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$$u(t) = P \sin \omega t$$

$$G(s) = \frac{K(s+z_1)(s+z_2)\cdots(s+z_n)}{(s+s_1)(s+s_2)\cdots(s+s_n)}$$

$$Y(s) = G(s) \cdot u(s),$$
  $u(s) = P \cdot \frac{\omega}{s^2 + \omega^2}$ 

$$Y(s) = G(s) \cdot \frac{P\omega}{s^2 + \omega^2}$$
$$= \frac{a}{s + j\omega} + \frac{\overline{a}}{s - j\omega} + \frac{b_1}{s + s_1} + \frac{b_2}{s + s_2} + \dots + \frac{b_n}{s + s_n}$$

$$y(t) = ae^{-j\omega t} + \overline{a}e^{j\omega t} + b_1e^{-s_1t} + b_2e^{-s_2t} + \dots + b_ne^{-s_nt}$$



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$$y(t) = ae^{-j\omega t} + \overline{a}e^{j\omega t}$$

$$a = G(s) \cdot \frac{P\omega}{s^2 + \omega^2} (s + j\omega) \Big|_{s=-j\omega} = -\frac{P}{2j} G(-j\omega)$$

$$\overline{a} = G(s) \cdot \frac{P\omega}{s^2 + \omega^2} (s - j\omega) \Big|_{s=j\omega} = \frac{P}{2j} G(j\omega)$$

$$G(j\omega) = G_x + jG_y$$

$$= |G(j\omega)| \cos \phi + j |G(j\omega)| \sin \phi$$

$$= |G(j\omega)| (\cos \phi + j \sin \phi) = |G(j\omega)| e^{j\phi}$$

Similarly,  $G(-j\omega) = |G(-j\omega)|e^{-j\phi}$ 

$$\Rightarrow a = -\frac{P}{2j}G(-j\omega) = -\frac{P}{2j}|G(j\omega)|e^{-j\phi}$$
$$\overline{a} = \frac{P}{2j}G(j\omega) = \frac{P}{2j}|G(j\omega)|e^{j\phi}$$





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$$\therefore y(t) = ae^{-j\omega t} + \overline{a}e^{j\omega t}$$

$$= -\frac{P}{2j} |G(j\omega)| e^{-j\phi} e^{-j\omega t} + \frac{P}{2j} |G(j\omega)| e^{j\phi} e^{j\omega t}$$

$$= |G(j\omega)| \frac{P}{2j} (e^{j(\omega t + \phi)} - e^{-j(\omega t + \phi)})$$

$$= |G(j\omega)| P \sin(\omega t + \phi)$$



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# **Frequency Response of First Order Systems**





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# **Frequency Response of Second Order Systems**





$$y(t) = A(j\omega)\sin(\omega t + \phi)$$

Magnitude ratio

$$M(\omega) = \left|\frac{y(t)}{r(t)}\right| = \left|G(j\omega)\right|$$

**Phase**  $\phi(j\omega) = \angle G(j\omega)$ 

 $\Rightarrow$  Frequency response





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$$\frac{Y(s)}{R(s)} = G(s) = \frac{\omega_n^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$$

$$\frac{Y(j\omega)}{R(j\omega)} = G(j\omega) = \frac{\omega_n^2}{(j\omega)^2 + 2\zeta\omega_n\omega j + \omega_n^2} = \frac{\omega_n^2}{\omega_n^2 - \omega^2 + 2\zeta\omega_n\omega j}$$

$$M(\omega) = \left|\frac{Y(j\omega)}{R(j\omega)}\right| = \left|\frac{\omega_n^2}{(\omega_n^2 - \omega^2)^2 + 4\zeta^2\omega_n^2\omega^2} \cdot (\omega_n^2 - \omega^2 - 2\zeta\omega_n\omega j)\right|$$

$$= \frac{\omega_n^2}{\sqrt{(\omega_n^2 - \omega^2)^2 + 4\zeta^2\omega_n^2\omega^2}} = \frac{1}{\sqrt{\left(1 - \frac{\omega_n^2}{\omega_n^2}\right)^2 + 4\zeta^2 \cdot \frac{\omega^2}{\omega_n^2}^2}}$$

$$Y(j\omega)$$

$$\frac{Y(j\omega)}{R(j\omega)} = M(j\omega) \angle G(j\omega)$$



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$$M^{2}(\omega) = \frac{1}{\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)^{2} + 4\zeta^{2} \cdot \frac{\omega^{2}}{\omega_{n}^{2}}}$$

$$\frac{dM^{2}(\omega)}{d\omega} = \frac{-\left[2\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)\left(-\frac{2\omega}{\omega_{n}^{2}}\right) + 8\zeta^{2} \cdot \frac{\omega}{\omega_{n}^{2}}\right]}{\left[\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)^{2} + 4\zeta^{2} \cdot \frac{\omega^{2}}{\omega_{n}^{2}}\right]^{2}} = 0$$

$$\Rightarrow 2\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right)\left(-\frac{2\omega}{\omega_{n}^{2}}\right) + 8\zeta^{2} \cdot \frac{\omega}{\omega_{n}^{2}} = 0, \qquad -\left(1 - \frac{\omega^{2}}{\omega_{n}^{2}}\right) + 2\zeta^{2} = 0$$

$$\therefore \omega = \omega_{n}\sqrt{1 - 2\zeta^{2}} = \omega_{m}, \qquad M_{m} = \frac{1}{2\zeta\sqrt{1 - \zeta^{2}}}$$



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# **Unit Step Response VS Frequency Response**



Unit step response

Frequency response







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### **Control System Design in Frequency Response**

1. large  $\omega_m \rightarrow$  fast time response

- 2.  $M_m$ ,  $M_p$  : function of damping ratio  $\zeta$ large  $M_m \rightarrow$  large  $M_p$
- 3. good damping characteristics  $1 < M_m < 1.4$
- 4. minimum effect of any undesirable noise





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### **Exponential Determination of Transfer Function**

$$r(t) \longrightarrow \begin{array}{c} Unknown \\ System \end{array} \longrightarrow y(t)$$

$$r(t) = \sin \omega_i t$$
;  $y(t) = M(\omega_i) \sin(\omega_i t - \phi)$ 

some frequency 
$$\omega_i \rightarrow \text{measure} M(\omega_i)$$
  
steady state





- Bode Plot (Logarithmic Plots)
  - i) Log Magnitude  $L_m G(j\omega) = 20 \log |G(j\omega)| dB$
  - ii) dB : decibel, Logarithm of the magnitude
- iii) 1 decade :  $1Hz \sim 10Hz$  frequency width  $2.5Hz \sim 25Hz$
- Drawing the Bode Plots

The reason of using logarithm – mathematical operation  $\times, \div \rightarrow +, -$ – typical three types : simple straight line asymtotic approximations



ex) 
$$G(s) = \frac{K(1+sT_{1})(1+sT_{2})}{s^{2}(1+sT_{3})(1+as+bs^{2})}$$
$$L_{m}G(j\omega) = L_{m}K + L_{m}(1+j\omega T_{1}) + L_{m}(1+j\omega T_{2})$$
$$-2L_{m}(j\omega) - L_{m}(1+j\omega T_{3}) - L_{m}\left\{1+aj\omega+b(j\omega)^{2}\right\}$$





iii) 
$$(1+j\omega T)$$
  
 $L_m(1+j\omega T) = 20\log|1+j\omega T| = 20\log\sqrt{1+\omega^2 T^2}$   
 $\omega T << 1 \rightarrow L_m(1+j\omega T) \cong 20\log 1 = 0dB$   
 $\omega T >> 1 \rightarrow L_m(1+j\omega T) \cong 20\log \omega T$   
 $= 20\log \omega + 20\log T$   
iv)  $\frac{1}{1+j\omega T}$   
 $L_m(1+j\omega T)^{-1} = -20\log|1+j\omega T| = -20\log\sqrt{1+\omega^2 T^2}$   
 $\omega T << 1 \rightarrow L_m(1+j\omega T)^{-1} \cong -20\log 1 = 0dB$   
 $\omega T >> 1 \rightarrow L_m(1+j\omega T)^{-1} \cong -20\log \omega T$   
 $= -20\log \omega - 20\log T$ 







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#### **Vibration Isolation in Rotating Systems**

Vibration due to rotating unbalance

(1)  $M\ddot{x} + b\dot{x} + kx = p(t)$ =  $mr\omega^2 \sin \omega t$ X(s) 1

$$\frac{A(s)}{P(s)} = \frac{1}{Ms^2 + bs + k} = G(s)$$





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### **Vibration Isolation in Rotating Systems**

(2) Frequency response

$$\frac{x(j\omega)}{p(j\omega)} = \frac{1}{-M\omega^2 + bj\omega + k} = G(j\omega)$$

$$x(t) = X \sin(\omega t + \phi) = |G(j\omega)| \sin(\omega t + \phi) \cdot mr\omega^2$$

$$\phi = -\tan^{-1} \frac{b\omega}{k - M\omega^2} = -\tan^{-1} \frac{2\zeta\omega/\omega_n}{1 - \omega^2/\omega_n^2}$$

$$|G(j\omega)| = \frac{1}{\sqrt{(k - M\omega^2)^2 + b^2\omega^2}}$$

$$= \frac{1/k}{\sqrt{(1 - \omega^2/\omega_n^2)^2 + (2\zeta\omega/\omega_n)^2}}$$

$$\frac{1}{k}$$

$$w range$$



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# **Vibration Isolation in Rotating Systems**

#### Vibration isolators



Reduce the magnitude of force transmitted from a machine to its foundation



**Suspensions** 

Reduce the magnitude of motion transmitted from a vibratory foundation to a system

Isolator i) - load supporting elements, spring.

- energy-dissipating elements, damper.
- ii) Synthetic rubber : both load supporting and energy dissipating



# Transmissibility

Transmissibility :

A measure of the reduction of transmitted force or motion afforded by an isolator.

 $TR(machine) = \frac{the force amplitude transmitted to the foundation}{the amplitude of the exciting force}$ 

 $TR(suspensions) = \frac{the \ vibration \ amplitude \ of \ the \ system}{the \ vibration \ amplitude \ of \ the \ foundation}$ 



### **Transmissibility for Force Excitation**

$$p(t) = mr\omega^{2} \sin \omega t = F_{0} \sin \omega t$$
$$f(t) = b\dot{x} + kx = F_{t} \sin (\omega t + \phi)$$
$$X(s) = \frac{1}{Ms^{2} + bs + k}P(s)$$
$$F_{t}(s) = (bs + k)X(s) = \frac{bs + k}{Ms^{2} + bs + k} \cdot P(s)$$

Frequency response

$$\frac{F_t(j\omega)}{p(j\omega)} = \frac{bj\omega + k}{-M\omega^2 + bj\omega + k} = \frac{(b/M)j\omega + (k/M)}{-\omega^2 + (b/M)j\omega + k/M}$$



# **Transmissibility for Force Excitation**

$$k/M = \omega_n^2, \quad b/M = 2\zeta\omega_n \quad \Rightarrow \frac{F_t(j\omega)}{p(j\omega)} = \frac{1+j(2\zeta\omega/\omega_n)}{1-\omega^2/\omega_n^2+j(2\zeta\omega/\omega_n)}$$

$$TR = \frac{F_t}{F_0} = \left|\frac{F_t(j\omega)}{p(j\omega)}\right| \qquad TR$$

$$= \frac{\sqrt{1+(2\zeta\omega/\omega_n)^2}}{\sqrt{(1-\omega^2/\omega_n^2)^2+(2\zeta\omega/\omega_n)^2}} \qquad TR$$

$$= \frac{\sqrt{1+(2\zeta\beta)^2}}{\sqrt{(1-\beta^2)^2+(2\zeta\beta)^2}}$$
Note,  $\beta = \sqrt{2}, \quad TR = 1$ 
for any  $\zeta$ 

$$\int_{0}^{1+(2\zeta\beta)} \frac{1}{\sqrt{(1-\beta^2)^2+(2\zeta\beta)^2}} \qquad \int_{0}^{1+(2\zeta\beta)} \frac{1}{\sqrt{(1-\beta^2)^2+(2\zeta\beta)^2}} \qquad \int_{0}^$$



# **Automobile Suspension Systems**



x(t)



Simplified model





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If b is small and  $\omega = \omega_n$ Then resonance – excessive vibration – extremely large force transmitted

 $\omega$  is very close to  $\omega_n$ , critical speed, => A dynamic vibration absorber







$$M \ddot{x} = -kx - b\dot{x} - k_a (x - y) + p(t)$$
  

$$m_a \ddot{y} = -k_a (y - x)$$
  
Laplace Transform  

$$(Ms^2 + bs + k + k_a) X(s) - k_a Y(s) = P(s)$$
  

$$(m_a s^2 + k_a) Y(s) - k_a X(s) = 0$$
  

$$\Rightarrow \frac{X(s)}{P(s)} = \frac{m_a s^2 + k_a}{(Ms^2 + bs + k + k_a)(m_a s^2 + k_a) - k_a^2}$$

**Frequency response** 

$$\frac{X(j\omega)}{P(j\omega)} = \frac{-m_a\omega^2 + k_a}{\left(-M\omega^2 + k + k_a\right)\left(-m_a\omega^2 + k_a\right) - k_a^2} \qquad \text{for small } b$$

The transmitted force

$$f(t) = kx + b\dot{x} \cong kx$$



$$|X(j\omega)| = \left| \frac{k_a - m_a \omega^2}{\left( -M\omega^2 + k + k_a \right) \left( k_a + m_a \omega^2 \right) - k_a^2} \right| |P(j\omega)|$$
$$= \left| \frac{mr\omega^2 \left( k_a - m_a \omega^2 \right)}{\left( -M\omega^2 + k + k_a \right) \left( k_a + m_a \omega^2 \right) - k_a^2} \right|$$

If we choose  $k_a$ ,  $m_a$  so that,

$$k_a - m_a \omega^2 = 0$$
  
then,  $|X(j\omega)| = 0$ ,  $|f(t)| = 0$  at this frequency.

If  $\omega_{exc} = \omega_n$ 

we use dynamic vibration absorbers.





Physically,  $f_{a}(t) \longrightarrow p(t) = P \sin \omega t$  M  $f_{a}(t) \longrightarrow p(t) = p(t)$  M

$$f_a(t) = -k_a(x - y) = k_a y = p(t), \qquad (at \ \omega_{exc} = \omega_n, \ x \cong 0)$$

Spring force  $k_a y$  cancels p(t).

$$\frac{Y(j\omega)}{P(j\omega)} = \frac{k_a}{\left(-M\omega^2 + k + k_a\right)\left(-m_a\omega^2 + k_a\right) - k_a^2}$$
  
If we choose  $m_a, k_a$  so that,  $k_a - m_a\omega^2 = 0$ 

then 
$$\frac{Y(j\omega)}{P(j\omega)} = -\frac{1}{k_a}$$
 magnitude ratio  $\left|\frac{1}{k_a}\right|$   
 $p(t) = \sin \omega t$   
 $y(t) = \frac{1}{k_a} P \sin(\omega t - 180^\circ) = -\frac{P}{k_a} \sin \omega t$  (phase  $-180^\circ$ )



# Seismograph



A device used to measure ground displacement during earthquakes

Equation of motion :  $m\ddot{x} = k(y-x) + b(\dot{y}-\dot{x})$ measurement : z = x - y

$$\Rightarrow m(\ddot{y}+\ddot{z})+b\dot{z}+kz=0$$

$$m\ddot{z} + b\dot{z} + kz = -m\ddot{y}$$

$$\Rightarrow Transfer Function : \frac{Z(s)}{Y(s)} = \frac{-ms^2}{ms^2 + bs + k} = \frac{-s^2}{s^2 + 2\zeta\omega_n s + \omega_n^2}$$
$$\left(\frac{k}{m} = \omega_n^2, \ \frac{b}{m} = 2\zeta\omega_n\right)$$



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# Seismograph



# Accelerometer



The system configuration is basically the same as the seismograph, but choice of undamped natural frequency is different.

Equation of motion : the same as the seismograph,  $m\ddot{z} + b\dot{z} + kz = -m\ddot{y}$ 

$$\Rightarrow Transfer Function : \frac{Z(s)}{s^2 Y(s)} = \frac{-m}{ms^2 + bs + k} = \frac{-1}{s^2 + 2\zeta \omega_n s + \omega_n^2}$$





# **Performance Characteristics**



$$Y = G(s)e$$

$$R - H(s)Y = e$$

$$\Rightarrow Y = G(s)(R - H(s)Y)$$

$$Y(1 + G(s)H(s)) = G(s)R$$

$$\Rightarrow \frac{Y(s)}{R(s)} = \frac{G(s)}{1 + G(s)H(s)} = G_0(s)$$



Consider, 2nd order system

$$\frac{Y(s)}{R(s)} = \frac{K}{s^2 + 2\zeta\omega_n s + \omega_n^2} = G_0(s)$$

Now, consider "time-response"-"pole" of the transfer function relations

$$Y(s) = \frac{1}{s} \frac{K}{s^2 + 2\zeta \omega_n s + \omega_n^2} \qquad (step input response)$$
$$= \frac{K/\omega_n^2}{s} + \frac{-(K/\omega_n^2)s - 2\zeta K/\omega_n}{s^2 + 2\zeta \omega_n s + \omega_n^2}$$

pole  $s^2 + 2\zeta \omega_n s + \omega_n^2 = 0$  $s = -\zeta \omega_n \pm \sqrt{\zeta^2 - 1} \omega_n$ 



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