

457.562 Special Issue on River Mechanics (Sediment Transport) .02 Review of Fluid Mechanics





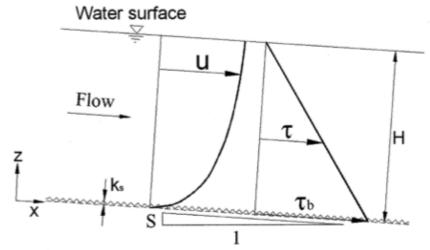
 Shear stress between adjacent layers in a simple parallel flow is det ermined by the viscosity and the velocity gradient



- If laminar flow is disturbed by an obstacle or roughness, then disturb ances must be damped by the molecular viscosity (stable)
- However, in turbulent flow, inertia overcomes the viscous force and disturbance grows and becomes random motion.
- Therefore, in turbulent motion rapid mixing of fluid particles during fl ow cause significant diffusion.
- What is the actual meaning of viscosity? Is there any other name of viscosity?



- Consider a steady, turbulent, uniform, open-channel flow having
 - H : a mean depth
 - U: a mean flow depth
 - -B: a mean width (much greater than depth)
 - S : a mean slope
 - $-k_s$: effective height
 - In a wide channel (*B/H* >>1)





- The roughness height (or effective height) will be proportional to a m ean size or diameter *D*.
- The bed shear stress

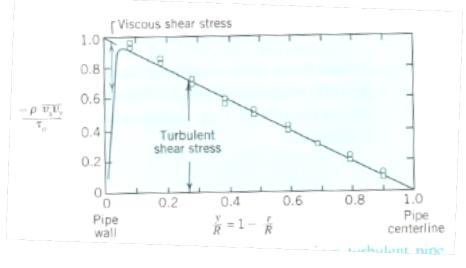
 $\tau_b = \rho g HS$ (ρ = water density, g=gravity)

- The above equation is the one-dimensional momentum conservation n equation
- Using the boundary shear stress, the shear velocity can be defined as

$$u_* = \sqrt{\tau_b / \rho}$$

- The sear velocity (the boundary shear stress)
 - provides a direct measure of the flow intensity
 - able to entrain and transport sediment particles.





- The above figure shows typical variation of total shear stress in a tur bulent pipe flow.
- Based on the linear relationship of shear profile (in the previous clas s), and the Prandtl relationship,

$$\tau = \tau_b \left(1 - \frac{z}{H} \right) = \rho l^2 \left(\frac{du}{dz} \right)^2$$



Near wall, *I=kz*, and mixing length in a pipe flow

$$l = kz \left(1 - \frac{z}{H}\right)^{1/2}$$

Then the previous equation can be written

$$\frac{\tau_b}{\rho \kappa^2 z^2} = \left(\frac{du}{dz}\right)^2$$
$$\frac{du}{dz} = \frac{\sqrt{\tau_b / \rho}}{\kappa z} = \frac{u_*}{\kappa z}$$
$$\frac{u}{u_*} = \frac{1}{\kappa} \ln\left(\frac{z}{z_0}\right) \qquad (1)$$

Only in a relatively thin layer (z/H=0.2)

u = time -averaged flow velocity at a distance z above the bed $z_0 =$ bed roughness length (distance above the bed where $u(z_0) = 0$) $\kappa =$ is von Karman's constant ~0.41.



Through intensive experiment works,

$$\frac{u}{u_*} = \frac{1}{\kappa} \ln \frac{u_* z}{v} + 5.5 \qquad (2)$$

• When the effective roughness height becomes by (1) and (2)

 $z_b = v / 9u_*$ Hydraulic smooth condition of flow

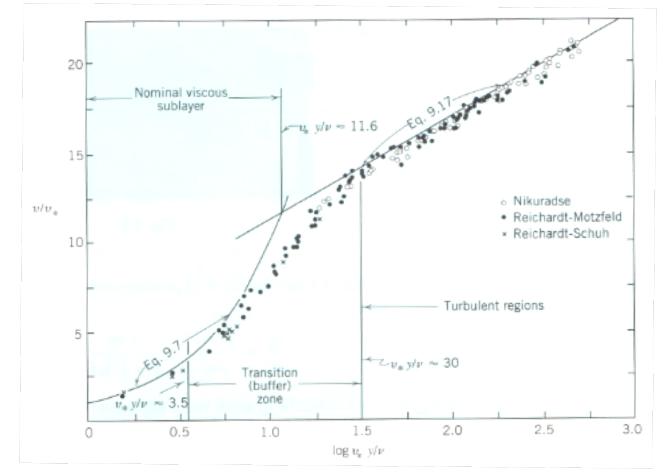
Very near the smooth wall (viscous sublayer), velocity is linear with depth

$$\frac{u}{u_*} = \frac{u_* z}{v} \qquad (\text{until } z = \delta_v)$$

Then two equations above need to match

$$\frac{u}{u_*} = \frac{u_*\delta_v}{v} = \frac{1}{\kappa} \ln \frac{u_*\delta_v}{v} + 5.5 \quad (at \ z = \delta_v)$$
$$\frac{u_*\delta_v}{v} = 11.6 \qquad \delta_v = 11.6 \frac{v}{u_*}$$





Velocity distribution near a smooth wall



- But, most boundaries in river flow are rough. Therefore we need to c onsider roughness height.
- If $k_s / \delta_v > 1$, then viscous sublayer will not exist.
- The corresponding logarithmic velocity profile is given by

$$\frac{u}{u_*} = \frac{1}{\kappa} \ln\left(\frac{z}{k_s}\right) + 8.5 = \frac{1}{\kappa} \ln\left(30\frac{z}{k_s}\right)$$

It follows that

 $z_0 = k_s / 30$ Hydraulically rough flow

Then what happens in the intermediate condition?



- 1. Flow velocity distribution : law of the wall
- Hydraulically rough flow

$$\frac{k_s}{\delta_v} \gg 1, \text{Re}_* \gg 11.6$$

- Hydraulically smooth flow $\frac{k_s}{\delta_{...}} \ll 1$, Re_{*} $\ll 11.6$
- Where "the Roghness Reynolds number"

$$\operatorname{Re}_{*} = \frac{u_{*}k_{s}}{v}, \quad (\operatorname{R}_{*c} = 11.6)$$

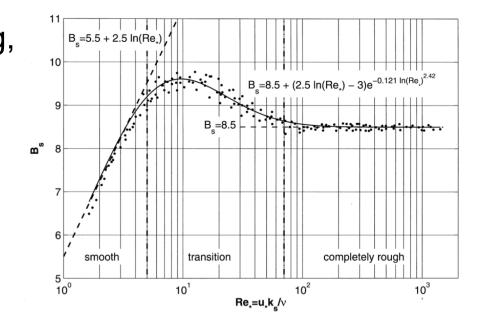
 Yalin (1992) proposed a way to represent both ranges and intermedi ate ranges as

$$\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln\left(\frac{z}{k_s}\right) + B_s$$
$$B_s = 8.5 + \left[2.5\ln(\operatorname{Re}_*) - 3\right] e^{-0.121\left[\ln\operatorname{Re}_*\right]^{2.42}}$$



- 1. Flow velocity distribution : law of the wall
 - Alternative way of writing, $\frac{\overline{u}}{u_*} = \frac{1}{\kappa} \ln \left(A_s \frac{z}{k_s} \right)$

$$A_s = e^{\kappa B_s}$$





2. Velocity-Defect and Log-Wakes laws

- Equation (1) requires some knowledge of the bed roughness charact eristics.
- An alternative formulation can be obtained if the flow depth H is intro duced as the relevant length scale.
- Assuming that the maximum flow velocity u_{max} takes places at the w ater surface z=H, then equation (1) can be manipulated to obtain the velocity-defect law (outer low of the wall).

$$\frac{u_{max} - u}{u_*} = -\frac{1}{\kappa} \ln \frac{z}{H}$$
(3)

 But, there were so many arguments and defects or log law working only in z/H<0.2. Nezu and Nakagawa added a wake function.

$$\frac{u}{u_*} = \frac{1}{\kappa} \ln\left(\frac{u_*z}{\nu}\right) + 5.5 + w\left(\frac{z}{H}\right)$$



2. Velocity-Defect and Log-Wakes laws

 W(z/H) is the wake function first proposed by Coles (1956) for turbul ent boundary-layer flows,

$$w\left(\frac{z}{H}\right) = \frac{2W_0}{\kappa}\sin^2\left(\frac{\pi}{2}\frac{z}{H}\right)$$

- W₀ is the Coles wake parameter, expressing the strength of the wak e function.
- Equation (3) can also be written in log-wake form

$$\frac{u_{max} - u}{u_*} = -\frac{1}{\kappa} \ln\left(\frac{z}{H}\right) + \frac{2W_0}{\kappa} \cos^2\left(\frac{\pi z}{2H}\right)$$

 Coles wake parameter varies depending on the condition and incras es with increasing sediment concentration raging from 0.191 to 0.86
 1.



2. Relations for channel resistance

- Most river flows are indeed hydraulically rough.
- The previous log law is used to obtain an appropriate expression for depth-averaged velocity U.

$$U = \frac{1}{H} \int_{0}^{H} u \, dz$$

$$\frac{U}{u_{*}} = \frac{1}{H} \int_{k_{s}}^{H} \left[\frac{1}{\kappa} \ln \left(\frac{z}{k_{s}} \right) + 8.5 \right] dz \qquad \text{(Avoiding } z=0 : \text{ singular point)}$$

$$\frac{U}{u_{*}} = \frac{1}{\kappa} \ln \left(\frac{H}{k_{s}} \right) + 6 = \frac{1}{\kappa} \ln \left(11 \frac{H}{k_{s}} \right)$$

This relation is known as Keulegan's resistance law for rough flow.

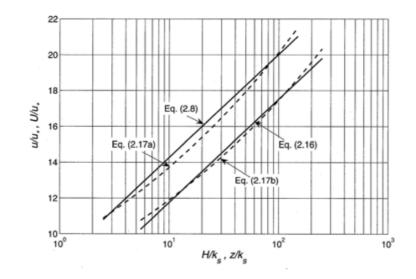


2. Relations for channel resistance

Manning-Strickler form.

$$\frac{u}{u_*} = \frac{1}{\kappa} \ln\left(30\frac{z}{k_s}\right) \cong 9.34 \left(\frac{z}{k_s}\right)^{1/6}$$
$$\frac{U}{u_*} = \frac{1}{\kappa} \ln\left(11\frac{H}{k_s}\right) \cong 8.1 \left(\frac{H}{k_s}\right)^{1/6}$$

Log - > power : was first presented by Keulegan.





2. Relations for channel resistance

Bed shear stress can be found as

$$\tau_{b} = \rho u_{*}^{2} = \rho C_{f} U^{2}$$

$$C_{f} = \left[\frac{1}{\kappa} \ln\left(11\frac{H}{k_{s}}\right)\right]^{-2} \quad or \quad \left[8.1\left(\frac{H}{k_{s}}\right)^{1/6}\right]^{-2}$$

In two-dimensional case

$$(\tau_{bs}, \tau_{bn}) = \rho C_f \left[U^2 + V^2 \right]^{1/2} (U, V)$$