $m \geq 4$  for large tokamak

Stable even lower  $m(\leq 2)$  for reactor tokamak

### 3) Microinstabilities

Kinetic approach - limited MHD approach

Anomalous transports

#### a. Two-stream instability

Particle bunching  $\rightarrow \tilde{E} \rightarrow$  bunching  $\uparrow \rightarrow$  unstable

#### b. Drift (or Universal) instability

Driven by  $\nabla p$  (or  $\nabla n$ ) in  $B$  (with  $k_\parallel \ll k_\perp$ )

$\rightarrow$  Excitation of drift waves with  $\omega/k_\perp = v_{De}$

Most unstable very short  $\lambda$  ( $k_\perp r_{Li} \approx 1$ )

Anomalous transport  $\rightarrow D_{Bohm}$

Stabilization : Good curvature (min-B), Shear, Finite  $\beta$

#### c. Trapped particle modes $\rightarrow$ cross-field diffusion $\uparrow$

Enhanced drift instability by trapped particle effects

Trapped electron mode (TEM)

Trapped ion mode (TIM)

Note)

No microinstability ( $\omega \geq \Omega_i$ )

due to velocity-space free energy (non-Maxwellian) in Tokamaks

### E. Neoclassical transports

$$\Gamma = -D \nabla n \approx -\frac{(\Delta r)^2}{\tau} \nabla n \quad : \text{Fick's law}$$

$$\mathbf{q} = -\kappa \nabla T \approx -\frac{(\Delta r)^2 n}{\tau_E} \nabla T \quad : \text{Fourier's law}$$

Thermal diffusivity :

$$\chi \equiv \frac{\kappa}{n} \approx \frac{(\Delta r)^2}{\tau_E} \approx D \Rightarrow \tau_E \approx \frac{a^2}{\chi}$$

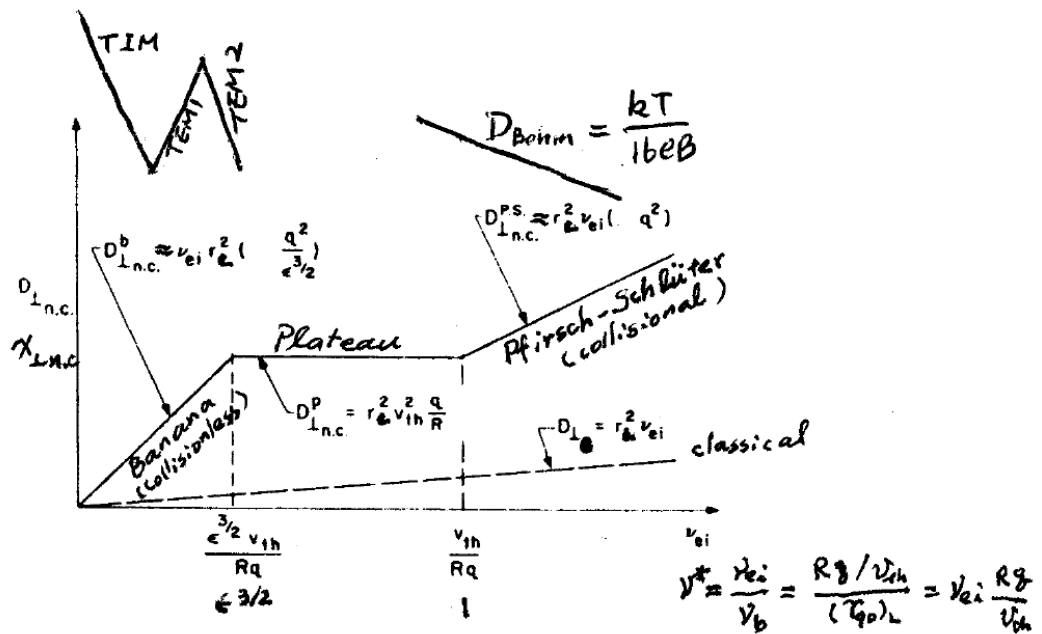
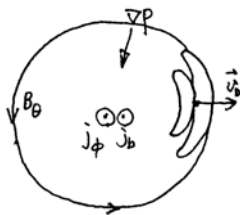


FIGURE 5.13. Neoclassical electron transport for tokamaks.

Bootstrap current - Toroidal current induced by radial NC diffusion in banana regime



→ Direct consequence of banana toroidal drift combined with collisional diffusion

$$j_b \approx -\frac{\epsilon^{1/2}}{B_\theta} \frac{\partial p}{\partial r}$$

$$\frac{I_{BS}}{I} = 0.67 \beta_p \sqrt{\epsilon}$$

→ Non-inductive current drive

Steady-state operation

Observed in TFTR, JT-60, JET (~89)

Comparison of NC results with experiments :

	$\tau$	$\tau_E$	$\chi$
ions	×(wrong)	×	Δ(may agree)
electrons	×	×	×(O(1)~O(2))



### F. Energy confinement time scaling ( $\tau_E \equiv E_{th} / P_{heat}$ )

Alcator scaling (75-82): OHT(Alcator A)

$$\tau_E = 5 \times 10^{-21} \overline{n_e} a^2$$

Neo-Alcator (82) : OHT(Alcator-C)  $\tau_{NA} = 1.92 \times 10^{-21} \overline{n_e} a R^2$

Goldston (84) : AHT(TFTR)  $\rightarrow$  L-mode (Low-confinement mode)

$$\tau_G = 3 \times 10^{-5} I_p R^{1.75} a^{-0.37} \kappa^{0.5} A_i^{0.5} / P_{aux}^{0.5}$$

ITER L-mode (89)

$$\tau_{ITER-89} = 0.048 I^{0.85} R^{1.2} a^{0.3} n_{20}^{0.1} B^{0.2} (A_{i\kappa_x} / P(MW))^{0.5}$$

ITER H-mode (90)  $\rightarrow$  H-mode (High-confinement mode)

$$\tau_{H-mode} \approx 2 \sim 3 \times \tau_{L-mode} \quad \text{1984 ASDEX}$$

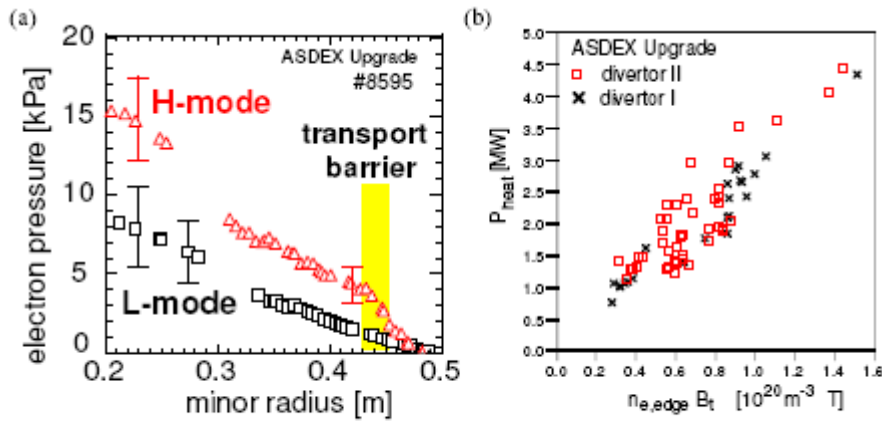
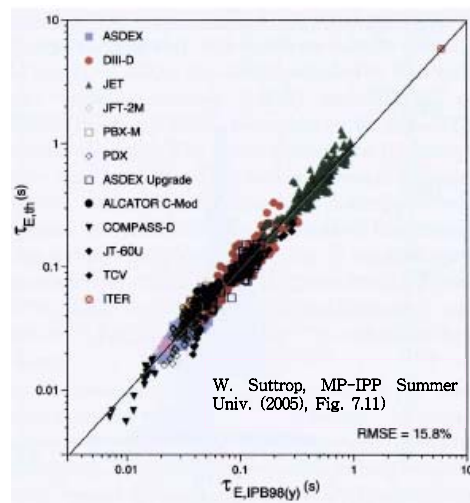


Figure 7.13: (a) Comparison of electron pressure profiles in L-mode and H-mode, (b) Experimental power threshold for barrier formation (W. Suttrop, MP-IPP Summer Univ. (2005))

### International confinement database for a large set of tokamaks

$$\tau_{E,IPB98(y)} = 0.0365 I(MA)^{0.97} B(T)^{0.08} P_{heat}(MW)^{-0.63} N_e(10^{19} m^{-3})^{0.41} M^{0.2} R(m)^{1.99} \epsilon^{0.23} \kappa^{0.67}$$



(cf) Different scaling laws caused by differences in geometry, impurity levels,  $\Delta B$ , additional heatings, start-up & operation

(cf) Murakami empirical scaling law for density limitation for OHT (1976):

$$\bar{n}_e (m^{-3}) < (1-2) \times 10^{20} \frac{B_\phi(T)}{R(m) q_a}$$

Greewald density limit:

$$\bar{n}/10^{20} = \kappa J(MA/m^2) = I_{MA}/(\pi ab)$$

Hugill-Diagram

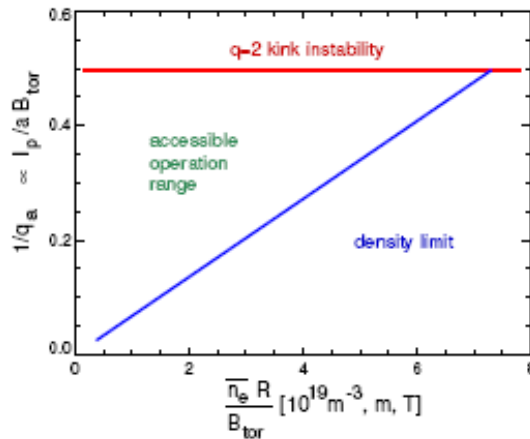


Figure 7.17: Hugill-Diagram showing the basic tokamak operational range. An upper limit to the plasma current is given by the onset of  $n = 1, m = 2$  kink instability for  $q = 2$  at the plasma edge. A density limit exists at  $\bar{n}_e \propto B_t / (Rq) \propto I_p / a^2$  (W. Suttrop, MP-IPP Summer Univ. (2005))

Advanced tokamak (AT) operation (Reversed shear by continuous bootstrap CD):

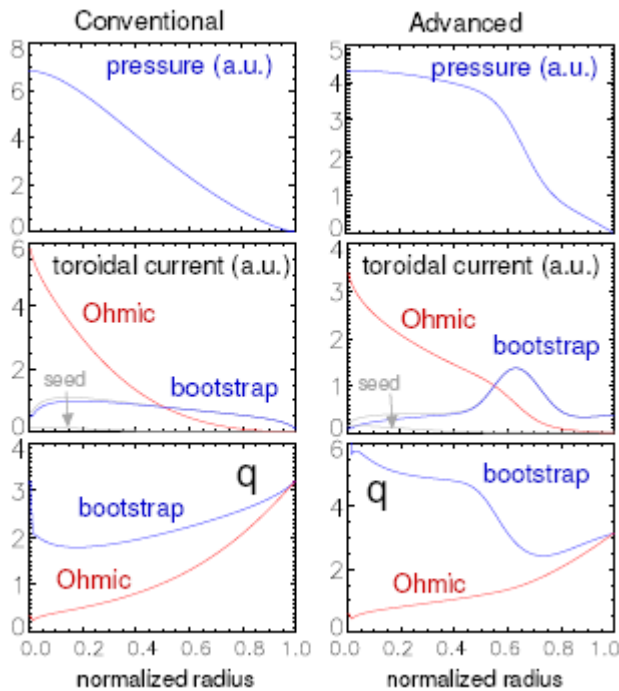


Figure 7.16: Profiles of plasma pressure  $p$ , current density  $j$  and safety factor  $q$  comparing conventional and advanced scenarios. A conventional plasma has a monotonously rising  $q$  profile. Flat or reversed reversed  $q$  profile can lead to a transport barrier. In an ideal advanced scenario, the resulting steep pressure gradient creates a bootstrap current that maintains the  $q$ -profile non-inductively in steady state. Ohmic and bootstrap contributions to  $j$  and  $q$  are shown separately. (W. Suttrop, MP-IPP Summer Univ. (2005))

**Tokamaks in the World with  $I_{\text{plasma}} \geq 1 \text{ MA}$**

Device	Institution/ Country	Geometry R (m) / a (m)	Elong- ation b/a	Plasma current $I_p$ (MA)	Toroid- al field $B_t$ (T)	Heating Power $P_{\text{add}}$ (MW)	P/R (MW/m)	Max.press. ratio $\beta_{\text{max}}$ (%)	Start of Operation
TFTR	PPPL/USA	2.5/0.85	1	2.5	5.2	32	13	1.7	1982 — 1998
JET	JET J.U./EU	2.96/1.2	1.8	7/ (6 div.)	3.5	>50	19.5	4.4	1983
DIII-D	GA/USA	1.67/0.67	2	3.5	2.2	20	12	5.1	1986
TORÉ SUPRA	EUR-CEA/ France	2.4/0.75	1	1.7	4.5	22	9.3	1.5	1988
T-15	Kurchatov/ Russ .Fed.	2.4/0.7	1	2	4	15	6.2	2.1	1988
FTU	EUR-ENEA/ Italy	0.93/0.31	1	$\leq 1.6$	8	$\leq 9$	8.5	1.9	1990
JT-60U	JAERI/ Japan	3.4/1.1	1.6	6	4.2	30	8.8	3.9	1991
ASDEX Upgrade	EUR-IPP/ Germany	1.65/0.5	1.6	$\leq 1.6$	3.9	$\leq 28$	$\leq 16$	2.5	1991
TCV	EUR-Suisse, CRPP/ Switzerland	0.88/0.24	3	<1.2	1.4	4.5	5.1	8.3	1992
ITER (project)	Europe, Japan, Russ. Fed., USA	8.1/2.8	1.6	21	5.7	100	39.5	3.2	(2008)
SNUT-79	SNU	0.65/0.15	1	0.12	3	OH			~85
KAIST-T (PRETEXT)	KAIST	0.53/0.15	1	0.08	10	OH			90
KT-1	KAERI	0.27/0.05	1	0.1	4	OH			86
KT-2	KAERI	1.4/0.25	1.8	0.5	3	ICRH/ECRH			(98)
KSTAR	KBRI	1.8/0.5	2.0	2.0	3.5	NBI/ICRH			2005