

전자물리특강: 유기반도체의 전자 구조

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Brief Review of QM

Time dependent Schrodinger equation

$$i\hbar \frac{\partial \Psi}{\partial t} = -\frac{\hbar^2}{2m} \nabla^2 \Psi + V\Psi$$

Time independent Schrodinger equation

$$-\frac{\hbar^2}{2m} \nabla^2 \psi + (V(x, y, z) - E)\psi(x) = 0$$

In spherical coordinates,

$$\int_0^{2\pi} \int_0^\pi \int_0^\infty |\psi(r, \theta, \phi)|^2 r^2 \sin \theta dr d\theta d\phi = 1$$

Particle in a wall

$$V = \infty \quad \left| \begin{array}{c} V = 0 \\ \hline 0 & x & L \end{array} \right| \quad V = \infty$$
$$-\frac{\hbar^2}{2m} \frac{d^2 \psi}{dx^2} - E \psi = 0 \quad E = \frac{\hbar^2 k^2}{2m}$$
$$\psi_n = \left(\frac{2}{L} \right)^{1/2} \sin \frac{n \pi x}{L}$$
$$E_n = \frac{\hbar^2}{2m} \left(\frac{n \pi}{L} \right)^2 = \frac{n^2 \pi^2 \hbar^2}{2m L^2}$$



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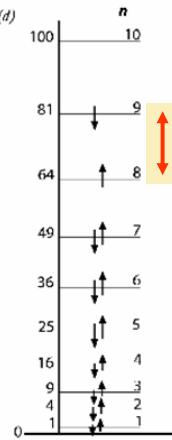
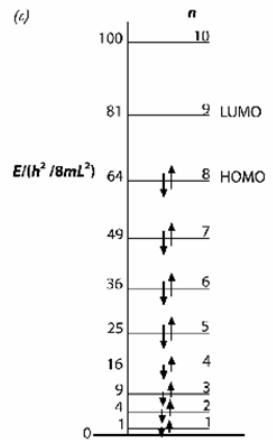
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Energy gap – simple quantum mechanics

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Polymer: Chain of N atoms separated by a (L=Na)
Electrons from p-orbitals fill the energy levels.

Particle in a box.



Energy levels

$$E_n = \frac{h^2 n^2}{8mL^2}$$

HOMO level

$$E(HOMO) = \frac{h^2}{8m(Na)^2} (N/2)^2$$

LUMO level

$$E(LUMO) = \frac{h^2}{8m(Na)^2} (N/2+1)^2$$

$$\text{Energy Gap } \Delta E = E(LUMO) - E(HOMO) = \frac{h^2}{8m(Na)^2} (N+1) \approx \frac{h^2}{8ma^2} \left(\frac{1}{N}\right)$$



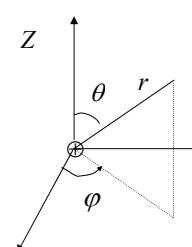
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Atomic orbitals

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Hydrogen atom

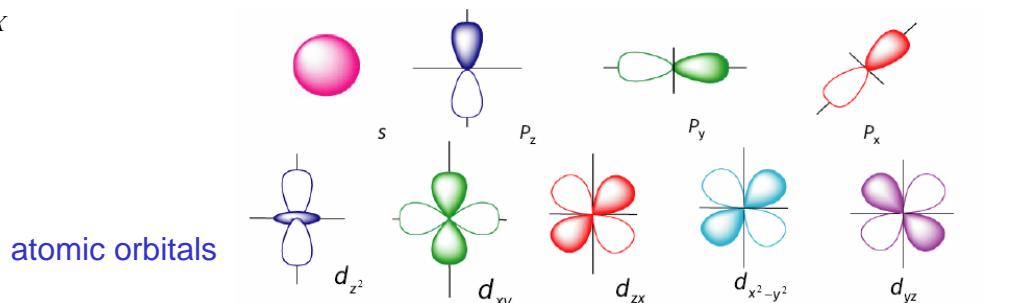


$$\nabla^2 \psi + \left(\frac{2m}{\hbar^2} \right) \{E - V(r)\} \psi = 0, \quad V(r) = -\frac{q}{r}$$

$$\frac{1}{r^2} \frac{\partial}{\partial r} \left(r^2 \frac{\partial}{\partial r} \right) \psi + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \left(\sin \theta \frac{\partial}{\partial \theta} \right) \psi + \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \psi + \frac{2m}{\hbar^2} \left(\frac{q^2}{r} + E \right) \psi = 0$$

$$\psi(r, \theta, \phi) = R(r) \Theta(\theta) \Phi(\phi)$$

X



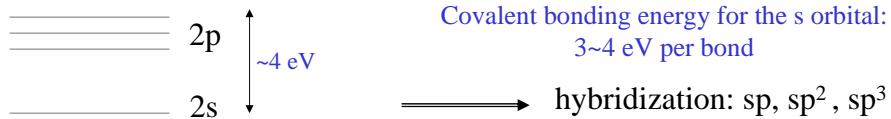
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Carbon and Hybridization

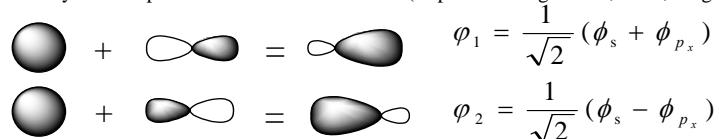
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Electronic Configuration of carbon: C=1s²2s²2p²



“In carbon, three possible hybridizations occur: sp, sp², and sp³; other group IV elements such as Si, Ge exhibit primarily sp³ hybridization. Carbon differs from Si and Ge insofar as carbon does not have inner atomic orbitals except for the spherical 1s orbitals, and the absence of nearby inner orbitals facilitates hybridizations involving only valence s and p orbitals for carbon.”

Ref. R. Saito *et al.* “Physical Properties of Carbon Nanotubes” (Imperial College Press, 1998, Singapore) p.5.



Orthonormality: $\langle \varphi_1 | \varphi_1 \rangle = 1$, $\langle \varphi_2 | \varphi_2 \rangle = 1$, $\langle \varphi_1 | \varphi_2 \rangle = 0$, $\langle \varphi_2 | \varphi_1 \rangle = 0$

sp hybridization



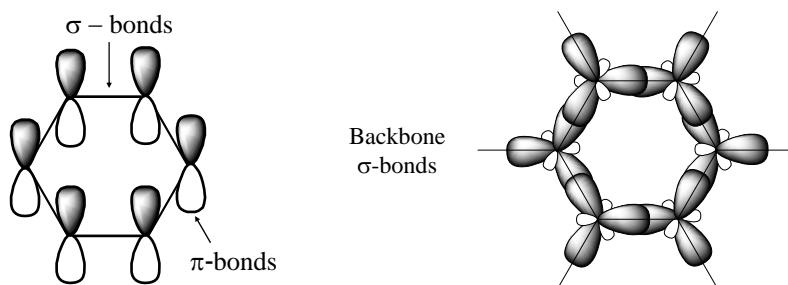
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Electronic Configuration of Benzene

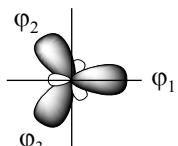
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Benzene: sp² (σ bonds) + p_z (π bonds)



sp² (σ bonds)

$$\varphi_1 = \frac{1}{\sqrt{3}} (\phi_s + \sqrt{2} \phi_{p_x})$$



$$\varphi_2 = \frac{1}{\sqrt{3}} (\phi_s - \frac{1}{\sqrt{2}} \phi_{p_x} + \frac{\sqrt{3}}{2} \phi_{p_y})$$

$$\varphi_3 = \frac{1}{\sqrt{3}} (\phi_s - \frac{1}{\sqrt{2}} \phi_{p_x} - \frac{\sqrt{3}}{2} \phi_{p_y})$$



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Huckel Molecular Orbital Theory

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Basic Assumption

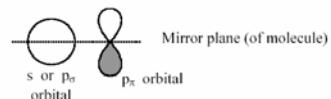
- The atomic orbitals contributing to the π -bonding are of a different symmetry to the atomic orbitals contributing to the σ -bonding and may be treated independently.
- Coulomb integrals for all the carbon atoms are assumed to be identical.
- All resonance integrals between directly-bonded atoms are assumed to be the same; whilst those between atoms that are not directly bonded are neglected.
- All overlap integrals representing the overlap of atomic orbitals centered on different atoms are neglected.

Ethylene C_2H_4

$$\psi = C_1\phi_1 + C_2\phi_2$$

$$\begin{vmatrix} H_{11} - ES_{11} & H_{12} - ES_{12} \\ H_{21} - ES_{21} & H_{22} - ES_{22} \end{vmatrix} = 0 \quad H_{ij} = \int \phi_i^* \hat{H} \phi_j d\tau \quad S_{ij} = \int \phi_i^* \phi_j d\tau$$

Assume $S_{ij} = \delta_{ij}$, $H_{ii} = \alpha$, $H_{ij} = \beta$ (neighboring atoms)



$$\begin{vmatrix} \alpha - E & \beta \\ \beta & \alpha - E \end{vmatrix} = 0$$

$$E = \alpha \pm \beta \quad (+ : \text{bonding}, - : \text{antibonding})$$



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Example: Ethylene Molecules

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Ethylene C_2H_4

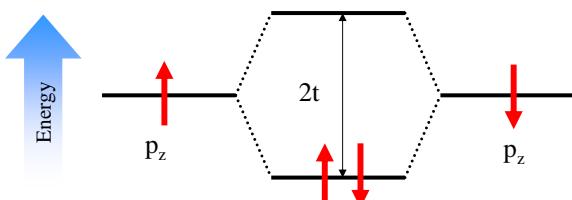
$$E = \alpha \pm \beta \quad (+ : \text{bonding}, - : \text{antibonding})$$

$$\begin{vmatrix} \alpha - E & \beta \\ \beta & \alpha - E \end{vmatrix} = 0$$

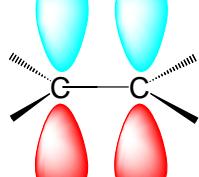
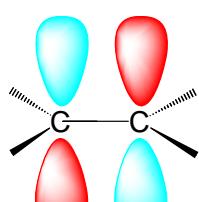
Ethylene C_2H_4

$$\Psi_{\pi^*} = \frac{1}{\sqrt{2}}(\varphi_1 - \varphi_2)$$

π^* (antibonding orbital)



$$\Psi_{\pi} = \frac{1}{\sqrt{2}}(\varphi_1 + \varphi_2)$$



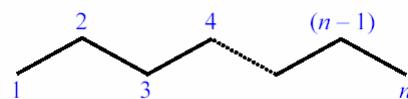
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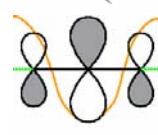
Example: Linear Conjugated Hydrocarbons

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General Solution (n -atom chain ; e.g. $C_nH_{(n+2)}$ conjugated polyenes)



$$c_s \propto \sin\left(\frac{\pi ks}{n+1}\right)$$



1 2 3

	1	2	3	\dots	n	
1	x	1	0	0	\dots	0
2	1	x	1			
3	0	1	x			
	0					
					x	1
n	0	\dots	\dots	\dots	1	x

= 0

n - total number of atoms in the conjugated chain

s - atom number (i.e. 1, 2, ..., n)

k - quantum number, identifying the MO (= 1, 2, ..., n)

$$E = \alpha + 2\beta \cos\left(\frac{\pi k}{n+1}\right)$$

$$(\alpha + 2\beta) \leq E \leq (\alpha - 2\beta)$$

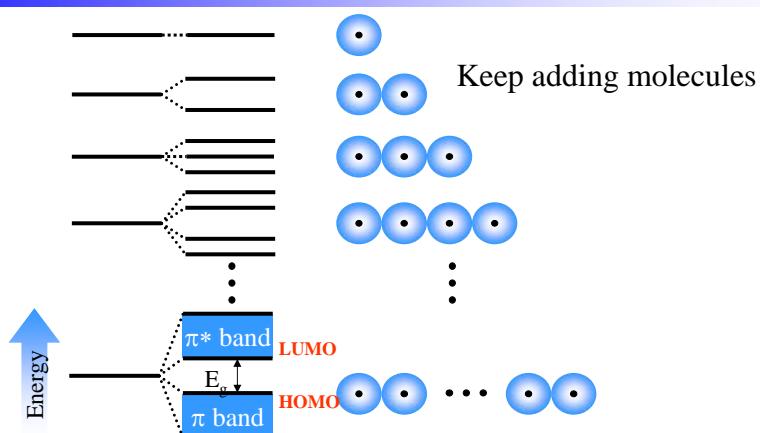


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Macromolecules (polymers) and solid: Formation of bands

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- The average energy separation between MOs decrease as n increases \rightarrow Formation of bands.
- The HOMO-LUMO separation, $\Delta E = EHOMO - ELUMO$, decreases as n increases.
- Energetically favored electron excitation is from the HOMO to the LUMO
- The interchain transfer integral (t) expresses the ease of charge transfer between two interacting chains.
- The larger the HOMO (LUMO) band-width (the magnitude of transfer integral), the higher the hole (electron) mobility.



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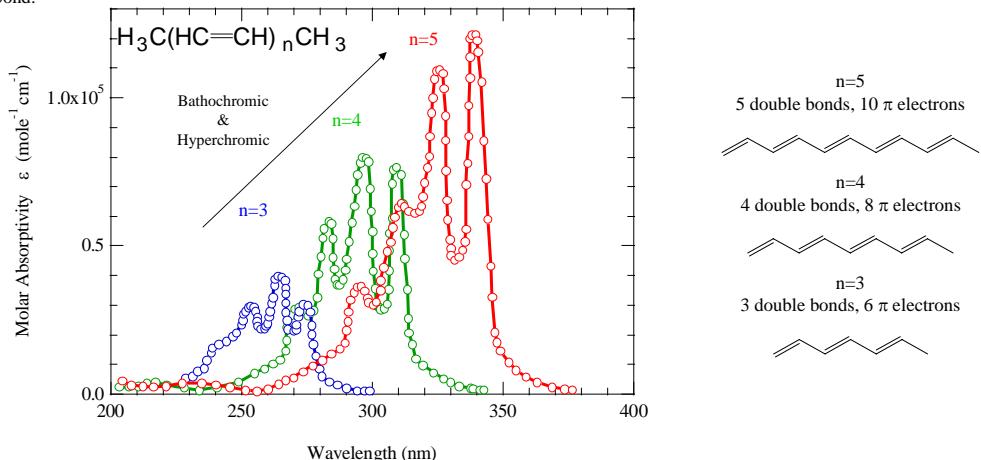
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Absorption spectra of polyenes: Dependence of conjugation

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Extending conjugation results in bathochromic and hyperchromic shifts in absorption.

The dimethylpolyene spectra show that each additional double bond in the conjugated π -electron system shifts the absorption maximum about 30 nm to longer wavelengths. Also, the molar absorptivity (ϵ) roughly doubles with each new conjugated double bond.



Ref. D. L. Pavia, G. M. Lampman, G. S. Kriz, Introduction to Spectroscopy, 2nd ed. (Saunders College Publishing, 1996), p.289



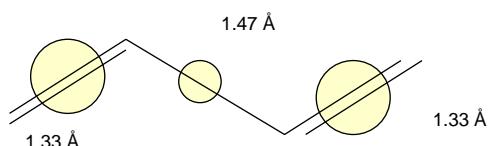
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Electron-Phonon coupling

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- Strong **electron-phonon coupling**:
 p -electron bond densities directly determine the geometric structure,
e.g., butadiene, C₄H₆.



- Vibronic progressions in optical absorption and photoluminescence spectra
- Stokes shift between optical absorption and emission.
- Polaron formation

* Charge transfer or low energy charge-excitations (via doping or photoexcitation) strongly affects the π -electron bond densities → geometric structure.

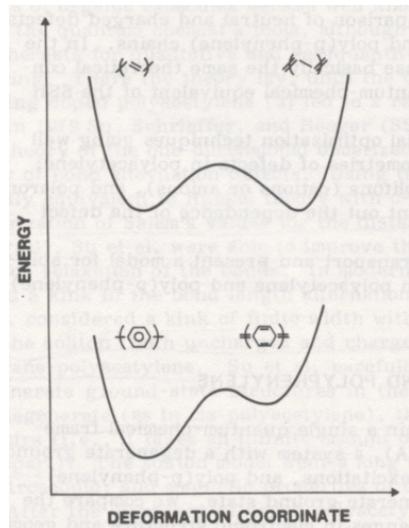


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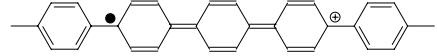
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Quasiparticles in Conjugated Polymers

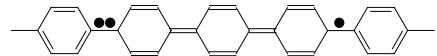
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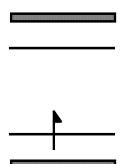
Potential Energy of Conjugated Polymers



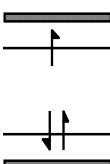
Positive polaron



Negative polaron



Positive polaron



Negative polaron

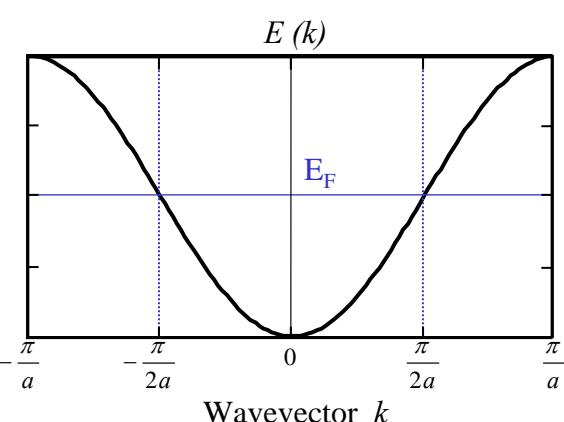
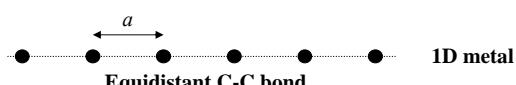


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1D infinite chain of carbons

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Tight-binding approximation: $E = E_{pz} - \beta - 2t \cos ka$



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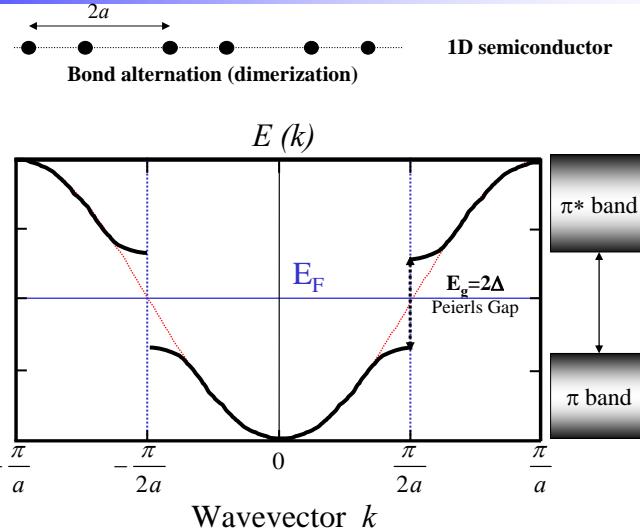
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Peierls instability → Energy Band gap

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Sir Rudolf E. Peierls
1907-1995



R.E. Peierls *Quantum Theory of Solids* (Oxford University, London Chap.5(1955).



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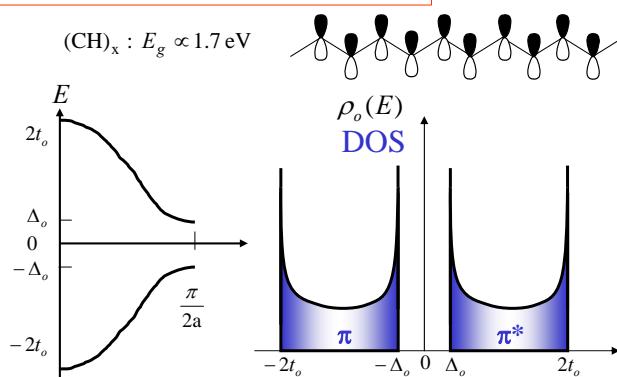
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Polyacetylene: SSH Hamiltonian

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$$\begin{aligned} H &= H_{kinetic} + H_{elastic} + H_{electronic} \\ &= \frac{M}{2} \sum_n \dot{u}_n^2 + \frac{K}{2} \sum_n (u_n - u_{n+1})^2 - \sum_{n,s} [t_0 + \alpha(u_n - u_{n+1})] (c_{n+1,s}^+ c_{n,s} + c_{n,s}^+ c_{n+1,s}) \end{aligned}$$

Peierls gap : $2\Delta = 16t_o \exp[-(1 + \frac{1}{2\lambda})]$, where $\lambda = \frac{2\alpha^2}{\pi t_o K}$.

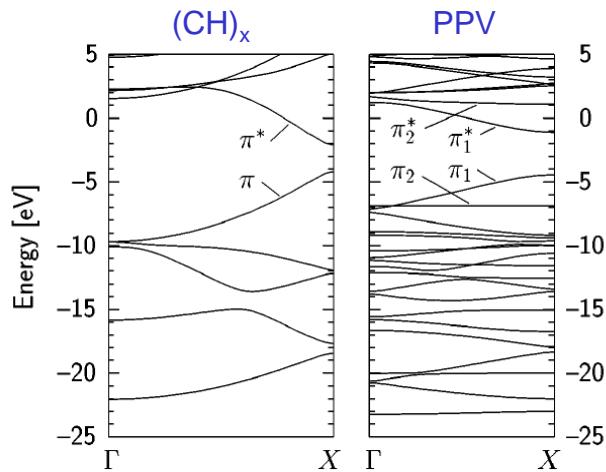


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Polyacetylene & PPV: Band structure

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M. Rohlfing and S. G. Louie, *Phys. Rev. Lett.* **82**, 1959 (1999)



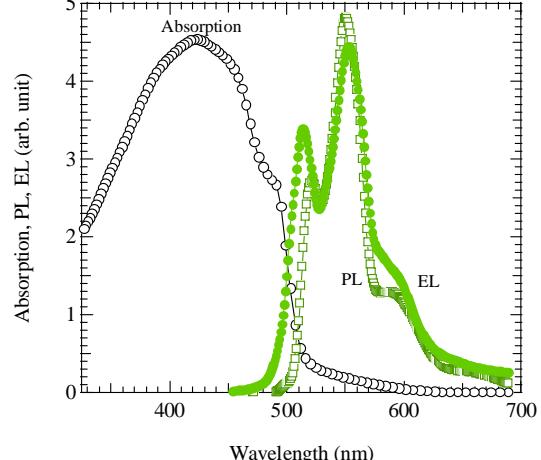
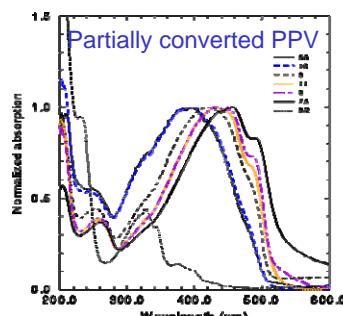
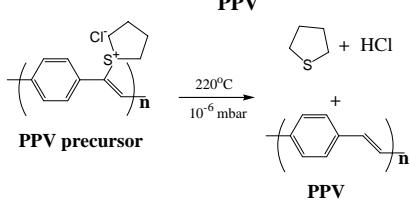
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Absorption and PL Spectra of Polymer: PPV

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Precursor route to poly(p-phenylenevinylene),
PPV

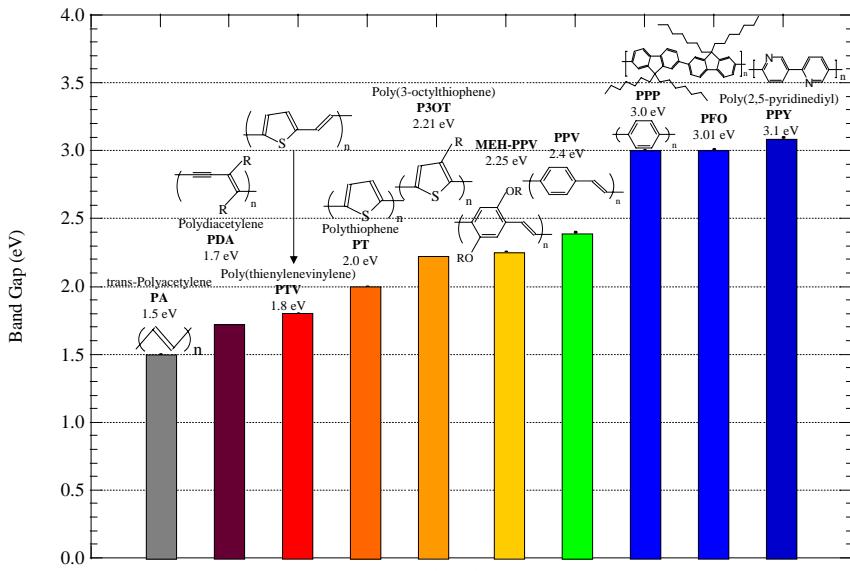


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π - π^* Energy Gap of Conjugated Polymers

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Ref. N. C. Greenham and R. H. Friend, Solid State Physics **49**, 1 (1995) p.13; A. Monkman et al., Phys. Rev. Lett. **6**, 1358 (2001)



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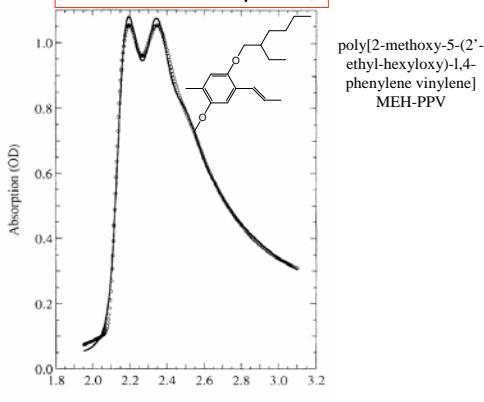
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DOS and Absorption

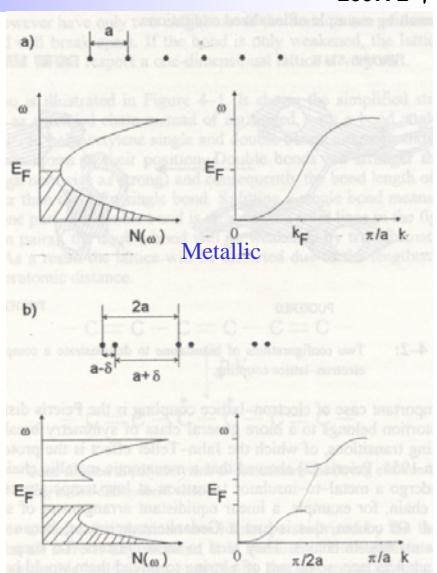
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Density of States (DOS):
Square-root singularity in 1-d.

$$D(E) = \frac{dN(E)}{dE} \propto \frac{1}{\sqrt{E - E_g}}$$



T. W. Hagler et al., Phys. Rev. B **51**, 14199 (1995)



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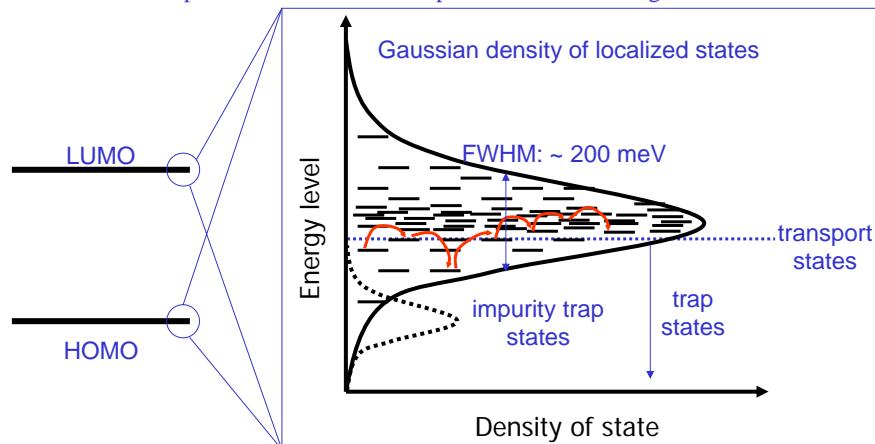
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Narrow band width and hopping transport

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- Weak van der Waals interaction between molecules

→ Each molecule keep its molecular levels independent of surrounding matrix in solid



Traps: Localized states below transport states



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