

Introduction to Nuclear Fusion

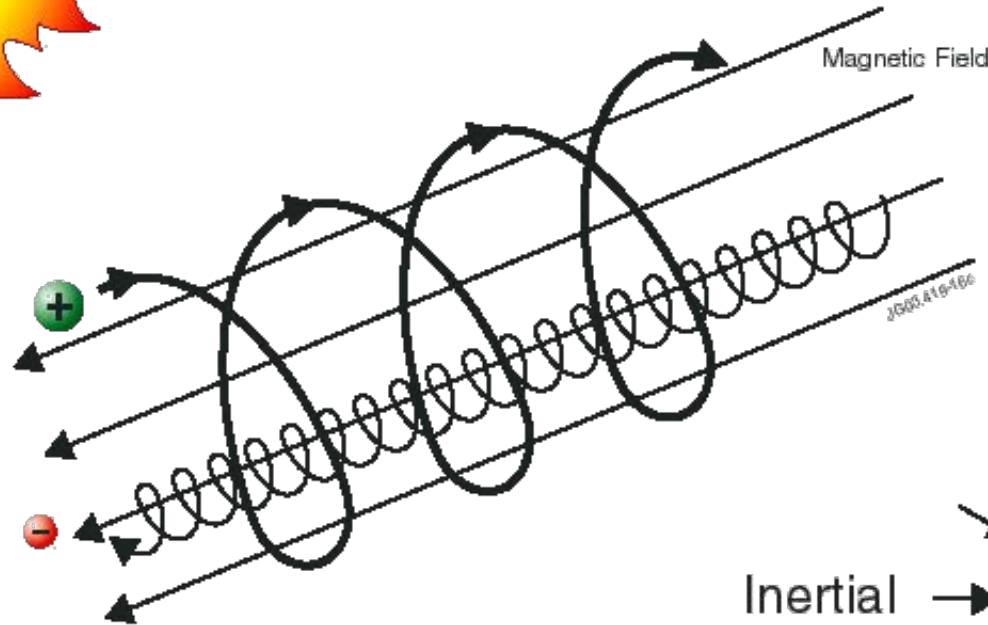
Prof. Dr. Yong-Su Na

To build a sun on earth



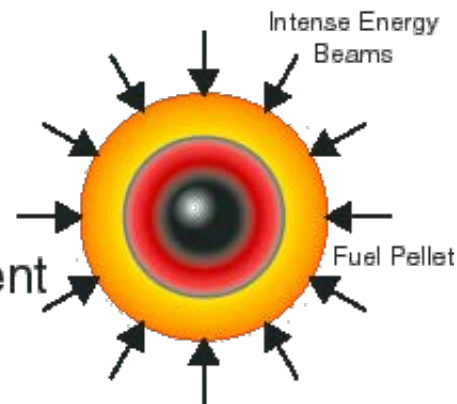
Gravitational
Confinement

Magnetic Confinement



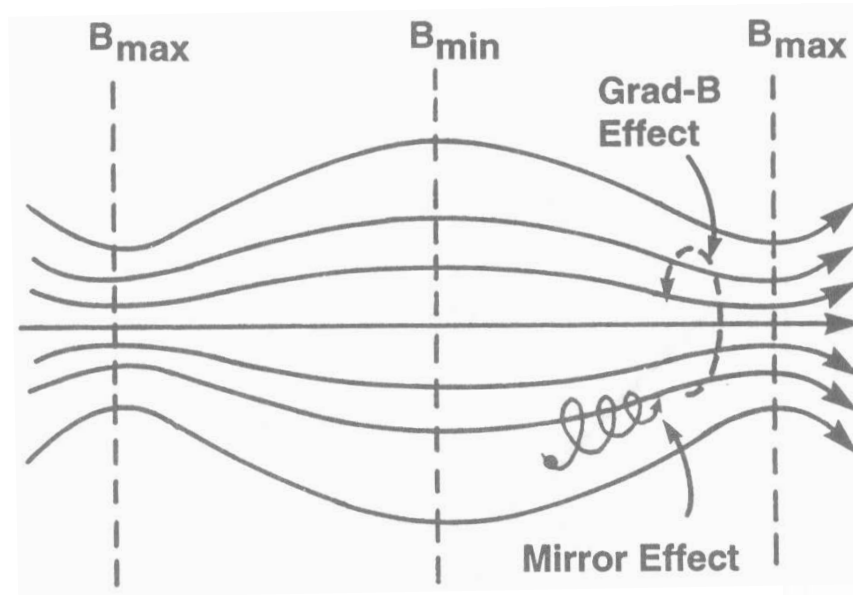
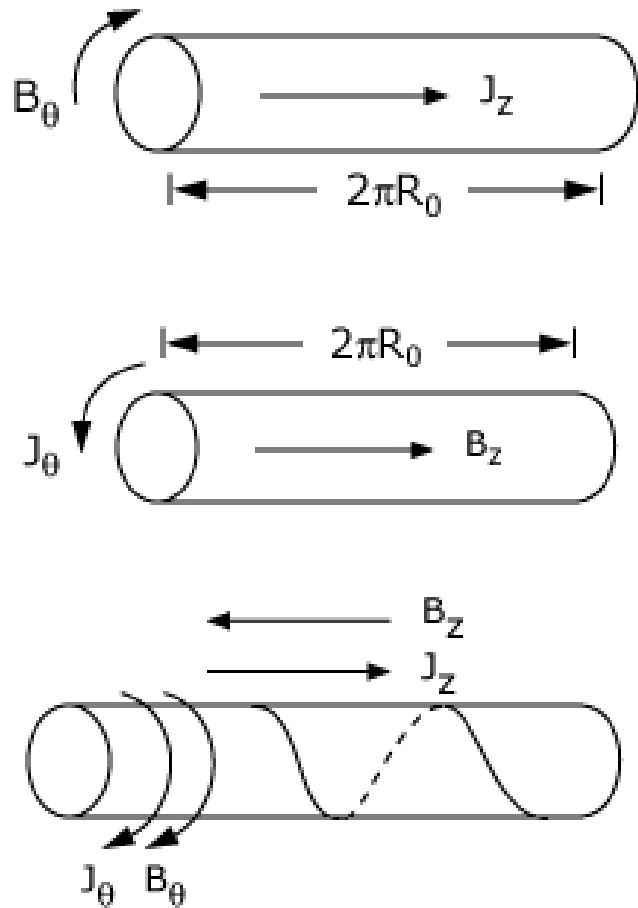
- Open magnetic confinement
- Closed magnetic confinement

Inertial
Confinement

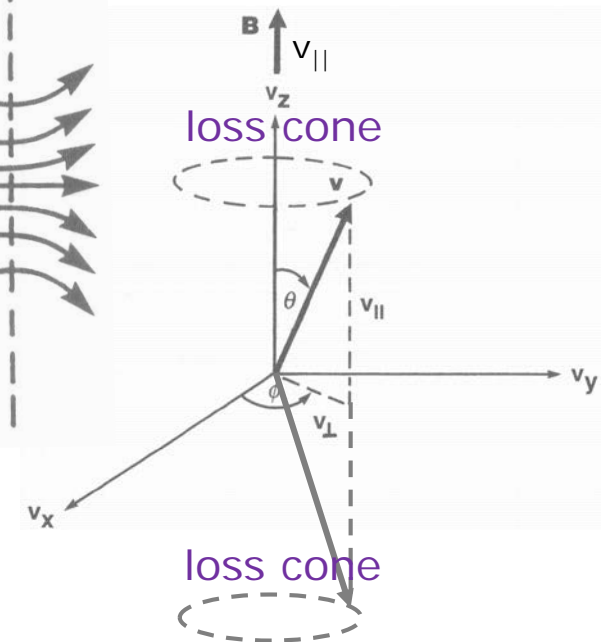


What is closed magnetic confinement?

Open Magnetic System

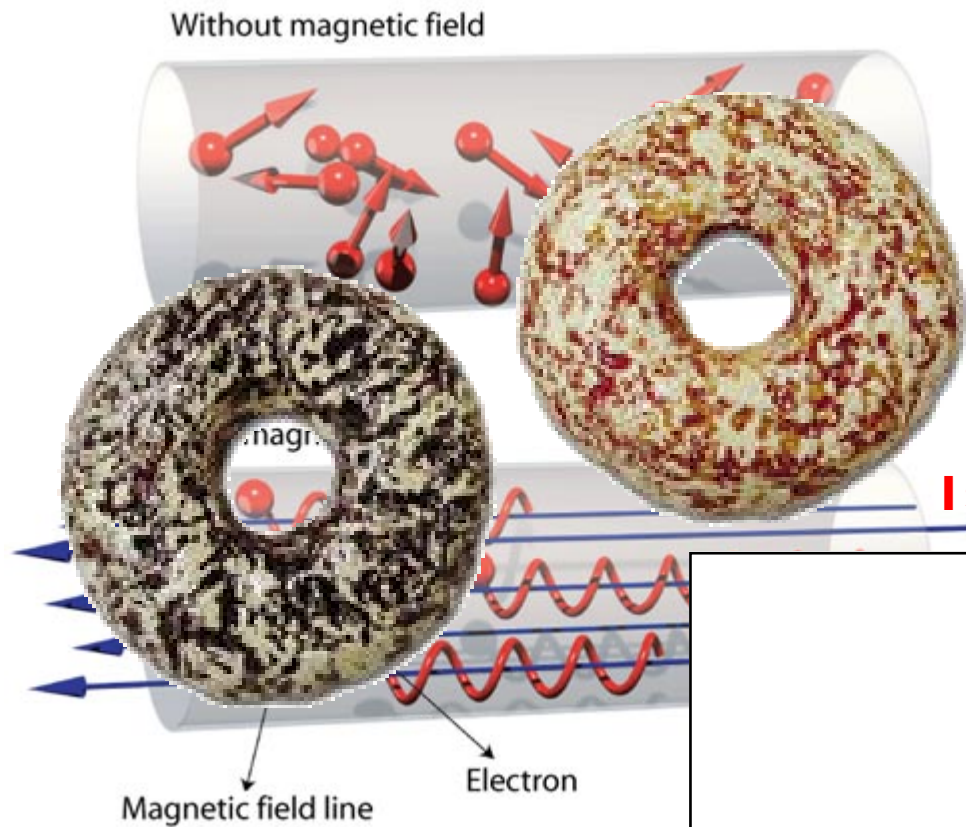


$$\sin^2 \theta \geq \frac{B_{\min}}{B_{\max}}$$



- Suffering from end losses

Open Magnetic System

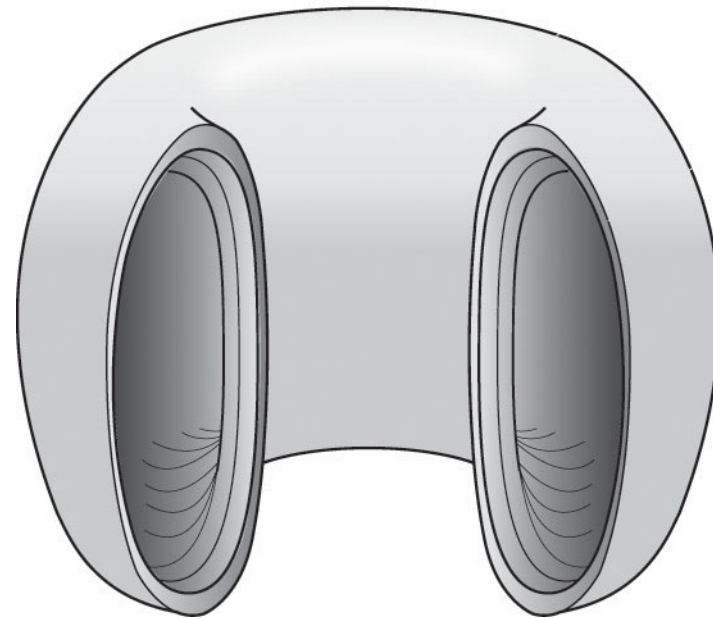
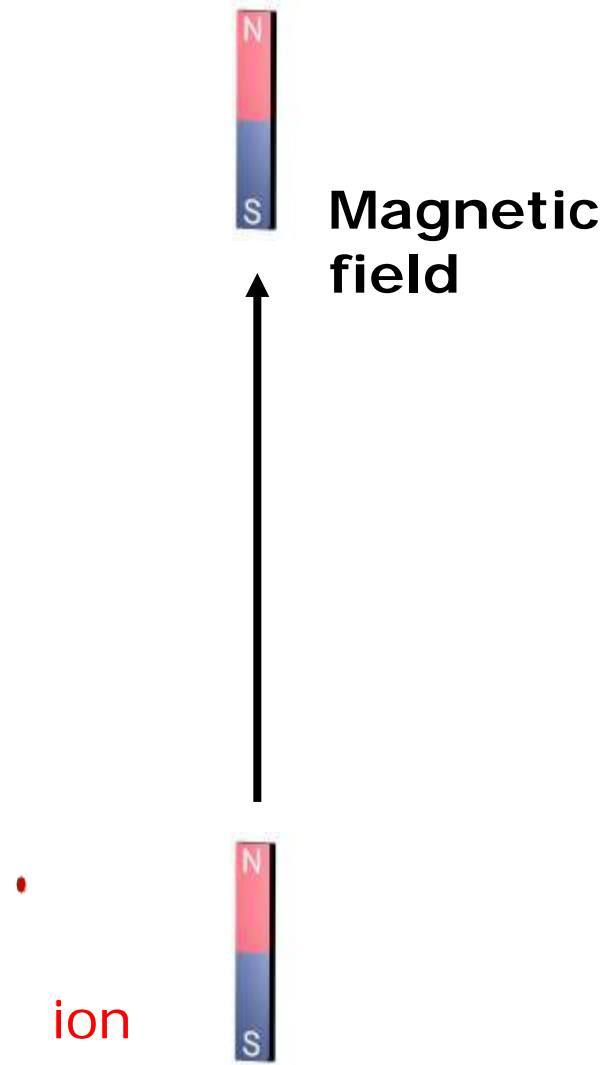


Is this motion realistic?



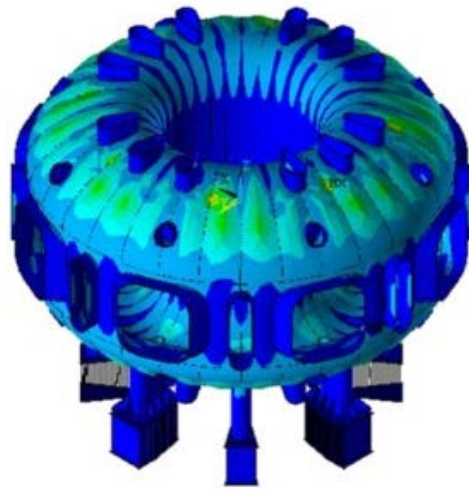
Magnetic field

Closed Magnetic System

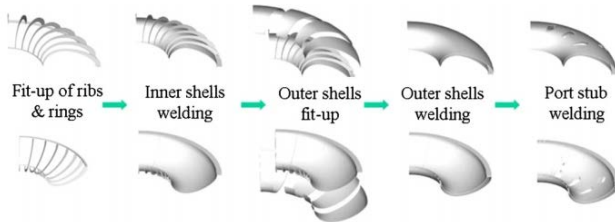


Donut-shaped vacuum vessel

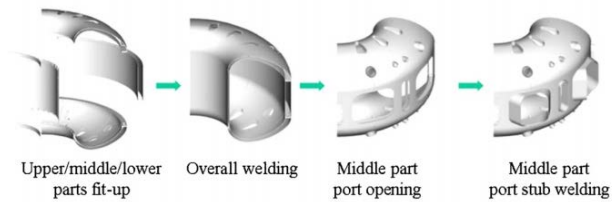
Closed Magnetic System



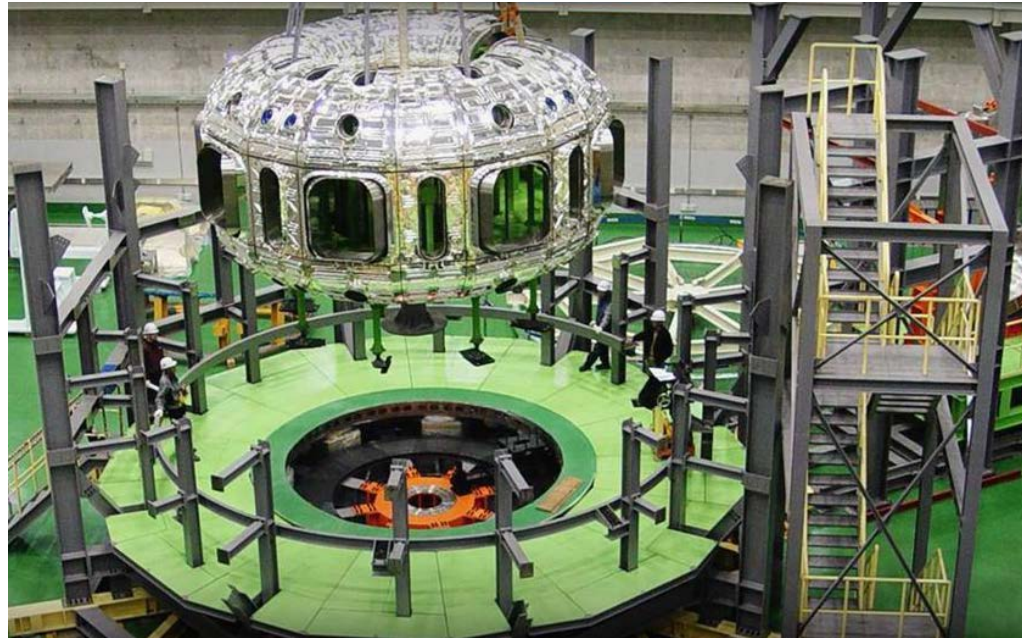
ANSYS 5.6
 FEB 27 2001
 11:10:55
 NODAL SOLUTION
 STEP=9999
 SINT (AVG)
 MIDDLE
 PowerGraphics
 EFACET=1
 AVRES=Max
 DMX =3.005
 SMN =.012972
 SMX =204.413
 .012972
 22.724
 45.435
 68.146
 90.857
 113.568
 136.279
 158.991
 181.702
 204.413



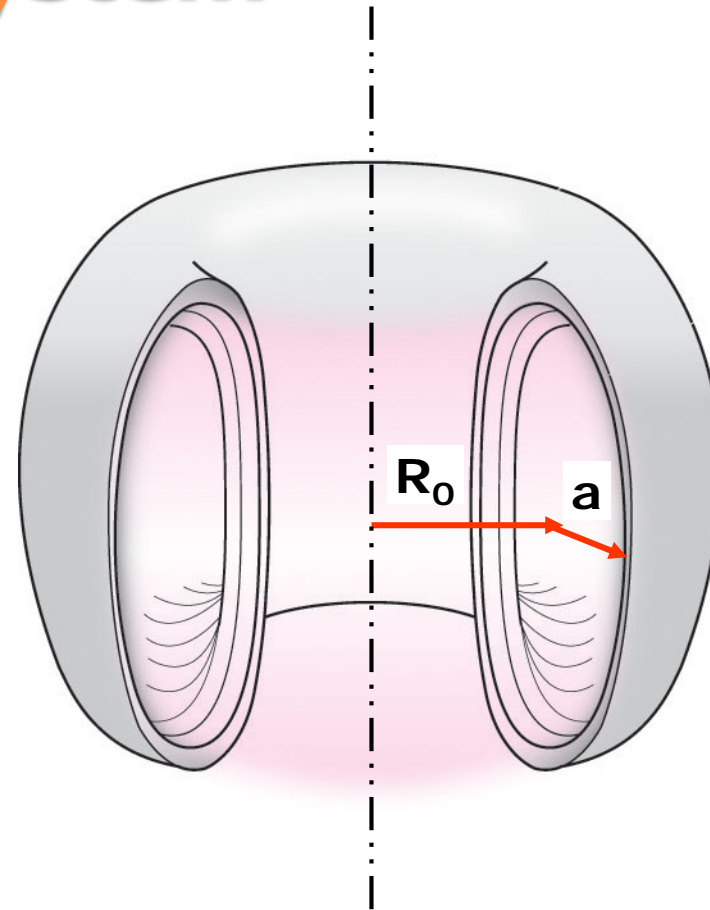
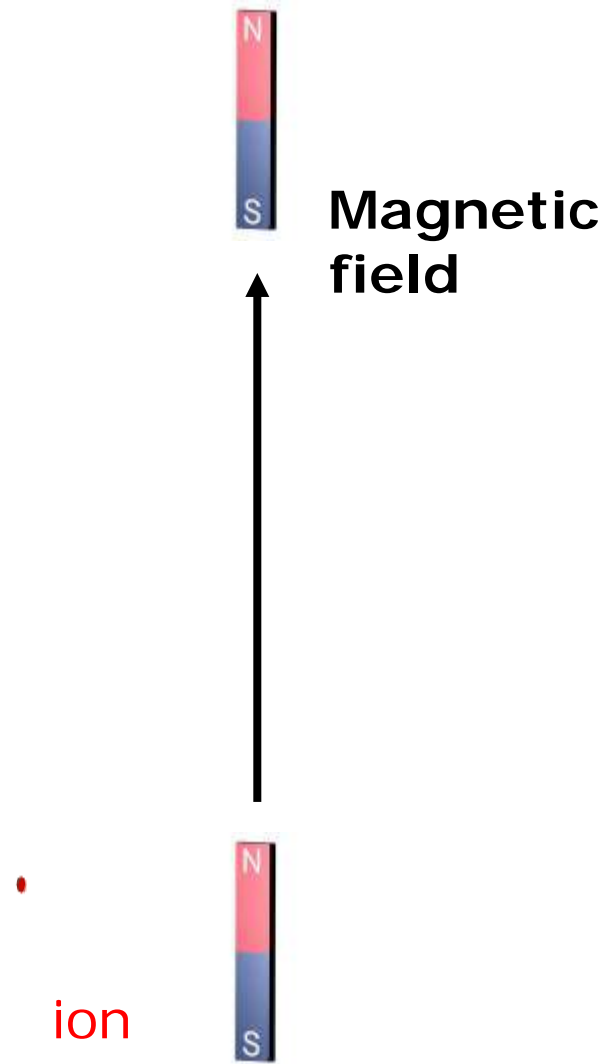
(a) Upper/lower parts fabrication procedures



(b) Quadrant fabrication procedure

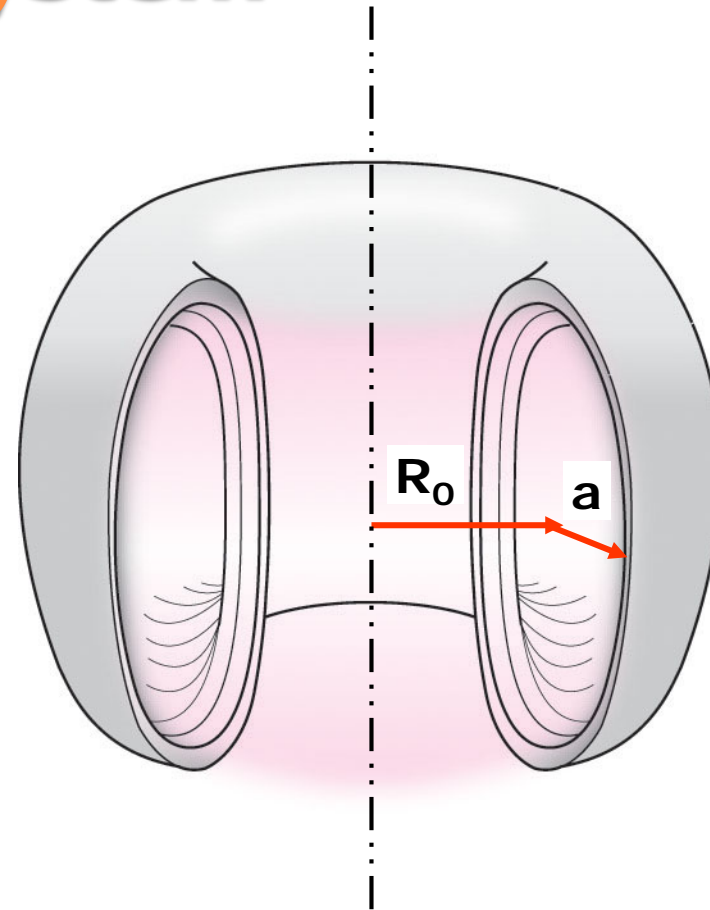
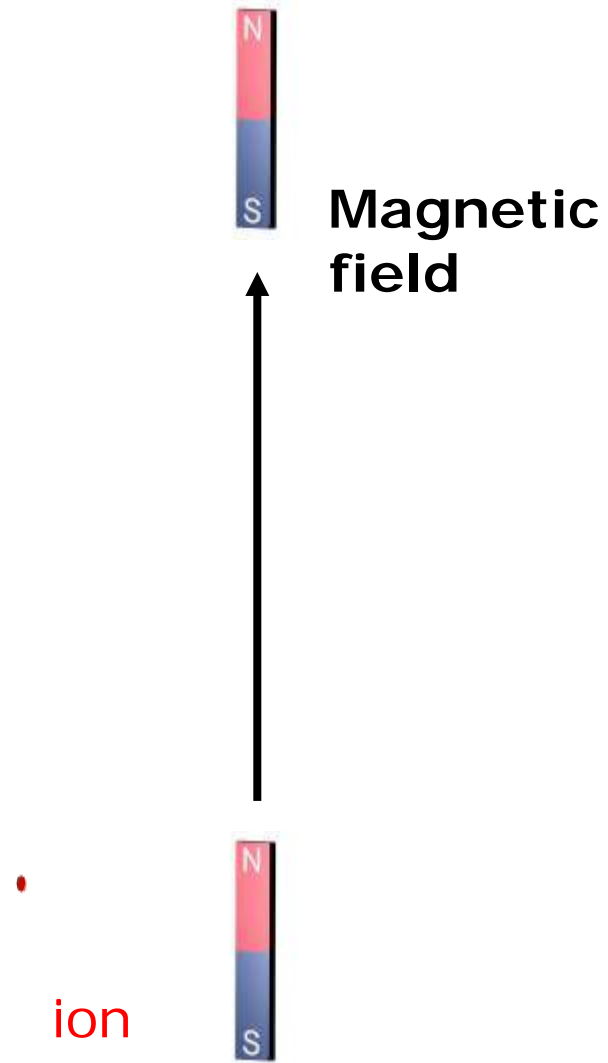


Closed Magnetic System



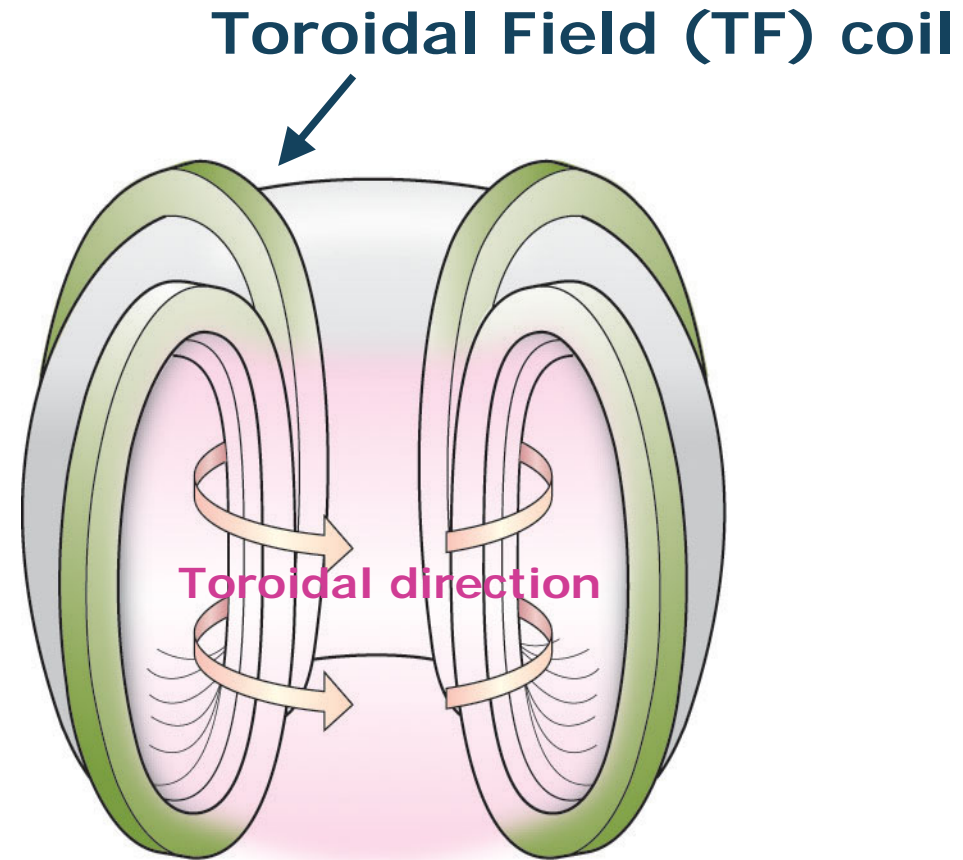
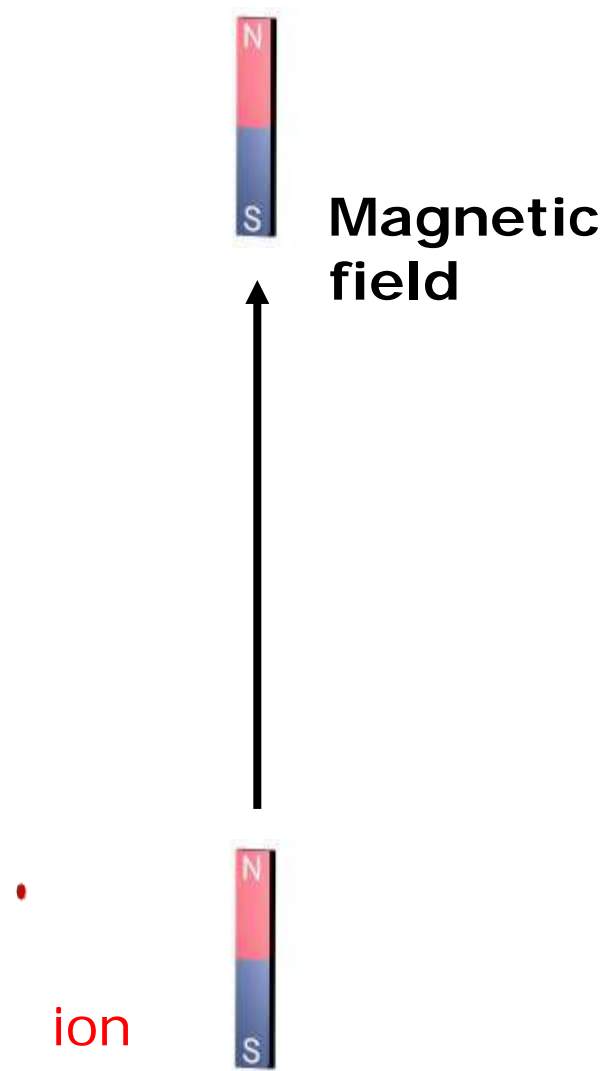
Plasma needs to be confined
 $R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$ in KSTAR

Closed Magnetic System



Plasma needs to be confined
 $R_0 = 6.2 \text{ m}$, $a = 2.0 \text{ m}$ in ITER

Closed Magnetic System



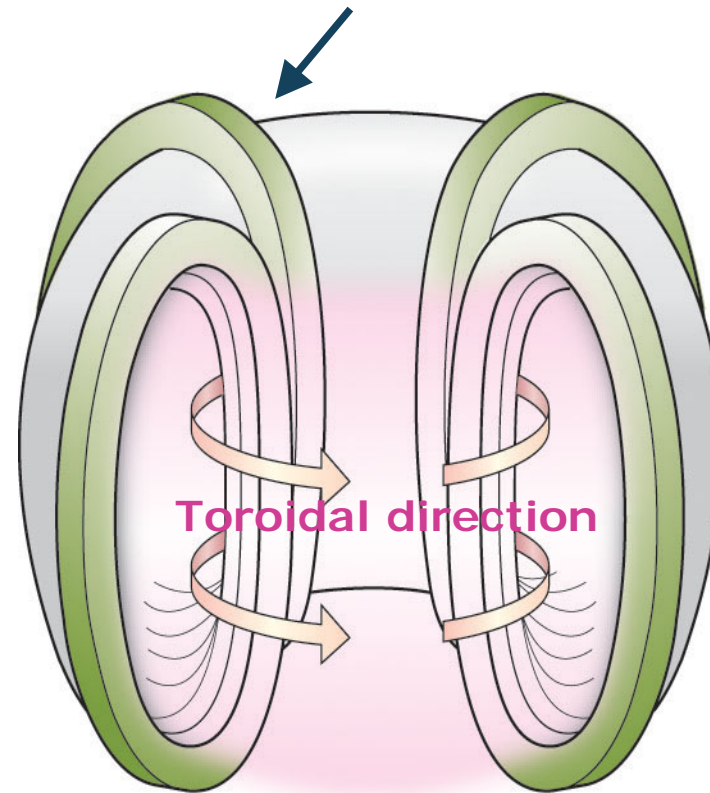
Applying toroidal magnetic field
3.5 T in KSTAR, 5.3 T in ITER

Closed Magnetic System



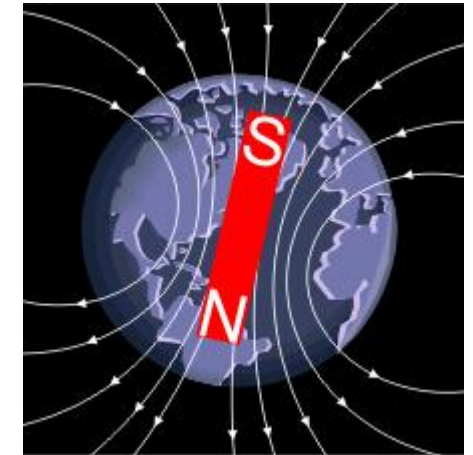
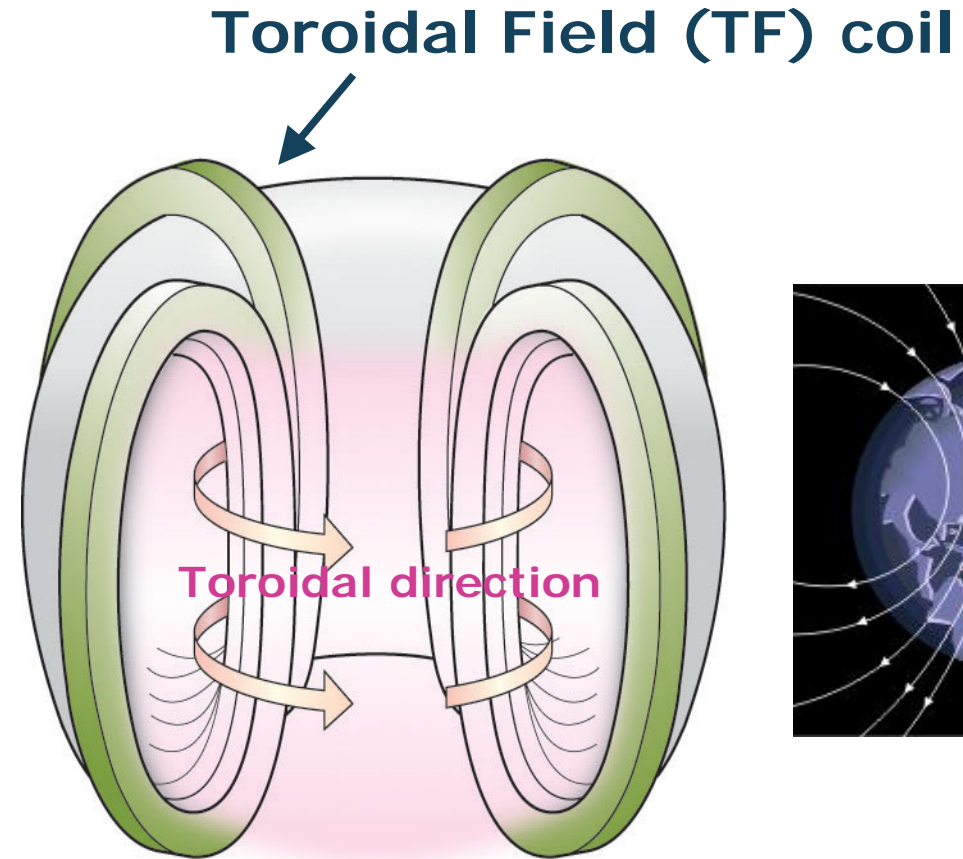
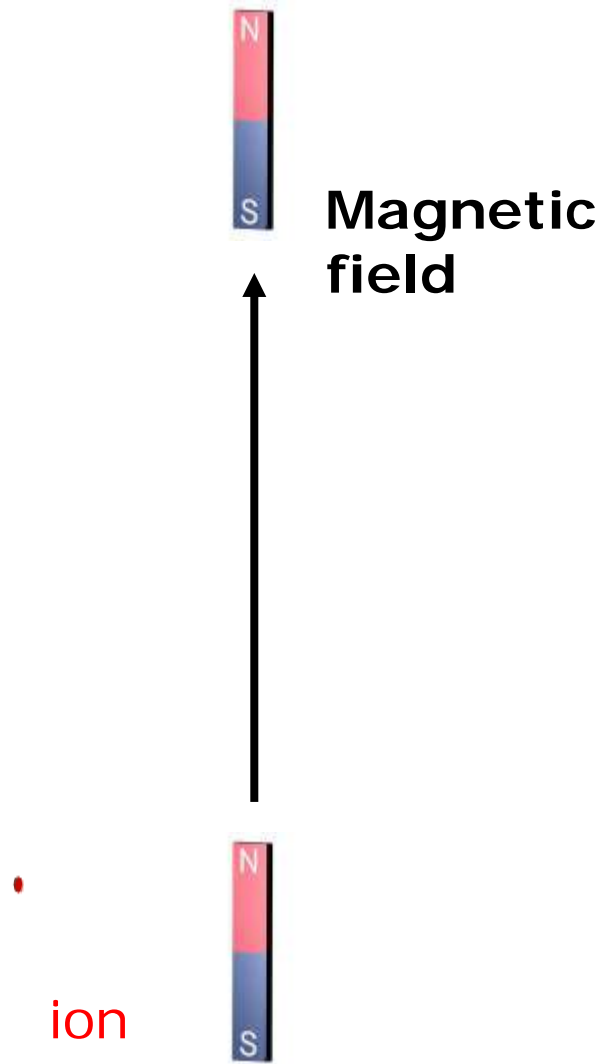
KSTAR

Toroidal Field (TF) coil



Applying toroidal magnetic field
3.5 T in KSTAR, 5.3 T in ITER

Closed Magnetic System



Magnetic field of earth?
0.5 Gauss = 0.00005 T

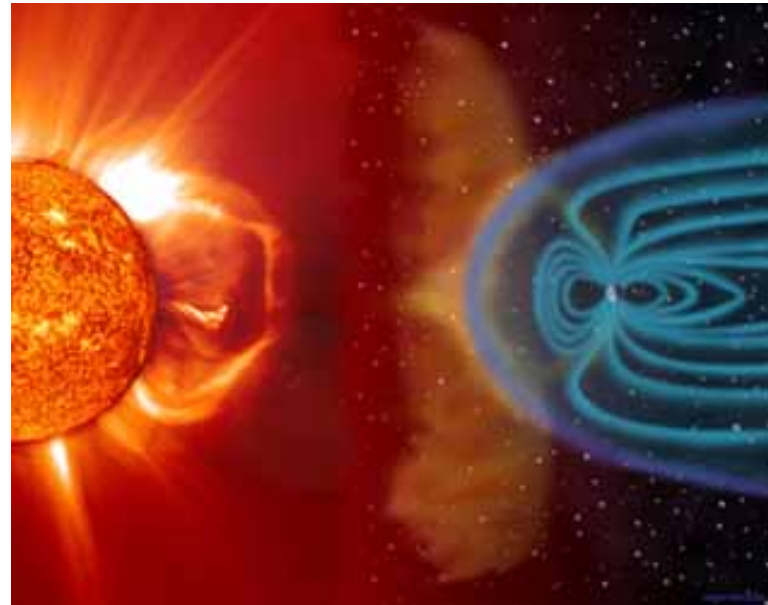
Closed Magnetic System



Magnetic field



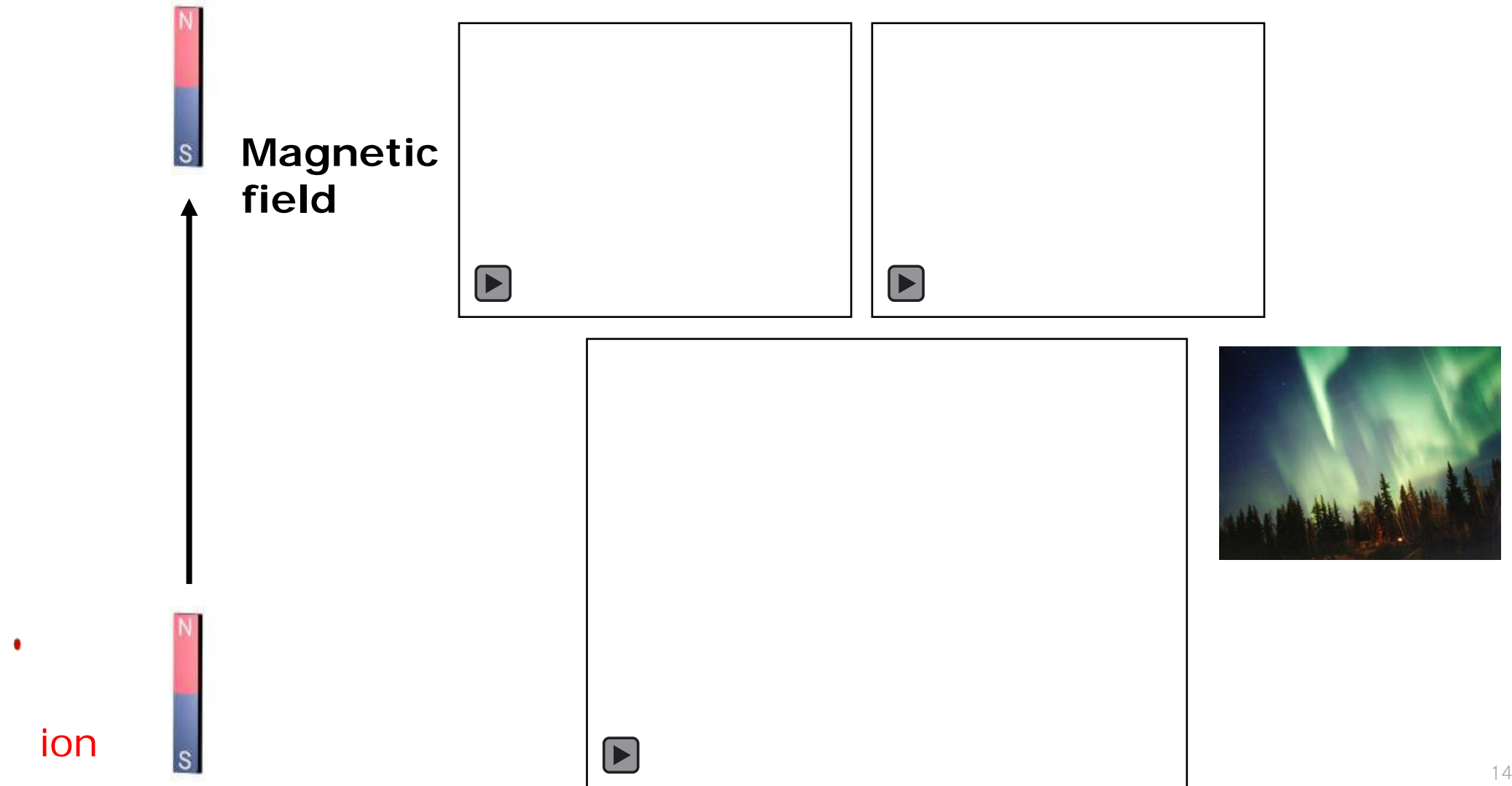
ion



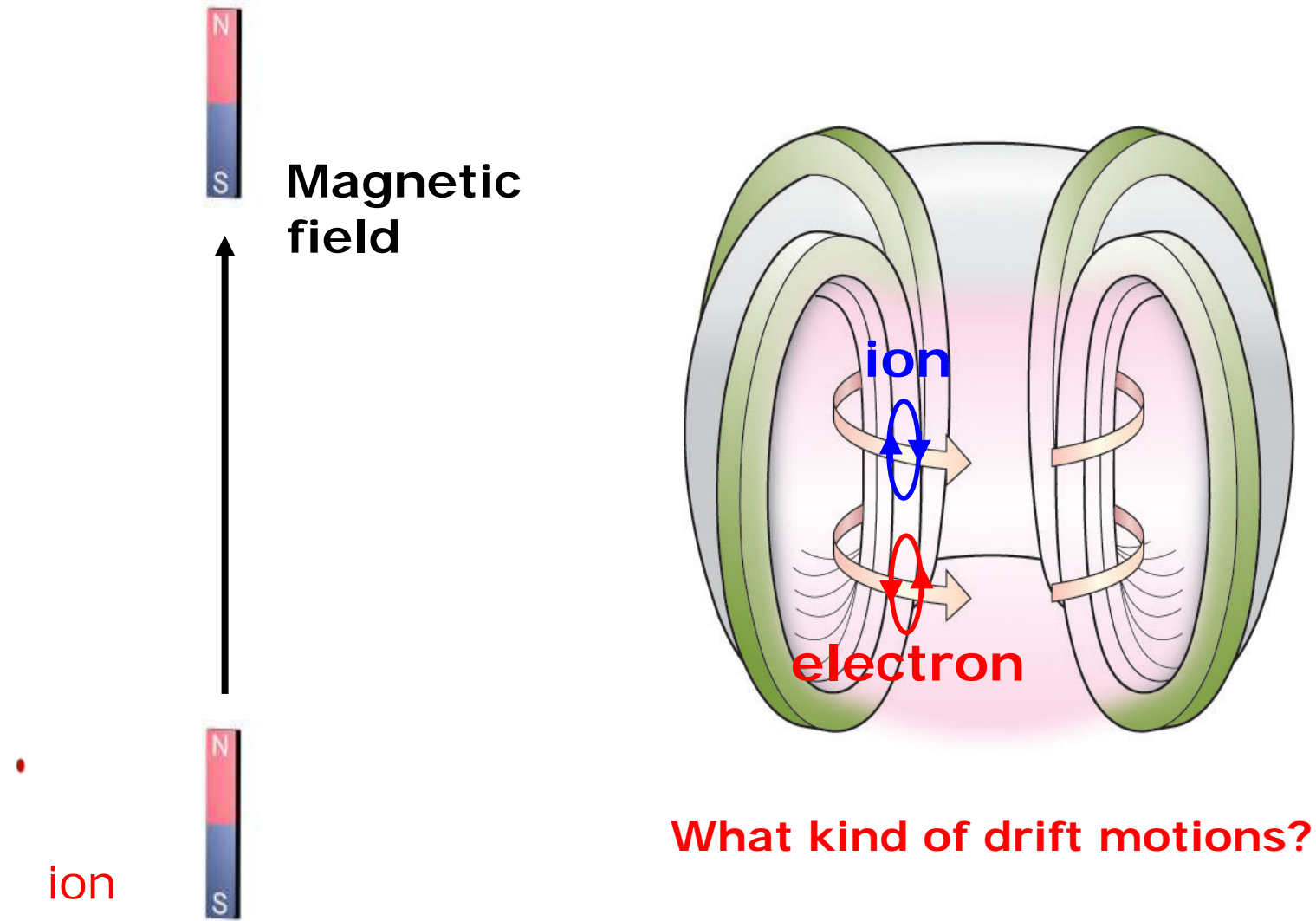
Magnetic field of earth?

0.5 Gauss = 0.00005 T

Closed Magnetic System

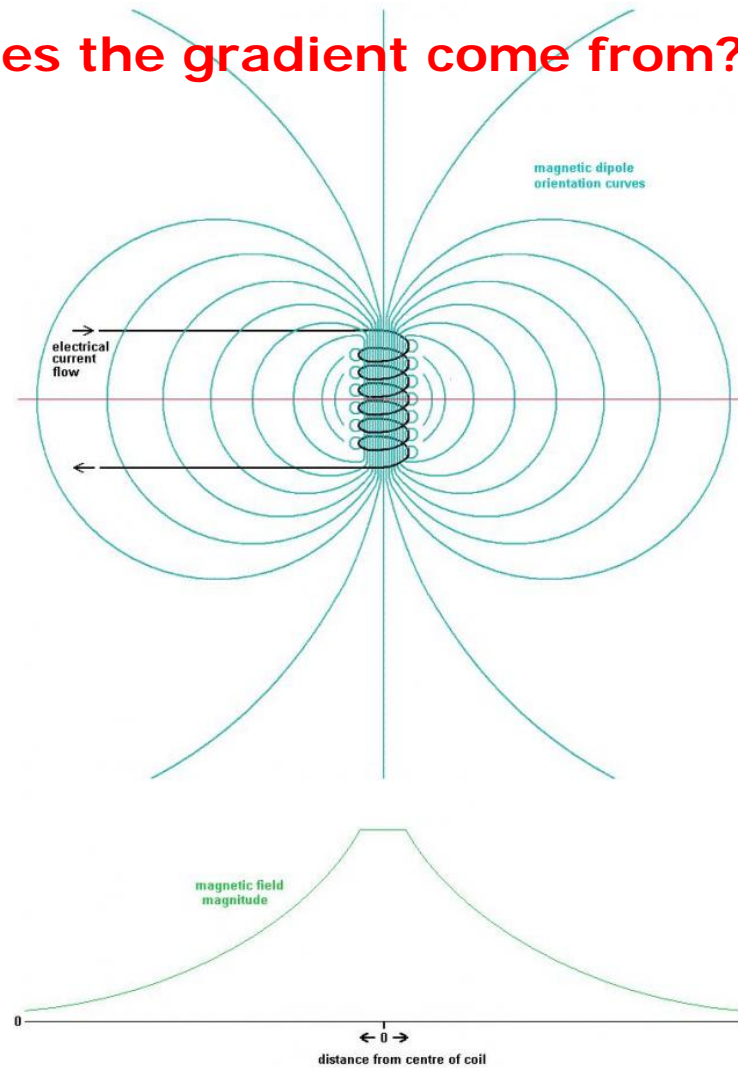
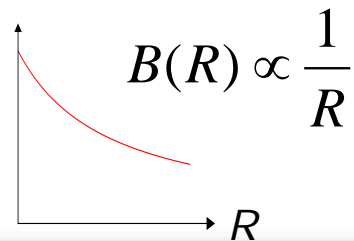
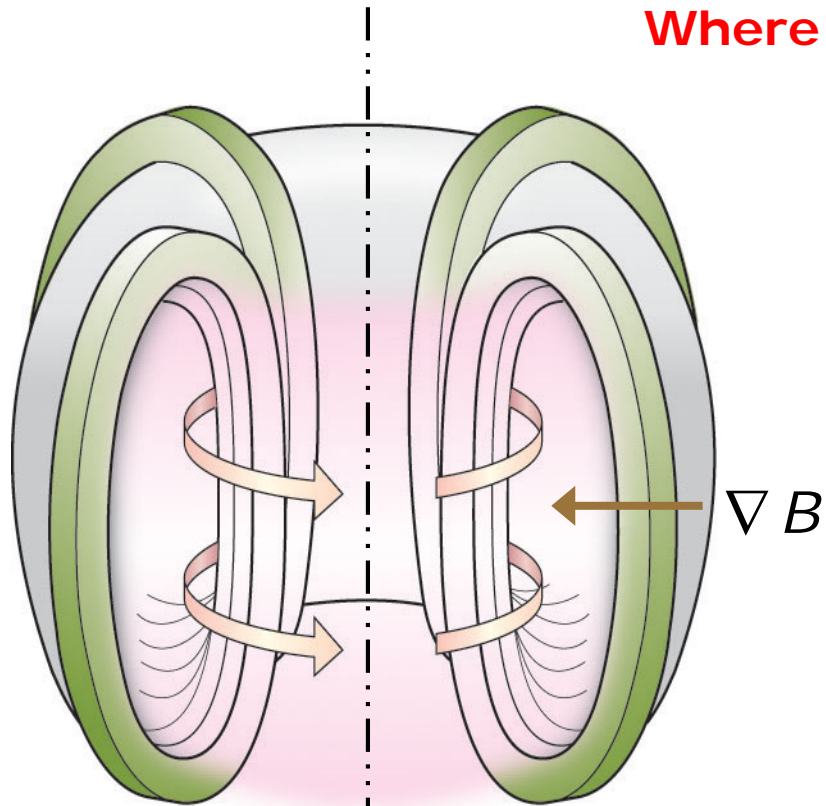


Closed Magnetic System



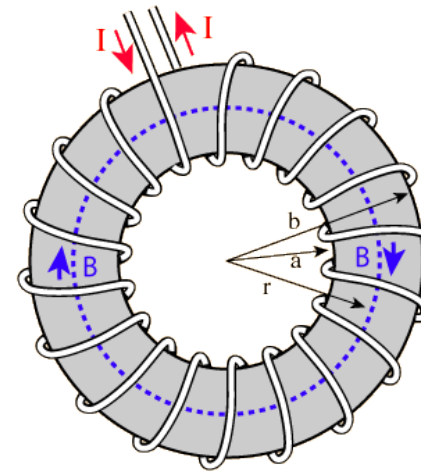
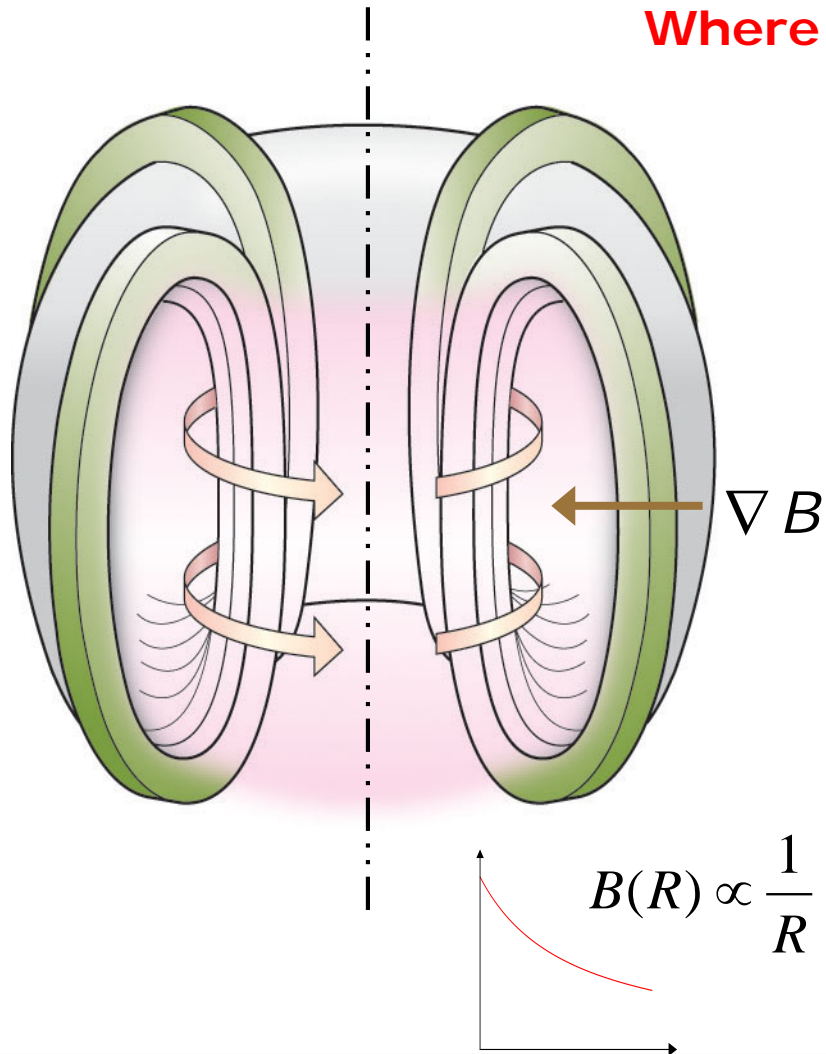
Closed Magnetic System

Where does the gradient come from?



Closed Magnetic System

Where does the gradient come from?

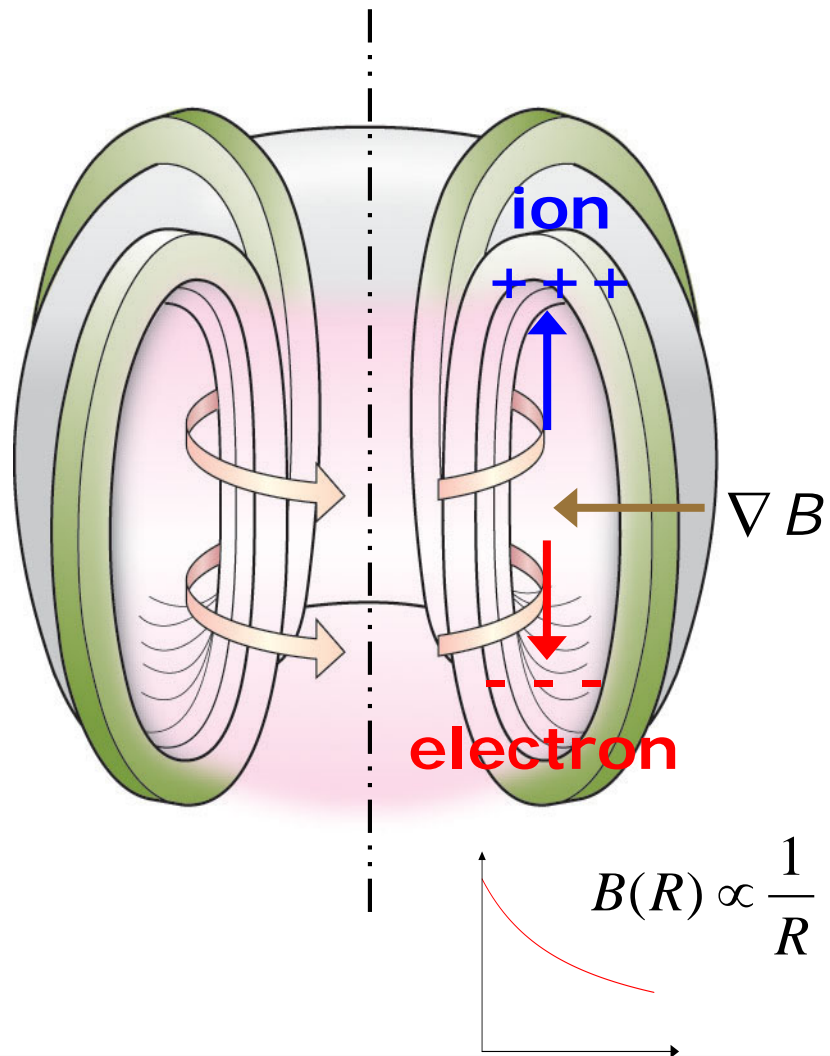


$$\nabla \times \mathbf{B} = \mu_0 \mathbf{J}$$

$$\oint \mathbf{B}_\phi \cdot d\mathbf{l} = \mu_0 N I_c$$

$$B_\phi(R) = \frac{\mu_0 N I_c}{2\pi R}$$

Closed Magnetic System

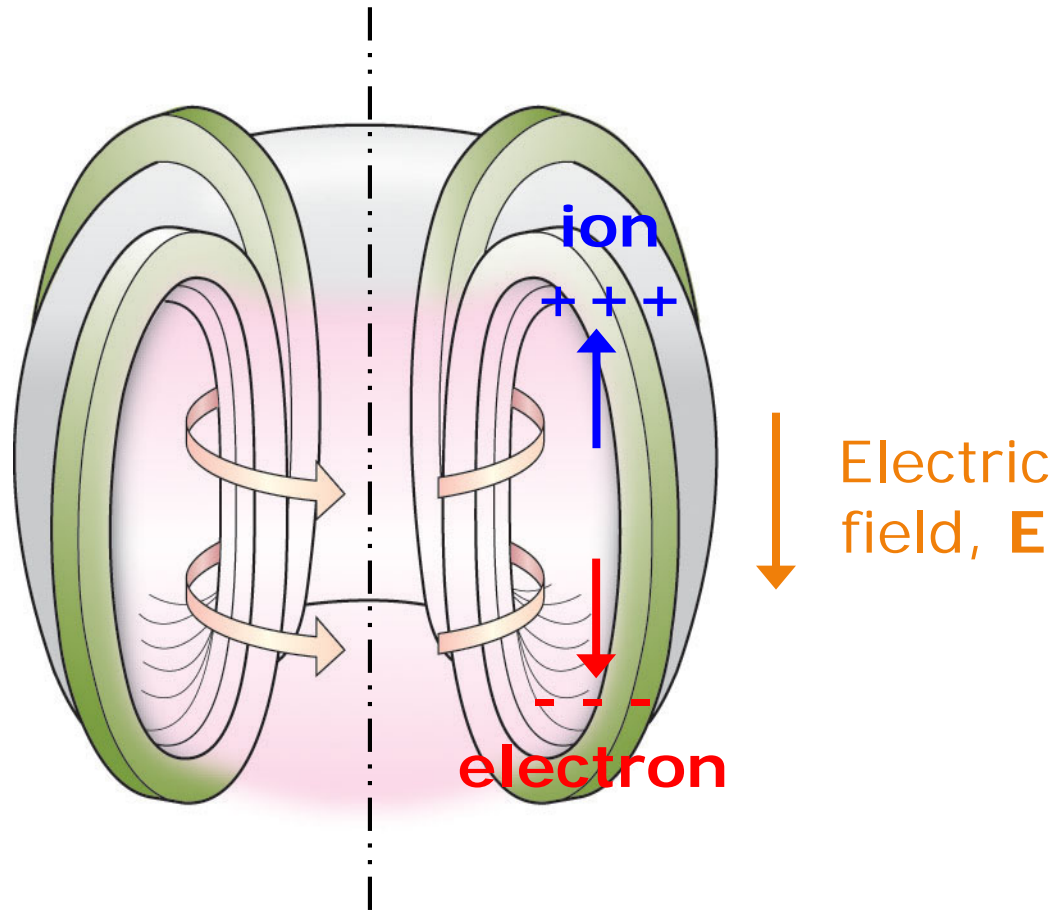


$$\mathbf{v}_{D,R} = \frac{mv_{\parallel}^2}{qB_0^2} \frac{\mathbf{R}_0 \times \mathbf{B}_0}{R^2}$$

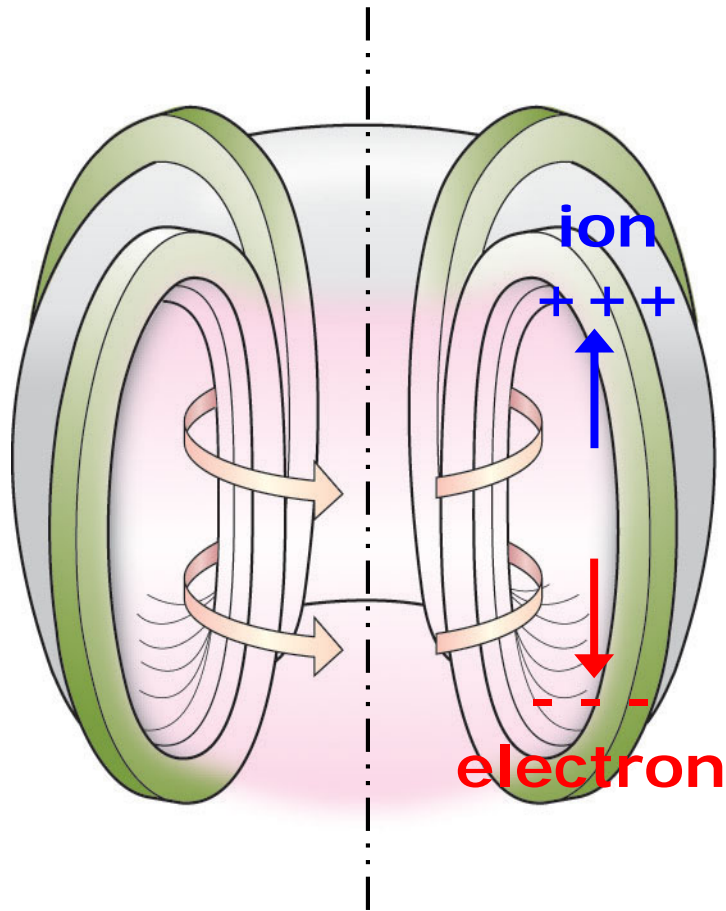
$$\begin{aligned} \mathbf{v}_{D,\nabla B} &= \pm \frac{1}{2} v_{\perp} r_L \frac{\mathbf{B} \times \nabla B}{B^2} \\ &= \frac{mv_{\perp}^2}{2qB} \frac{\mathbf{B} \times \nabla B}{B^2} \end{aligned}$$

$$\mathbf{v}_D = \frac{m}{q} \frac{1}{R_0 B_{\phi}(R_0)} \left[v_{\parallel}^2 + \frac{v_{\perp}^2}{2} \right] \mathbf{e}_z$$

Closed Magnetic System



Closed Magnetic System

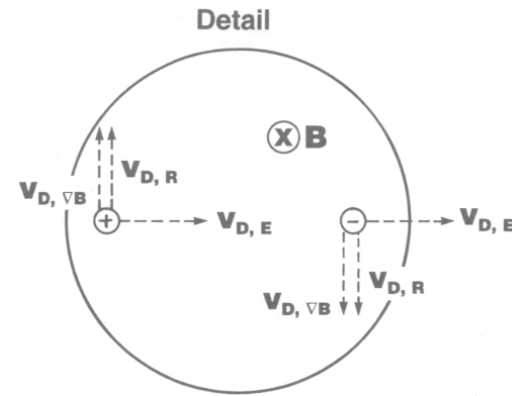


➔ **E × B** drift

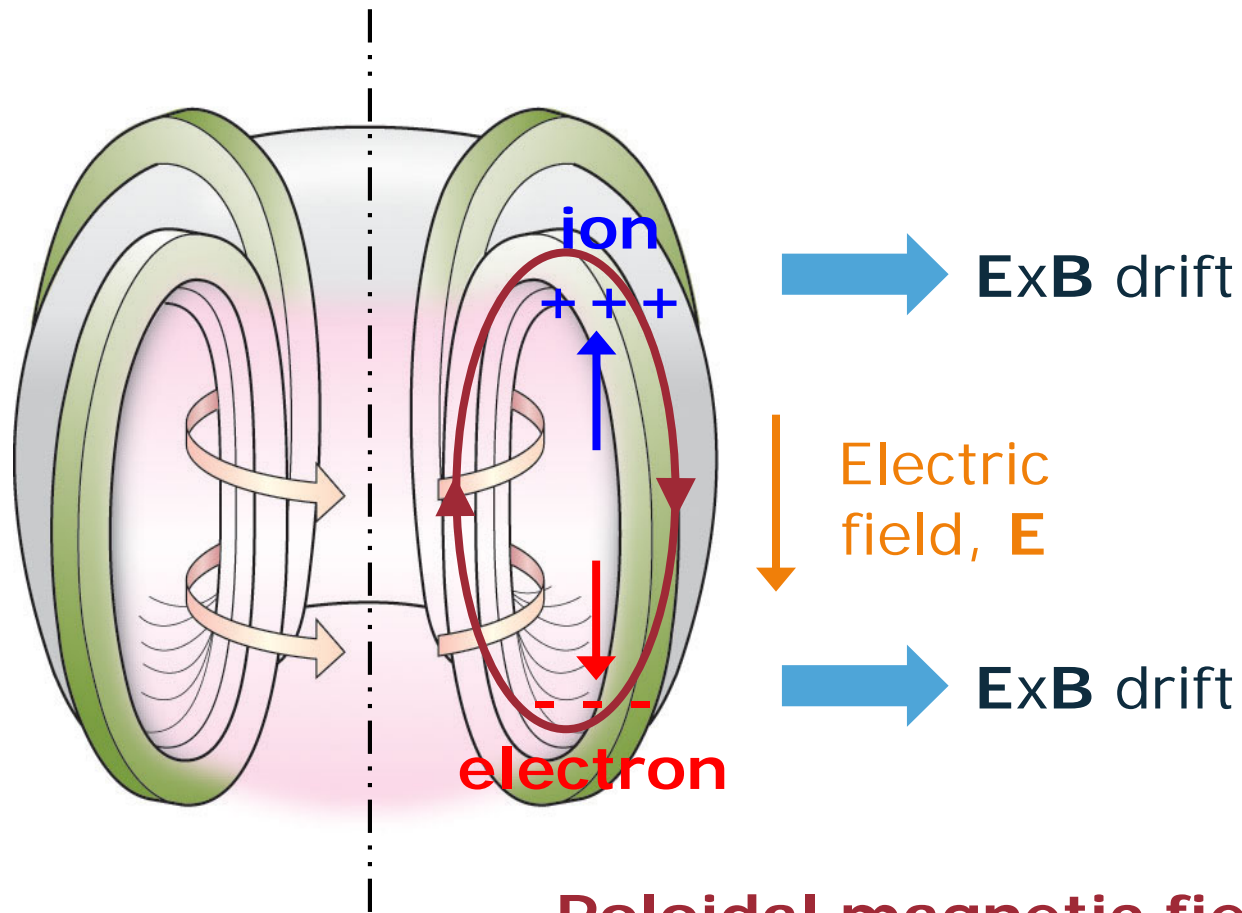
↓ Electric field, **E**

➔ **E × B** drift

$$\mathbf{v}_{D,E} = \frac{\mathbf{E} \times \mathbf{B}}{B^2} = \frac{E}{B_\phi(R_0)} \cdot \frac{\mathbf{R}}{R_0}$$



Closed Magnetic System

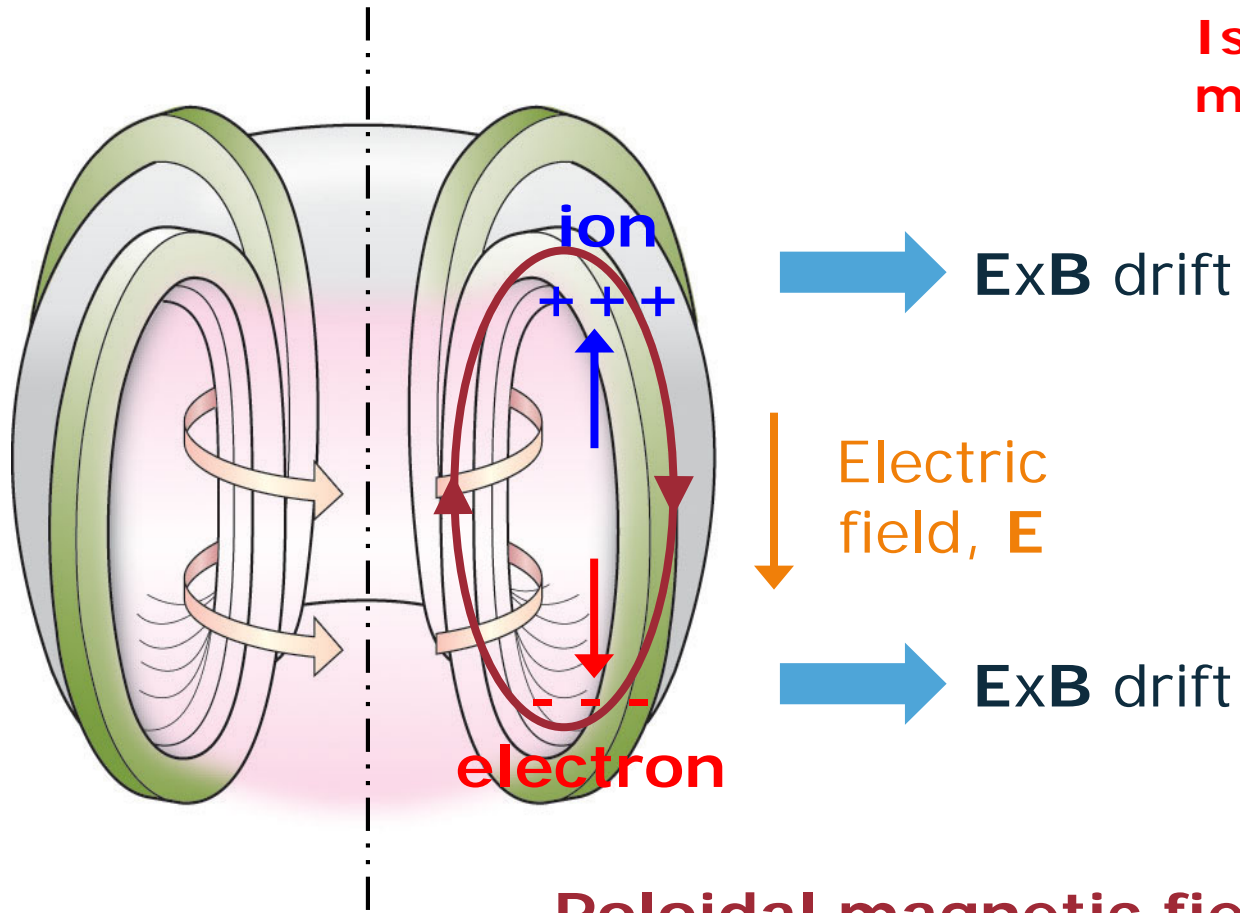


Poloidal magnetic field required
Tokamak .VS. Stellarator

What is a tokamak?

Closed Magnetic System

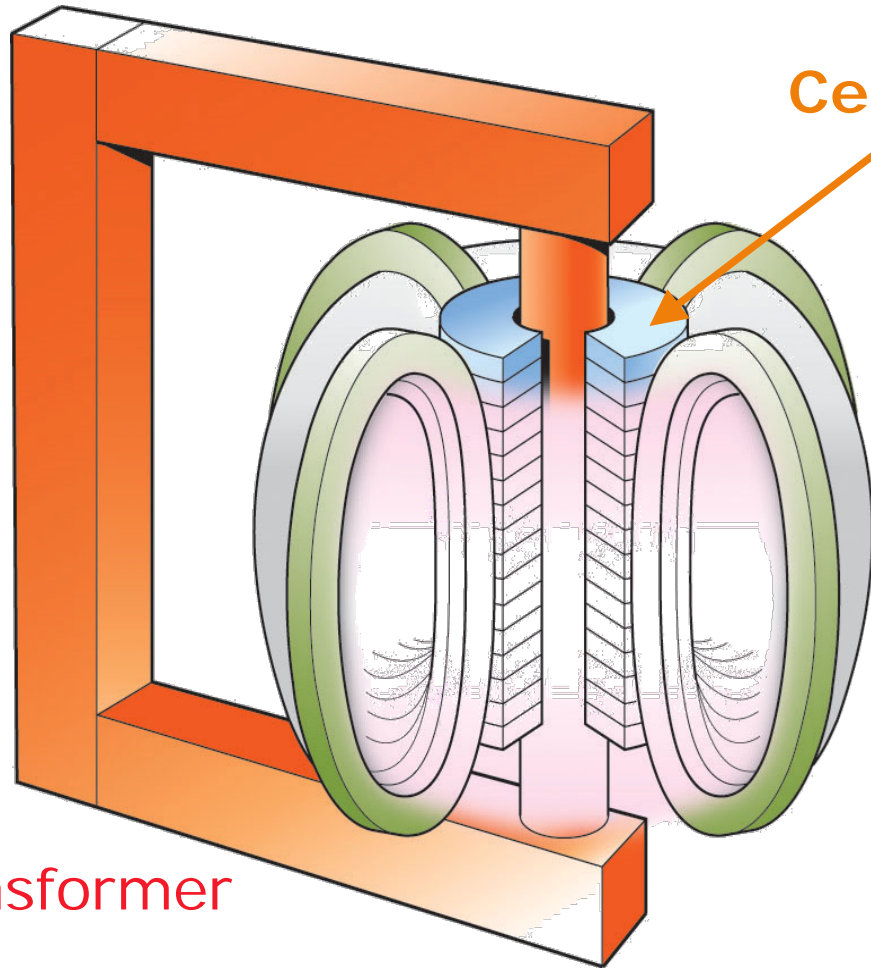
Is the plasma current driven if toroidal magnetic field is applied?



Poloidal magnetic field required

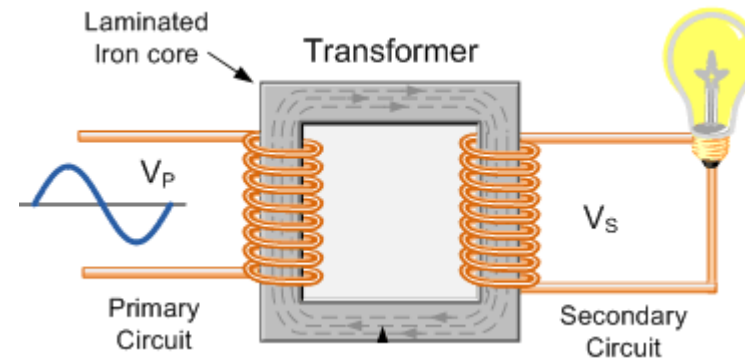
How to drive plasma current?

Tokamak



Central Solenoid (CS)

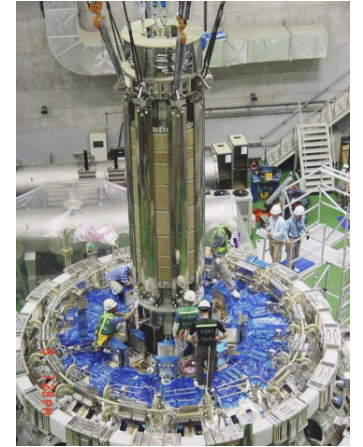
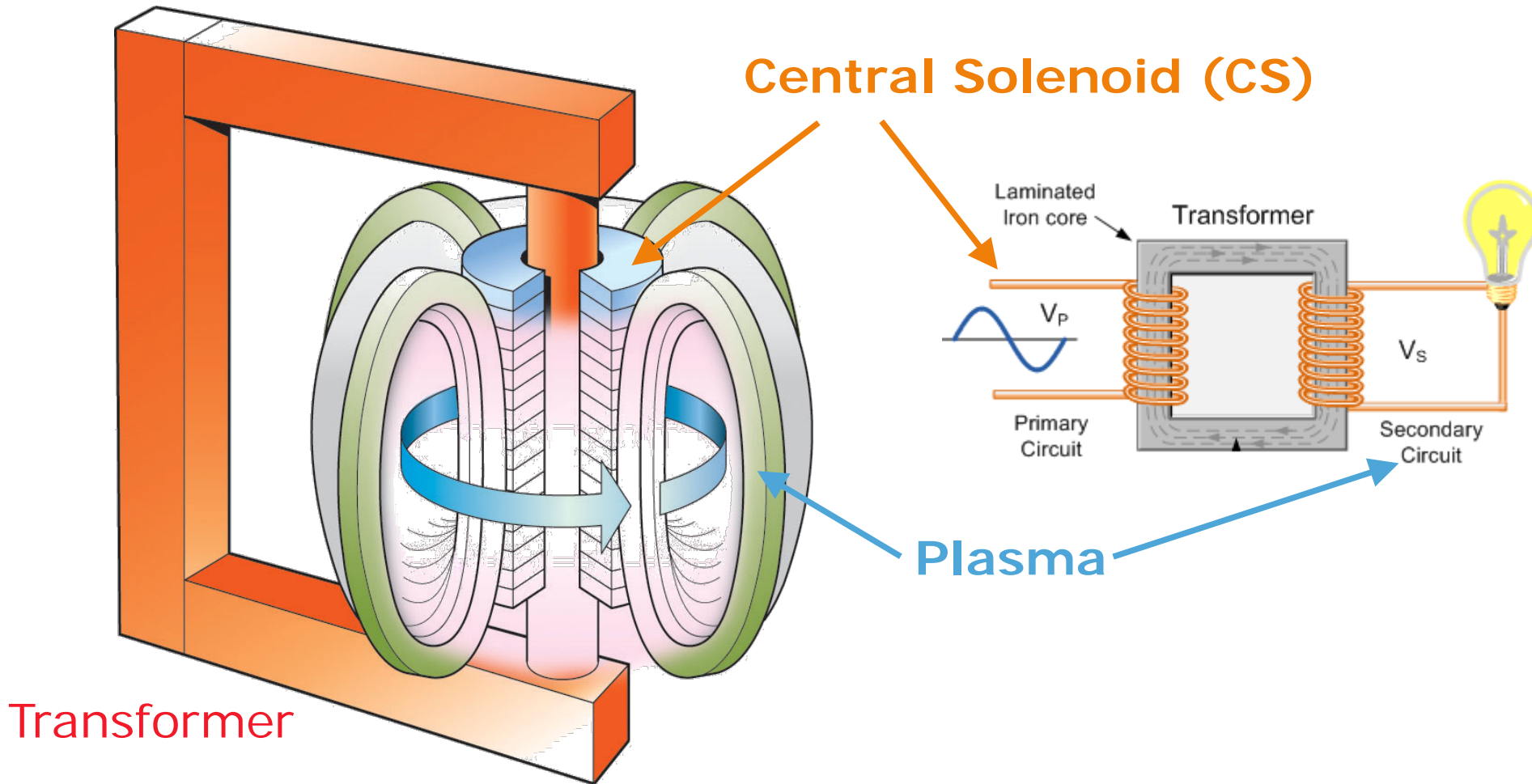
Transformer



Faraday's law

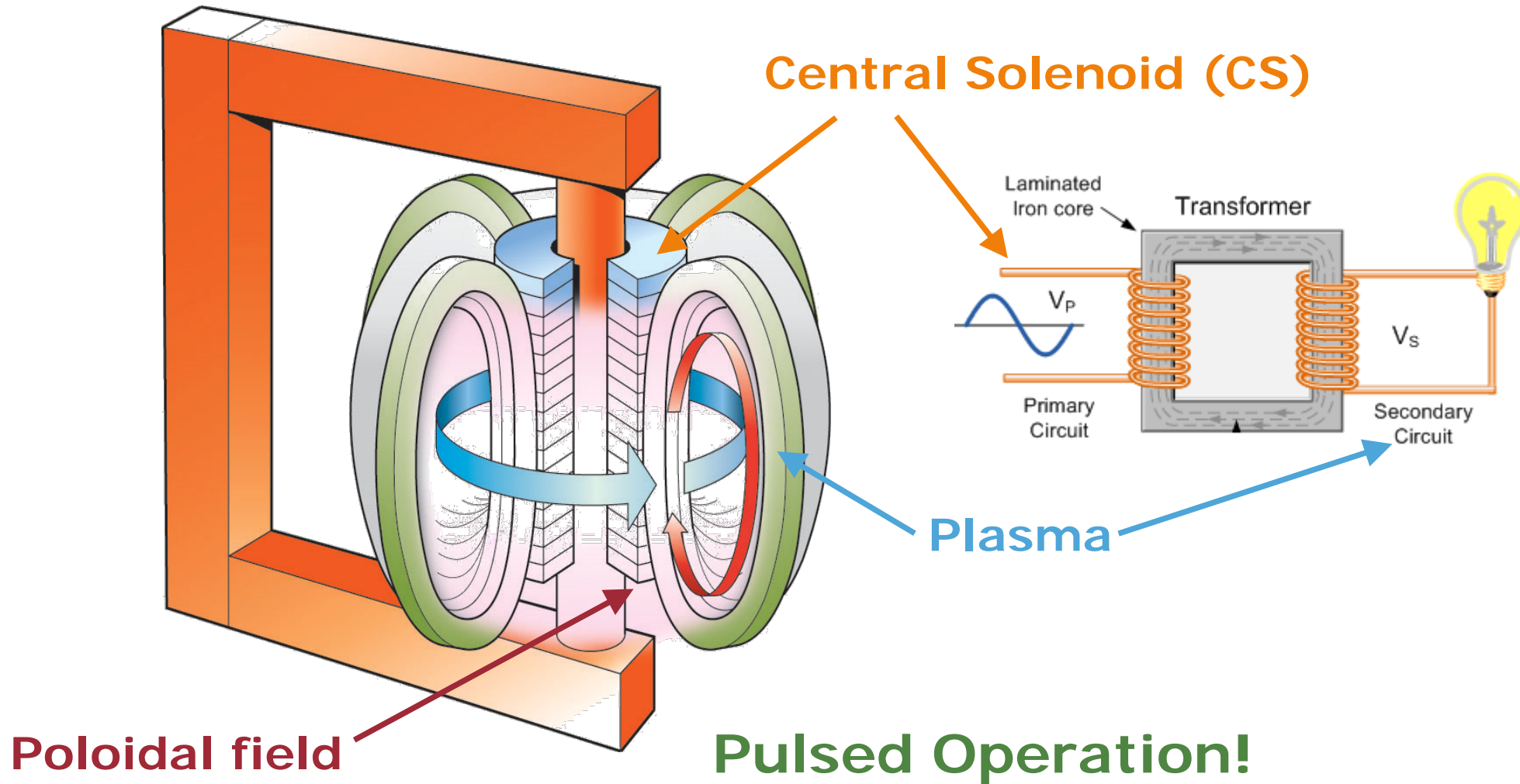
$$\mathcal{V} = -\frac{d}{dt} \int_S \vec{B} \cdot d\vec{S}$$

Tokamak



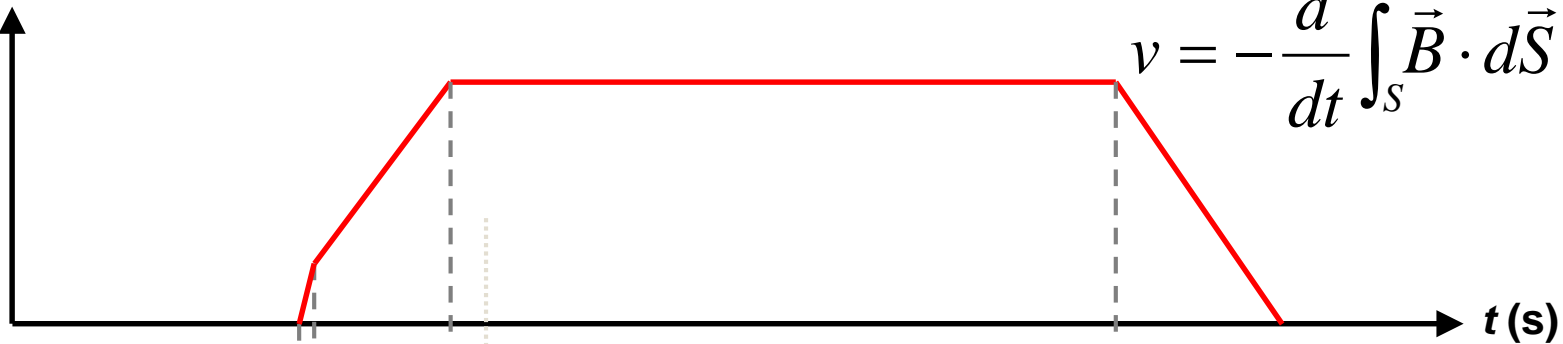
KSTAR

Tokamak



Pulsed Operation

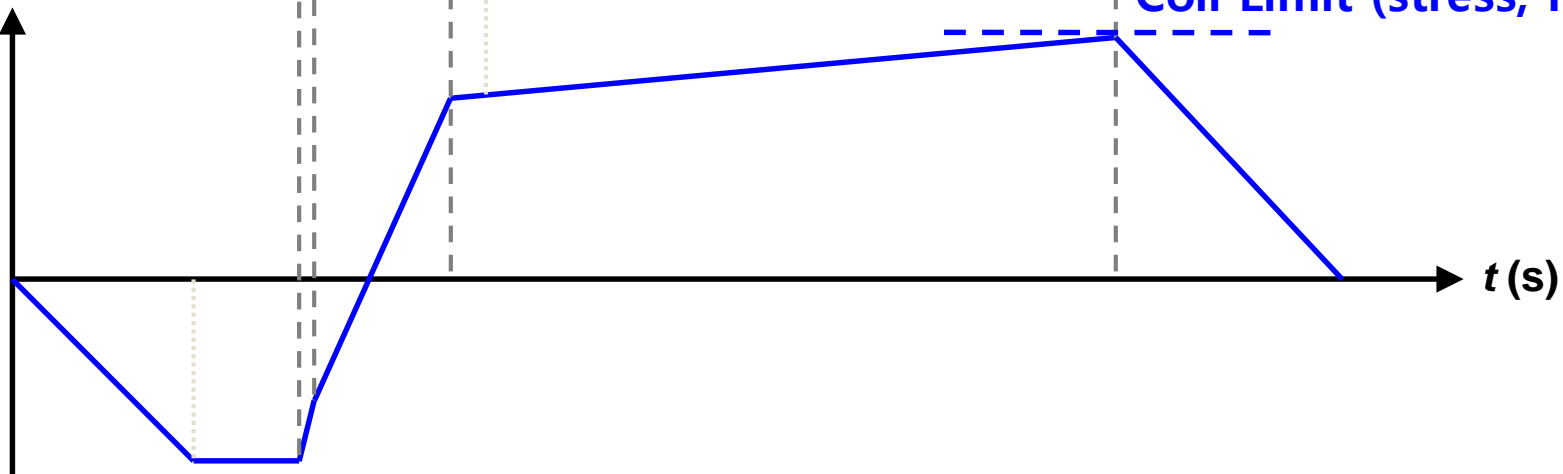
Plasma Current
(MA)



Faraday's law

$$v = -\frac{d}{dt} \int_s \vec{B} \cdot d\vec{S}$$

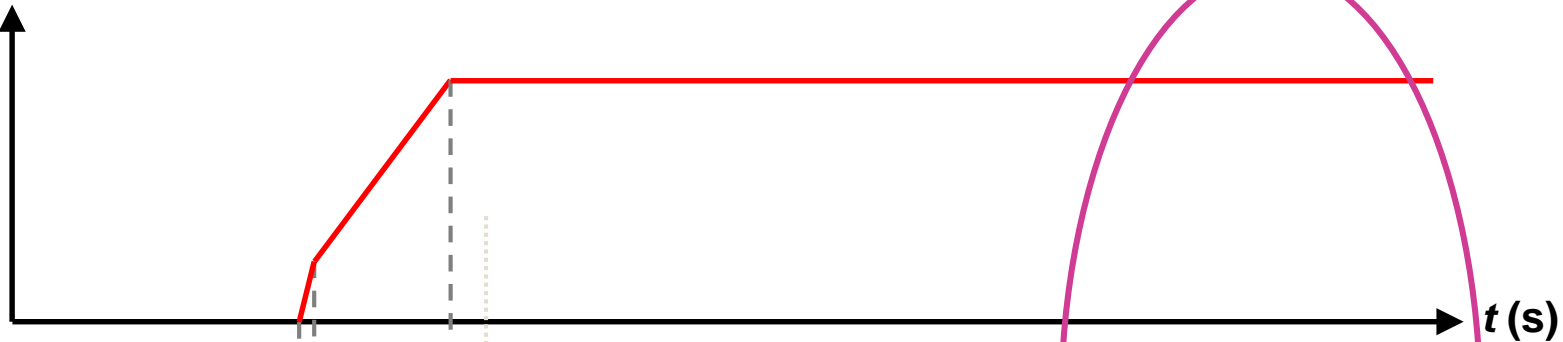
CS Coil Current
(kA)



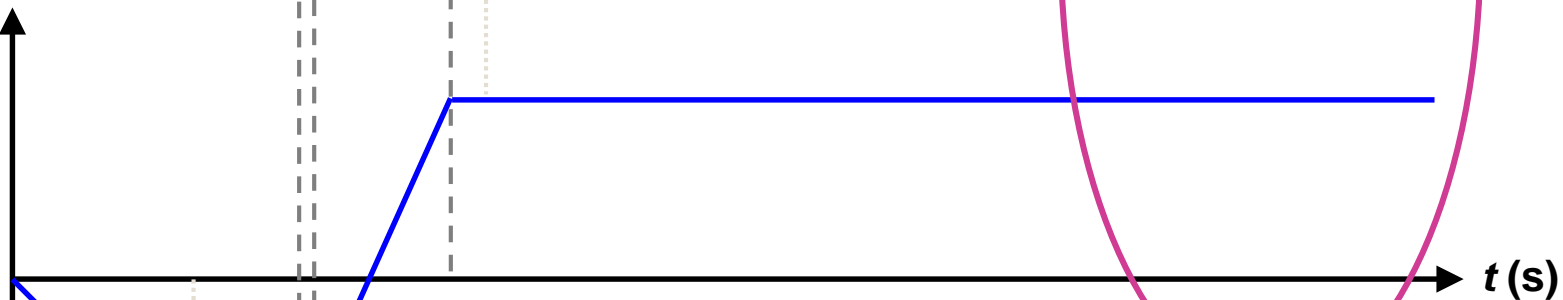
Inherent drawback of Tokamak!

Steady-State Operation

Plasma Current
(MA)



CS Coil Current
(kA)

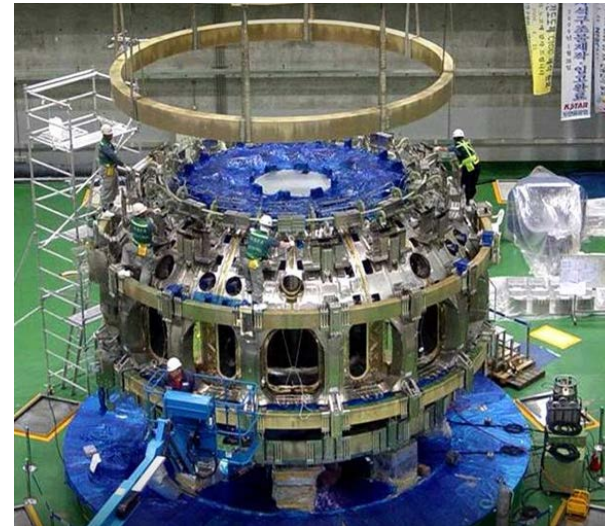
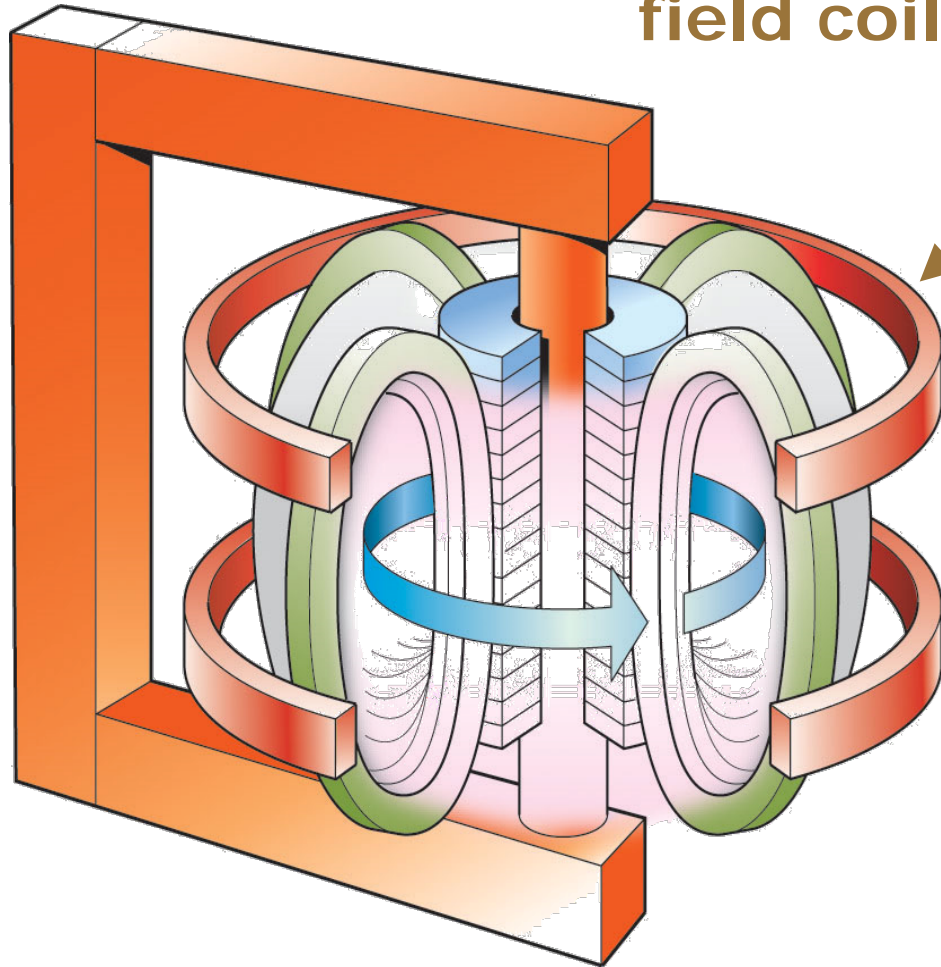


$$d/dt \sim 0$$

Steady-state operation
by self-generated and externally driven current

Tokamak

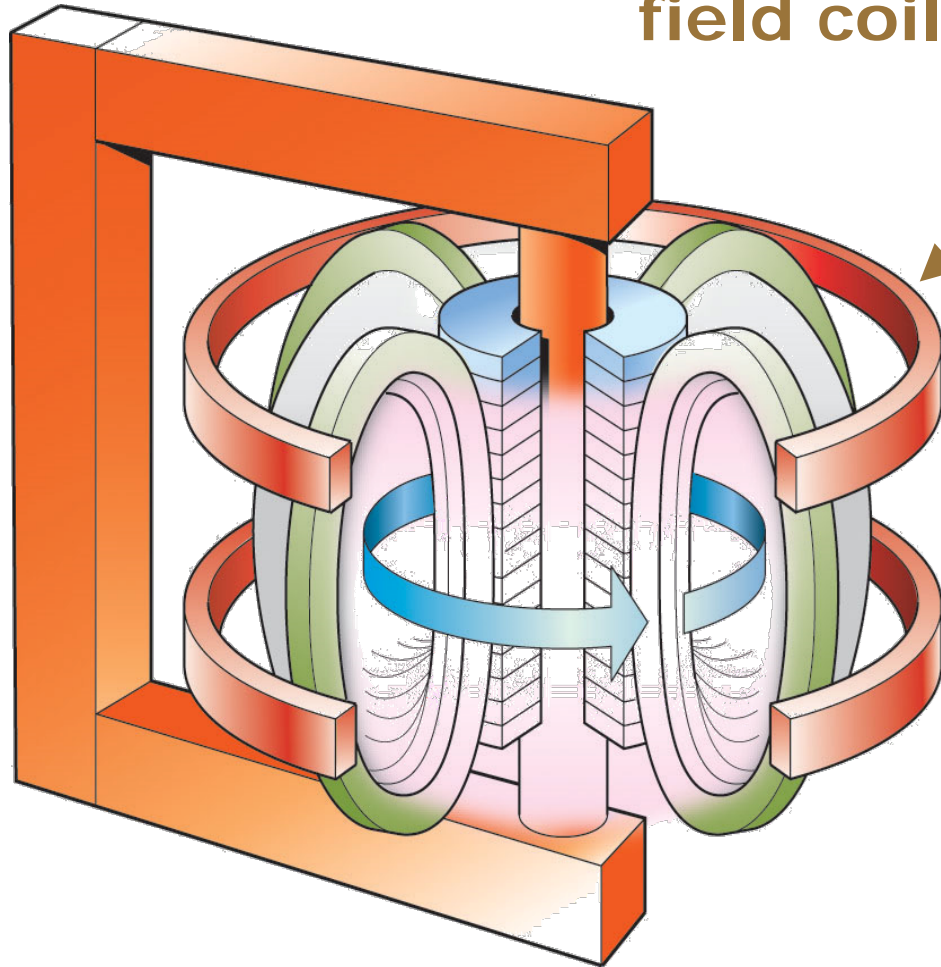
Adding vertical (equilibrium)
field coils (PF: Poloidal Field)



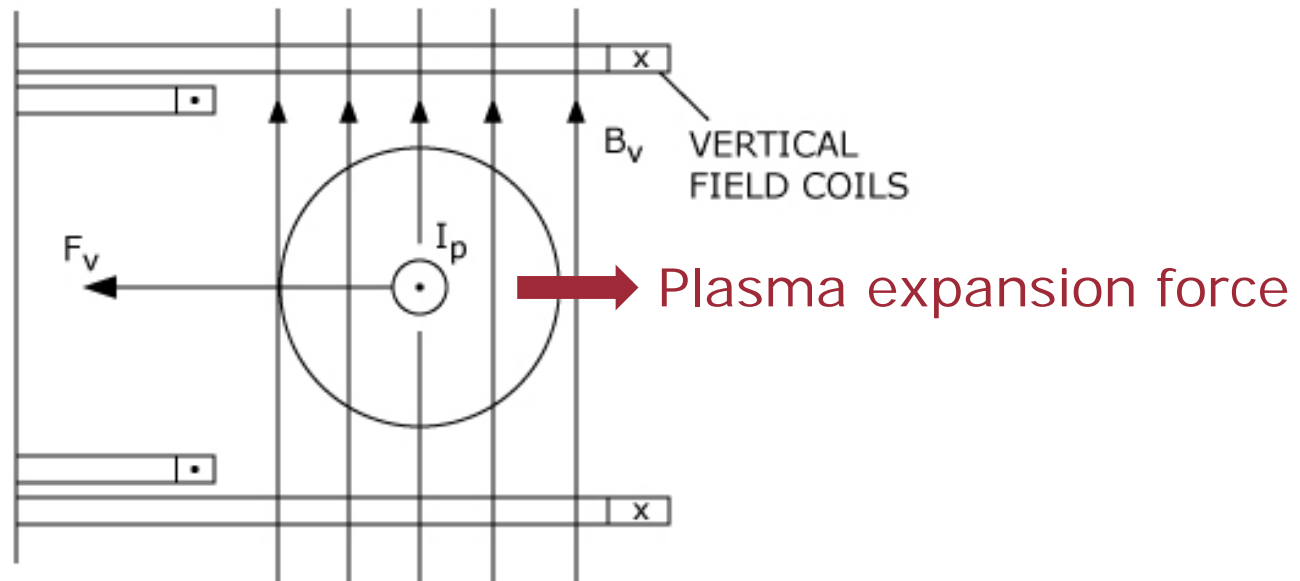
KOTAR

Tokamak

Adding vertical (equilibrium)
field coils (PF: Poloidal Field)

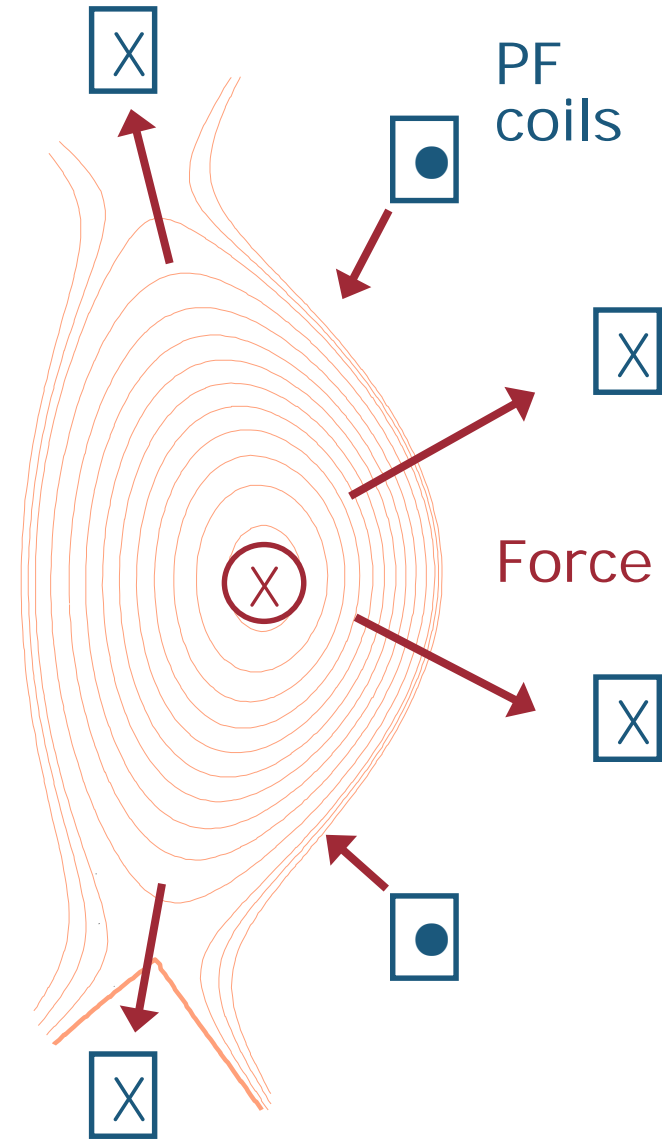
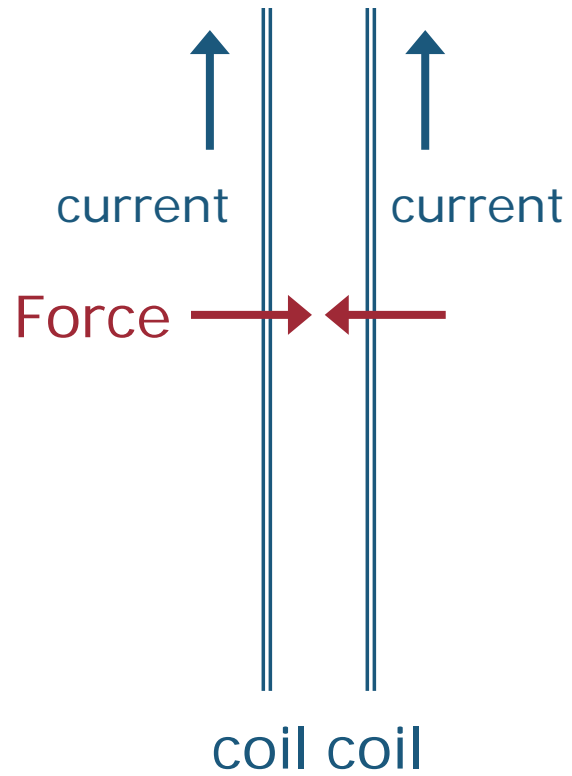


Tokamak



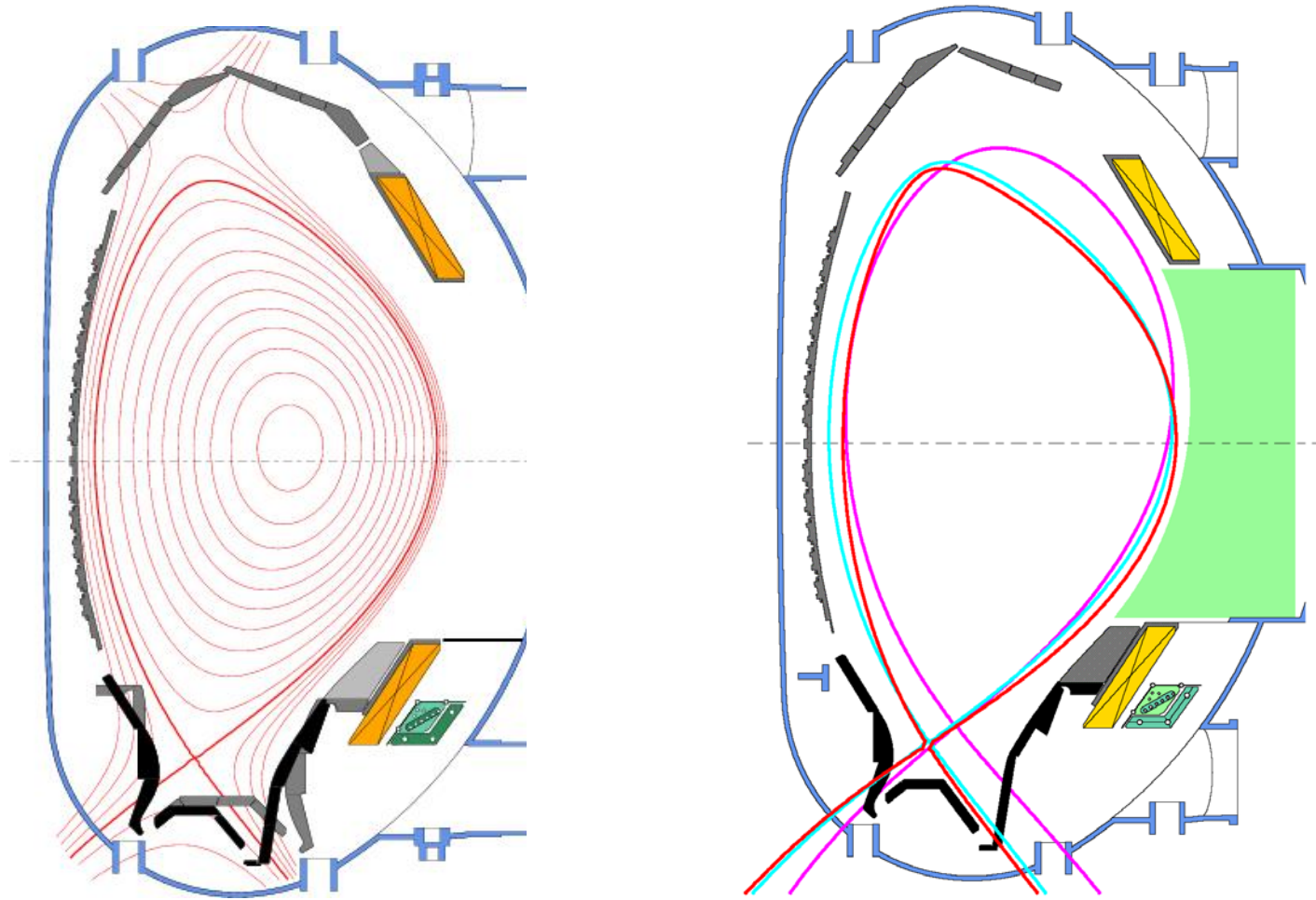
**Force balance by vertical field coils:
Plasma positioning**

Tokamak



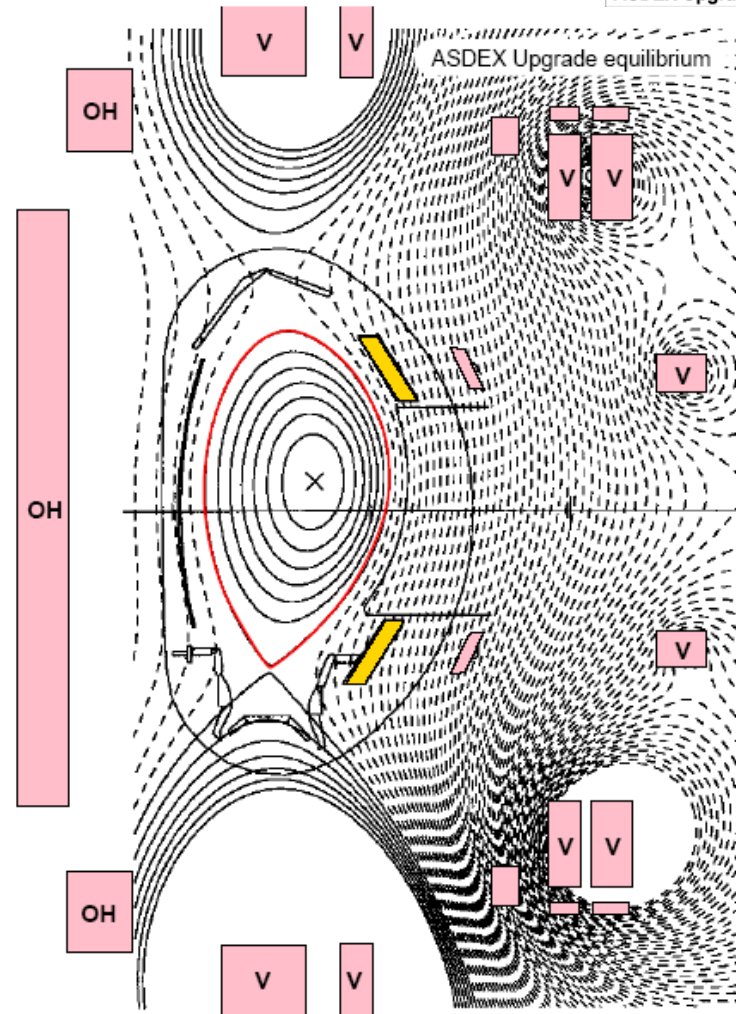
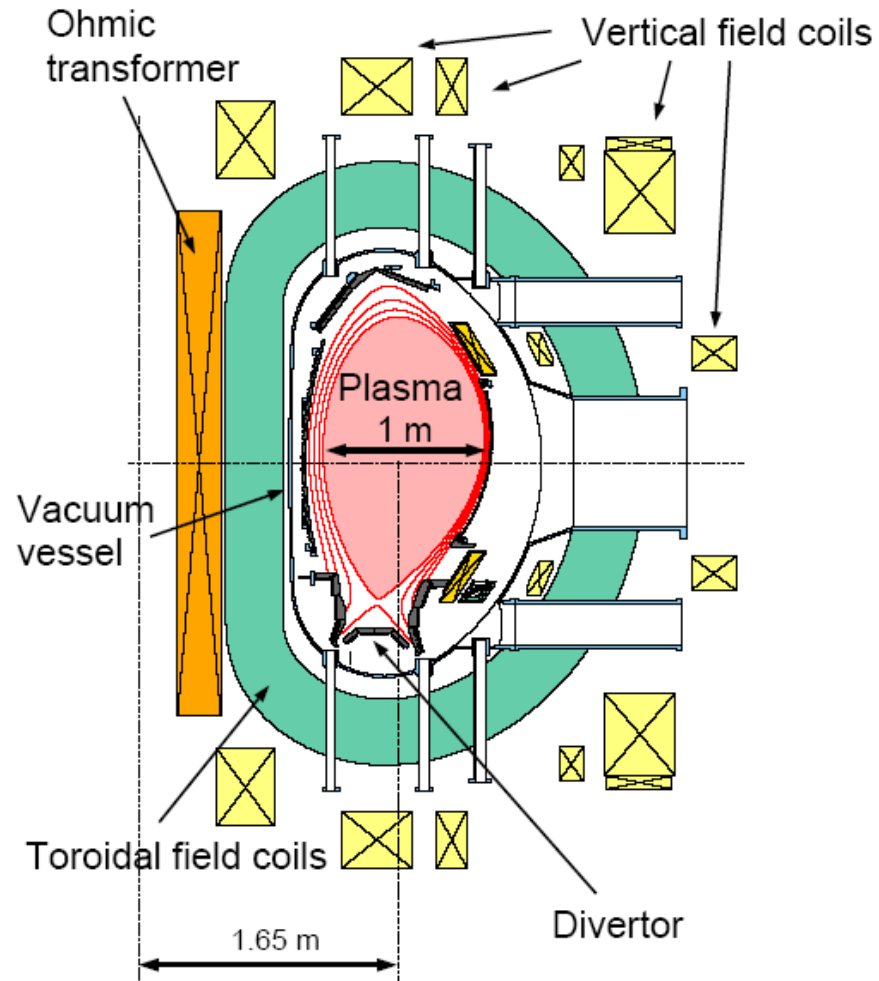
Plasma shaping by PF coils

Tokamak



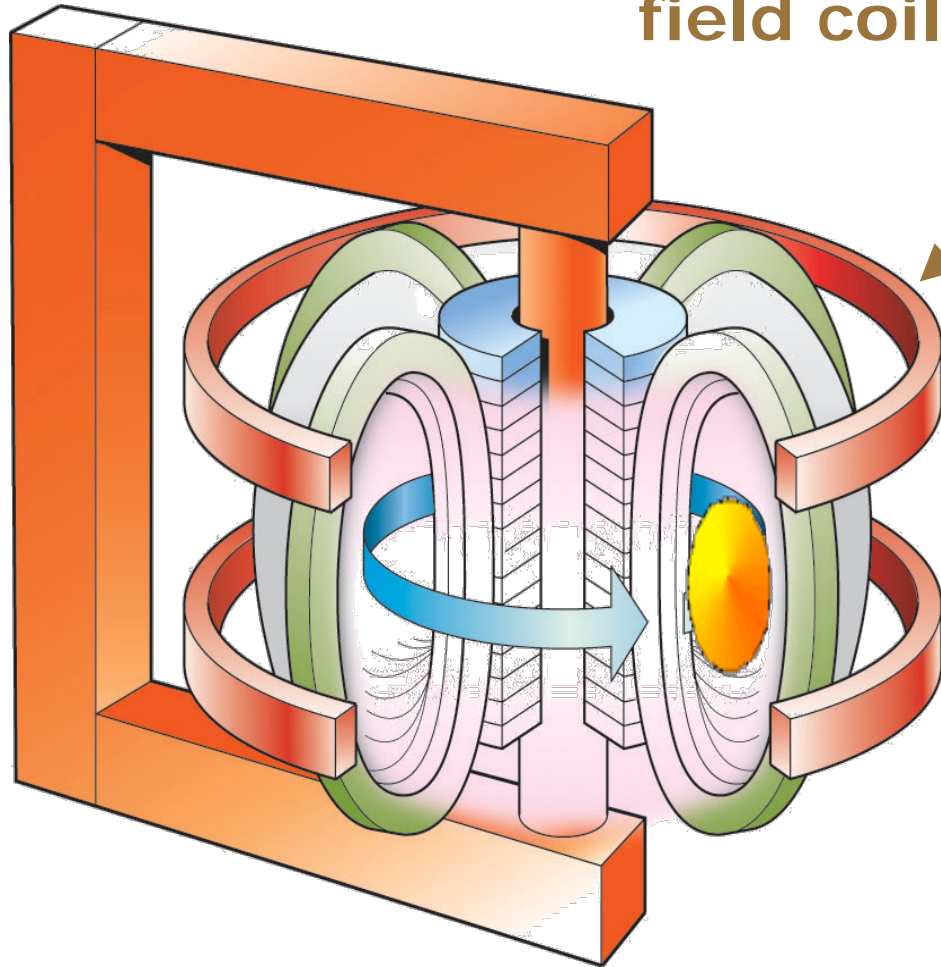
The plasma shape can be modified by PF coil currents.

Tokamak



Tokamak

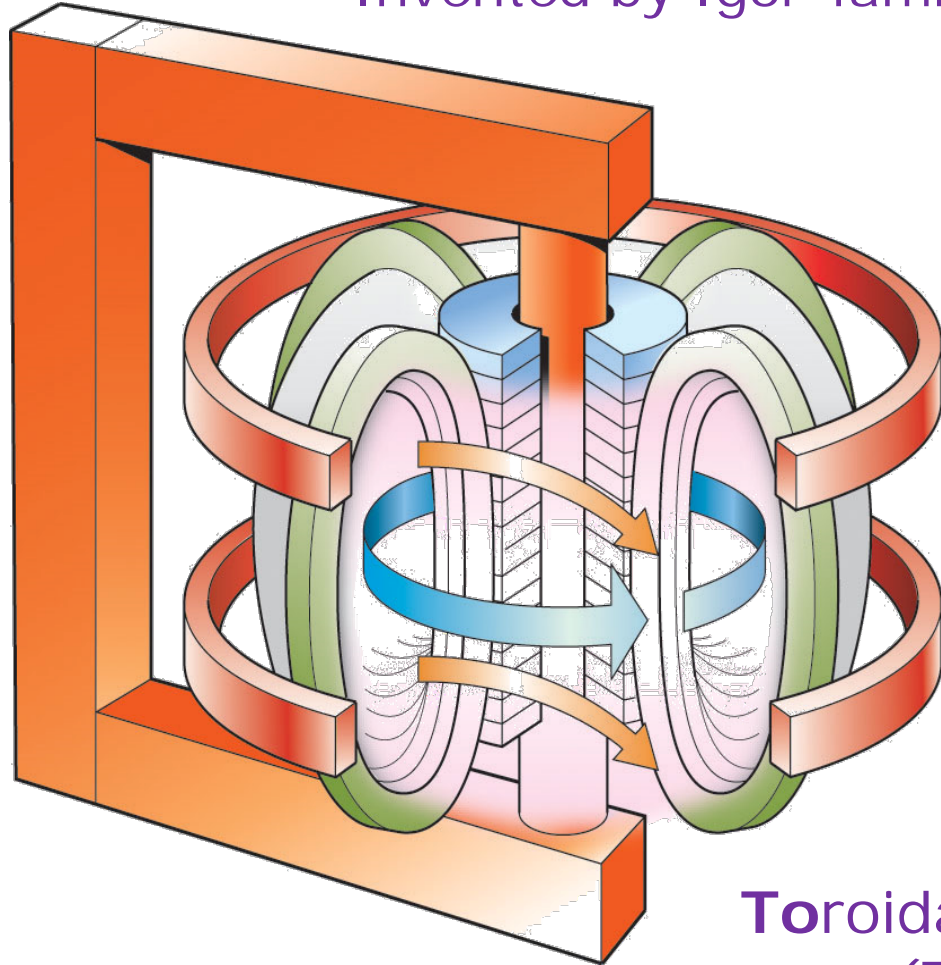
Adding vertical (equilibrium)
field coils (PF: Poloidal Field)



Plasma positioning & shaping by PF coils

Tokamak

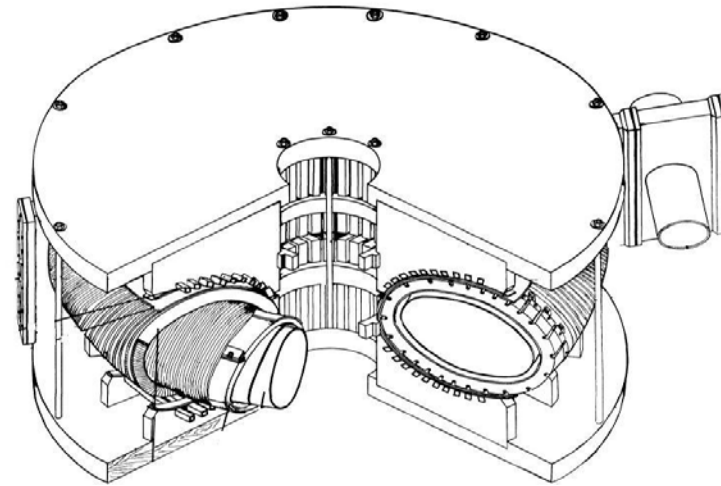
Invented by Igor Tamm and Andrei Sakharov in 1952



WIKIPEDIA

Toroidalnaja kamera magnitnaja katushka
(Toroidal chamber magnetic coil)

Tokamak



Cutaway of the Toroidal Chamber in
Artsimovitch's Paper *Research on
Controlled Nuclear Fusion in the USSR*



Toroidalnaja kamera magnitnaja katushka
(Toroidal chamber magnetic coil)

1958 IAEA FEC, Geneva, Switzerland



T1: The world's first tokamak,
Kurchatov Institute, Moscow Russia

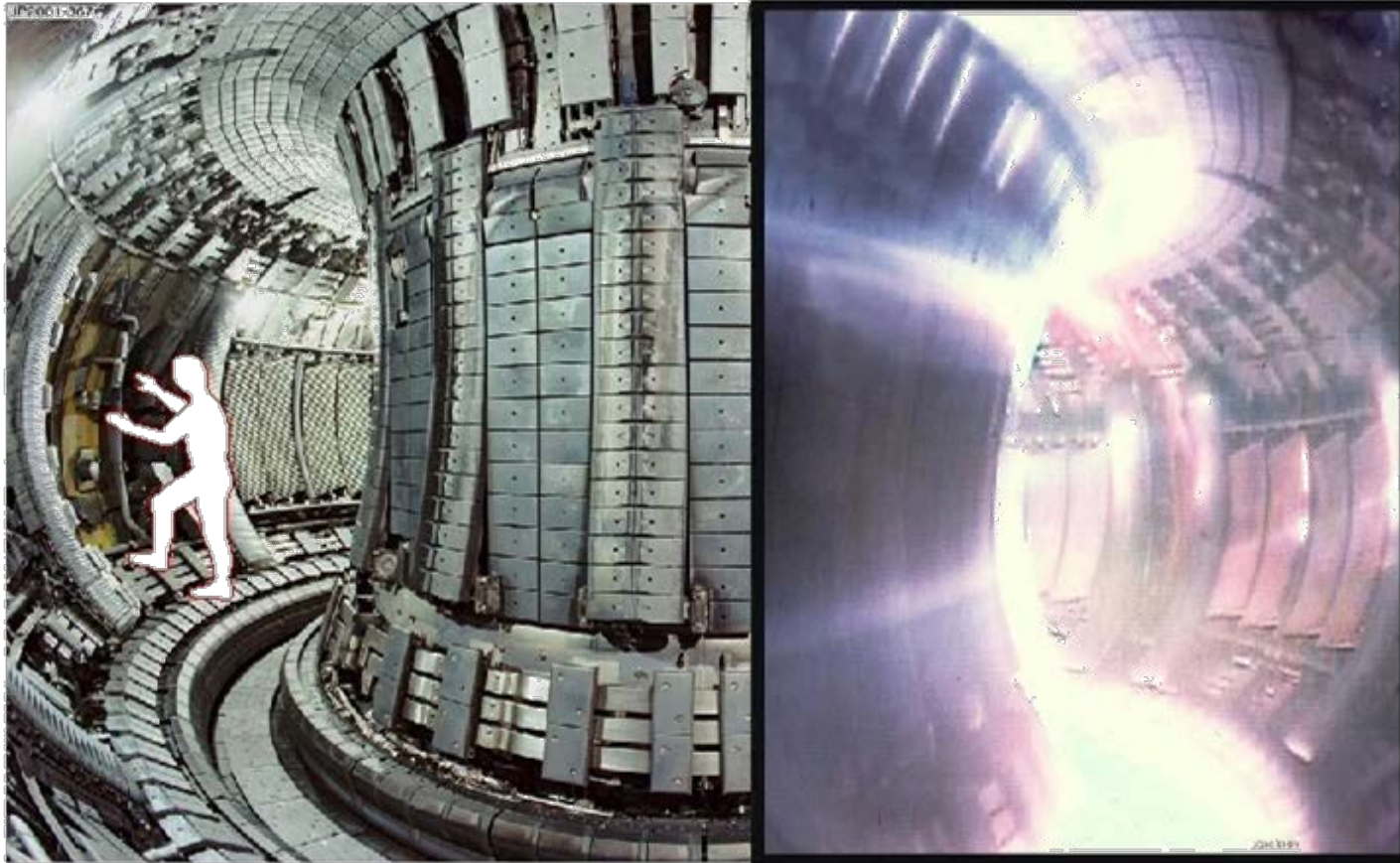
Tokamak

JET (Joint European Torus): $R_0 = 3 \text{ m}$, $a = 0.9 \text{ m}$, 1983-today



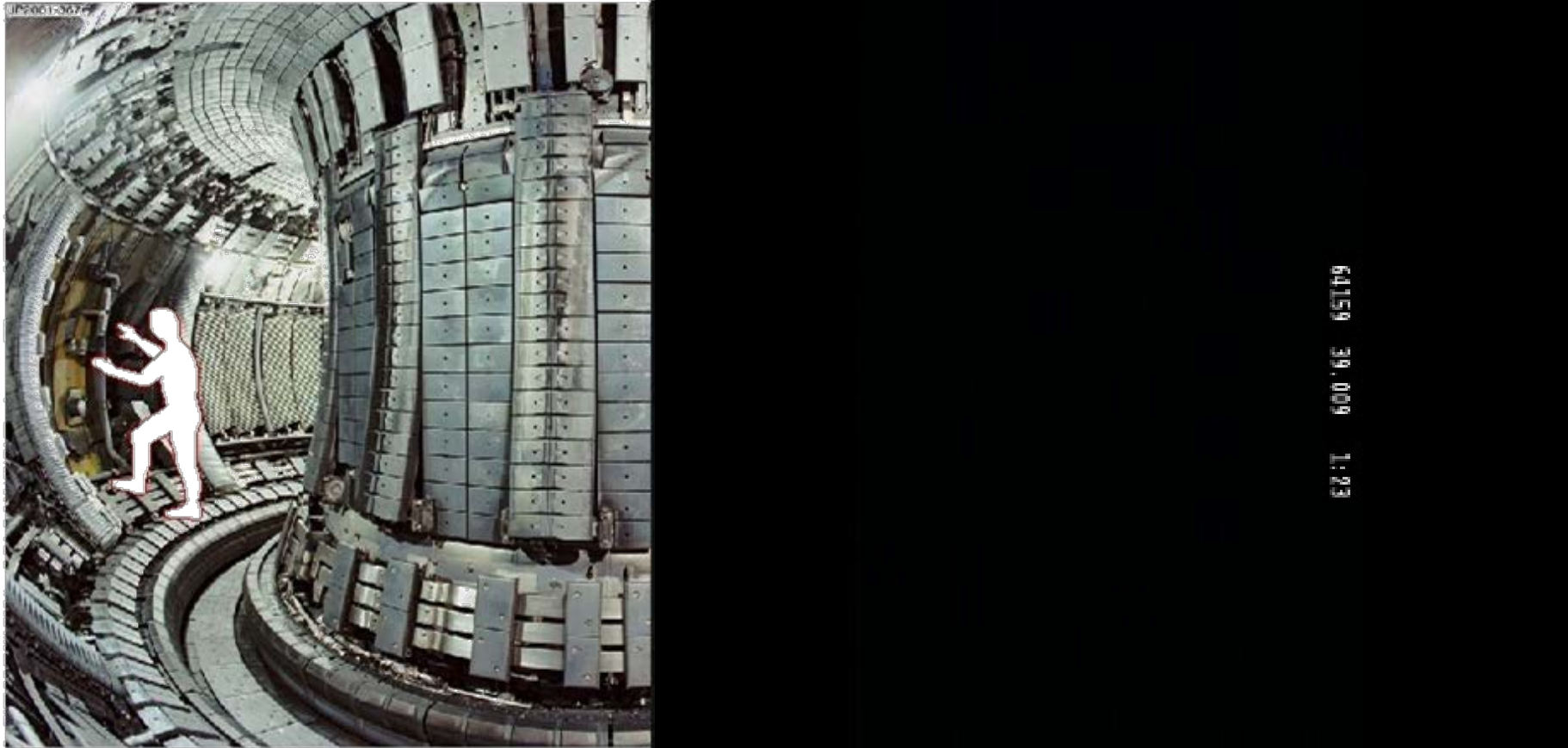
Tokamak

JET (Joint European Torus): $R_0 = 3$ m, $a = 0.9$ m, 1983-today



Tokamak

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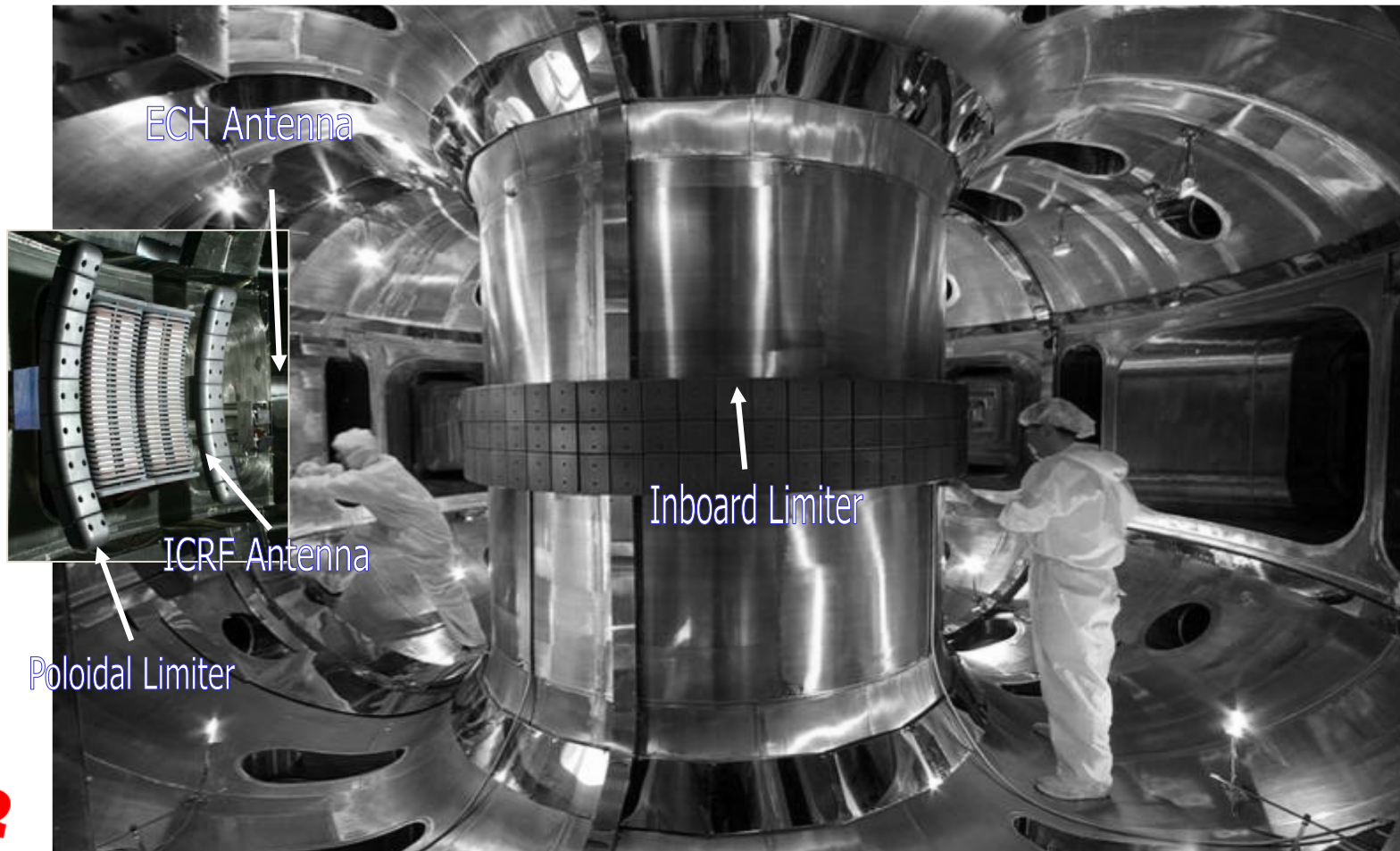


Tokamak



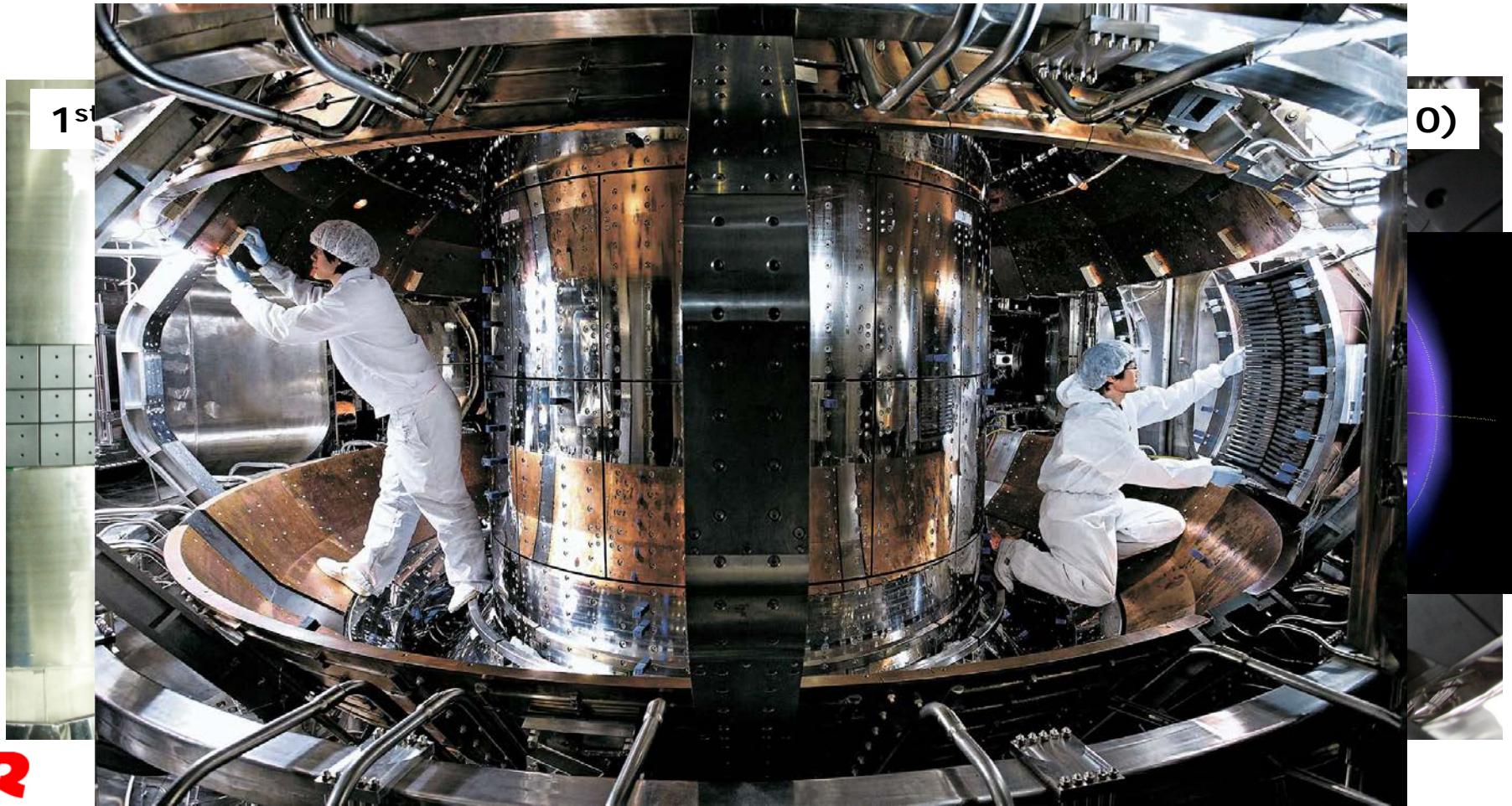
Tokamak

KSTAR (Korea Superconducting Tokamak Advanced Research):
 $R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$, 2007-today



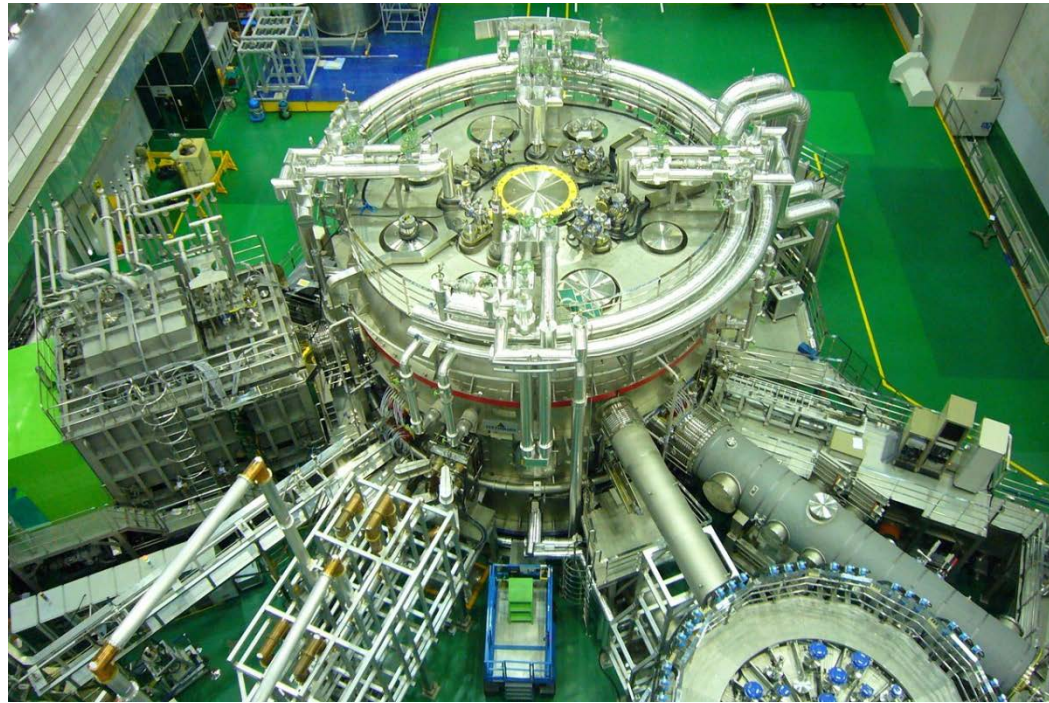
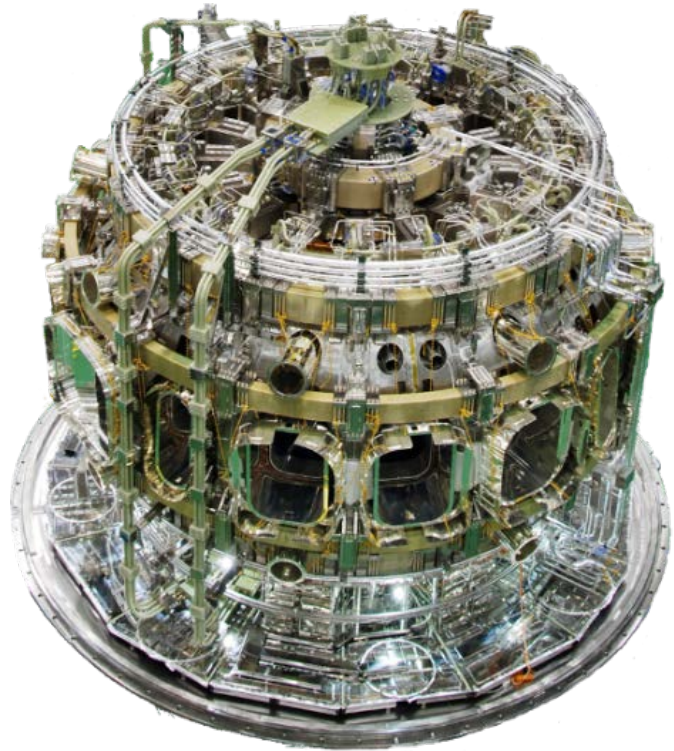
Tokamak

KSTAR (Korea Superconducting Tokamak Advanced Research):
 $R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$, 2007-today



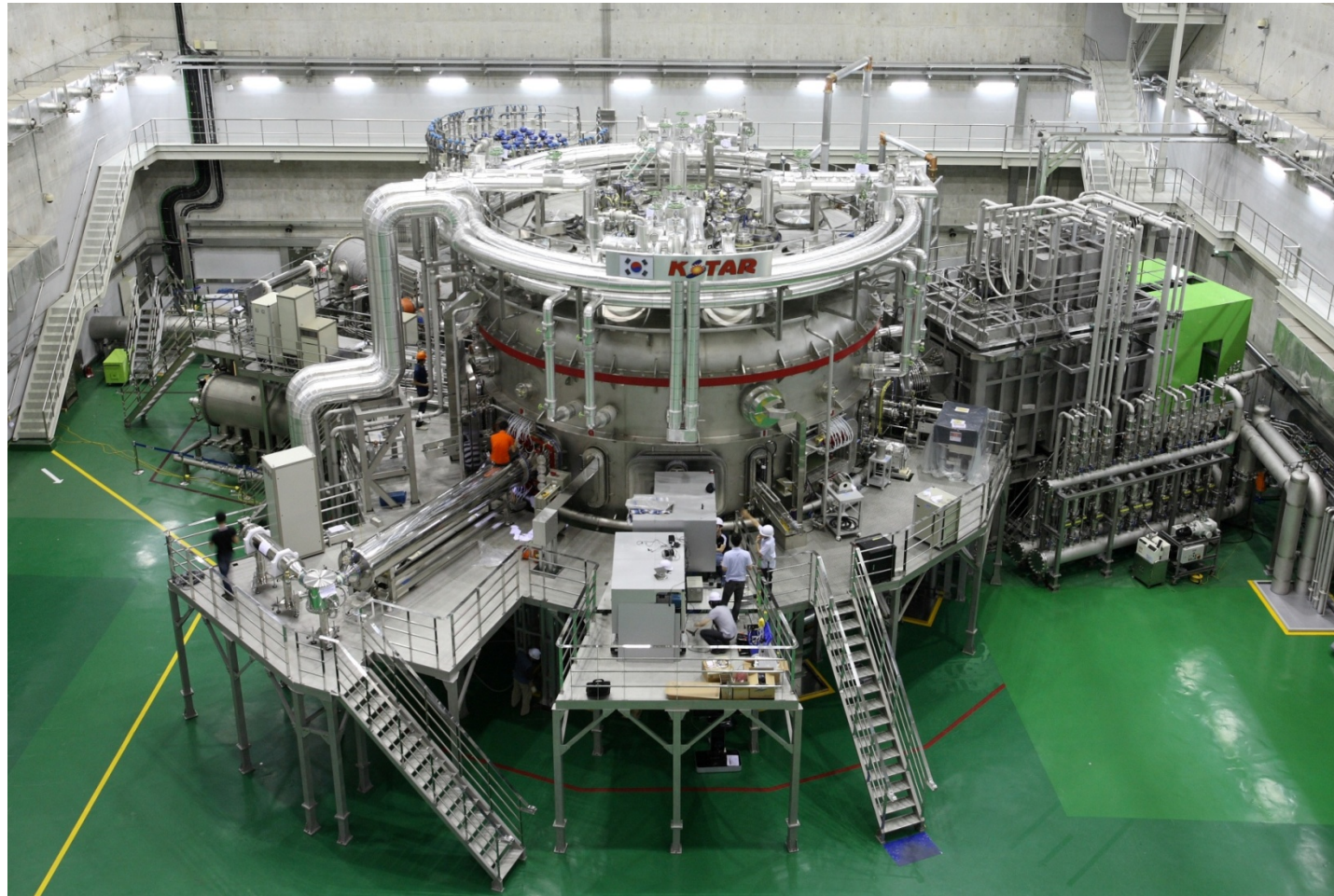
Tokamak

KSTAR (Korea Superconducting Tokamak Advanced Research):
 $R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$, 2007-today



Tokamak

KSTAR (Korea Superconducting Tokamak Advanced Research):
 $R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$, 2007-today



Tokamak

KSTAR (Korea Superconducting Tokamak Advanced Research):

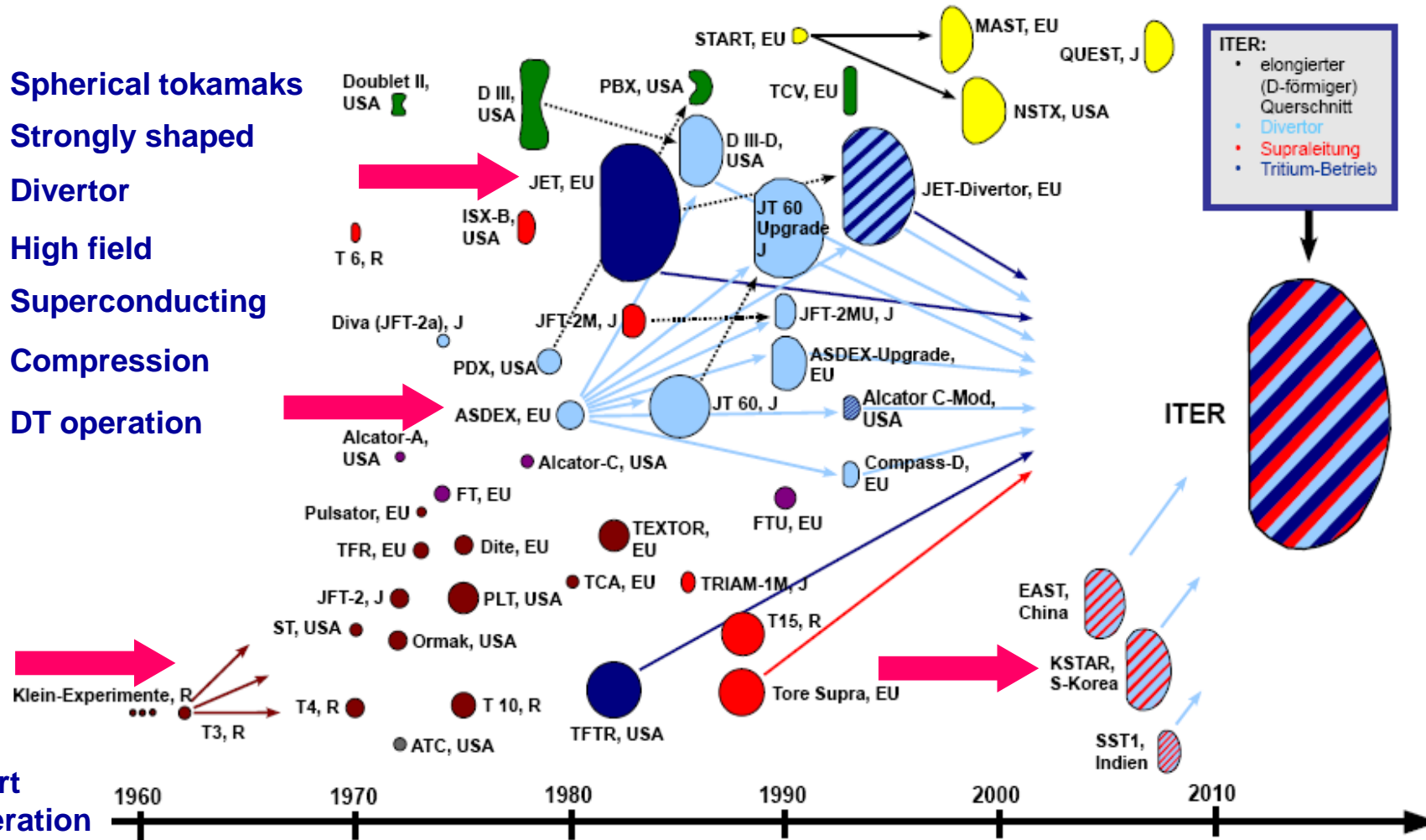
$R_0 = 1.8 \text{ m}$, $a = 0.5 \text{ m}$, 2007-today

KSTAR 1st plasma

Analyse the KSTAR 1st plasma

Tokamak

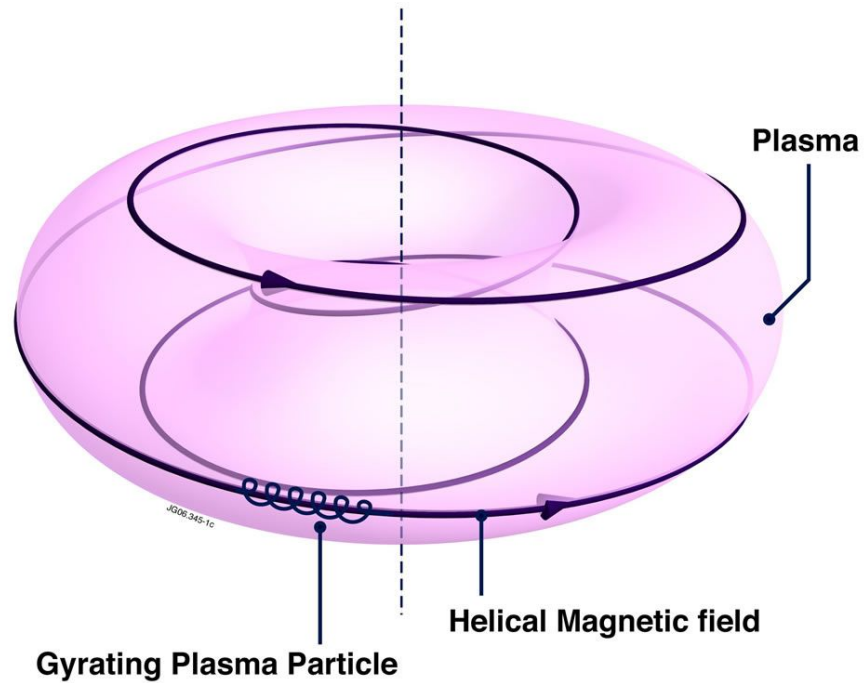
-  Spherical tokamaks
-  Strongly shaped
-  Divertor
-  High field
-  Superconducting
-  Compression
-  DT operation



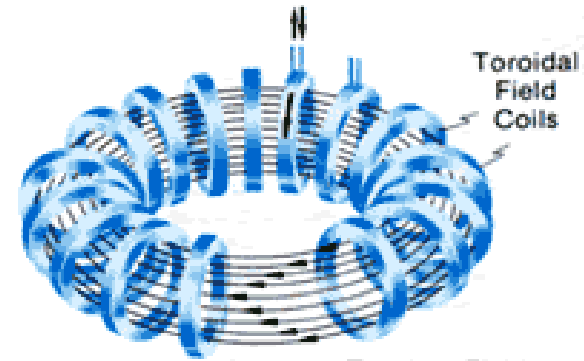
Basic tokamak variables

How to characterize a tokamak?

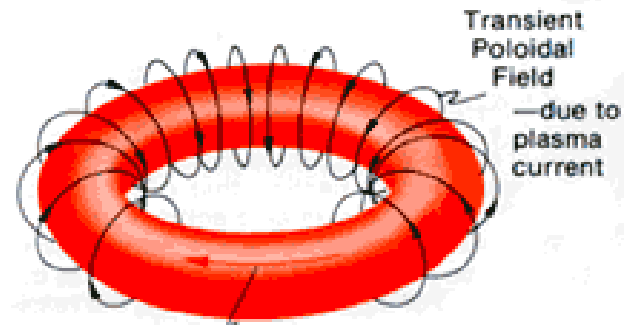
Tokamak



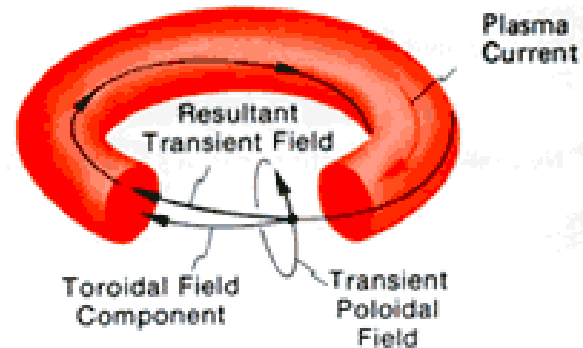
Relatively Constant Electric Current



Constant Toroidal Field

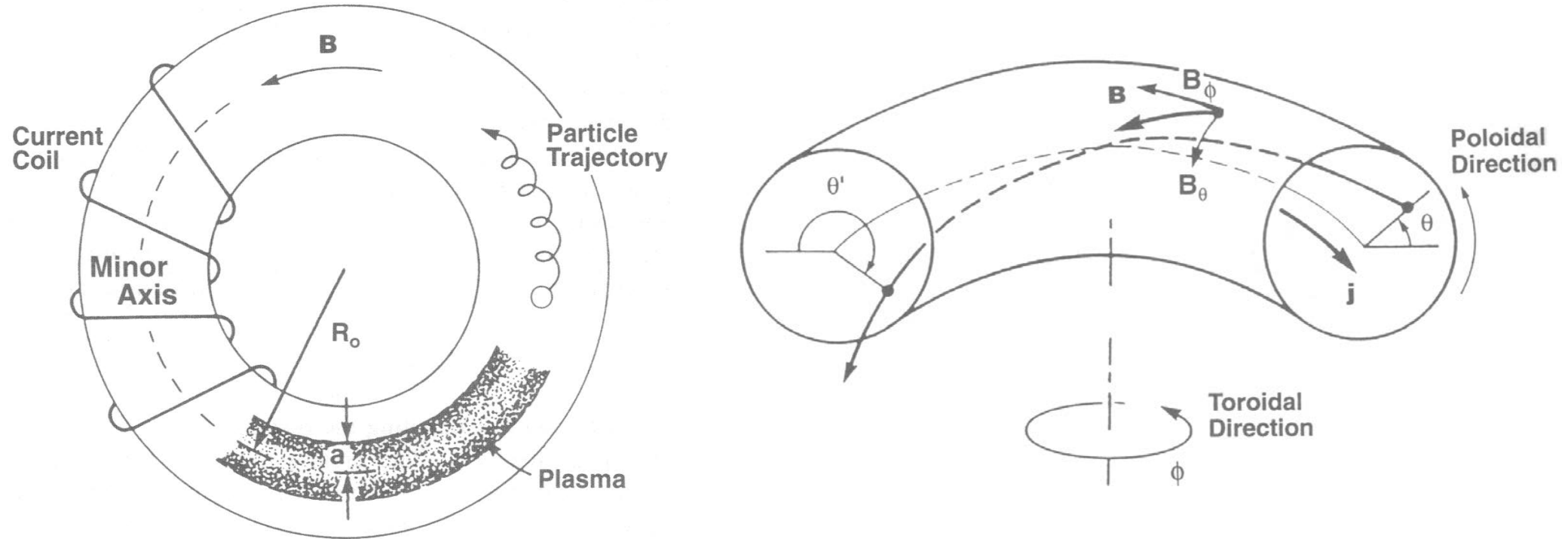


Transient Plasma Current



Basic Tokamak Variables

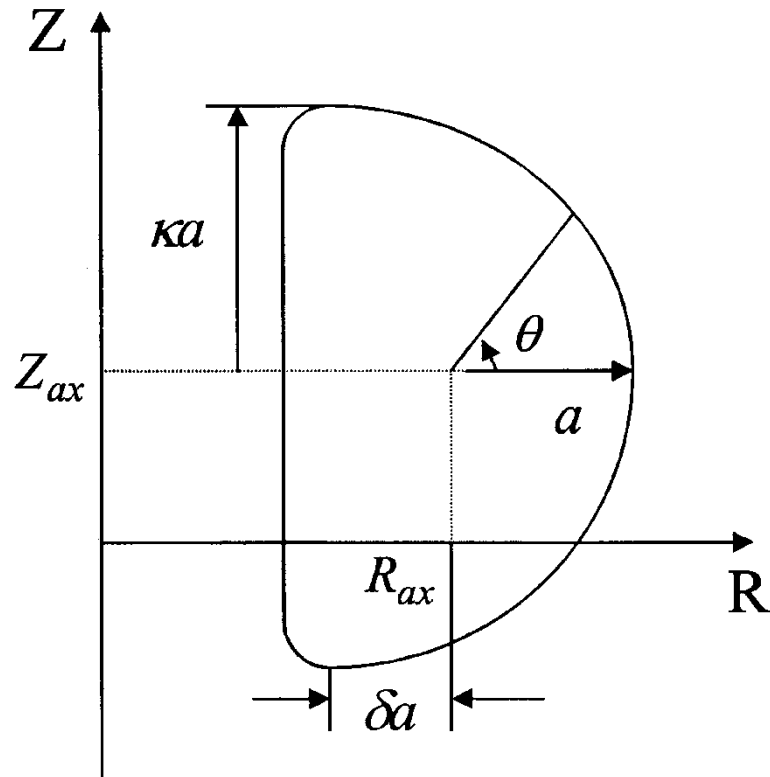
- Cylindrical and local coordinates for a tokamak



- Major radius: R_0 , minor radius: a
- Aspect ratio: $R_0/a \sim 3-5$
ex) KSTAR: 3.6, ITER: 3.1

Basic Tokamak Variables

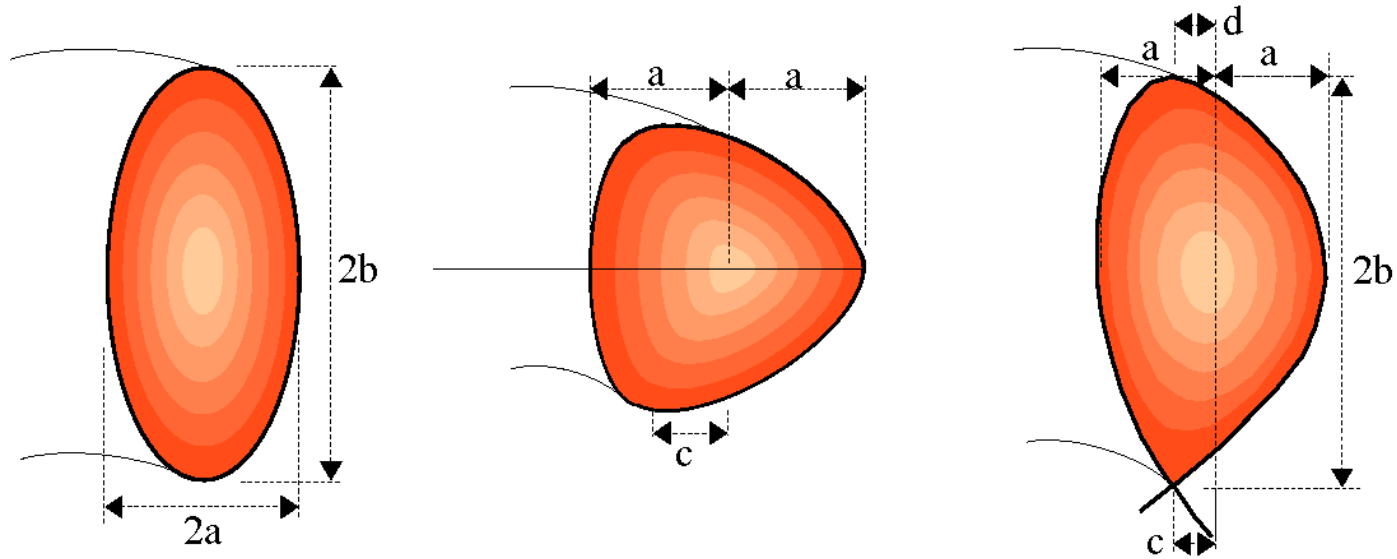
- Plasma equilibrium parameters



- Elongation: κ
- Triangularity: δ

Basic Tokamak Variables

- Plasma equilibrium parameters

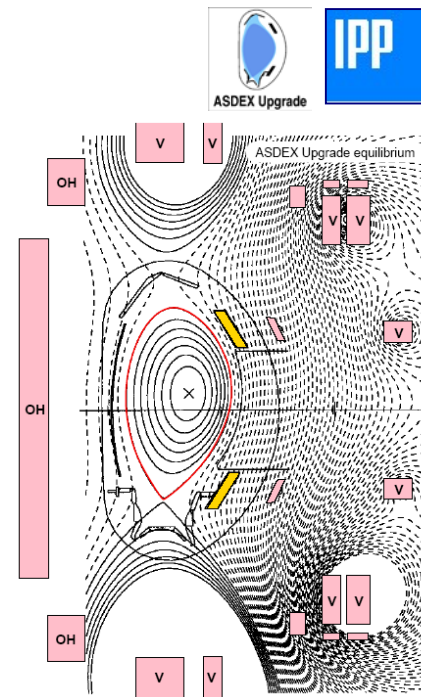
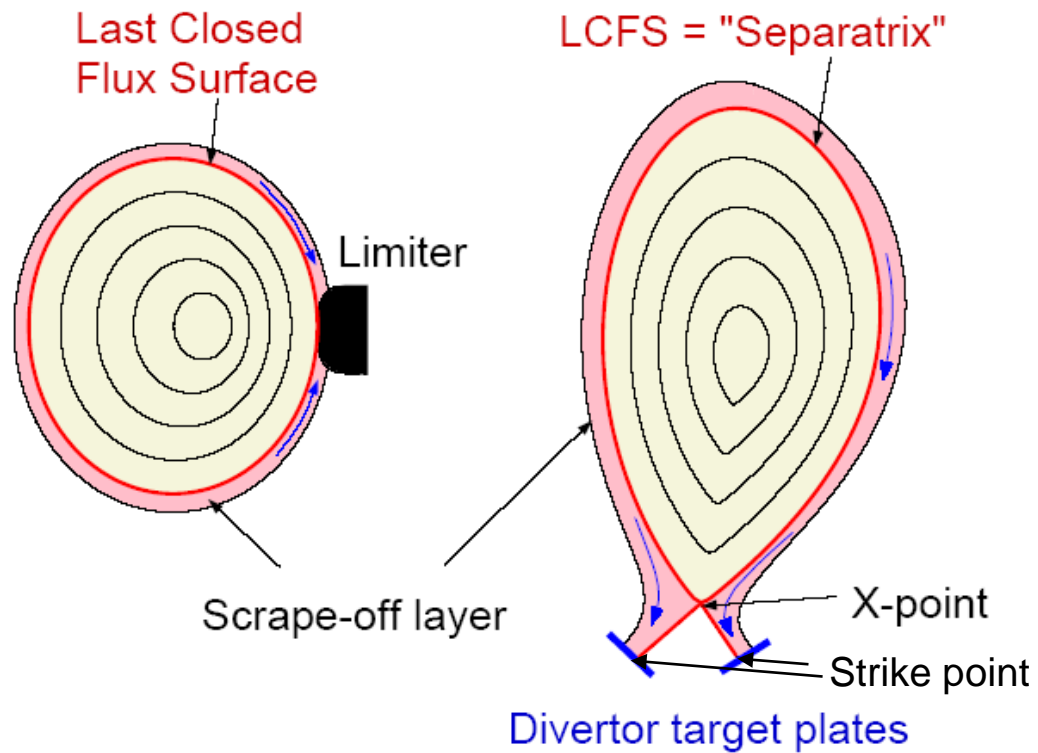


$$\kappa = \frac{b}{a}$$

$$\delta = \frac{c+d}{2a}$$

Basic Tokamak Variables

- Separation of plasma from wall by a limiter and a divertor

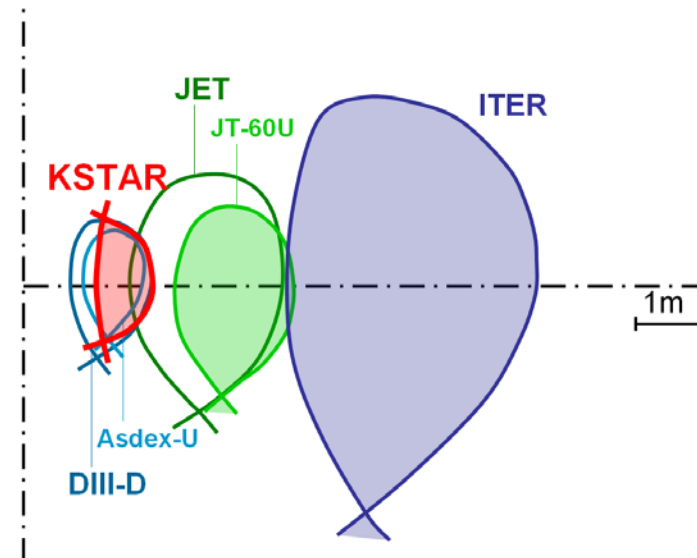


Basic Tokamak Variables

- Plasma equilibrium parameters

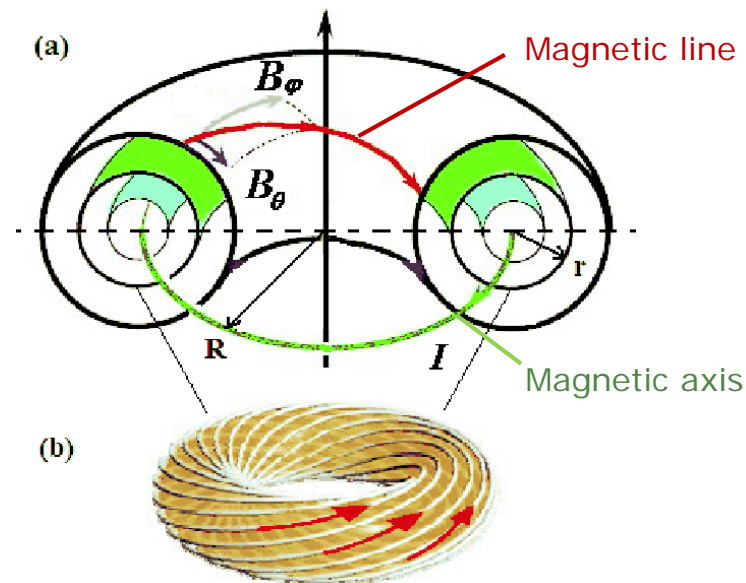
Parameters	KSTAR	ITER
Major Radius, R_0	1.8 m	6.2 m
Minor Radius, a	0.5 m	2.0 m
Plasma Current, I_p	2.0 MA	15 MA
Elongation, κ_x	2.0	1.85
Triangularity, δ_x	0.8	0.5
Toroidal Field, B_0	3.5 T	5.3 T
Pulse Length	300 s	500 s
Fuel	H, D	D, T

- Plasma shape



Magnetic Flux Surface

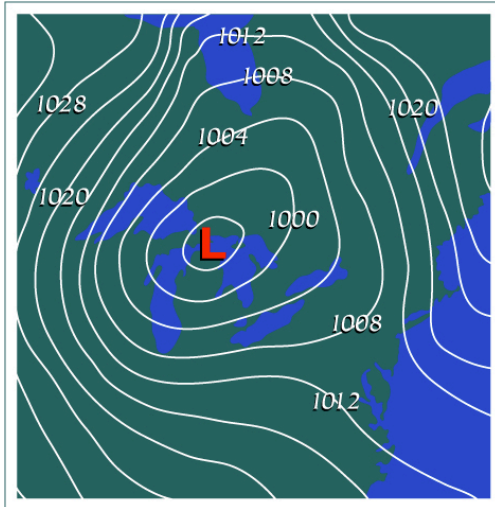
- In a tokamak configurations with confined plasmas the magnetic lines lie on a set of nested toroidal surfaces called flux surfaces.
- Pressure is constant along a magnetic field line.
- Magnetic lines lie in surfaces of constant pressure.
- Flux surfaces are surfaces of constant pressure.
- The current lines lie on surfaces of constant pressure.



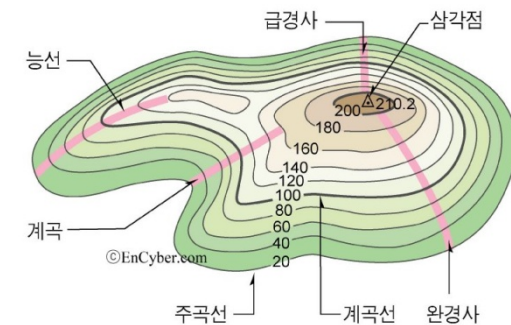
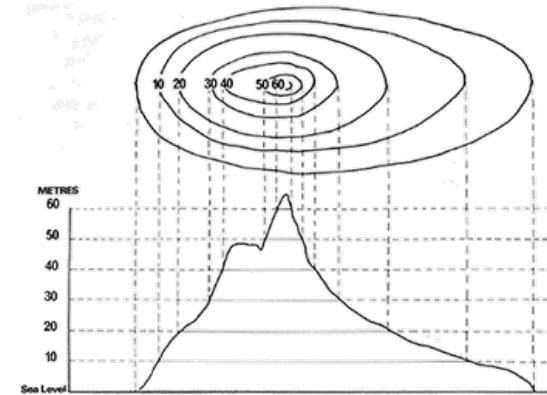
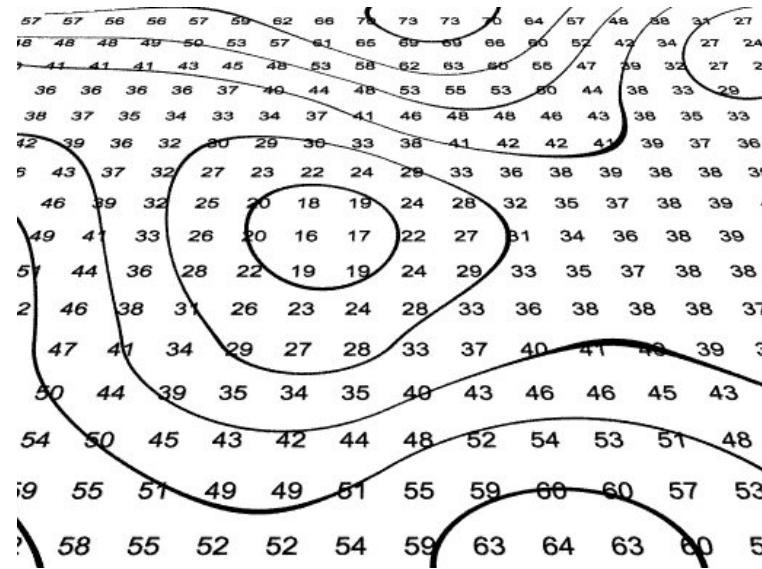
$$\vec{J} \times \vec{B} = \nabla p$$

$$\vec{B} \cdot \nabla p = 0 \quad \vec{J} \cdot \nabla p = 0$$

Magnetic Flux Surface



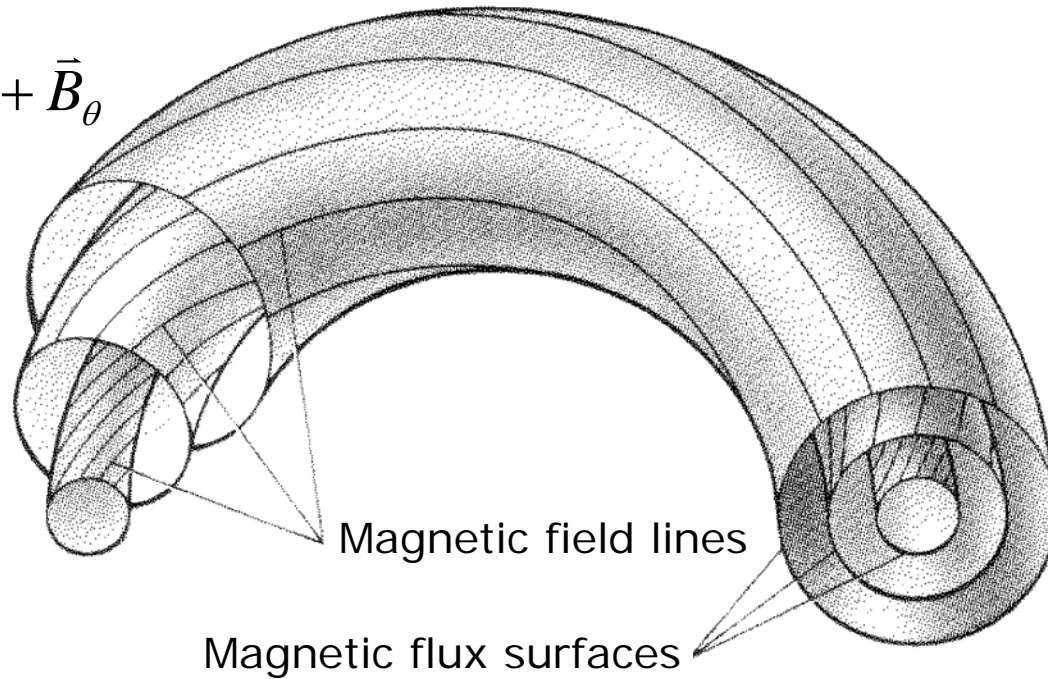
Isobar: $p = \text{const.}$



Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit

$$\vec{B} = \vec{B}_\phi + \vec{B}_\theta$$



- Rotational transform:
 $\Delta\theta$? when $\Delta\Phi = 2\pi$

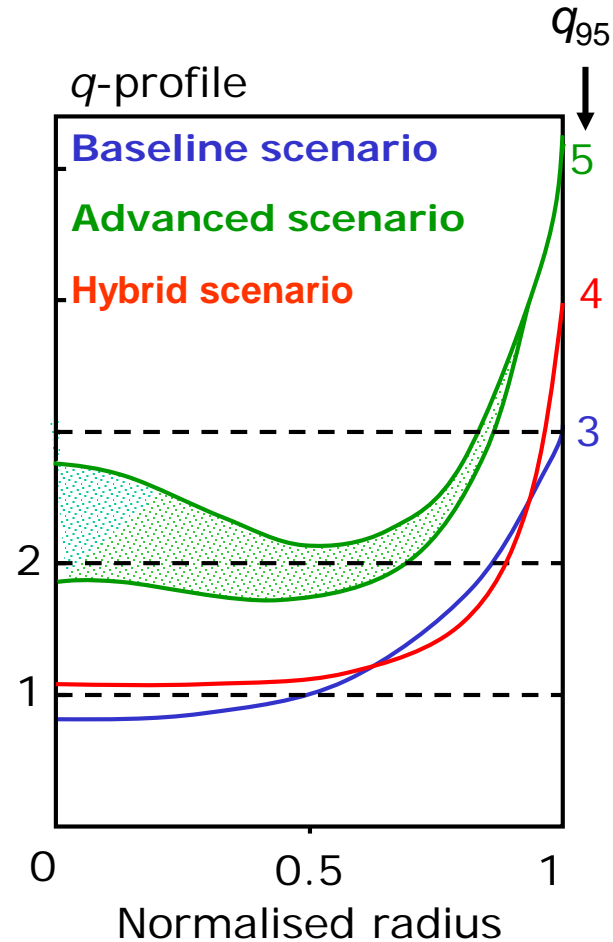
$$t = \frac{\Delta\theta}{\Delta\phi} = \frac{\frac{B_\theta}{r}}{\frac{B_\phi}{2\pi R}} = \frac{2\pi B_\theta}{r B_\phi}$$

$$\leftarrow \frac{R d\phi}{B_\phi} = \frac{r d\theta}{B_\theta}$$

$$q = \frac{\text{number of toroidal windings}}{\text{number of poloidal windings}} = \frac{2\pi}{t} = \frac{r B_\phi}{R B_\theta}$$

Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit



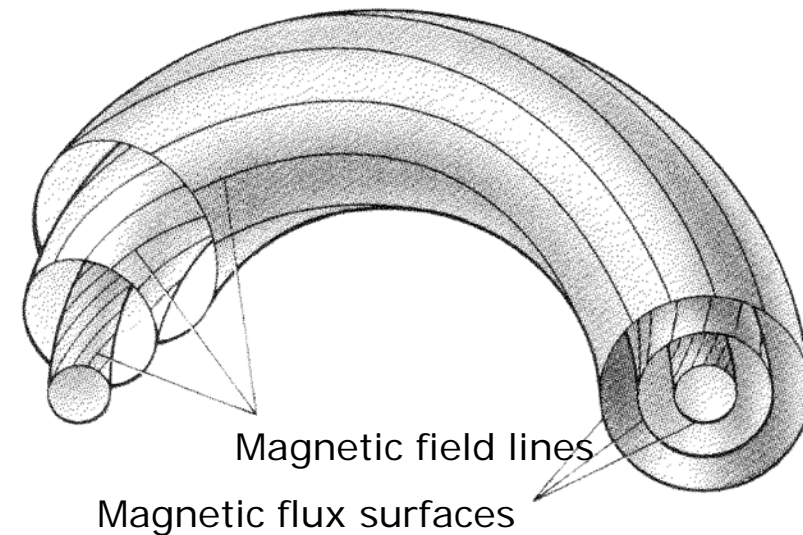
Basic Tokamak Variables

- Magnetic Shear

- Measuring the change in pitch angle of a magnetic field line from one flux surface to the next
- Playing an important role in stabilizing MHD instabilities, particularly those driven by the pressure gradient:

A perturbation aligned with $\mathbf{B}(r)$ will, at a point with increased minor radial distance $r+dr$, encounter field lines at a different angle which again will vary as the perturbation grows to another distance $r+dr'$. Any helically resonant instabilities are thus radially localised.

$$s(r) \equiv \frac{r}{q} \frac{dq}{dr}$$



Tokamak equilibrium

Tokamak Equilibrium

$$\nabla p = \vec{J} \times \vec{B}$$

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\nabla \cdot \vec{B} = 0$$

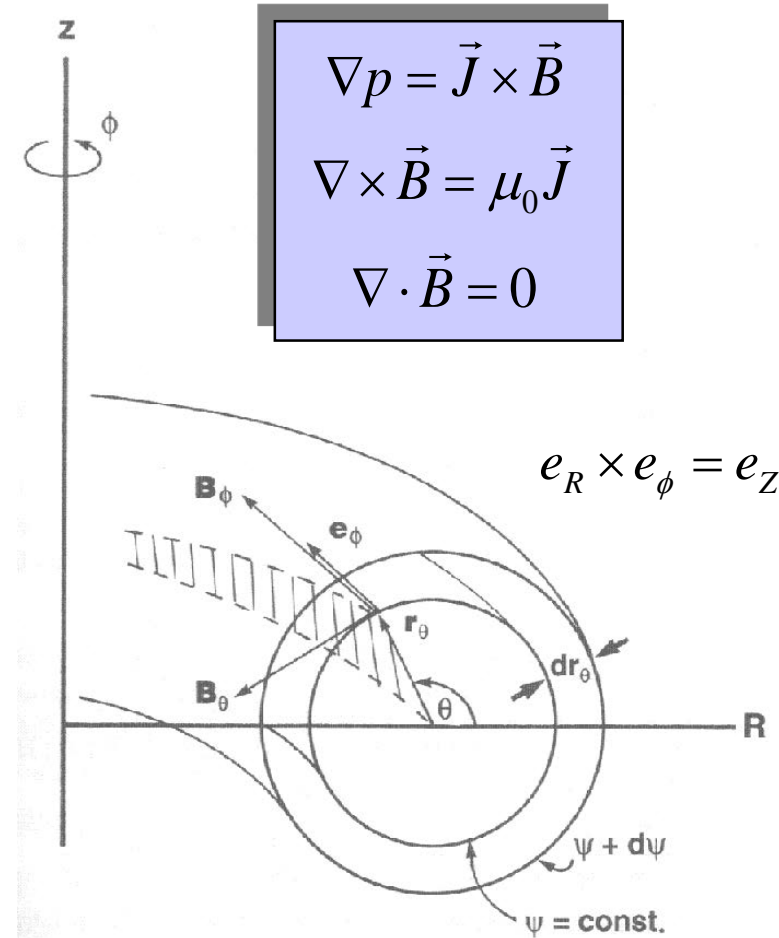
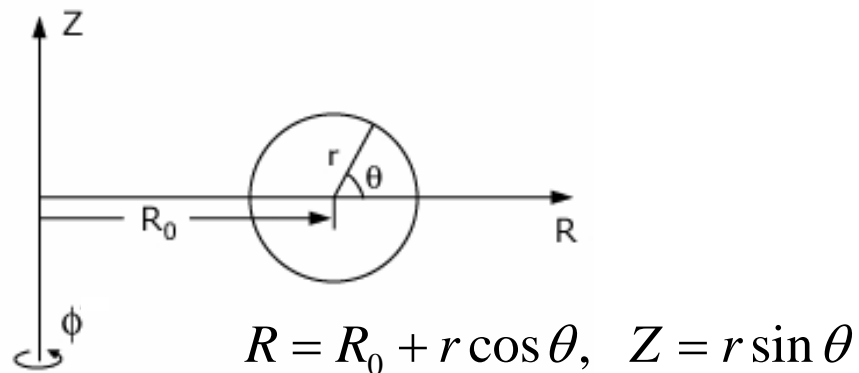
- Consider the axisymmetric torus, the simplest, multi-dimensional configuration
- We shall derive the Grad–Shafranov equation for axisymmetric equilibria.
- This provides a complete description of toroidal equilibrium:
radial pressure balance, toroidal force balance,
 β limits, q -profiles, etc.

Tokamak Equilibrium

- The Grad-Shafranov Equation

$$\Delta^* \psi = -\mu_0 R^2 \frac{dp}{d\psi} - F \frac{dF}{d\psi}$$

- Nonlinear
- Partial differential equation
- Grad and Rubin (1958), Shafranov (1960)
- Toroidal axisymmetric $\partial/\partial\phi = 0$



Tokamak Equilibrium

- The Grad-Shafranov Equation

Sequence of solution of the MHD equilibrium equations

1. The $\nabla \cdot \mathbf{B} = 0$
2. Ampere's law: $\mu_0 \mathbf{J} = \nabla \times \mathbf{B}$
3. The momentum equation: $\mathbf{J} \times \mathbf{B} = \nabla p$

$$\vec{J} \times \vec{B} = \nabla p$$

$$\nabla p = (\vec{J}_\phi + \vec{J}_p) \times (\vec{B}_\phi + \vec{B}_p)$$

$$\nabla p = \vec{J} \times \vec{B}$$

$$\nabla \times \vec{B} = \mu_0 \vec{J}$$

$$\nabla \cdot \vec{B} = 0$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

- Momentum equation $\vec{J} \times \vec{B} = \nabla p$

$$\nabla p = (\vec{J}_\phi + \vec{J}_p) \times (\vec{B}_\phi + \vec{B}_p)$$

- aim: to express each term with ψ and F $\psi = \frac{\Psi_p}{2\pi}, F = \frac{\mu_0 I_p}{2\pi}$

$$\nabla \cdot \vec{B} = 0 \quad \longrightarrow \quad \vec{B}_p = -\frac{1}{R} \vec{e}_\phi \times \nabla \psi$$

$$\mu_0 \vec{J}_p = \nabla \times \vec{B}_\phi \quad \longrightarrow \quad \vec{J}_p = -\frac{1}{\mu_0 R} \vec{e}_\phi \times \nabla F, \quad B_\phi = \frac{F}{R}$$

$$\mu_0 \vec{J}_\phi = (\nabla \times \vec{B})_\phi \quad \longrightarrow \quad J_\phi = -\frac{1}{\mu_0 R} \Delta^* \psi$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

- The $\nabla \cdot \mathbf{B}$ Equation

In cylindrical coordinates for toroidal axisymmetric fields ($\partial/\partial\Phi=0$)

$$\nabla \cdot \vec{B} = 0 \quad \frac{1}{R} \frac{\partial(RB_R)}{\partial R} + \frac{1}{R} \frac{\partial B_\phi}{\partial \phi} + \frac{\partial B_Z}{\partial Z} = 0$$

$$\vec{B} \cdot \nabla \psi = 0 \quad B_R \frac{\partial \psi}{\partial R} + B_Z \frac{\partial \psi}{\partial Z} = 0$$

Magnetic flux
surface

$$B_R = -\frac{1}{R} \frac{\partial \psi}{\partial Z}, \quad B_Z = \frac{1}{R} \frac{\partial \psi}{\partial R}$$

ψ : Stream function
for the poloidal
magnetic field

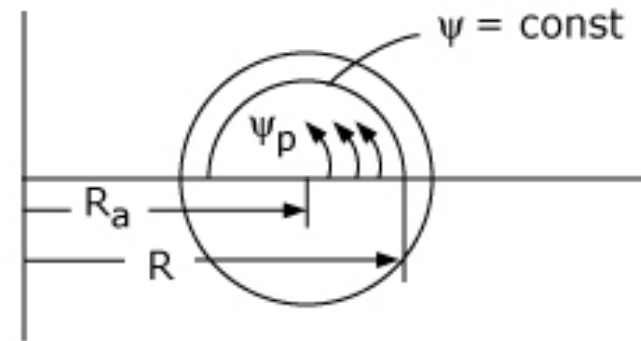
Tokamak Equilibrium

- The Grad-Shafranov Equation

- The stream function ψ is closely related to the poloidal flux in the plasma.

$$\begin{aligned}\psi_p &= \int \vec{B}_p \cdot d\vec{A} \\ &= \int_0^{2\pi} d\phi \int_{R_a}^R R dR B_z(R, Z=0) \\ &= \int_{R_a}^R 2\pi R \frac{1}{R} \frac{\partial \psi}{\partial R} dR \quad \leftarrow B_z = \frac{1}{R} \frac{\partial \psi}{\partial R} \\ &= 2\pi [\psi(R, 0) - \psi(R_a, 0)] = 2\pi \psi\end{aligned}$$

Poloidal flux on axis is zero



Tokamak Equilibrium

- The Grad-Shafranov Equation

- The $\nabla \cdot \mathbf{B}$ Equation

$$\nabla \cdot \vec{B} = 0 \quad \frac{1}{R} \frac{\partial(RB_R)}{\partial R} + \frac{1}{R} \frac{\partial B_\phi}{\partial \phi} + \frac{\partial B_Z}{\partial Z} = 0$$

$$\vec{B} \cdot \nabla \psi = 0 \quad B_R \frac{\partial \psi}{\partial R} + B_Z \frac{\partial \psi}{\partial Z} = 0$$

Magnetic flux surface

ψ : Stream function for the poloidal magnetic field

$$B_R = -\frac{1}{R} \frac{\partial \psi}{\partial Z}, \quad B_Z = \frac{1}{R} \frac{\partial \psi}{\partial R}$$

$$\vec{B} = B_\phi \vec{e}_\phi + \vec{B}_p$$

$$2\pi d\psi = 2\pi(\psi + d\psi) - 2\pi\psi = \int_{A+\delta A} \vec{B}_p \cdot d\vec{A} - \int_A \vec{B}_p \cdot d\vec{A} = \int_{\delta A} \vec{B}_p \cdot d\vec{A} \approx B_p 2\pi R dr$$

$$RB_p = \frac{d\psi}{dr} = |\nabla \psi| \quad \leftarrow 2\pi\psi = \psi_p = \int \vec{B}_p \cdot d\vec{A}$$

$$\vec{B}_p = -\frac{1}{R} \vec{e}_\phi \times \nabla \psi$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

- Ampere's law

$$\mu_0 \vec{J}_p = \nabla \times \vec{B}_\phi$$

$$\int_{\partial A} \nabla \times \vec{B}_\phi \cdot d\vec{A} = \oint \vec{B}_\phi \cdot d\vec{l} = 2\pi d(RB_\phi) = 2\pi d\psi \frac{\partial(RB_\phi)}{\partial\psi}$$

$$\int_{\partial A} \mu_0 \vec{J}_p \cdot d\vec{A} = \mu_0 J_p 2\pi R dr$$

$$\Rightarrow \frac{\partial(RB_\phi)}{\partial\psi} = \frac{\mu_0 J_p R}{(d\psi/dr)}$$

$$\Rightarrow \mu_0 J_p R = \frac{\partial F}{\partial\psi} \frac{\partial\psi}{\partial r} = \frac{\partial F}{\partial\psi} |\nabla\psi| = \frac{\partial F}{\partial r} = |\nabla F(\psi)| \quad \leftarrow F(\psi) \equiv RB_\phi$$

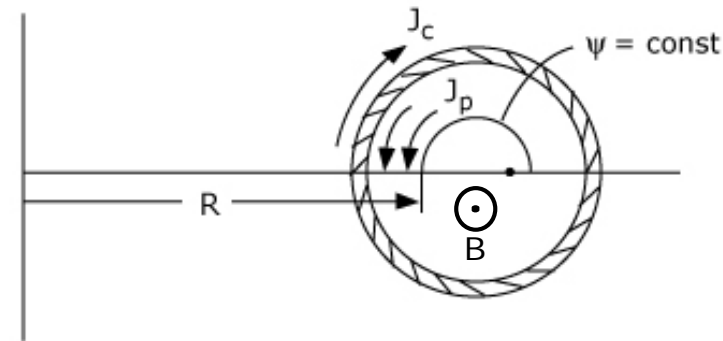
$$\vec{J}_p = -\frac{1}{\mu_0 R} \vec{e}_\phi \times \nabla F, \quad B_\phi = \frac{F}{R}$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

- Interpretation of $F(\psi)$

$$\begin{aligned}
 I_p &= \int \vec{J}_p \cdot d\vec{A} \\
 &= \int_0^{2\pi} d\phi \int_0^R R dR J_z(R, Z=0) \\
 &= \int_0^{2\pi} d\phi \int_0^R R dR \frac{1}{\mu_0 R} \frac{\partial}{\partial R} (R B_\phi) \\
 &= \int_0^R 2\pi R \frac{1}{\mu_0 R} \frac{\partial F}{\partial R} dR \quad \leftarrow \quad J_z = \frac{1}{\mu_0} \frac{1}{R} \frac{\partial}{\partial R} (R B_\phi), \quad F(\psi) \equiv R B_\phi \\
 &= 2\pi [F(R,0) - F(0,0)] / \mu_0 \quad R=0 \rightarrow F(\psi) = 0 \\
 &= 2\pi F(\psi) / \mu_0
 \end{aligned}$$



$I_p(\psi)$ is the total poloidal current and the toroidal field coils passing through the circle $\psi(R, 0) = \text{const}$.

Tokamak Equilibrium

- The Grad-Shafranov Equation
 - Ampere's law

$$\mu_0 \vec{J} = \nabla \times \vec{B}$$

$$J_\phi = \frac{1}{\mu_0} \left(\frac{\partial B_R}{\partial Z} - \frac{\partial B_Z}{\partial R} \right) = -\frac{1}{\mu_0} \left(\frac{\partial}{\partial R} \left(\frac{1}{R} \frac{\partial \psi}{\partial R} \right) + \frac{1}{R} \frac{\partial^2 \psi}{\partial Z^2} \right)$$
$$\equiv -\frac{1}{\mu_0 R} \Delta^* \psi \quad \leftarrow B_R = -\frac{1}{R} \frac{\partial \psi}{\partial Z}, \quad B_Z = \frac{1}{R} \frac{\partial \psi}{\partial R}$$

$$\Delta^* \psi \equiv R \frac{\partial}{\partial R} \left(\frac{1}{R} \frac{\partial \psi}{\partial R} \right) + \frac{\partial^2 \psi}{\partial Z^2} \quad \text{elliptic operator}$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

- Momentum equation

$$\vec{J} \times \vec{B} = \nabla p$$

$$\nabla p = (\vec{J}_\phi + \vec{J}_p) \times (\vec{B}_\phi + \vec{B}_p)$$

$$= \left(\vec{J}_\phi - \frac{1}{\mu_0 R} \vec{e}_\phi \times \nabla F \right) \times \left(\vec{B}_\phi - \frac{1}{R} \vec{e}_\phi \times \nabla \psi \right)$$

$$= -\vec{J}_\phi \times \left(\frac{1}{R} \vec{e}_\phi \times \nabla \psi \right) - \left(\frac{1}{\mu_0 R} \vec{e}_\phi \times \nabla F \right) \times \vec{B}_\phi$$

$$= \frac{\vec{J}_\phi}{R} \nabla \psi - \frac{\vec{B}_\phi}{\mu_0 R} \nabla F \quad \leftarrow A \times (B \times C) = B(A \cdot C) - C(A \cdot B)$$

Symmetry

$$\vec{J}_p = -\frac{1}{\mu_0 R} \vec{e}_\phi \times \nabla F$$

$$\vec{B}_p = -\frac{1}{R} \vec{e}_\phi \times \nabla \psi$$

$$B_\phi = \frac{F(\psi)}{R}$$

$$J_\phi = -\frac{1}{\mu_0 R} \Delta^* \psi$$

$$J_\phi = R \frac{dp}{d\psi} + \frac{B_\phi}{\mu_0} \frac{dF}{d\psi} = R \frac{dp}{d\psi} + \frac{F(\psi)}{\mu_0 R} \frac{dF}{d\psi} = -\frac{1}{\mu_0 R} \Delta^* \psi$$

Tokamak Equilibrium

- The Grad-Shafranov Equation

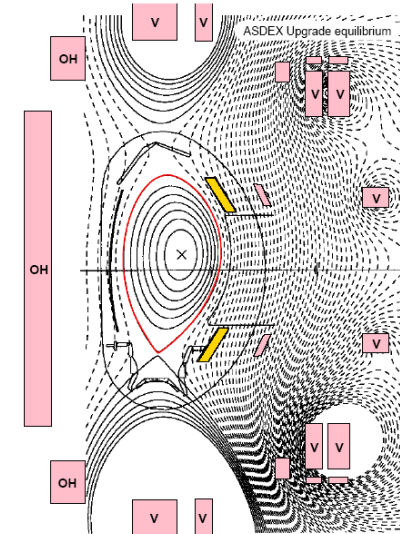
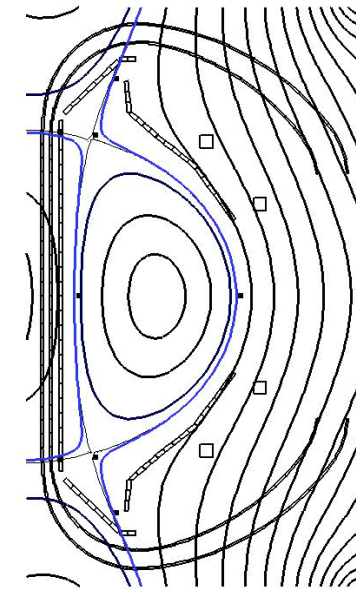
$$\Delta^* \psi = -\mu_0 R^2 \frac{dp}{d\psi} - F \frac{dF}{d\psi}$$

↑
 $\mathbf{J}_\phi \times \mathbf{B}_p$:
 strength
 of \mathbf{B}_p

↑
 ∇p :
 plasma load
 on a magnetic
 flux surface

↑
 $\mathbf{J}_p \times \mathbf{B}_\phi$:
 strength
 of \mathbf{B}_ϕ

- BCs: provided by the transformer-induced poloidal magnetic field outside the plasma
- In practice, the G-S equation is solved numerically to find the geometrical location of the magnetic surfaces and the radial distribution of the axial current density in a way that is consistent with the experimentally measured pressure profiles (p) and the externally applied field (F).



References

Lesch, Astrophysics, IPP Summer School (2008)