# Mechanics and Design 

Chapter 1<br>Vectors and Tensors

School of Mechanical and Aerospace Engineering
Seoul National University

- Vectors, Vector Additions, etc.

Convention:

$$
\mathbf{a}=a_{x} \mathbf{i}+a_{y} \mathbf{j}+a_{z} \mathbf{k}
$$

Base vectors:

$$
\mathbf{i}, \mathbf{j}, \mathbf{k} \text { or } \mathbf{e}_{1}, \mathbf{e}_{2}, \mathbf{e}_{3}
$$

## Indicial notation:

$$
\begin{gathered}
a_{i}=\mathbf{a} \cdot \mathbf{e}_{\mathbf{i}} \\
\mathbf{a}=a_{i} \mathbf{e}_{\mathbf{i}}=a_{1} \mathbf{e}_{\mathbf{1}}+a_{2} \mathbf{e}_{2}+a_{3} \mathbf{e}_{3}
\end{gathered}
$$

- Vectors, Vector Additions, etc.


## Summation Convention:

Repeated index

$$
\begin{aligned}
a_{i} a_{i} & =a_{1} a_{1}+a_{2} a_{2}+a_{3} a_{3} \\
a_{k k} & =a_{11}+a_{22}+a_{33} \\
a_{i j} a_{i j} & =a_{11} a_{11}+a_{12} a_{12}+a_{13} a_{13} \\
& +a_{21} a_{21}+a_{22} a_{22}+a_{23} a_{23} \\
& +a_{31} a_{31}+a_{32} a_{32}+a_{33} a_{33}
\end{aligned}
$$

## - Scalar Product and Vector Product

## Scalar Product

The scalar product is defined as the product of the two magnitudes times the cosine of the angle between the vectors

$$
\begin{equation*}
\mathbf{a} \cdot \mathbf{b}=a b \cos \theta \tag{1-1}
\end{equation*}
$$

From the definition of eq. (1-1), we immediately have

$$
\begin{aligned}
(m \mathbf{a}) \cdot(n \mathbf{b}) & =m n(\mathbf{a} \cdot \mathbf{b}) \\
\mathbf{a} \cdot \mathbf{b} & =\mathbf{b} \cdot \mathbf{a} \\
\mathbf{a} \cdot(\mathbf{b}+\mathbf{c}) & =\mathbf{a} \cdot \mathbf{b}+\mathbf{a} \cdot \mathbf{c}
\end{aligned}
$$

The scalar product of two different unit base vectors defined above is zero, since $\cos \left(90^{\circ}\right)=0$, that is,

$$
\begin{aligned}
& \mathbf{i} \cdot \mathbf{j}=\mathbf{j} \cdot \mathbf{k}=\mathbf{k} \cdot \mathbf{i}=0 \\
& \mathbf{i} \cdot \mathbf{i}=\mathbf{j} \cdot \mathbf{j}=\mathbf{k} \cdot \mathbf{k}=1
\end{aligned}
$$

- Scalar Product and Vector Product

Then the scalar product becomes

$$
\begin{aligned}
\mathbf{a} \cdot \mathbf{b} & =\left(a_{x} \mathbf{i}+a_{y} \mathbf{j}+a \mathbf{k}\right) \cdot\left(b_{x} \mathbf{i}+b_{y} \mathbf{j}+b_{z} \mathbf{k}\right) \\
& =a_{x} b_{x}+a_{y} b_{y}+a_{z} b_{z}
\end{aligned}
$$

or in an indicial notation, we have

$$
\mathbf{a} \cdot \mathbf{b}=a_{i} b_{i}
$$

Vector Product (or cross product)

$$
\mathbf{c}=\mathbf{a} \times \mathbf{b}
$$

defined as a vector c perpendicular to both $\mathbf{a}$ and $\mathbf{b}$ in the sense that makes $\mathbf{a}, \mathbf{b}$, c a right-handed system. Magnitude of vector c is given by

$$
c=a b \sin (\theta)
$$

## - Scalar Product and Vector Product

In terms of components, vector product can be written as

$$
\begin{aligned}
\mathbf{a} \times \mathbf{b} & =\left(a_{x} \mathbf{i}+a_{y} \mathbf{j}+a_{z} \mathbf{k}\right) \times\left(b_{x} \mathbf{i}+b_{y} \mathbf{j}+b_{z} \mathbf{k}\right) \\
& =\left|\begin{array}{ccc}
\mathbf{i} & \mathbf{j} & \mathbf{k} \\
a_{x} & a_{y} & a_{z} \\
b_{x} & b_{y} & b_{z}
\end{array}\right| \\
& =\left(a_{y} b_{z}-a_{z} b_{y}\right) \mathbf{i}+\left(a_{z} b_{x}-a_{x} b_{z}\right) \mathbf{j}+\left(a_{x} b_{y}-a_{y} b_{x}\right) \mathbf{k}
\end{aligned}
$$



The vector product is distributive as

- Fig. 1.1 Definition of vector product

$$
\mathbf{a} \times(\mathbf{b}+\mathbf{c})=(\mathbf{a} \times \mathbf{b})+(\mathbf{a} \times \mathbf{c})
$$

but it is not associative as

$$
\mathbf{a} \times(\mathbf{b} \times \mathbf{c}) \neq(\mathbf{a} \times \mathbf{b}) \times \mathbf{c}
$$

## - Scalar Product and Vector Product

## Scalar Triple Product

$$
(\mathbf{a} \times \mathbf{b}) \cdot \mathbf{c}=\mathbf{a} \cdot(\mathbf{b} \times \mathbf{c})=\left|\begin{array}{lll}
a_{x} & a_{y} & a_{z} \\
b_{x} & b_{y} & b_{z} \\
c_{x} & c_{y} & c_{z}
\end{array}\right|
$$

Permutation Symbol $e_{m n r}$

$$
e_{m n r}=\left\{\begin{array}{c}
0 \text { when any two indices are equal } \\
+1 \text { when }(m, n, r) \text { is even purmutation of }(1,2,3) \\
-1 \text { when }(m, n, r) \text { is odd purmutation of }(1,2,3)
\end{array}\right.
$$

where
even permutations are $(1,2,3),(2,3,1)$, and $(3,1,2)$
odd permutations are $(1,3,2),(2,1,3)$, and $(3,2,1)$.
Using the permutation symbol, the vector product can be represented by

$$
\mathbf{a} \times \mathbf{b}=e_{p q r} a_{q} b_{r} \mathbf{i}_{p}
$$

## - Scalar Product and Vector Product

Kronecker Delta $\delta_{i j}$

$$
\delta_{p q}=\left\{\begin{array}{l}
1 \text { if } p=q \\
0 \text { if } p \neq q
\end{array}\right.
$$

Some examples are given below:

$$
\begin{gathered}
\delta_{i i}=3 \\
\delta_{i j} \delta_{i j}=\delta_{i i}=3 \\
u_{i} \delta_{i j}=u_{i}=u_{j} \\
T_{i j} \delta_{i j}=T_{i i}
\end{gathered}
$$

## - Rotation of Axes, etc

## Change of Orthogonal Basis

- Vector is independent of a coordinate system.
- However, the components of a vector change when the coordinate system changes.

Let $x_{i}$ and $\bar{X}_{i}$ be the two coordinate systems, as shown in Fig. 1-2, which have the same origins.
Also, let the orientation of the two coordinate systems is given by the direction cosines as

|  | $\overline{\mathbf{i}}_{1}$ | $\overline{\mathbf{i}}_{2}$ | $\overline{\mathbf{i}}_{3}$ |
| :---: | :---: | :---: | :---: |
| $\mathbf{i}_{1}$ | $a_{1}^{1}$ | $a_{2}^{1}$ | $a_{3}^{1}$ |
| $\mathbf{i}_{2}$ | $a_{1}^{2}$ | $a_{2}^{2}$ | $a_{3}^{2}$ |
| $\mathbf{i}_{3}$ | $a_{1}^{3}$ | $a_{2}^{3}$ | $a_{3}^{3}$ |



- Fig. 1.2 Coordinate transformation
- Rotation of Axes, etc

Under this system,

$$
\overline{\mathbf{i}_{\mathrm{r}}}=a_{r}^{s} \mathbf{i}_{\mathrm{s}} \quad \text { and } \quad \mathbf{i}_{\mathrm{s}}=a_{r}^{s} \overline{\mathbf{i}_{\mathbf{i}}}
$$

Since the unit vectors are orthogonal, we have

$$
\overline{\mathbf{i}}_{\mathrm{p}} \cdot \overline{\mathbf{i}}_{q}=\delta_{p q} \quad \text { and } \quad \mathbf{i}_{\mathrm{p}} \cdot \mathbf{i}_{q}=\delta_{p q}
$$

From these

$$
\begin{aligned}
& \overline{\mathbf{i}}_{\mathrm{p}} \cdot \overline{\mathbf{i}}_{\mathrm{q}}= a_{p}^{s} a_{q}^{r} \mathbf{i}_{\mathrm{s}} \cdot \mathbf{i}_{\mathrm{r}}=a_{p}^{s} a_{q}^{r} \delta_{r s}=a_{p}^{s} a_{q}^{s}=\delta_{p q} \\
& \text { i.e., } \quad a_{p}^{s} a_{q}^{s}=\delta_{p q}
\end{aligned}
$$

Similarly, we get

$$
a_{r}^{m} a_{r}^{n}=\delta_{m n}
$$

## 2D coodinate transformation

$$
\left[\begin{array}{l}
u_{x} \\
u_{y}
\end{array}\right]=\left[\begin{array}{cc}
\cos \theta & -\sin \theta \\
\sin \theta & \cos \theta
\end{array}\right]\left[\begin{array}{l}
\bar{u}_{x} \\
\bar{u}_{y}
\end{array}\right] \quad \text { or } \quad \mathbf{u}=A \overline{\mathbf{u}}
$$

The inverse of the above transformation

## equation becomes

$$
\left[\begin{array}{l}
\bar{u}_{x} \\
\bar{u}_{y}
\end{array}\right]=\left[\begin{array}{cc}
\cos \theta & \sin \theta \\
-\sin \theta & \cos \theta
\end{array}\right]\left[\begin{array}{l}
u_{x} \\
u_{y}
\end{array}\right] \quad \text { or } \overline{\mathbf{u}}=A^{-1} \mathbf{u}=A^{T} \mathbf{u}
$$



Coordinate Transformation of Vector

$$
\begin{aligned}
& \text { Let, } \quad A=\left[\begin{array}{lll}
a_{1}^{1} & a_{2}^{1} & a_{3}^{1} \\
a_{1}^{2} & a_{2}^{2} & a_{3}^{2} \\
a_{1}^{3} & a_{2}^{3} & a_{3}^{3}
\end{array}\right] \text { Then, } \mathbf{V}=A \overline{\mathbf{v}} \text { or } \overline{\mathbf{v}}=A^{T} \mathbf{v} \\
& A=\left[\begin{array}{lll}
\overline{\mathbf{a}}_{1} & \overline{\mathbf{a}}_{2} & \overline{\mathbf{a}}_{3}
\end{array}\right]
\end{aligned}
$$

Note that

$$
A^{T} A=A A^{T}=1
$$

## Second Order Tensors

The second order tensor may be expressed by tensor product or open product of two vectors. Let

$$
\mathbf{a}=a_{1} \mathbf{e}_{1}+a_{2} \mathbf{e}_{2}+a_{3} \mathbf{e}_{3} \quad \text { and } \quad \mathbf{b}=b_{1} \mathbf{e}_{1}+b_{2} \mathbf{e}_{2}+b_{3} \mathbf{e}_{3}
$$

$$
\begin{aligned}
& \mathrm{T}=\mathbf{a b}=\mathbf{a} \otimes \mathbf{b}=\left(a_{1} \mathbf{e}_{1}+a_{2} \mathbf{e}_{2}+a_{3} \mathbf{e}_{3}\right) \otimes\left(b_{1} \mathbf{e}_{1}+b_{2} \mathbf{e}_{2}+b_{3} \mathbf{e}_{3}\right) \\
&= a_{1} b_{1} \mathbf{e}_{1} \otimes \mathbf{e}_{1}+a_{1} b_{2} \mathbf{e}_{1} \otimes \mathbf{e}_{2}+a_{1} b_{3} \mathbf{e}_{1} \otimes \mathbf{e}_{3} \\
&+a_{2} b_{1} \mathbf{e}_{2} \otimes \mathbf{e}_{1}+a_{2} b_{2} \mathbf{e}_{2} \otimes \mathbf{e}_{2}+a_{2} b_{3} \mathbf{e}_{2} \otimes \mathbf{e}_{3} \\
&+a_{3} b_{1} \mathbf{e}_{3} \otimes \mathbf{e}_{1}+a_{3} b_{2} \mathbf{e}_{3} \otimes \mathbf{e}_{2}+a_{3} b_{3} \mathbf{e}_{3} \otimes \mathbf{e}_{3} \\
& \text { Or } \\
& \mathrm{T}= T_{11} \mathbf{e}_{1} \otimes \mathbf{e}_{1}+T_{12} \mathbf{e}_{1} \otimes \mathbf{e}_{2}+T_{13} \mathbf{e}_{1} \otimes \mathbf{e}_{3}+ \\
&+T_{21} \mathbf{e}_{2} \otimes \mathbf{e}_{1}+T_{22} \mathbf{e}_{2} \otimes \mathbf{e}_{2}+T_{23} \mathbf{e}_{2} \otimes \mathbf{e}_{3} \\
&+T_{31} \mathbf{e}_{3} \otimes \mathbf{e}_{1}+T_{32} \mathbf{e}_{3} \otimes \mathbf{e}_{2}+T_{33} \mathbf{e}_{3} \otimes \mathbf{e}_{3}
\end{aligned}
$$

Here, we may consider $\mathbf{e}_{i} \otimes \mathbf{e}_{j} \neq \mathbf{e}_{j} \otimes \mathbf{e}_{i}$ as a base of the second order tensor, and $T_{i j}$ is a component of the second order tensor T .

It can be seen that the second order tensor map a vector to another vector, that is,

$$
\begin{aligned}
\mathbf{u}=T \cdot \mathbf{v}= & \left(T_{11} \mathbf{e}_{1} \otimes \mathbf{e}_{1}+T_{12} \mathbf{e}_{1} \otimes \mathbf{e}_{2}+T_{13} \mathbf{e}_{1} \otimes \mathbf{e}_{3}+\right. \\
& +T_{21} \mathbf{e}_{2} \otimes \mathbf{e}_{1}+T_{22} \mathbf{e}_{2} \otimes \mathbf{e}_{2}+T_{23} \mathbf{e}_{2} \otimes \mathbf{e}_{3} \\
& \left.+T_{31} \mathbf{e}_{3} \otimes \mathbf{e}_{1}+T_{32} \mathbf{e}_{3} \otimes \mathbf{e}_{2}+T_{33} \mathbf{e}_{3} \otimes \mathbf{e}_{3}\right) \cdot\left(v_{1} \mathbf{e}_{1}+v_{2} \mathbf{e}_{2}+v_{3} \mathbf{e}_{3}\right) \\
= & \left(T_{11} v_{1}+T_{12} v_{2}+T_{13} v_{3}\right) \mathbf{e}_{1}+\left(T_{21} v_{1}+T_{22} v_{2}+T_{23} v_{3}\right) \mathbf{e}_{2} \\
& +\left(T_{31} v_{1}+T_{32} v_{2}+T_{33} v_{3}\right) \mathbf{e}_{3} \\
= & T_{i j} v_{j} \mathbf{e}_{i}
\end{aligned}
$$

Symmetric Tensors and Skew tensors

Symmetric tensor

$$
T_{i j}=T_{j i}
$$

Skew or
Antisymmetric tensor $\longrightarrow T_{i j}=-T_{j i}$

## Rotation of axes, change of tensor components

Let $u_{i}=T_{i p} v_{p}$ and $\bar{u}_{i}=\bar{T}_{i p} \bar{v}_{p}$
Where $u_{i}$ and $\bar{u}_{i}$ are the same vector but decomposed into two different coordinate systems, $x_{i}$ and $\bar{x}_{i}$. The same applies to $v_{i}$ and $\bar{v}_{i}$. Then by the transformation of vector, we get

$$
\begin{aligned}
\bar{u}_{i}=a_{i}^{j} u_{j} & =a_{i}^{j} T_{j q} v_{q} \\
& =a_{i}^{j} T_{j q} a_{p}^{q} \bar{v}_{p}
\end{aligned}
$$

Therefore,

$$
\bar{T}_{i p}=a_{i}^{j} a_{p}^{q} T_{j q}
$$

In matrix form, we have,

$$
\bar{T}=A^{T} T A \quad \text { or, } \quad T=A \bar{T} A^{T}
$$

## - Rotation of Axes, etc

Scalar product of two tensors

$$
\begin{aligned}
& \mathrm{T}: \mathrm{U}=T_{i j} U_{i j} \\
& \mathrm{~T} \cdot \cdot \mathrm{U}=T_{i j} U_{j i}
\end{aligned}
$$

If we list all the terms of tensor product, we get

$$
\begin{aligned}
\mathrm{T}: \mathrm{U}= & T_{i j} U_{i j} \\
= & T_{11} U_{11}+T_{12} U_{12}+T_{13} U_{13} \\
& +T_{21} U_{21}+T_{22} U_{22}+T_{23} U_{23} \\
& +T_{31} U_{31}+T_{32} U_{32}+T_{33} U_{33}
\end{aligned}
$$

The product of two second-order tensors
$\mathrm{T} \cdot \mathrm{U}$
$(\mathrm{T} \cdot \mathrm{U}) \cdot \mathrm{v}=\mathrm{T} \cdot(\mathrm{U} \cdot \mathrm{v})$
If $\mathrm{P}=\mathrm{T} \cdot \mathrm{U}$, then $P_{i j}=T_{i k} U_{k j}$

## - Rotation of Axes, etc

## The Trace

## Definition : $\operatorname{tr}(\mathrm{T})=T_{k k}$

Note that
(1) $\mathrm{A} \cdot \mathrm{B}=\operatorname{tr}(\mathrm{A} \cdot \mathrm{B})$
(2) $\mathrm{A}: \mathrm{B}=\operatorname{tr}\left(\mathrm{A} \cdot \mathrm{B}^{T}\right)=\operatorname{tr}\left(\mathrm{A}^{\mathrm{T}} \cdot \mathrm{B}\right)$
(3) $\operatorname{tr}(\mathrm{A} \cdot \mathrm{B})=\operatorname{tr}(\mathrm{B} \cdot \mathrm{A})$
(4) $\operatorname{tr}(\mathrm{A} \cdot \mathrm{B} \cdot \mathrm{C})=\operatorname{tr}(\mathrm{B} \cdot \mathrm{C} \cdot \mathrm{A})=\operatorname{tr}(\mathrm{C} \cdot \mathrm{A} \cdot \mathrm{B}) \quad$ Cyclic property of trace

Proof $>$
(1) $\mathrm{A} \cdot \mathrm{B}=\operatorname{tr}(\mathrm{A} \cdot \mathrm{B})$

$$
\text { Let, } \mathrm{P}=\mathrm{A} \cdot \mathrm{~B} \text {, i.e., } P_{i j}=A_{i k} B_{k j}
$$

$$
\text { Then, } \operatorname{tr}(\mathrm{P})=P_{i i}=A_{i k} B_{k i}=\mathrm{A} \cdot \cdot \mathrm{~B}
$$

(2) $\mathrm{A}: \mathrm{B}=\operatorname{tr}\left(\mathrm{A} \cdot \mathrm{B}^{T}\right)=\operatorname{tr}\left(\mathrm{A}^{\mathrm{T}} \cdot \mathrm{B}\right)$

$$
\begin{aligned}
& \text { Let, } \mathrm{P}=\mathrm{A} \cdot \mathrm{~B}^{\mathrm{T}}, \text { i.e., } \quad P_{i j}=A_{i k} B_{j k} \\
& \text { Then, } \operatorname{tr}\left(\mathrm{A} \cdot \mathrm{~B}^{T}\right)=P_{i i}=A_{i k} B_{i k}=\mathrm{A}: \mathrm{B}
\end{aligned}
$$

## - Review of Elementary Matrix Concepts

## Eigenvectors and eigenvalues of a matrix

Linear transformation, $y=M x$, associates, to each point $P\left(x_{1}, x_{2}, x_{3}\right)$, another point $Q\left(y_{1}, y_{2}, y_{3}\right)$. Also, it associates to any other point $\left(r x_{1}, r x_{2}, r x_{3}\right)$ on the line OP , another point ( $r y_{1}, r y_{2}, r y_{3}$ ) on the line OQ.
So we may consider the transformation to be a transformation of the line OP into the line OQ.

Now any line transformed into itself is called an eigenvector of the matrix M. That is,

$$
\mathbf{M x}=\lambda \mathbf{x}=\lambda \mathbf{I} \mathbf{x}
$$

A nontrivial solution will exist if and only if the determinant vanishes:

$$
|(\mathbf{M}-\lambda \mathbf{I})|=0
$$

Note that $3 \times 3$ determinant is expanded, it will be a cubic polynominal equation with real coefficients. The roots of this equation are called the eigenvalues of the matrix. Upon solving the equation, some of the roots could be complex numbers.


Linear Transformation
$M$ maps $\mathbf{v}$ into the same vector $\mathbf{v}$.


## - Review of Elementary Matrix Concepts

## A real symmetric matrix has only real eigenvalues.

If there were a complex root $\lambda$, then its complex conjugate $\bar{\lambda}$ is also a root. Therefore,

$$
\begin{aligned}
\mathbf{M} \mathbf{x} & =\lambda \mathbf{x} \\
\overline{\mathbf{M}} \mathbf{x} & =\bar{\lambda} \mathbf{x}
\end{aligned}
$$

These equations can be written as

$$
\begin{aligned}
\overline{\mathbf{x}}^{T} \mathbf{M x} & =\lambda \overline{\mathbf{x}}^{T} \mathbf{x} \\
\mathbf{x}^{T} \mathbf{M} \overline{\mathbf{x}} & =\bar{\lambda} \mathbf{x}^{T} \overline{\mathbf{x}}
\end{aligned}
$$

Note that $\overline{\mathbf{x}}^{T} \mathbf{x}=x_{k} \bar{x}_{k}=\mathbf{x}^{T} \overline{\mathbf{x}}$ and since M is symmetric, we get

$$
\begin{aligned}
\overline{\mathbf{x}}^{T} \mathbf{M x} & =M_{i j} \bar{x}_{i} x_{j} \\
& =M_{j i} \bar{x}_{j} x_{i} \quad \text { (by interchanging the dummy indices) } \\
& =M_{i j} x_{i} \bar{x}_{j}=\mathbf{x}^{T} \mathbf{M} \overline{\mathbf{x}} \text { (by symmetry of M) }
\end{aligned}
$$

## - Review of Elementary Matrix Concepts

Subtracting, we get

$$
(\lambda-\bar{\lambda}) \overline{\mathbf{x}}^{T} \mathbf{x}=0
$$

Since $x$ is nontrivial, $\overline{\mathbf{x}}^{T} \mathbf{x} \neq 0$. Therefore, we should have

$$
\lambda=\bar{\lambda}
$$

So that $\lambda$ must be real.

> We can obtain the eigenvector associated to each eigenvalue by substituting each eigenvalue into the matrix equation.
> When eigenvalues are all distinct : ??
> Two of the eigenvalues are equal : ??
> All of the eigenvalues are equal : ??

