

2019 Fall

Introduction to Materials Science and Engineering

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Contents for previous class

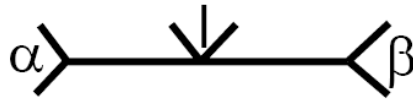
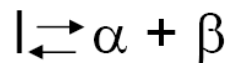
Chapter 11 Phase Diagrams Summary

- **Phase diagrams** are useful tools to determine:
 - the number and types of phases present,
 - the **composition** of each phase,
 - and the weight fraction of each phasegiven the temperature and composition of the system.
- The microstructure of an alloy depends on
 - its composition, and
 - whether or not cooling rate allows for maintenance of equilibrium.
- Important phase diagram phase transformations include eutectic, eutectoid, and peritectic.

Review of Invariant Binary Reactions

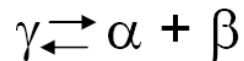
Eutectic Type

Eutectic



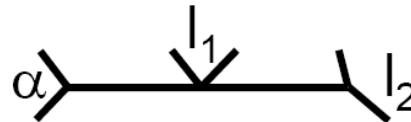
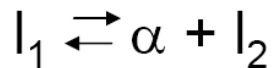
Al-Si, Fe-C

Eutectoid



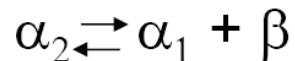
Fe-C

Monotectic



Cu-Pb

Monotectoid



Al-Zn, Ti-V

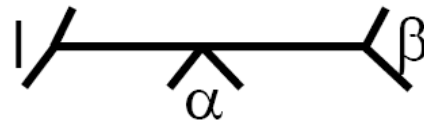
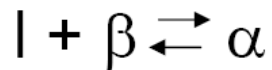
On cooling one phase going to two phases

Metatectic reaction: $\beta \leftrightarrow L + \alpha$ **Ex. Co-Os, Co-Re, Co-Ru** 3

Review of Invariant Binary Reactions

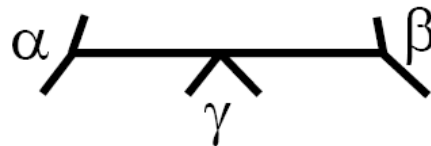
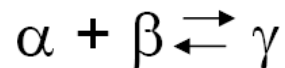
Peritectic Type

Peritectic



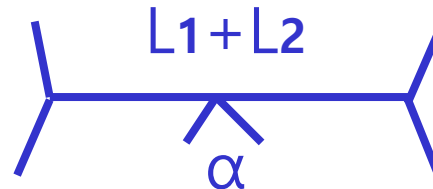
Fe-C

Peritectoid



Cu-Al

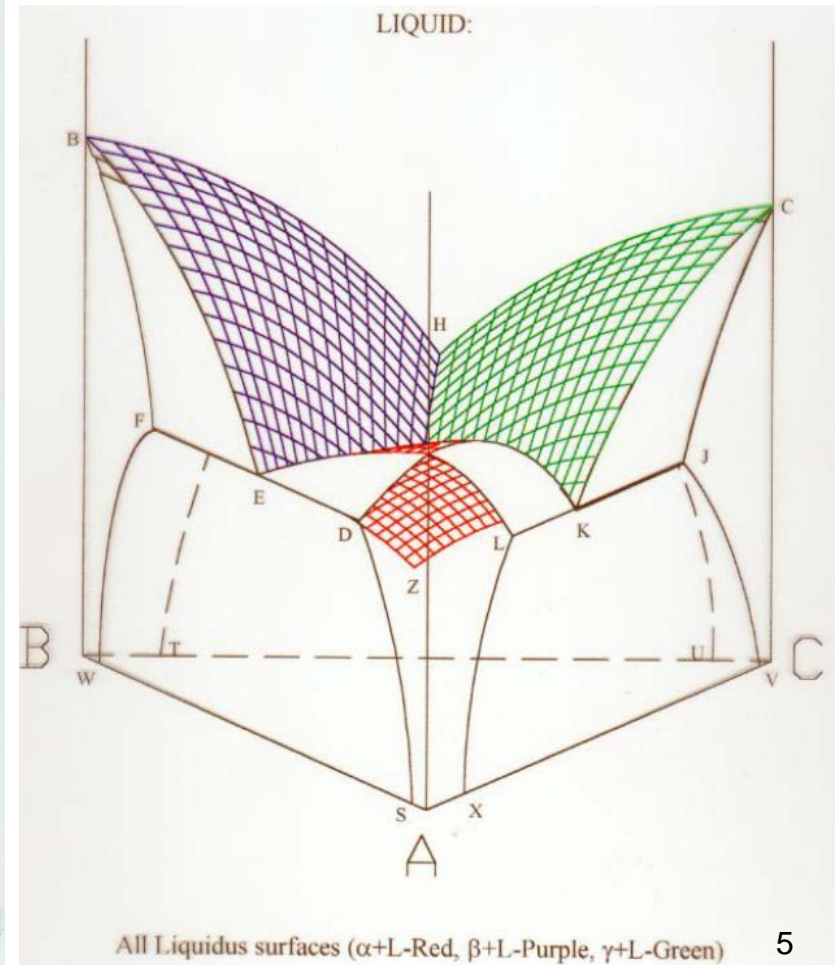
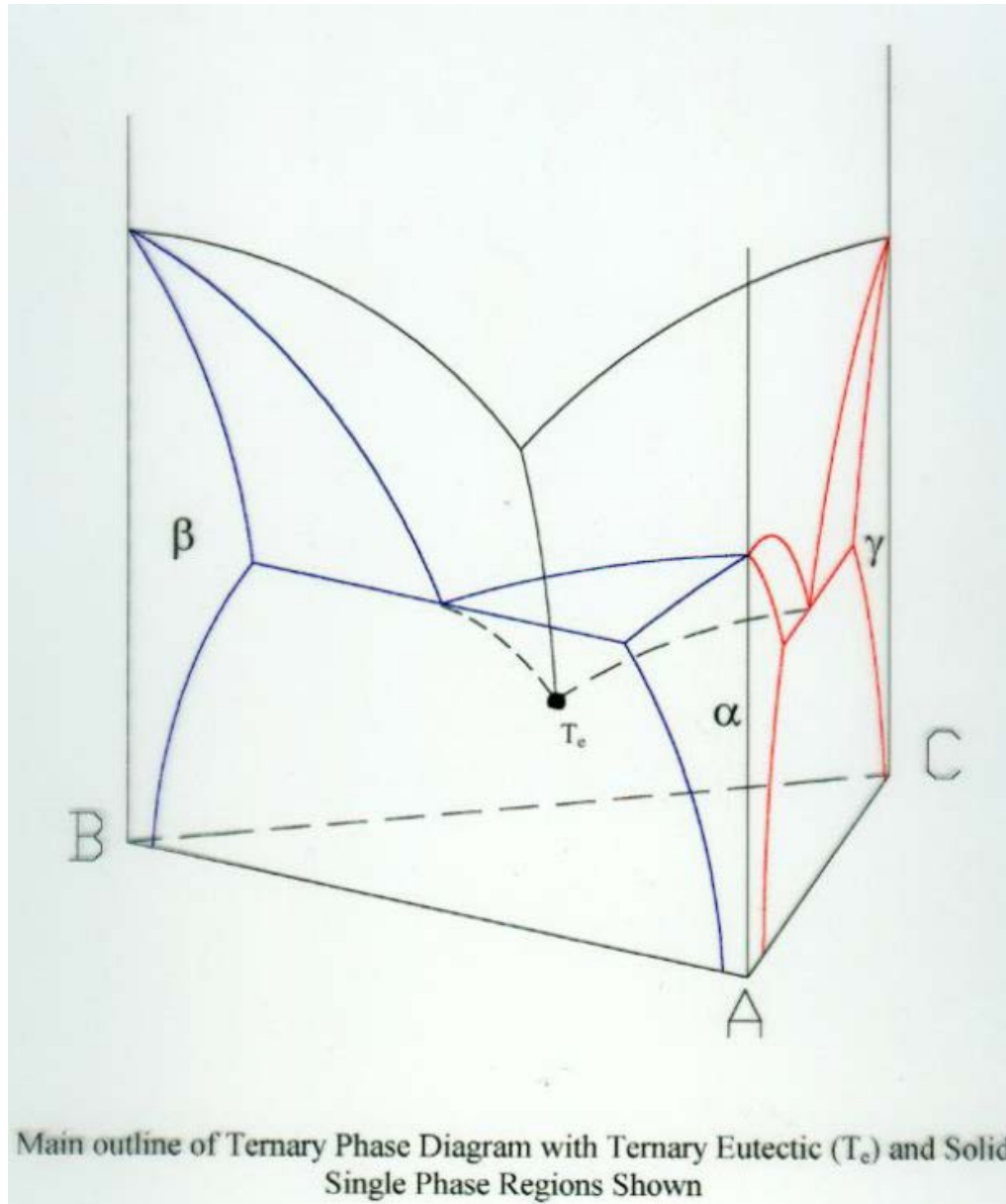
Syntectic reaction



K-Zn, Na-Zn,
K-Pb, Pb-U, Ca-Cd

On cooling two phases going to one phase

Ternary Eutectic System (with Solid Solubility)



< Quaternary phase Diagrams >

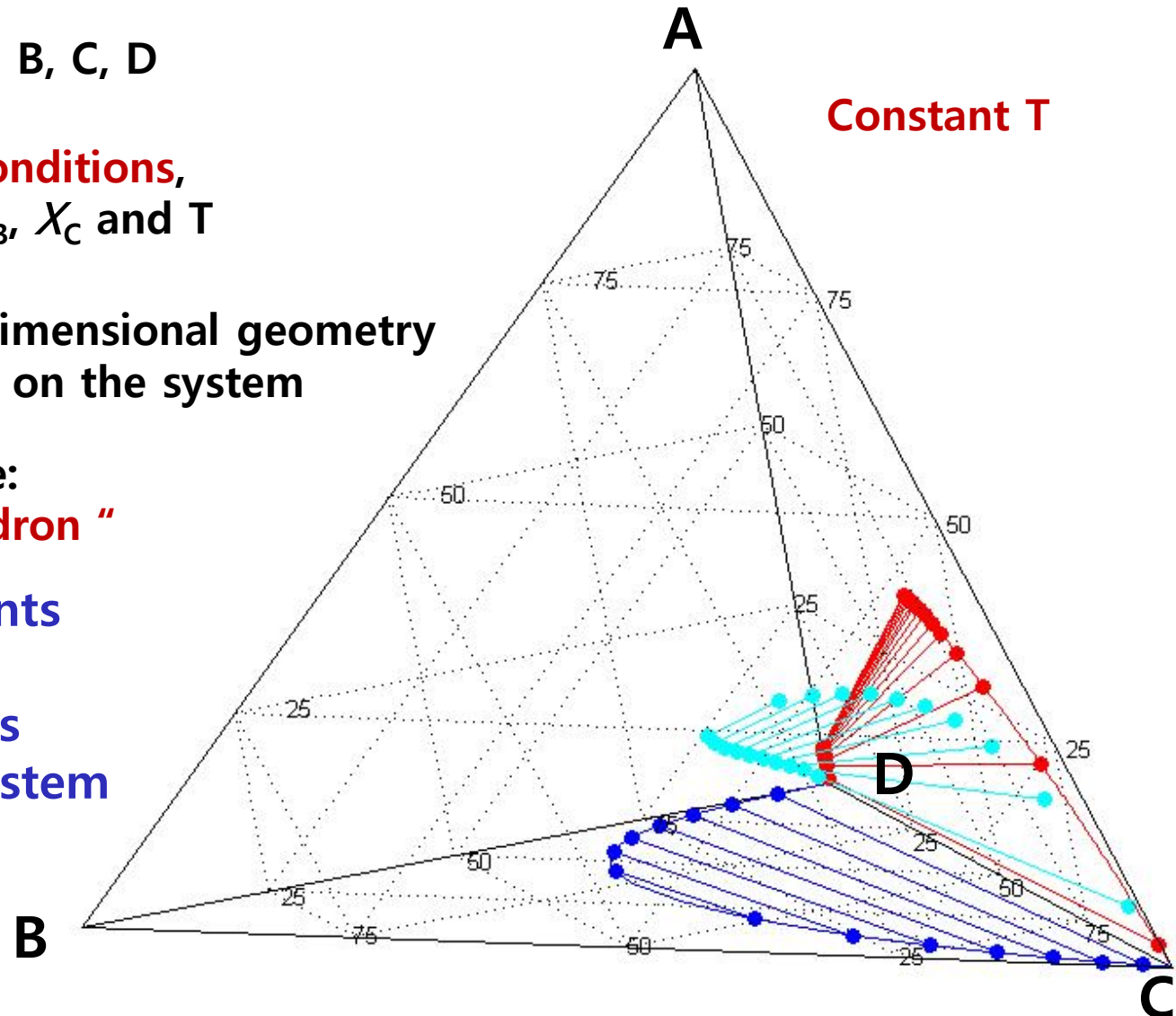
Four components: A, B, C, D

Assuming **isobaric conditions**,
Four variables: X_A , X_B , X_C and T

A difficulty of four-dimensional geometry
→ further restriction on the system

Most common figure:
" **equilateral tetrahedron** "

4 pure components
6 binary systems
4 ternary systems
A quaternary system

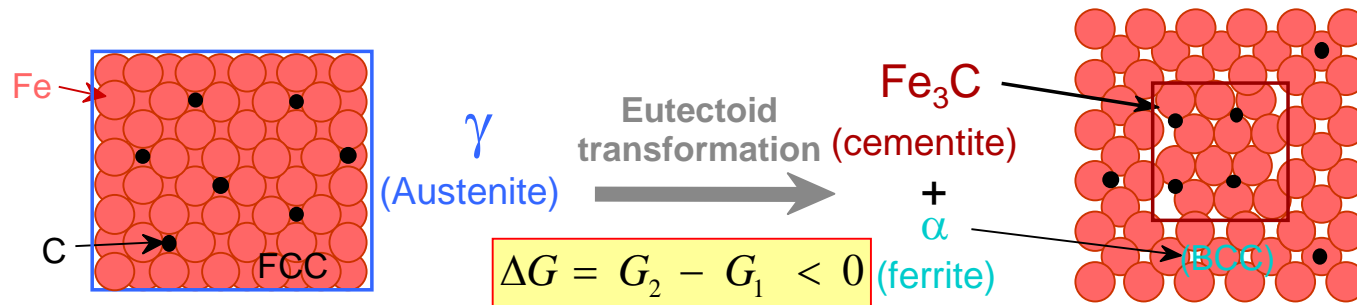


Contents for today's class

Chapter 12. Phase Transformations

ISSUES TO ADDRESS...

- Transforming one phase into another takes **time**.



- How does the rate of transformation depend on **time and T** ?
- How can we slow down the transformation so that we can **engineering non-equilibrium structures**?
- Are the **mechanical properties of non-equilibrium structures** better?

Contents_Phase transformation course

**Background
to understand
phase
transformation**

(Ch1) Thermodynamics and Phase Diagrams

(Ch2) Diffusion: Kinetics

(Ch3) Crystal Interface and Microstructure

**Representative
Phase
transformation**

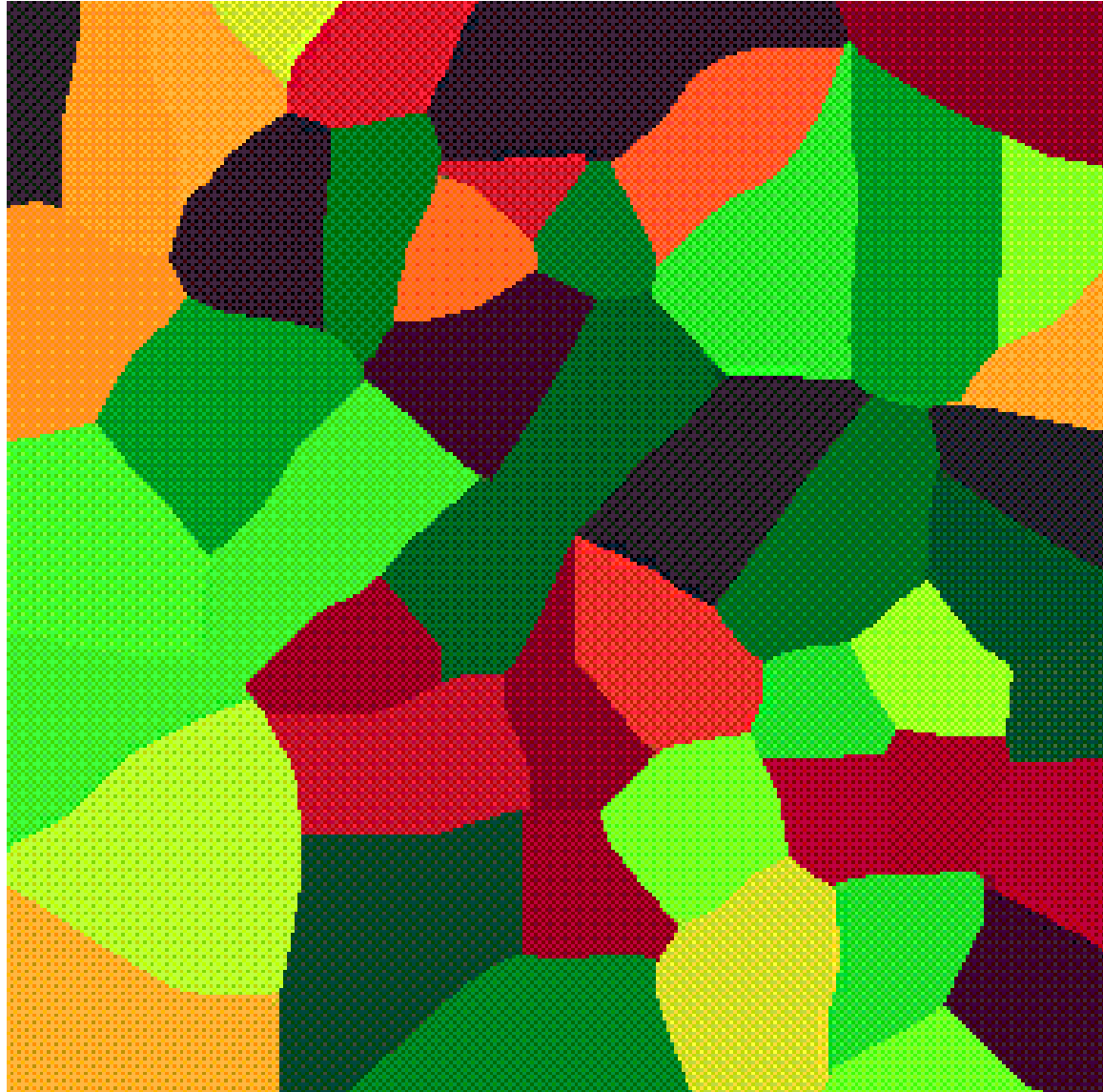
(Ch4) Solidification: Liquid \rightarrow Solid

(Ch5) Diffusional Transformations in Solid: Solid \rightarrow Solid

(Ch6) Diffusionless Transformations: Solid \rightarrow Solid

Phase Transformations

Solidification: Liquid \longrightarrow Solid



4 Fold Anisotropic Surface Energy/2 Fold Kinetics, Many Seeds

Solidification: **Liquid** \longrightarrow **Solid**

- casting & welding
- single crystal growth
- directional solidification
- rapid solidification

4.1. Nucleation in Pure Metals

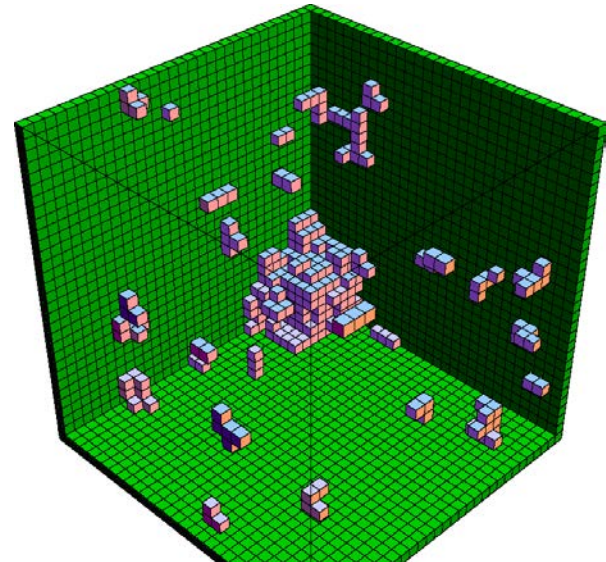
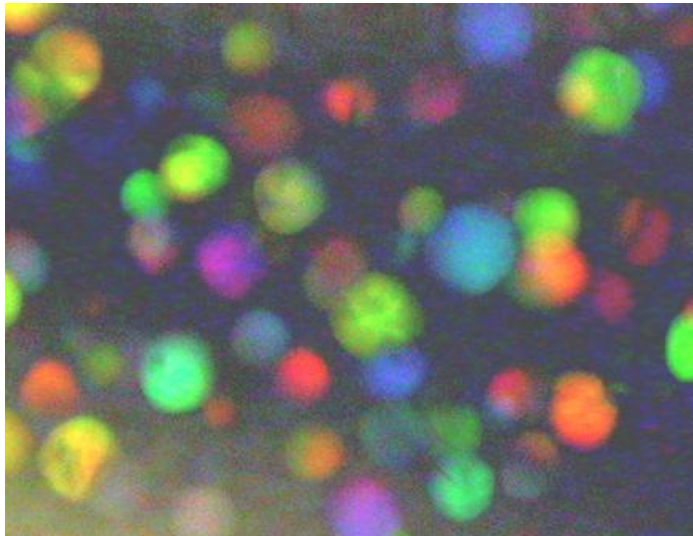
$$T_m : G_L = G_S$$

- Undercooling (supercooling) for nucleation: 250 K ~ 1 K

<Types of nucleation>

- **Homogeneous nucleation**

- **Heterogeneous nucleation**



4.1.1. Homogeneous Nucleation

Driving force for solidification

$$G^L = H^L - TS^L$$

$$G^S = H^S - TS^S$$

$$\Delta G = \Delta H - T \Delta S$$

$$L : \Delta H = H^L - H^S$$

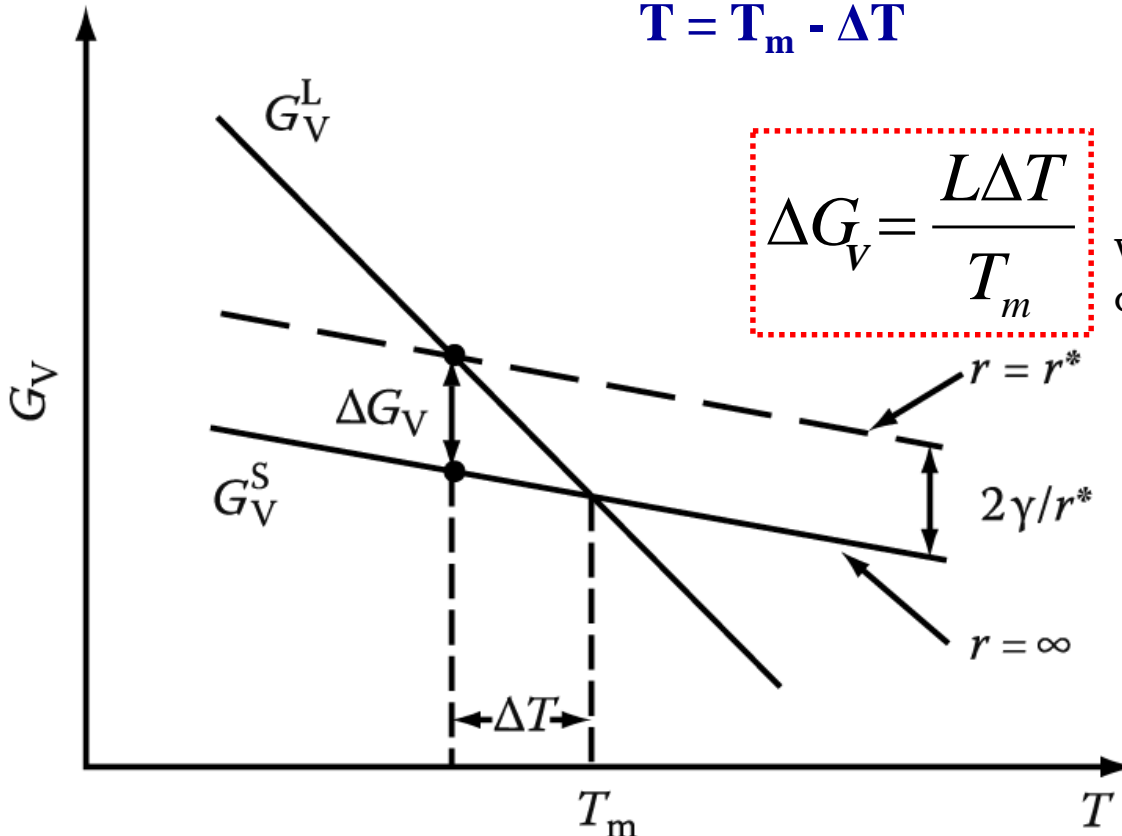
(Latent heat)

$$T = T_m - \Delta T$$

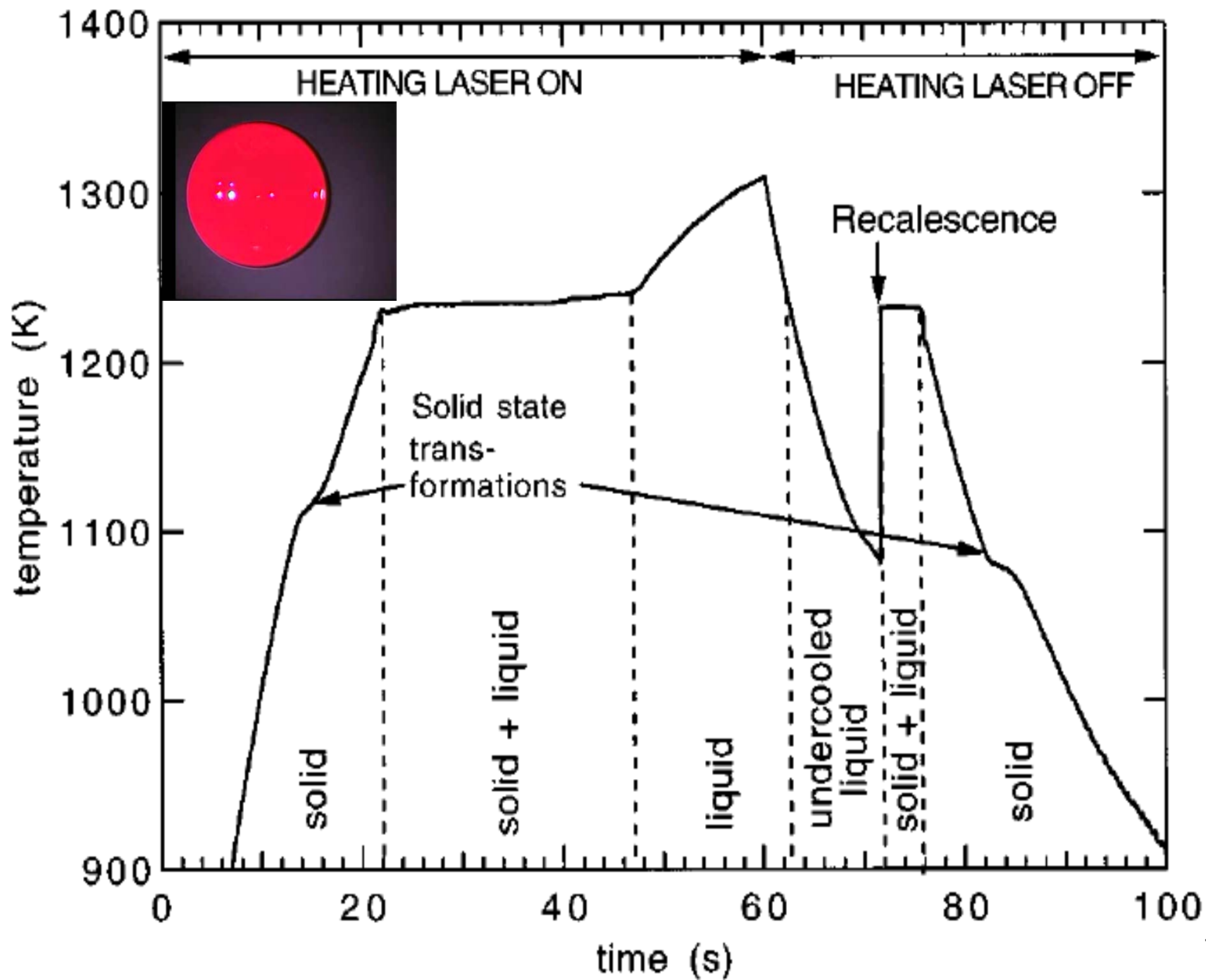
$$\Delta G = 0 = \Delta H - T_m \Delta S$$

$$\Delta S = \Delta H / T_m = L / T_m$$

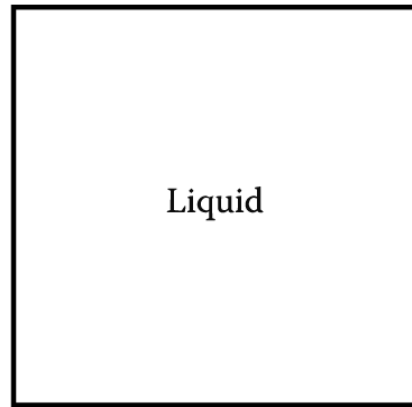
$$\Delta G = L - T(L/T_m) \approx (L\Delta T) / T_m$$



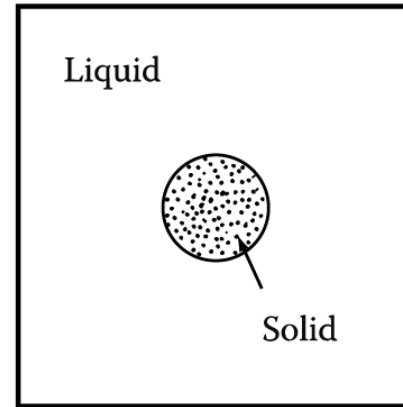
Variation of free energy per unit volume obtained from undercooling (ΔT)



4.1.1. Homogeneous Nucleation



(a) G_1



(b) $G_2 = G_1 + \Delta G$

$$G_1 = (V_S + V_L)G_V^L$$

$$G_2 = V_S G_V^S + V_L G_V^L + A_{SL} \gamma_{SL}$$

G_V^S, G_V^L : free energies per unit volume

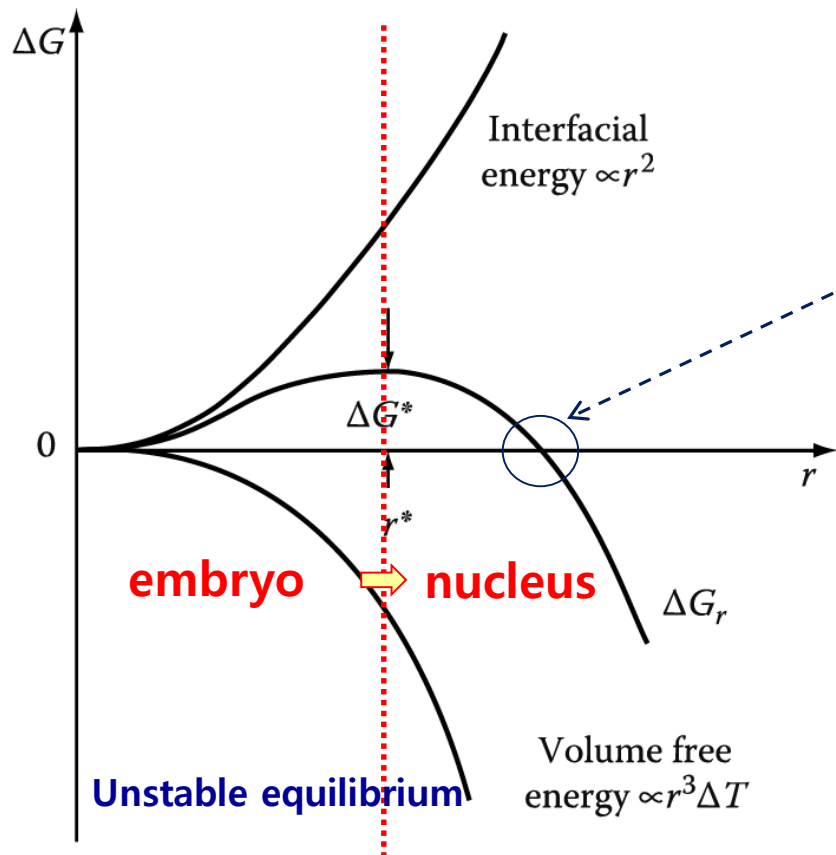
$$\Delta G = G_2 - G_1 = -V_S (G_V^L - G_V^S) + A_{SL} \gamma_{SL}$$

for spherical nuclei (isotropic) of radius : r

$$\Delta G_r = -\frac{4}{3} \pi r^3 \Delta G_V + 4\pi r^2 \gamma_{SL}$$

Volume (Bulk) Free Energy – stabilizes the nuclei (releases energy) Surface Free Energy- destabilizes the nuclei (it takes energy to make an interface)

Calculation of critical radius, r^*



$$\Delta G_r = -\frac{4}{3} \pi r^3 \Delta G_V + 4\pi r^2 \gamma_{SL}$$

Why r^* is not defined by $\Delta G_r = 0$?

$r < r^*$: **unstable** (lower free E by reduce size)

$r > r^*$: **stable** (lower free E by increase size)

r^* : critical nucleus size

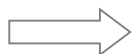
$$r^* \implies dG=0$$

nuclei $< r^*$ shrink; nuclei $> r^*$ grow (to reduce energy)

Gibbs-Thompson Equation

Critical ΔG of nucleation at r^*

$$r^* = \frac{2\gamma_{SL}}{\Delta G_V}$$



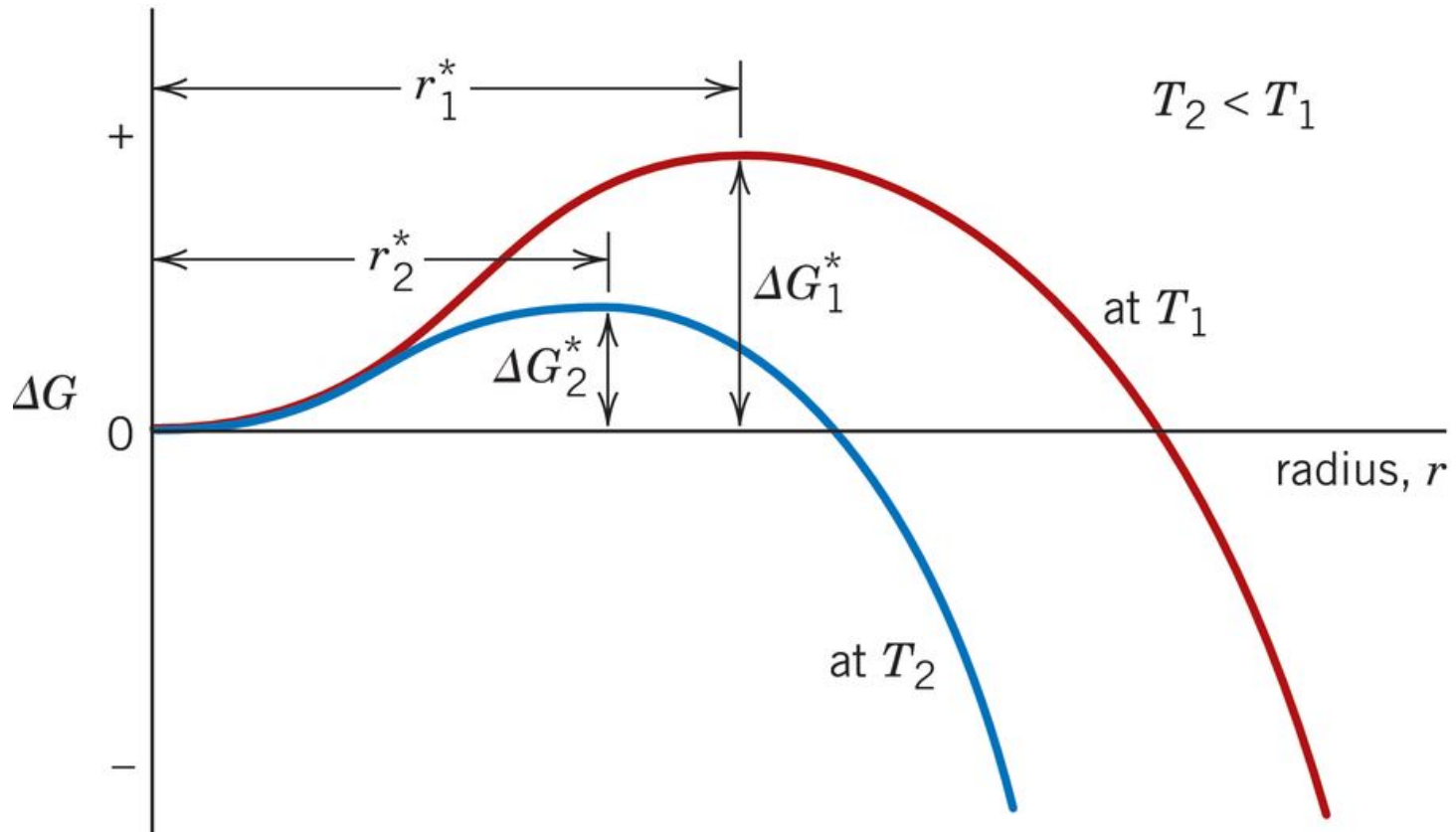
$$\Delta G^* = \frac{16\pi\gamma_{SL}^3}{3(\Delta G_V)^2} = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3L_V^2} \right) \frac{1}{(\Delta T)^2}$$

$$\Delta G_r = -\frac{4}{3} \pi r^3 \Delta G_V + 4\pi r^2 \gamma_{SL}$$

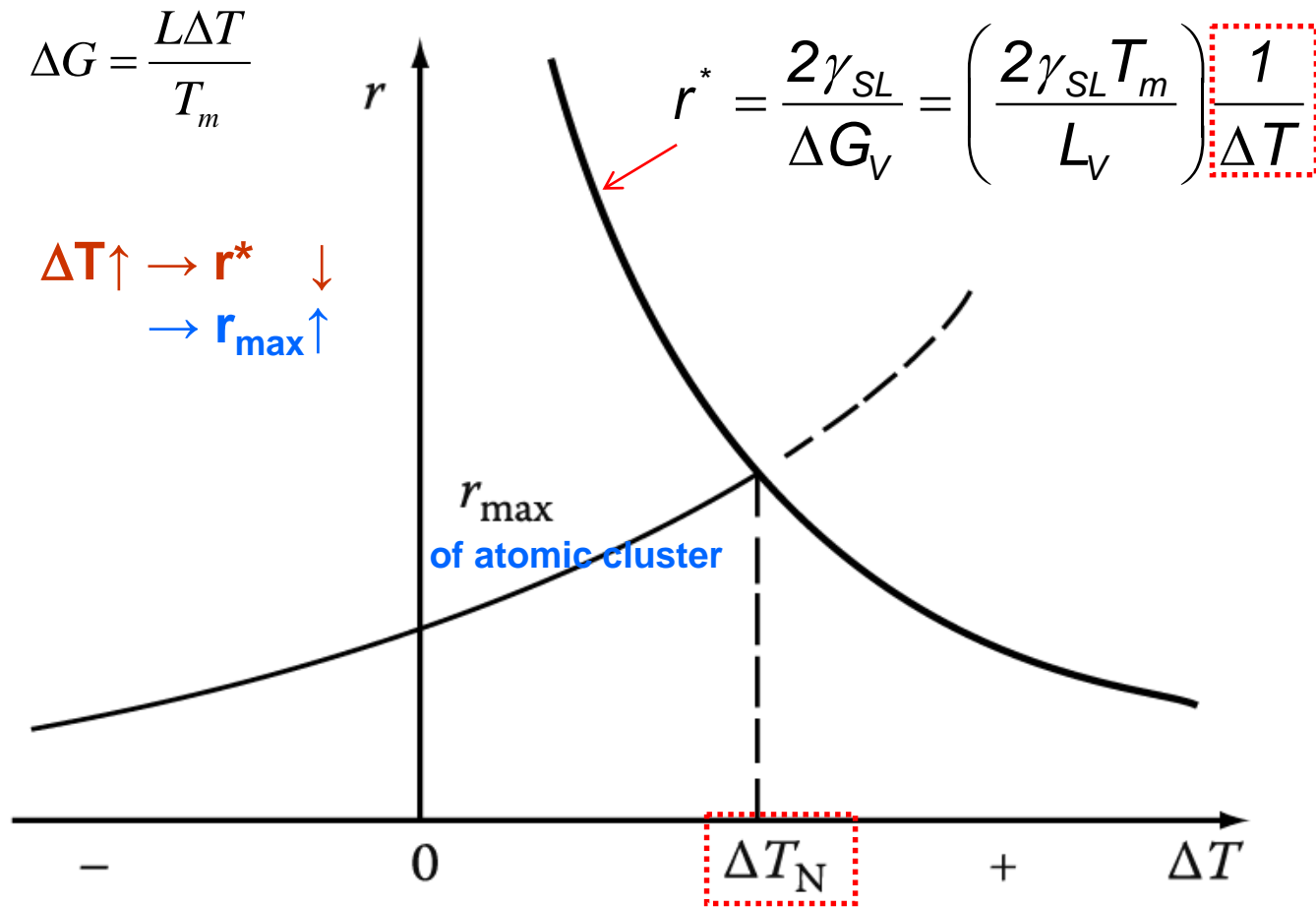
$$\Delta G_V = \frac{L\Delta T}{T_m}$$

$$r^* = \frac{2\gamma_{SL}}{\Delta G_V} = \left(\frac{2\gamma_{SL}T_m}{L_V} \right) \frac{1}{\Delta T}$$

$$\Delta G^* = \frac{16\pi\gamma_{SL}^3}{3(\Delta G_V)^2} = \left(\frac{16\pi\gamma_{SL}^3T_m^2}{3L_V^2} \right) \frac{1}{(\Delta T)^2}$$



The creation of a critical nucleus ~ thermally activated process



ΔT_N is **the critical undercooling** for homogeneous nucleation.

Fig. 4.5 The variation of r^* and r_{\max} with undercooling ΔT

The number of clusters with r^* at $\Delta T < \Delta T_N$ is negligible.

4.1.2. The homogeneous nucleation rate - kinetics

How fast solid nuclei will appear in the liquid at a given undercooling?

C_0 : atoms/unit volume

C^* : # of clusters with size of C^* (critical size) 임계핵 크기의 cluster 수

$$C^* = C_0 \exp\left(-\frac{\Delta G_{\text{hom}}^*}{kT}\right) \quad \text{clusters / m}^3$$

The addition of one more atom to each of these clusters will convert them into stable nuclei. 한 개 원자 추가로 확산시 핵생성

Homogeneous
Nucleation rate

$$N_{\text{hom}} = f_0 C_0 \exp\left(-\frac{\Delta G_{\text{hom}}^*}{kT}\right) \quad \text{nuclei / m}^3 \cdot \text{s}$$

$f_0 \sim 10^{11} \text{ s}^{-1}$: frequency \propto vibration frequency energy of diffusion in liquid surface area (const.)

$$\Delta G^* = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3L_v^2} \right) \frac{1}{(\Delta T)^2}$$

$C_0 \sim$ typically 10^{29} atoms/m³

$$N_{\text{hom}} \approx 1 \text{ cm}^{-3} \text{ s}^{-1} \quad \text{when } \Delta G^* \sim 78 kT$$

Reasonable nucleation rate

Homogeneous Nucleation in Solids

Concentration of Critical Size Nuclei per unit volume

$$C^* = C_0 \exp(-\Delta G^* / kT)$$

C_0 : number of atoms
per unit volume in the parent phase

Homogeneous Nucleation Rate

If each nucleus can be made supercritical at a rate of f per second,

$$N_{\text{hom}} = f C^*$$

$$f = \omega \exp(-\Delta G_m / kT)$$

f depends on how frequently a critical nucleus can receive an atom from the α matrix.

$\omega \propto$ vibration frequency, area of critical nucleus

ΔG_m : activation energy for atomic migration

$$N_{\text{hom}} = \omega C_0 \exp\left(-\frac{\Delta G_m}{kT}\right) \exp\left(-\frac{\Delta G^*}{kT}\right)$$

: This eq. is basically same with eq (4.12) except considering temp. dependence of f .

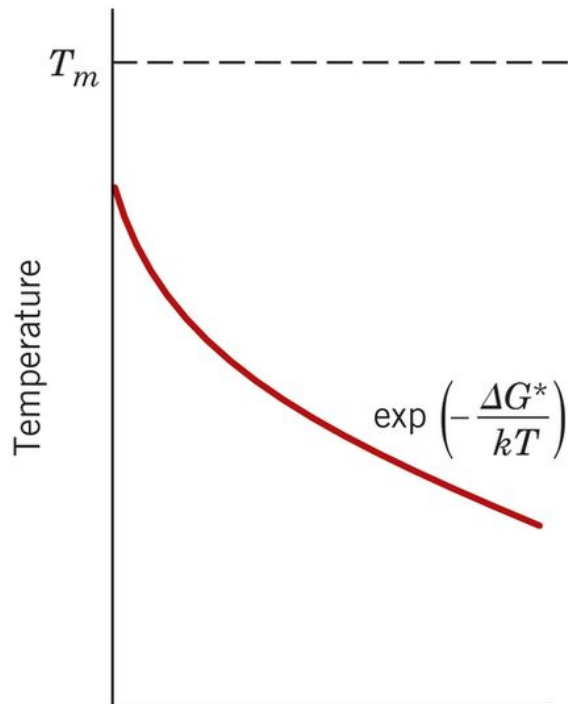
Homogeneous
Nucleation rate

$$N_{\text{hom}} = f_0 C_0 \exp\left(-\frac{\Delta G_{\text{hom}}^*}{kT}\right)$$

nuclei / m³·s

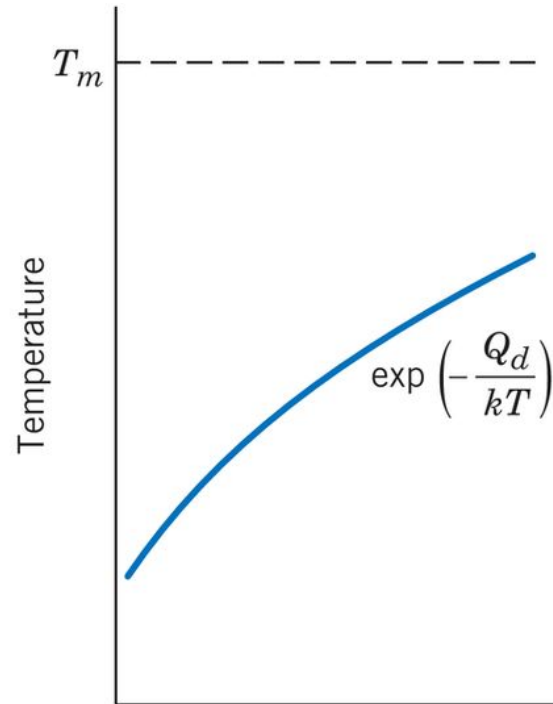
4.1.2. The homogeneous nucleation rate - kinetics

$$N_{\text{hom}} = \omega C_0 \exp\left(-\frac{\Delta G_m}{kT}\right) \exp\left(-\frac{\Delta G^*}{kT}\right) = Q_d$$



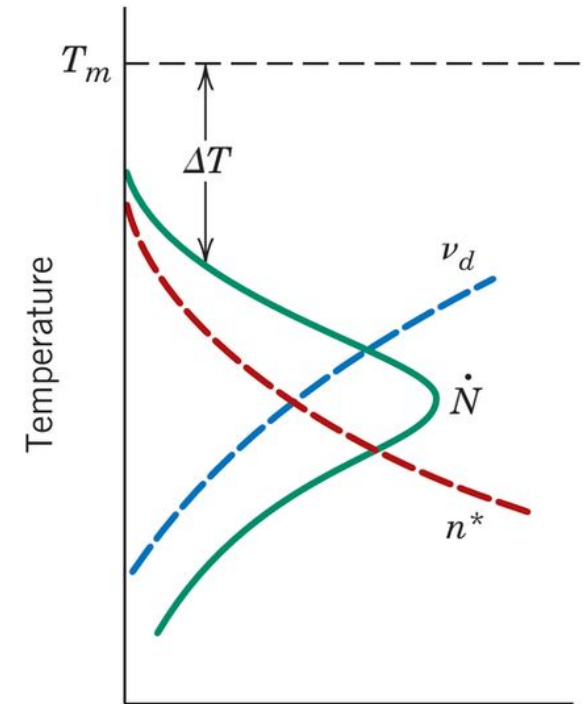
Number of stable nuclei, n^*

(a)



Frequency of attachment, ν_d

(b)



n^*, ν_d, \dot{N}

(c)

4.1.2. The homogeneous nucleation rate - kinetics

$$N_{\text{hom}} \approx f_0 C_o \exp\left\{-\frac{A}{(\Delta T)^2}\right\}$$

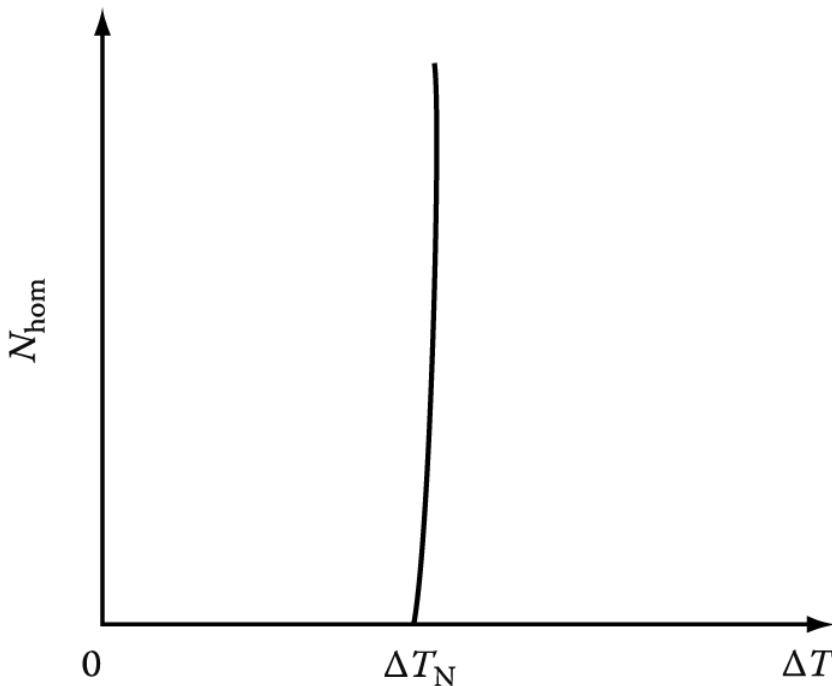
where $A = \frac{16\pi\gamma_{SL}^3 T_m^2}{3L_v^2 kT}$

A = relatively insensitive to Temp.

How do we define ΔT_N ?

$$N_{\text{hom}} \sim \frac{1}{\Delta T^2}$$

Changes by orders of magnitude from essentially zero to very high values over a very narrow temperature range



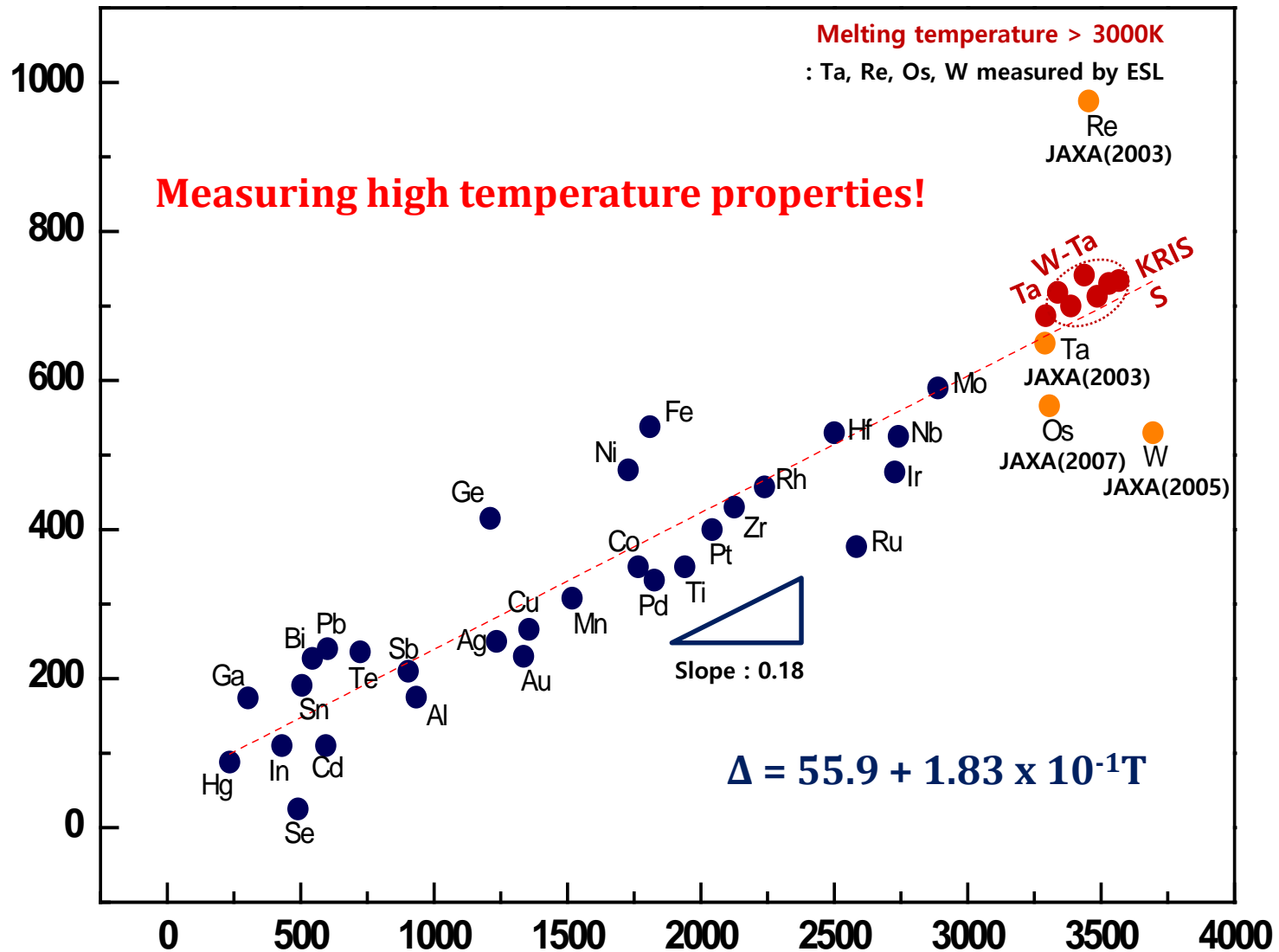
→ **critical value for detectable nucleation**

- critical supersaturation ratio
- critical driving force
- critical supercooling

→ **for most metals, $\Delta T_N \sim 0.2 T_m$ (i.e. $\sim 200\text{K}$)**

Fig. 4.6 The homogeneous nucleation rate as a function of undercooling ΔT . ΔT_N is **the critical undercooling** for homogeneous nucleation.

Maximum undercooling vs. Melting temperature



* **Copper** Homogeneous nucleation

$$\Delta T = 230 \text{ K} \rightarrow r^* \sim 10^{-7} \text{ cm} < 4 * (\text{Diameter of Cu atom})$$

If nucleus is spherical shape,

$$V = 4.2 * 10^{-21} \text{ cm}^3 \sim 360 \text{ atoms} (\because \text{one Cu atom } 1.16 * 10^{-23} \text{ cm}^3)$$

“Typically in case of metal” $\Delta T^* \sim 0.2 T_E / \sigma_{SL} \sim 0.4 L$

→ r^* (critical nucleus for homogeneous nucleation) of metal ~ 200 atoms

But, if cluster radius \sim (only 4 * atom diameter),

“no spherical shape”

(large deviation from spherical shape) →

→ **Possible structure for the critical nucleus of Cu**
: bounded only by {111} and {100} plane

- σ_{SL} may vary with the crystallographic nature of the surface.
- The faces of this crystal are close to their critical size for 2D nucleation at the critical temp for the nucleus as a whole.

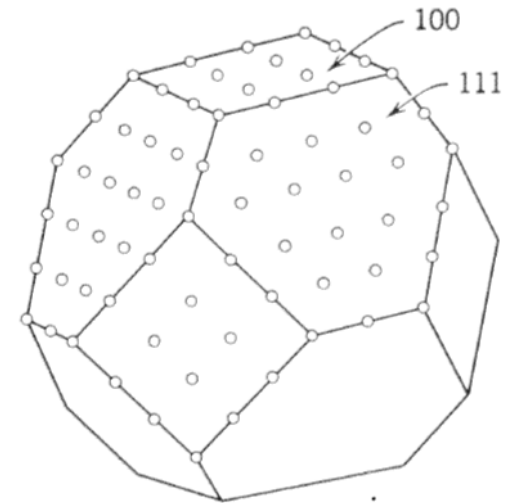


Fig. 3.11. Possible structure for the critical nucleus. (From B. Chalmers, *Physical Metallurgy*, John Wiley and Sons, New York, 1959, p. 246.)

Real behavior of nucleation: metal $\Delta T_{\text{bulk}} < \Delta T_{\text{small drop}}$

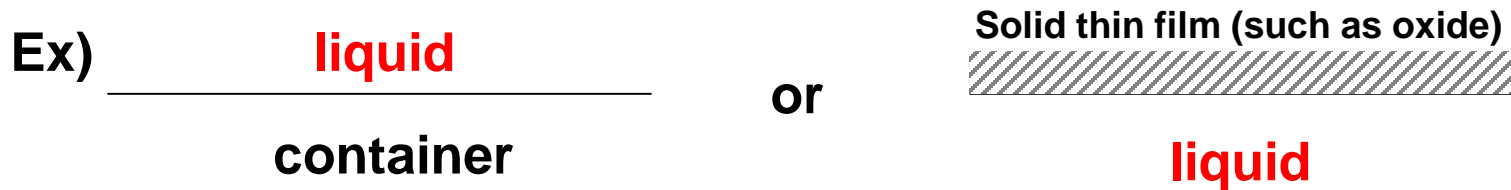
Under suitable conditions, liquid nickel can be undercooled (or supercooled) to 250 K below T_m (1453°C) and held there indefinitely without any transformation occurring.



Normally undercooling as large as 250 K are not observed.

The nucleation of solid at undercooling of only ~ 1 K is common.

The formation of a nucleus of critical size can be catalyzed by a suitable surface in contact with the liquid. → “Heterogeneous Nucleation”



Why this happens? What is the underlying physics?

Which equation should we examine?

$$\Delta G^* = \frac{16\pi\gamma_{SL}^3}{3(\Delta G_V)^2} = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3 L_V^2} \right) \frac{1}{(\Delta T)^2}$$

$$N_{\text{hom}} = f_0 C_o \exp\left(-\frac{\Delta G_{\text{hom}}^*}{kT}\right)$$

4.1.3. Heterogeneous nucleation

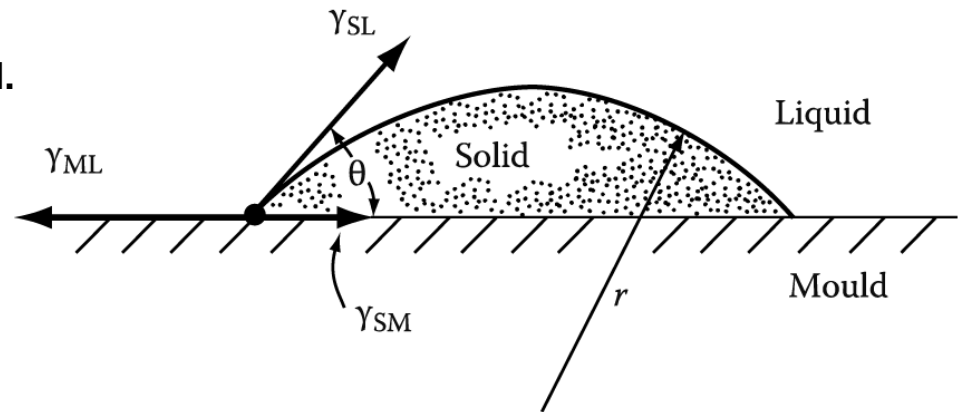
From
$$\Delta G^* = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3L_v^2} \right) \frac{1}{(\Delta T)^2}$$

Nucleation becomes easy if $\gamma_{SL} \downarrow$ by forming nucleus from mould wall.

Fig. 4.7 Heterogeneous nucleation of spherical cap on a flat mould wall.

$$\gamma_{ML} = \gamma_{SL} \cos \theta + \gamma_{SM}$$

$$\cos \theta = (\gamma_{ML} - \gamma_{SM}) / \gamma_{SL}$$



$$\Delta G_{het} = -V_S \Delta G_V + A_{SL} \gamma_{SL} + A_{SM} \gamma_{SM} - A_{SM} \gamma_{ML}$$

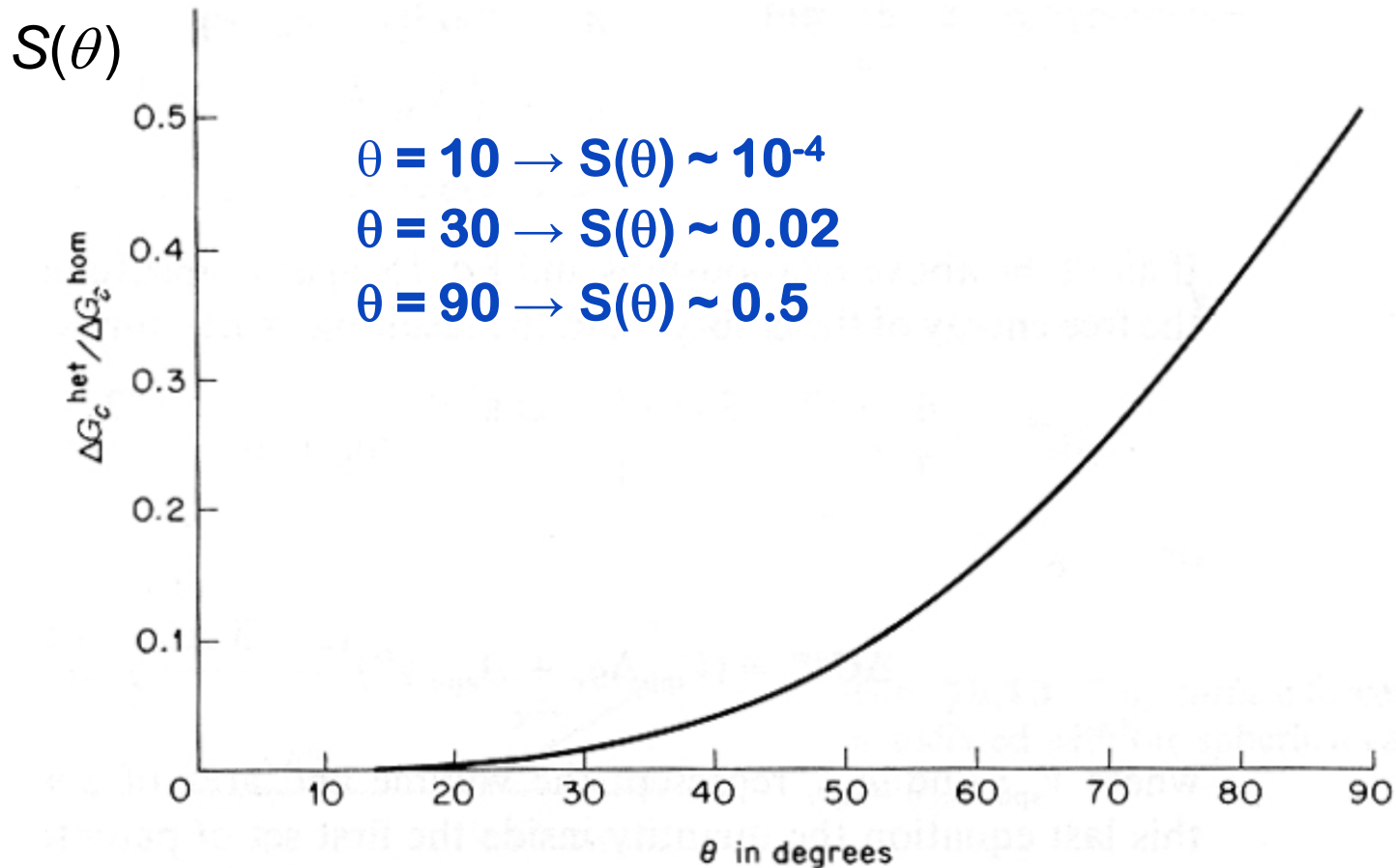
In terms of the wetting angle (θ) and the cap radius (r) (Exercies 4.6)

$$\Delta G_{het} = \left\{ -\frac{4}{3} \pi r^3 \Delta G_V + 4\pi r^2 \gamma_{SL} \right\} S(\theta)$$

where $S(\theta) = (2 + \cos \theta)(1 - \cos \theta)^2 / 4$

$S(\theta)$ has a numerical value ≤ 1 dependent only on θ (the shape of the nucleus)

$$\Delta G_{het}^* = S(\theta) \Delta G_{hom}^* \quad \Rightarrow \quad r^* = \frac{2 \gamma_{SL}}{\Delta G_V} \quad \text{and} \quad \Delta G^* = \frac{16 \pi \gamma_{SL}^3}{3 \Delta G_V^2} \cdot S(\theta)$$



$S(\theta)$ has a numerical value ≤ 1 dependent only on θ (the shape of the nucleus)

$$\Delta G_{het}^* = S(\theta) \Delta G_{hom}^*$$



$$r^* = \frac{2 \gamma_{SL}}{\Delta G_V} \quad \text{and} \quad \Delta G^* = \frac{16\pi\gamma_{SL}^3}{3\Delta G_V^2} \cdot S(\theta)$$

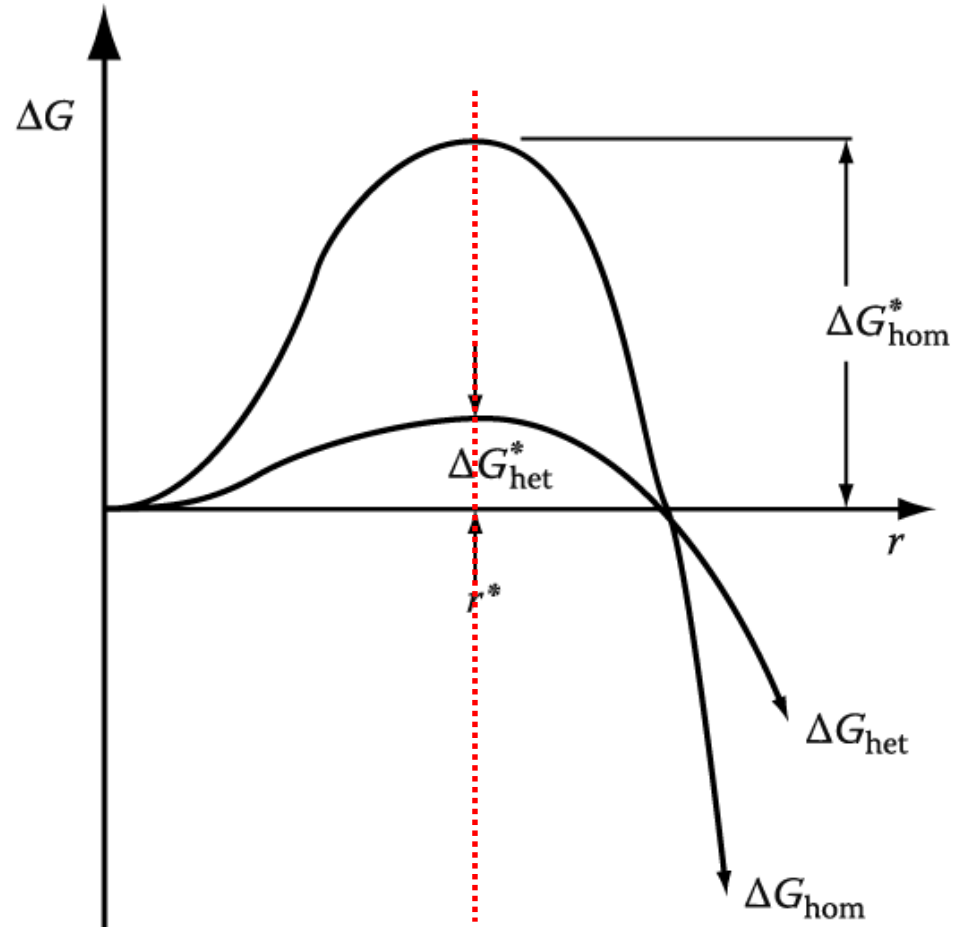
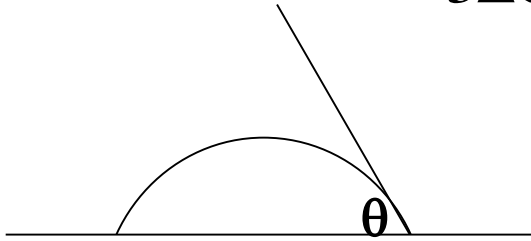


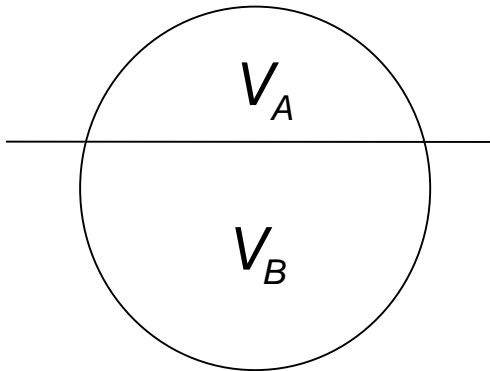
Fig. 4.8 The excess free energy of solid clusters for homogeneous and heterogeneous nucleation. Note r^* is independent of the nucleation site.

Barrier of Heterogeneous Nucleation

$$\Delta G^* = \frac{16\pi\gamma_{SL}^3}{3\Delta G_V^2} \cdot S(\theta) = \frac{16\pi\gamma_{SL}^3}{3\Delta G_V^2} \cdot \frac{(2 - 3\cos\theta + \cos^3\theta)}{4}$$



$$\Delta G_{het}^* = S(\theta)\Delta G_{hom}^*$$



$$\Delta G_{sub}^* = \Delta G_{homo}^* \left(\frac{2 - 3\cos\theta + \cos^3\theta}{4} \right)$$

$$\frac{V_A}{V_A + V_B} = \frac{2 - 3\cos\theta + \cos^3\theta}{4} = S(\theta)$$

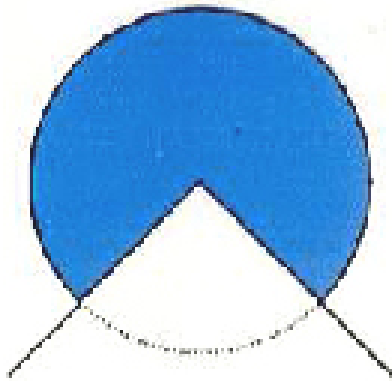
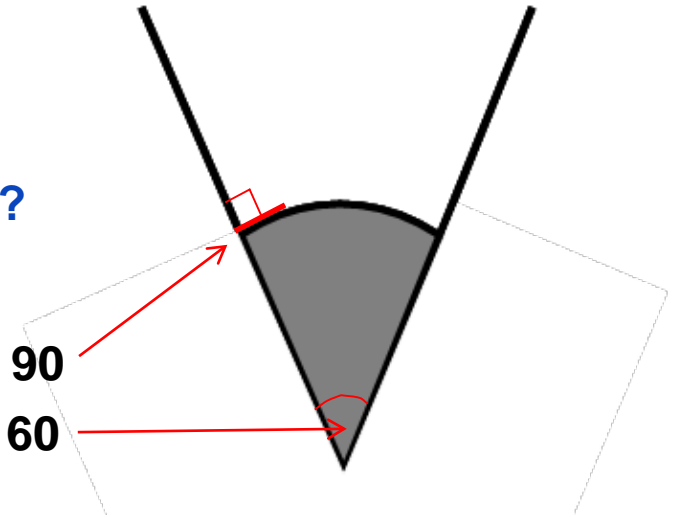
How about the nucleation at the crevice or at the edge?

Nucleation Barrier at the crevice

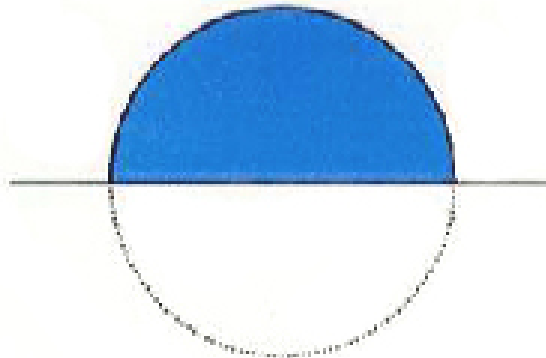
What would be the shape of nucleus and the nucleation barrier for the following conditions?

$$\frac{1}{6} \Delta G_{\text{homo}}^*$$

contact angle = 90
groove angle = 60



$$\frac{3}{4} \Delta G_{\text{homo}}^*$$

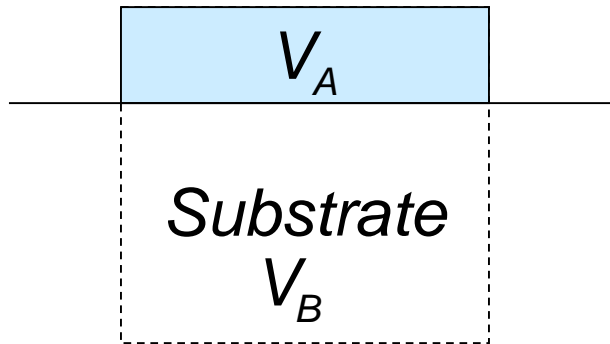


$$\frac{1}{2} \Delta G_{\text{homo}}^*$$

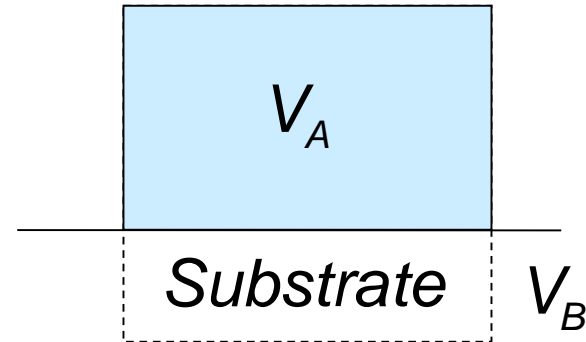


$$\frac{1}{4} \Delta G_{\text{homo}}^*$$

How do we treat the non-spherical shape?



Good Wetting



Bad Wetting

$$\Delta G_{sub}^* = \Delta G_{homo}^* \left(\frac{V_A}{V_A + V_B} \right)$$

Effect of good and bad wetting on substrate

The Effect of ΔT on ΔG^*_{het} & ΔG^*_{hom} ?

$$\Delta G^* = \frac{16\pi\gamma_{SL}^3}{3(\Delta G_V)^2} = \left(\frac{16\pi\gamma_{SL}^3 T_m^2}{3 L_v^2} \right) \frac{1}{(\Delta T)^2}$$

$$\Rightarrow \Delta G^* = \frac{16\pi\gamma_{SL}^3}{3\Delta G_V^2} \cdot S(\theta)$$

n_1 atoms in contact with the mold wall

$$n^* = n_1 \exp\left(-\frac{\Delta G^*_{het}}{kT}\right)$$

$$N_{het} = f_1 C_1 \exp\left(-\frac{\Delta G^*_{het}}{kT}\right)$$

Plot ΔG^*_{het} & ΔG^*_{hom} vs ΔT and N vs ΔT .

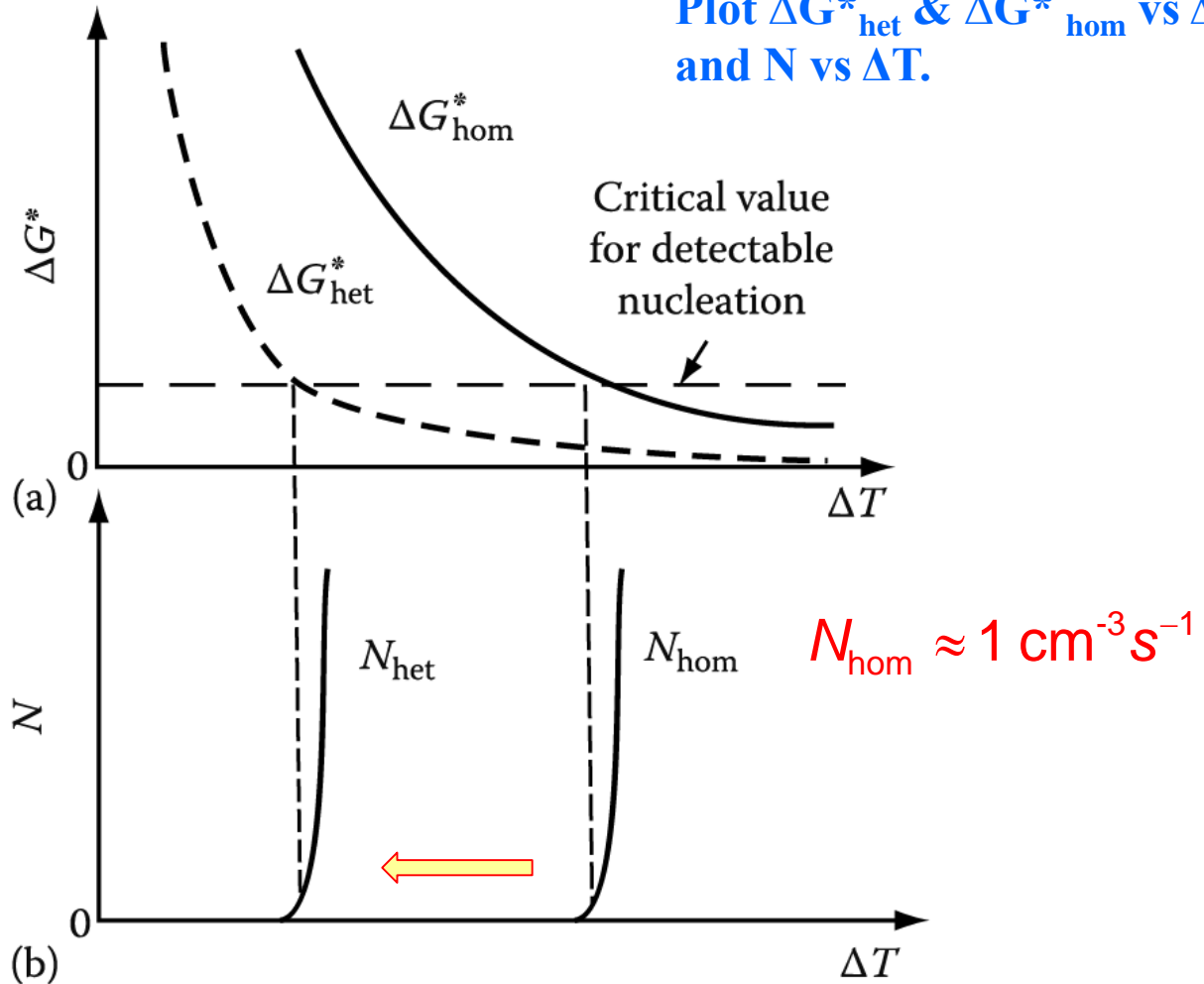
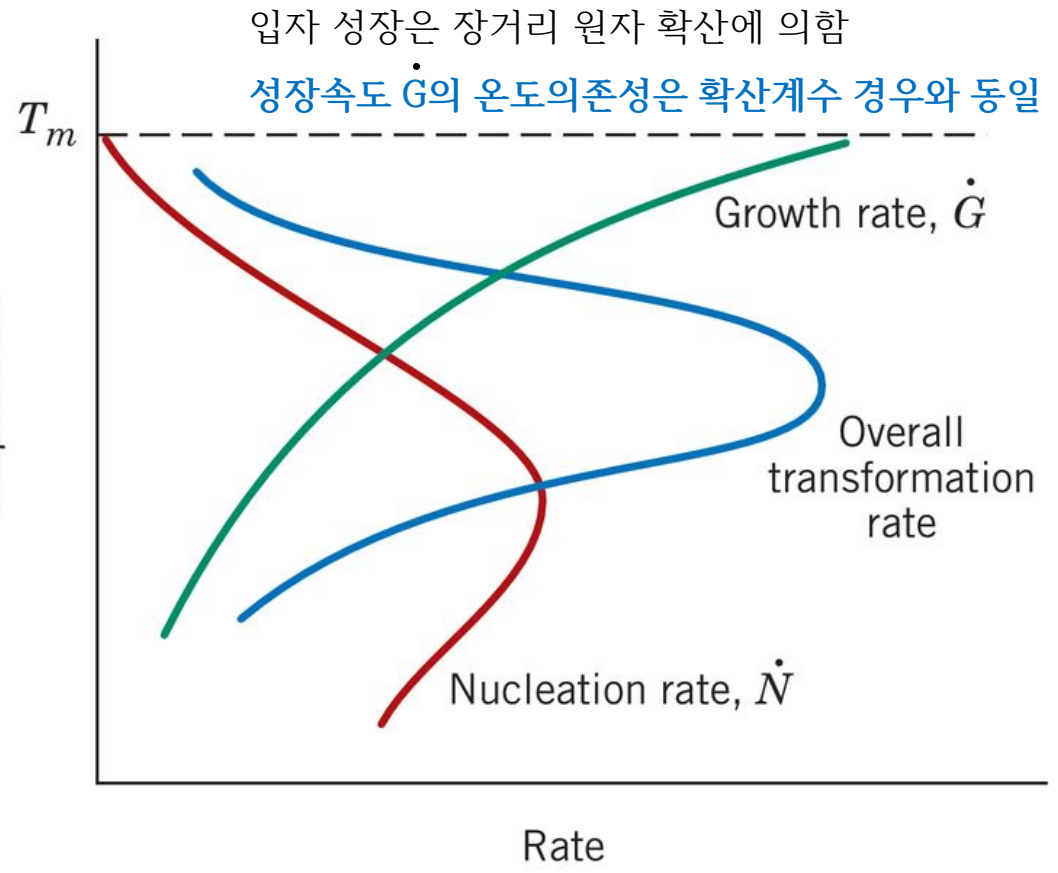
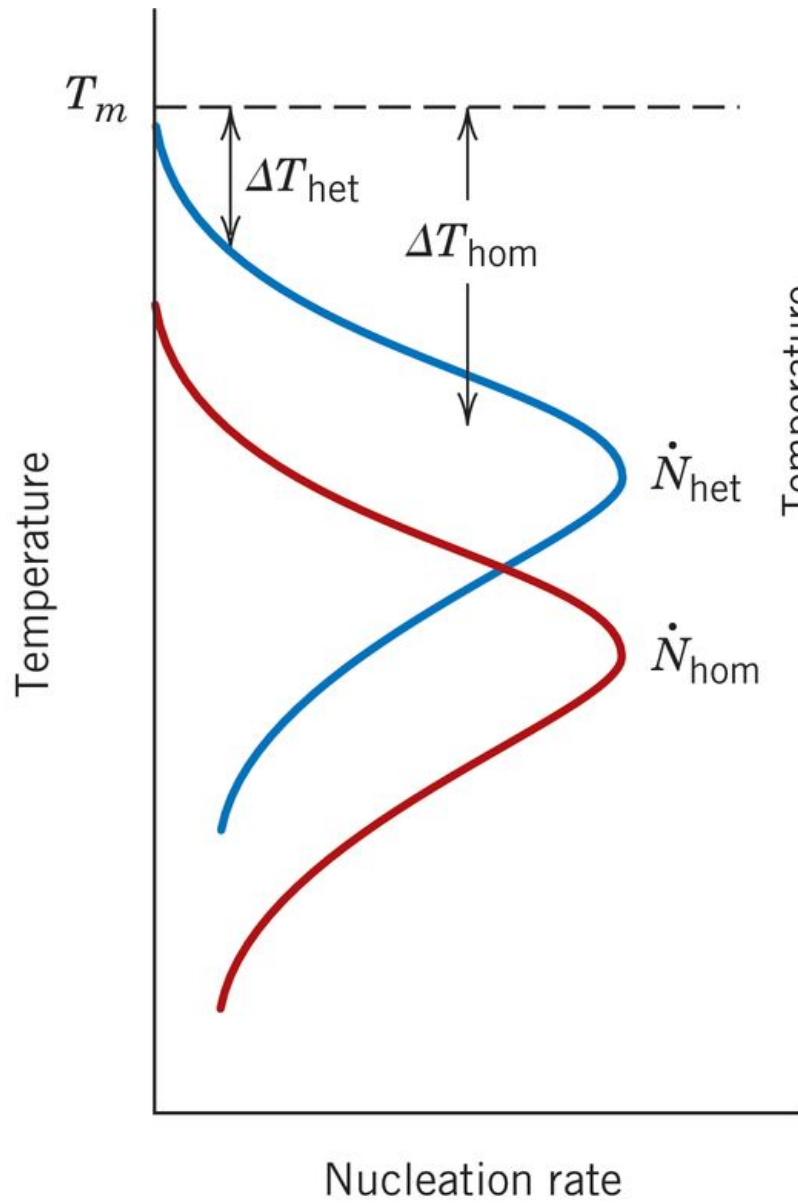


Fig. 4.9 (a) Variation of ΔG^* with undercooling (ΔT) for homogeneous and heterogeneous nucleation. (b) The corresponding nucleation rates assuming the same critical value of ΔG^*



3.7 The Nucleation of Melting

Although nucleation during solidification usually requires some undercooling, melting invariably occurs at the equilibrium melting temperature even at relatively high rates of heating.

Because, melting can apparently, start at crystal surfaces without appreciable superheating.

Why?

$$\gamma_{SL} + \gamma_{LV} < \gamma_{SV} \quad (\text{commonly})$$

In the case of gold,

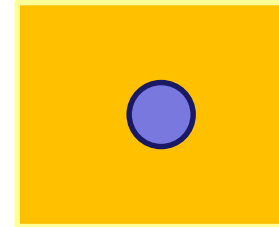
| | | |
|---------------|--------------|--------------|
| γ_{SL} | solid-liquid | 132 ergs/cm |
| γ_{LV} | liquid-vapor | 1128 ergs/cm |
| γ_{SV} | solid-vapor | 1400 ergs/cm |



In general, wetting angle = 0 \Rightarrow No superheating required!

Melting and Crystallization are Thermodynamic Transitions

Solidification: Liquid \longrightarrow Solid



<Thermodynamic>

- Interfacial energy $\Rightarrow \Delta T_N$

Liquid

T_m Undercooled Liquid

Solid

No superheating required!

- Interfacial energy \Rightarrow No ΔT_N

$$\gamma_{SL} + \gamma_{LV} < \gamma_{SV}$$

Melting: Liquid \longleftarrow Solid



5.4 Overall Transformation Kinetics – TTT Diagram

If isothermal transformation,

The fraction of Transformation as a function of Time and Temperature

$$\rightarrow f(t, T)$$

Plot f vs $\log t$.

- isothermal transformation
- $f \sim$ volume fraction of β at any time; $0 \sim 1$

Plot the fraction of transformation (1%, 99%) in T-log t coordinate.

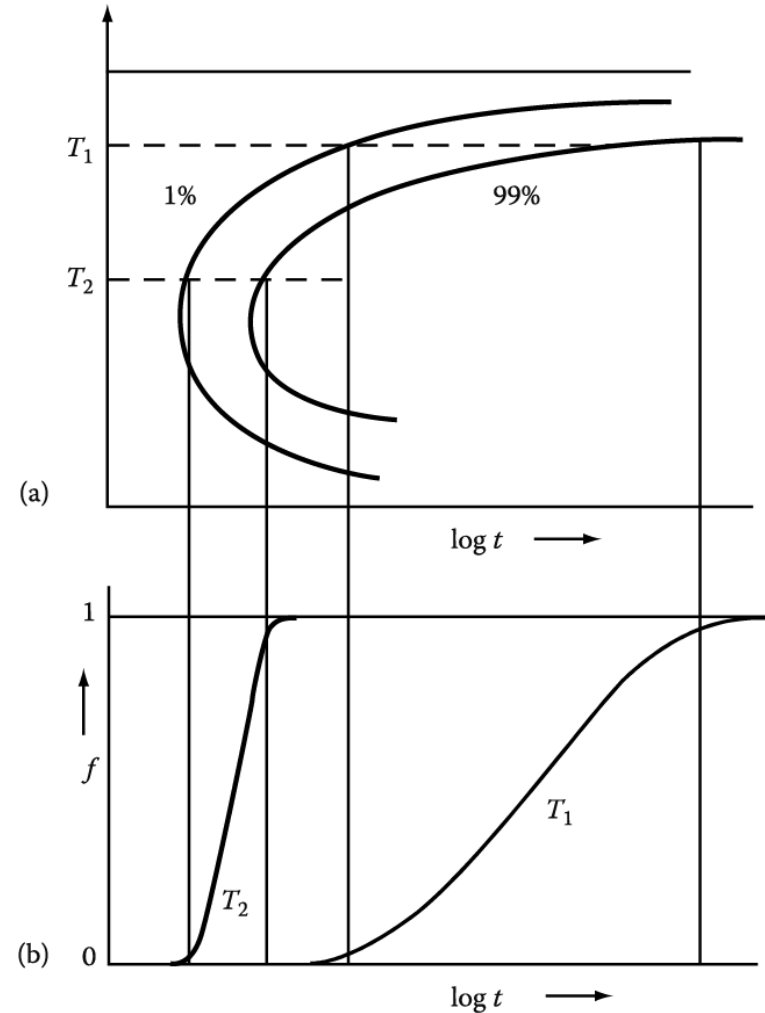


Fig. 5.23 The percentage transformation versus time for different transformation temperatures.

Transformation Kinetics

- Avrami proposed that for a three-dimensional nucleation and growth process kinetic law

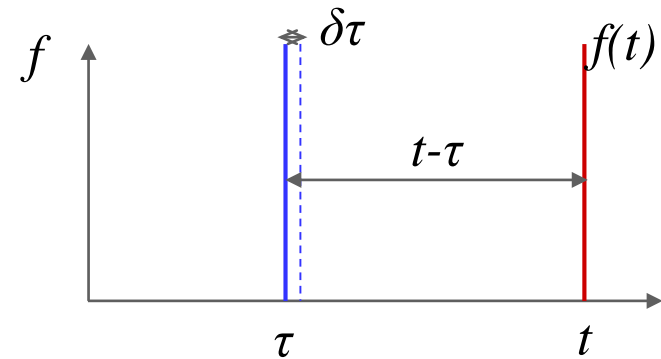
$$f = 1 - \exp(-kt^n) \quad \text{Johnson-Mehl-Avrami equation}$$

$$f : \text{volume fraction transformed} = \frac{\text{Volume of new phase}}{\text{Volume of specimen}}$$

- Assumption :
 - ✓ reaction produces by nucleation and growth
 - ✓ nucleation occurs randomly throughout specimen
 - ✓ reaction product grows rapidly until impingement

Constant Nucleation Rate Conditions

- Nucleation rate (I) is **constant**.
- Growth rate (v) is constant.
- No compositional change



$$df_e = \frac{\left(\begin{array}{l} \text{Vol. of one particle nucleated} \\ \text{during } d\tau \text{ measured at time } t \end{array} \right) \times \left(\begin{array}{l} \text{number of nuclei} \\ \text{formed during } d\tau \end{array} \right)}{\text{Volume of specimen}}$$

$$df_e = \frac{\frac{4}{3} \pi [v(t-\tau)]^3 \times (IV_0 d\tau)}{V_0}$$

$$f_e(t) = \int_0^t I \cdot \frac{4}{3} \pi [v(t-\tau)]^3 d\tau$$

$$= I \cdot \frac{4}{3} \pi v^3 \left[-\frac{1}{4} (t-\tau)^4 \right]_0^t = \frac{1}{3} \pi I v^3 t^4$$

$$V = \frac{4}{3} \pi r^3 = \frac{4}{3} \pi (vt)^3$$

$$V' = \frac{4}{3} \pi v^3 (t-\tau)^3$$

- do not consider impingement & repeated nucleation

- only true for $f \ll 1$

As time passes the β cells will eventually impinge on one another and the rate of transformation will decrease again.

Constant Nucleation Rate Conditions

- consider impingement + repeated nucleation effects

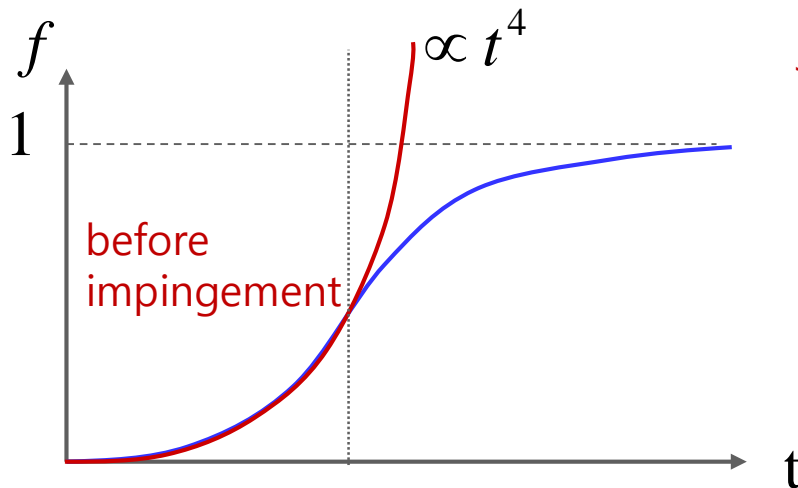
$$df = (1-f)df_e \longrightarrow df_e = \frac{df}{1-f}$$

$$f_e = -\ln(1-f)$$

$$f(t) = 1 - \exp(-f_e(t)) = 1 - \exp\left(-\frac{\pi}{3} I v^3 t^4\right)$$

* Short time:
 $1 - \exp(z) \sim z$ ($z \ll 1$)

* Long time:
 $t \rightarrow \infty, f \rightarrow 1$



Johnson-Mehl-Avrami Equation

$$f = 1 - \exp(-kt^n)$$

k : T sensitive $f(I, v)$ $-\frac{\pi}{3} I v^3$
 n : 1 ~ 4 (depend on nucleation mechanism)

Growth controlled.

Nucleation-controlled.

If no change of nucleation mechanism during phase transformation, n is not related to T .

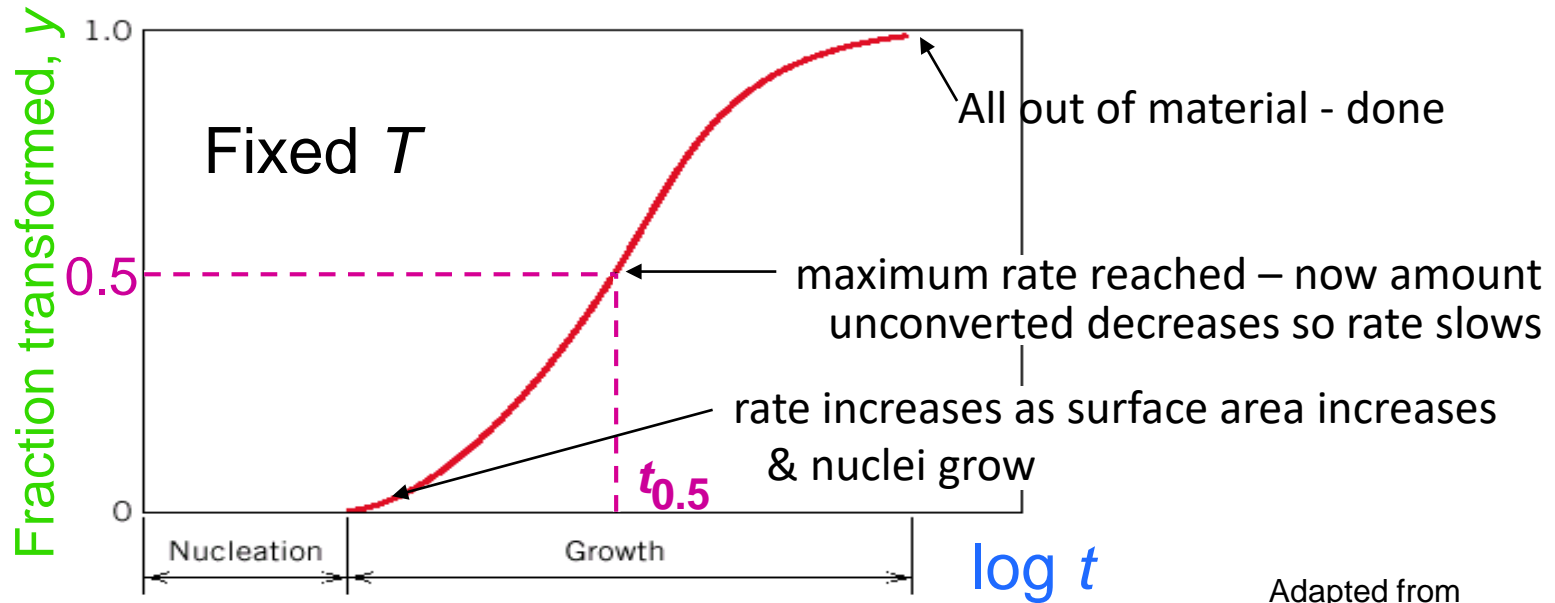
i.e. 50% transform

$$\text{Exp}(-0.7) = 0.5$$

$$kt_{0.5}^n = 0.7 \quad t_{0.5} = \frac{0.7}{k^{1/n}} \quad \frac{\pi}{3} I v^3 \quad \Rightarrow \quad t_{0.5} = \frac{0.9}{N^{1/4} v^{3/4}}$$

Rapid transformations are associated with (large values of k), or (rapid nucleation and growth rates)

Rate of Phase Transformation



Adapted from
Fig. 10.10,
Callister 7e.

Avrami rate equation $\rightarrow y = 1 - \exp(-kt^n)$

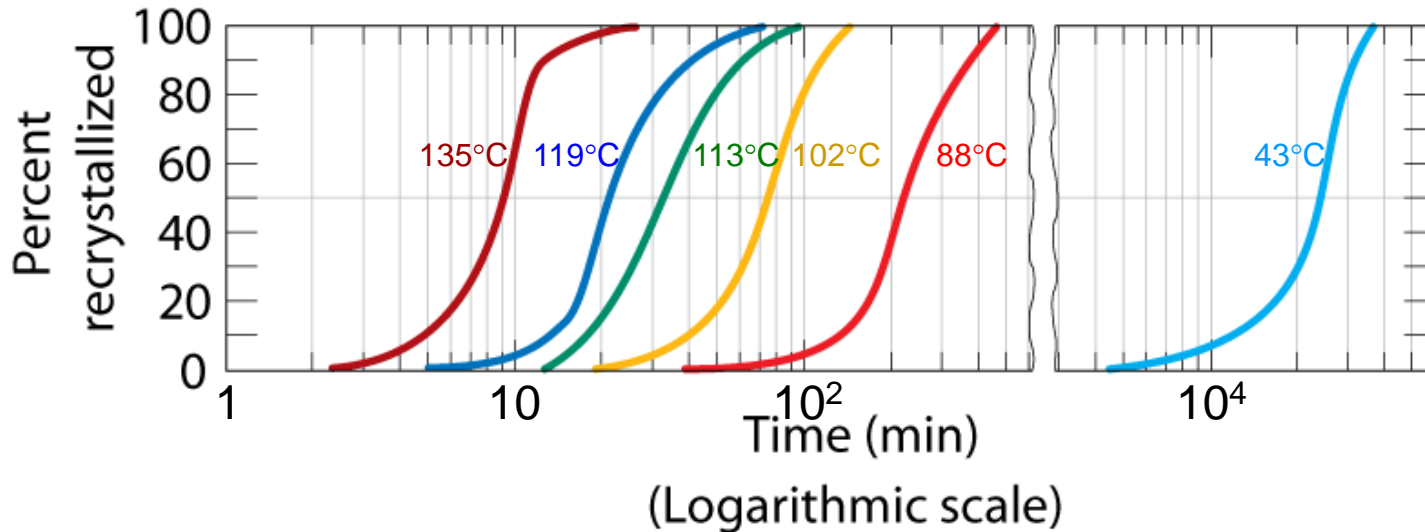
fraction transformed

time

– k & n fit for specific sample

By convention $\text{rate} = 1 / t_{0.5}$

Rate of Phase Transformations



Adapted from Fig. 10.11, *Callister 7e*. (Fig. 10.11 adapted from B.F. Decker and D. Harker, "Recrystallization in Rolled Copper", *Trans AIME*, **188**, 1950, p. 888.)

- In general, rate increases as $T \uparrow$

$$r = 1/t_{0.5} = A e^{-Q/RT}$$

- R = gas constant
- T = temperature (K)
- A = preexponential factor
- Q = activation energy

Arrhenius
expression

- r often small: equilibrium not possible!

Time-Temperature-Transformation Curves (TTT)

- How much time does it take at any one temperature for a given fraction of the liquid to transform (nucleate and grow) into a crystal?

- $f(t, T) \sim \pi I(T) \mu(T)^3 t^4 / 3$

where f is the fractional volume of crystals formed, typically taken to be 10^{-6} , a barely observable crystal volume.

Nucleation rates

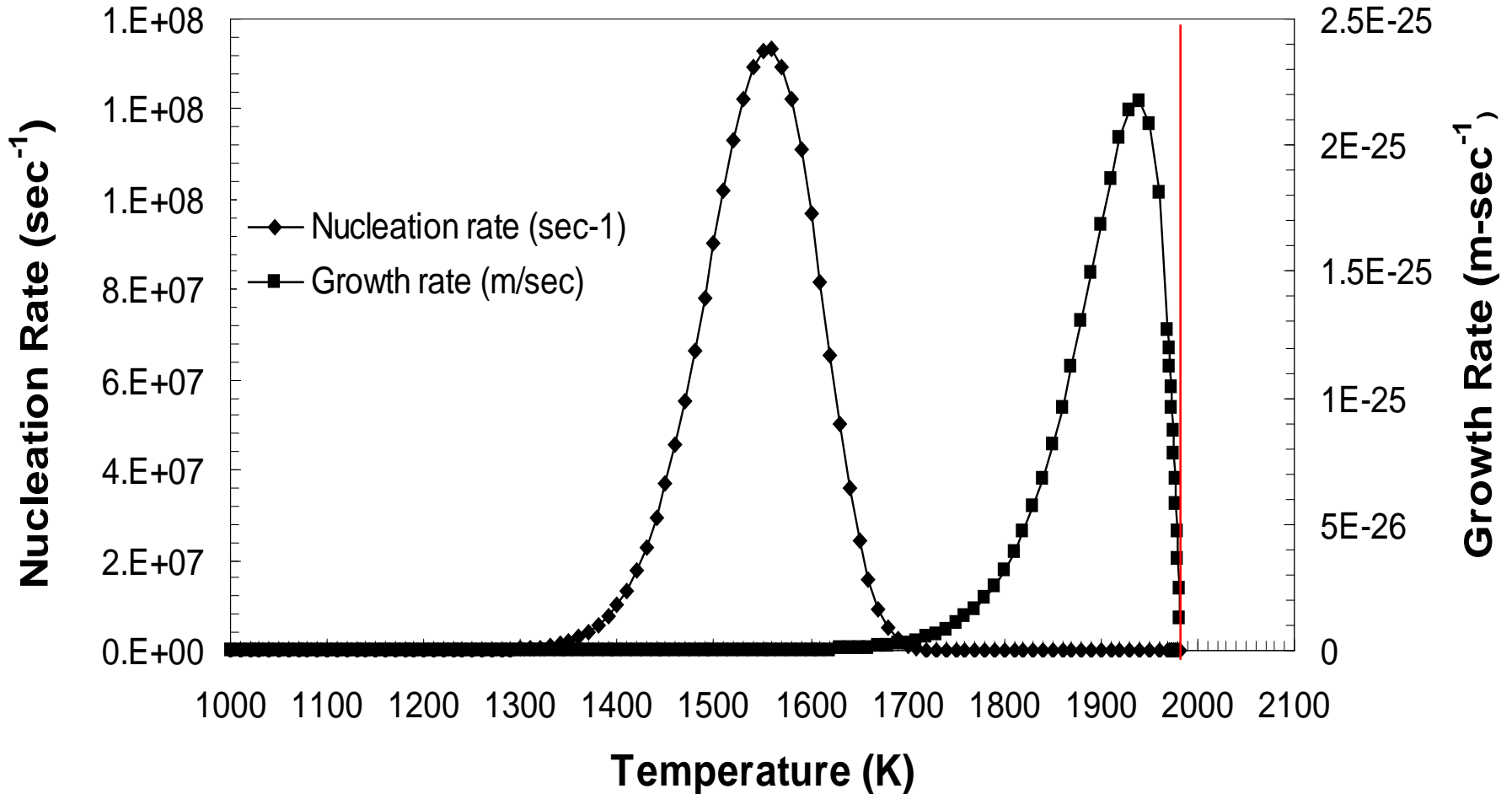
$$I = n \nu \exp \left\{ \left(\frac{16\pi\Delta H_{cryst}}{81RT} \right) \left(\frac{T_m}{\Delta T} \right)^2 \right\} \exp \left\{ \frac{-\Delta E_D}{RT} \right\}$$

Growth rates

$$\mu(T) = \left(\frac{fRT}{3N\pi a^2 \eta(T)} \right) \left(1 - \exp \left[\left(\frac{\Delta H_m}{RT} \right) \left(\frac{\Delta T}{T_m} \right) \right] \right)$$

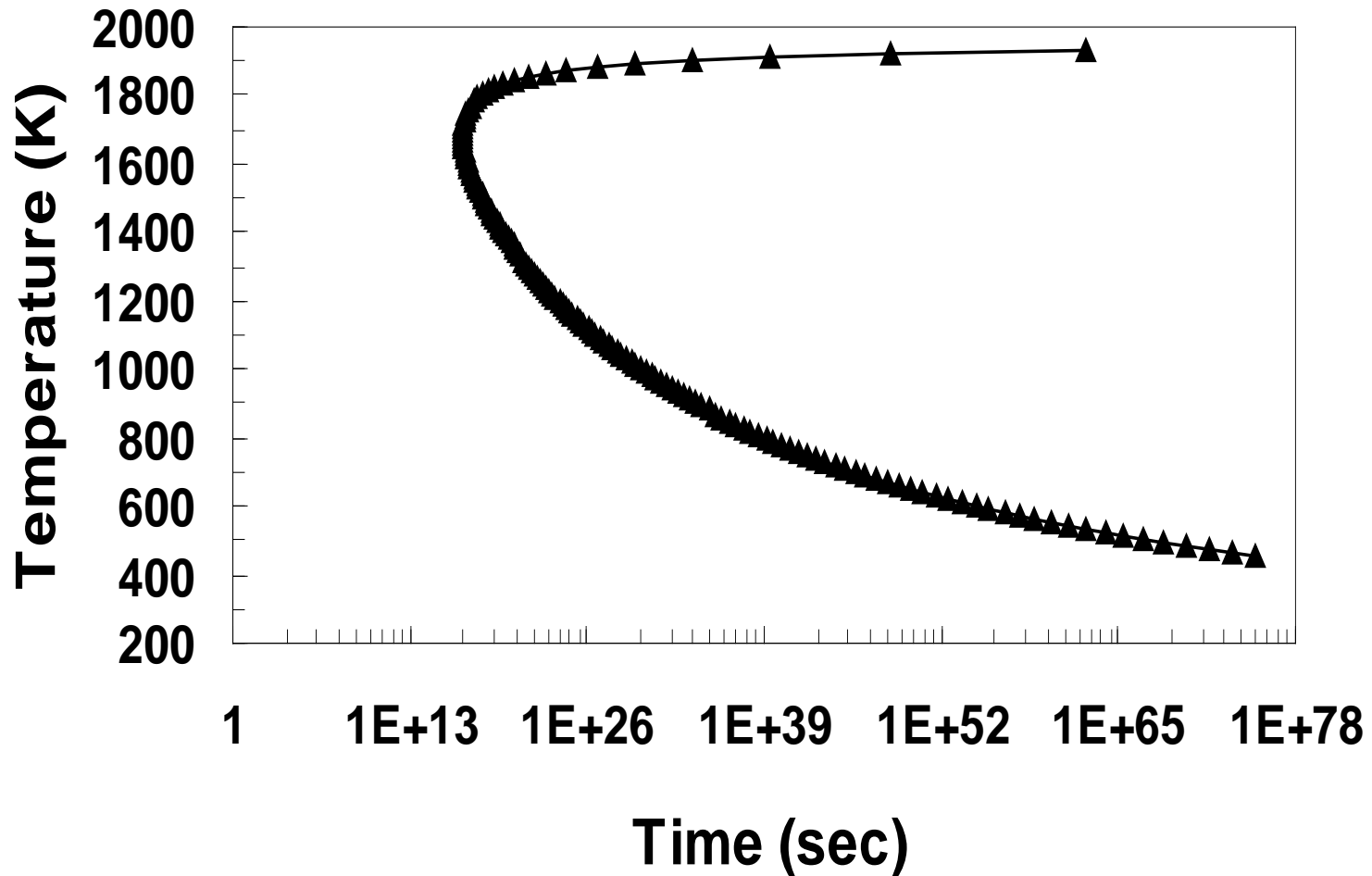
Nucleation and Growth Rates

Nucleation and Growth for Silica



Time Transformation Curves for Silica

T-T-T Curve for Silica



* Time-Temperature-Transformation diagrams

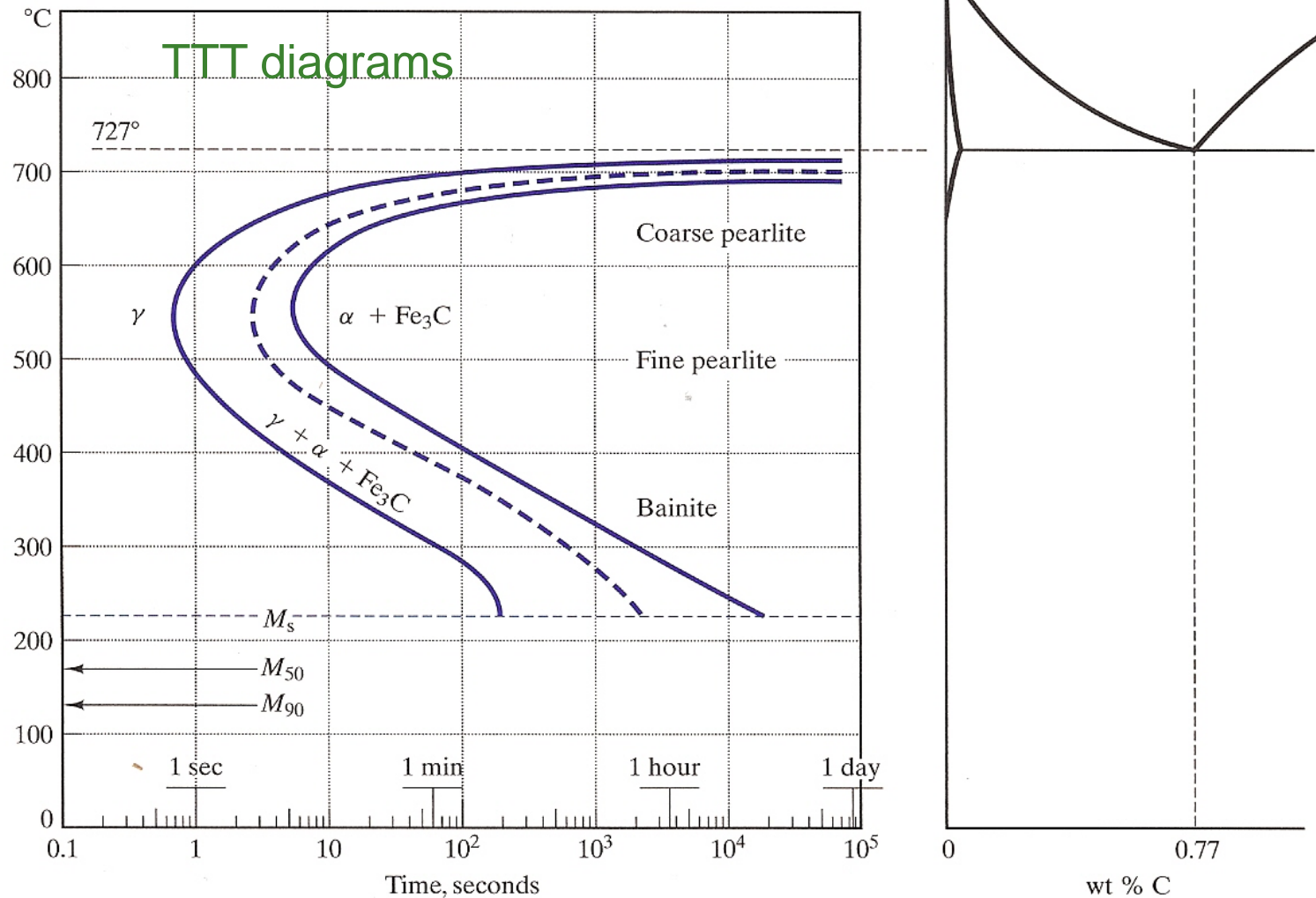


FIGURE 10.11 A more complete TTT diagram for eutectoid steel than was given in Figure 10.7. The various stages of the time-independent (or diffusionless) martensitic transformation are shown as horizontal lines. M_s represents the start, M_{50} represents 50% transformation, and M_{90} represents 90% transformation. One hundred percent transformation to martensite is not complete until a final temperature (M_f) of -46°C .

* Continuous Cooling Transformation diagrams

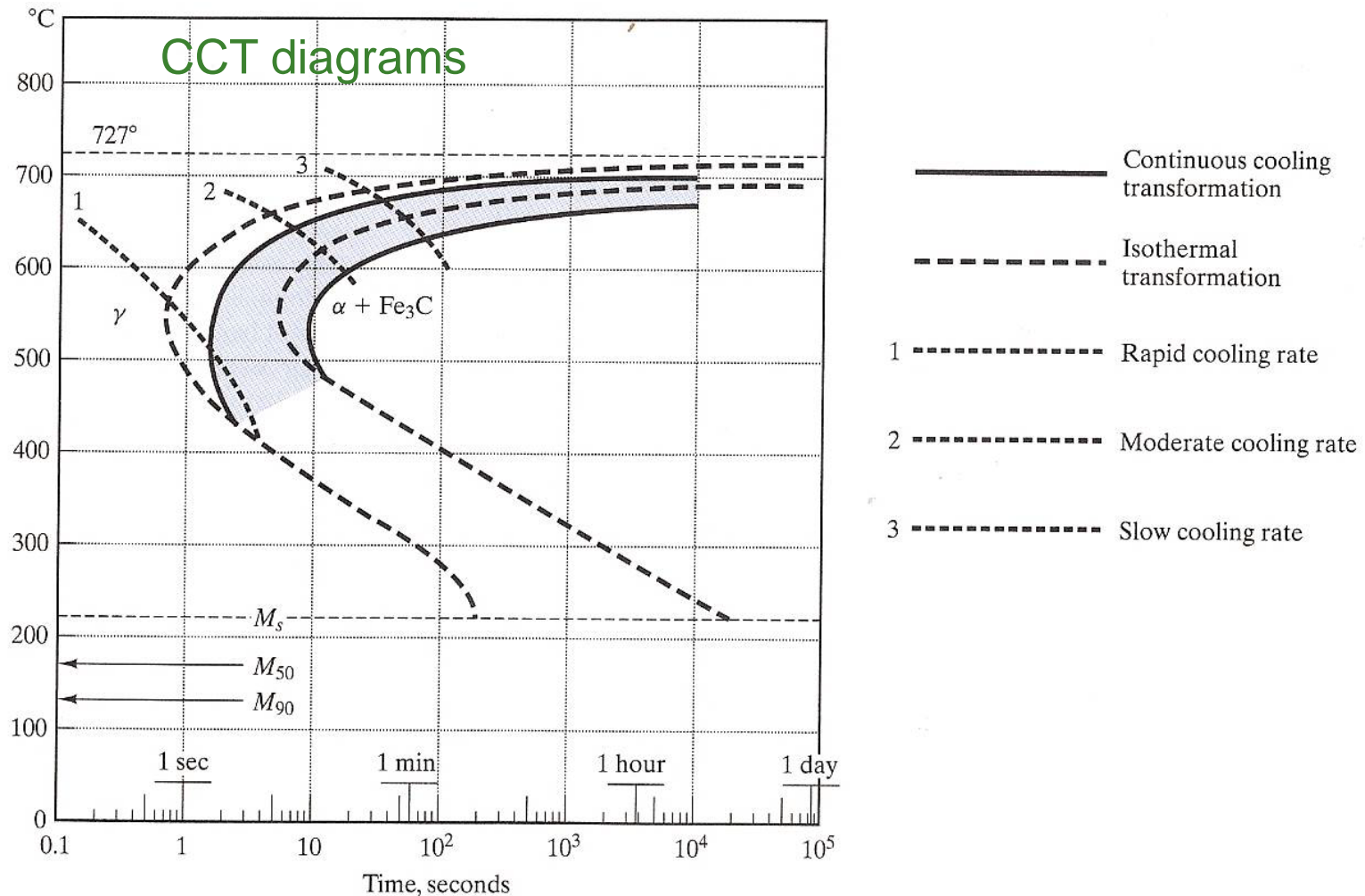


FIGURE 10.14 A continuous cooling transformation (CCT) diagram is shown superimposed on the isothermal transformation diagram of Figure 10.11. The general effect of continuous cooling is to shift the transformation curves downward and toward the right. (After Atlas of Isothermal Transformation and Cooling Transformation ⁴⁴Diagrams, American Society for Metals, Metals Park, OH, 1977.)