

## 6. Diffusionless Transf.

### 6.1. Charac. of Diffusionless Transf.

- the formation of M appears to be a random process
- M( $\alpha'$ ) in the shape of a lens and spans an entire grain  
the p of plates  $\approx f(\gamma/\text{grain size})$

① the transformed region coherent w/ the surrounding  $\gamma$ .

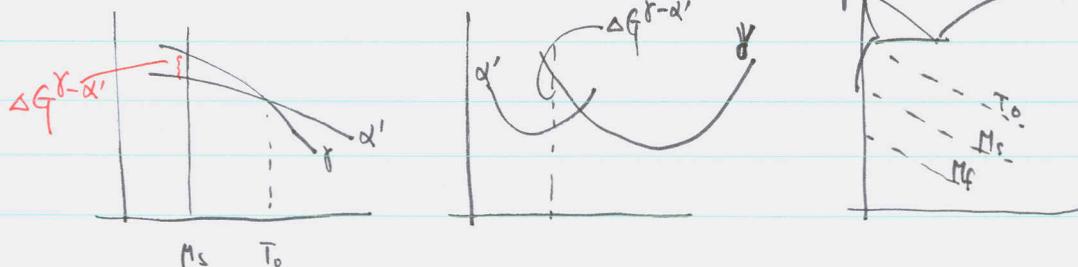
② the intersection of  $\alpha'$  w/ the surface - continuous

- Fig 6.2 ③ happens quick ( $10^{-7}$  s to span in a grain)

$\alpha'/\gamma$  interface speed  $\sim$  sound. in the solid.

cf. M of Fe-Ni exhibits isothermal growth. growth: indep.  
of thermal activi

- The M. formation starts at  $M_s$  : the temp. associated w/ a d.f.



- Increasing C contents decreases the  $M_s$  temp.  $M_f$  : finish Some retained  $\gamma$  (10-15%) may be not 100% Mart. due to the high elastic stressed <sup>int</sup> bet'n M. plates.

Like in eq.(1.17) the driving force for the nnd. of M at Ms.

$$\Delta G^{r-a'} = \Delta H^{r-a'} \frac{(T_0 - Ms)}{T_0}$$

Table 6.1  $\rightarrow$   $\Delta H$  information.

340 J/m

The ordered alloys show a rel. small  $\Delta T$ . (Fe-Pt 24% ordered)  
disordered Fe-Pt 2400 J/m.

### 6.1.1. The solid soln of Carbon in Fe.

- In an fcc lattice stuc. two possible positions for interstitial atoms.  
tetrahedral, octahedral.

Fig. 6.4.  $d_a = 0.225 D$      $d_6 = 0.414 D$     D: parent Fe dia.  $2.52\text{\AA}$   
 $(0.568\text{\AA})$      $(1.044\text{\AA})$

+ the dia. of carbon  $1.54\text{\AA} \rightarrow$  considerable distortion  
 $\hookrightarrow$  octahedral sites. occupied  $\uparrow$

- In a bcc stuc. 3 possible octahedral sites ( $\frac{1}{2}[100], \frac{1}{2}[010], \frac{1}{2}[001]$ )  
 $\hookrightarrow$  6 tetra .. in each unit cell.  
 $d_a = 0.291 D$      $d_6 = 0.155 D$     D:  $2.866\text{\AA}$ .

Note! ① although there's more 'free' space in bcc than the c-packed lattices, the larger # of possible interstitial sites (9) cause less space available / interst. than for fcc.

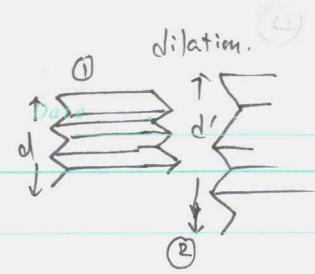
②  $d_c < d_a$ , still C, N prefer octa. sites very few  
 $\rightarrow$  causes considerable distortion.  $(\frac{1}{2}00)$  sites

distortion from  $\text{fcc}$  to  $\text{bct}$  (body centered tet.)  
at  $-100^\circ\text{C}$  the c/a ratio of the bct

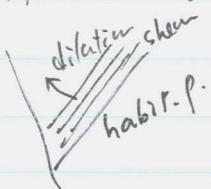
Fig 6.5(c)  $c/a = 1.005 + 0.045 (\text{w/o C})$

## 6.2 Martensite Xtallography

- Xtalliphic dependence, M grows, within a grain, in a limited # of orientations.
  - ex. Fe-alloys, the orientation variants and plate morphology chosen
    - dep. on. alloy content (C or Ni)
- The irrational nature of the growth planes of high C or Ni Martensites: Subject of much discussion. Reasons?
  - ① if M grows at speeds ~ that of sound, it needs highly mobile disl. interface.
    - a. ~~then~~, have to explain why the high mobility of an interface moving on  $\gamma$  planes<sup>which</sup> are not always asst'd w/  $\perp$  glide.
    - b. ~~the habit plane of M is observed to be macroscopically undistorted~~  
 $\hookrightarrow$  a plane common to both  $\gamma$  & M.
    - c. ~~the absence of plastic def~~  $\rightarrow$  the absence of plastic def in the form of discontinuity at the surface indicates that the shape strain does not cause any noticeable rotation of the habit plane.
  - $\rightarrow$  Coherence bet'n M and  $\gamma$ , this would cause additional displacement
- 1 For the habit plane undistorted, the M. transf. by a homo. shear // to the h.p.
- 2  $\gamma \rightarrow \alpha'$  entails  $\Delta V \uparrow$ . (4% expansion),  $\therefore$  the dilatation must occur normal to the habit plane (normal to the lens).
- 3 Then, can the bct M. lattice struc. be generated by simple shear?



the undistorted habit plane similar to twinning

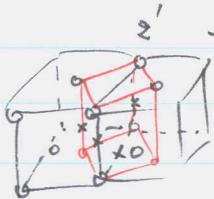


- Assume that
  - ① the equivalent macro. shape change in M formaz is a twinning shear // to the h.p. (twinning p)
  - ② + a simple uniaxial tensile dilatation ⊥ to the h.p. - strain called an invariant plane strain  $\Sigma_3 = 0$

### 6.2-1. The Bain model of the fcc $\rightarrow$ bct transf.

$\rightarrow$  twin plan

- Bain showed  $\gamma \rightarrow \alpha'$  (bcc) w/ the min. of atom movement &  $\Sigma$  in the parent lattice.



$\alpha'$  (bcc) can be achieved by

a) 20% contraction in  $\hat{z}'$  direction

b) 12% expansion in the x, y dire.

carbon sites in  $\hat{z}'$  of the bcc at  $\frac{1}{2}(\bar{1}00)$ .

$\rightarrow$  not exactly match w/ the position in fcc octa sites: the position change of C during transf.

- Min. atom movement: Bain distortion. Results in

$$(111)_\gamma \rightarrow (011)_{\alpha'}$$

$$[\bar{1}01]_\gamma \rightarrow [\bar{1}\bar{1}1]_{\alpha'} \quad (\text{k-s relation}) \quad \left. \begin{array}{l} \text{orientation } 50^\circ \text{ off.} \\ \end{array} \right.$$

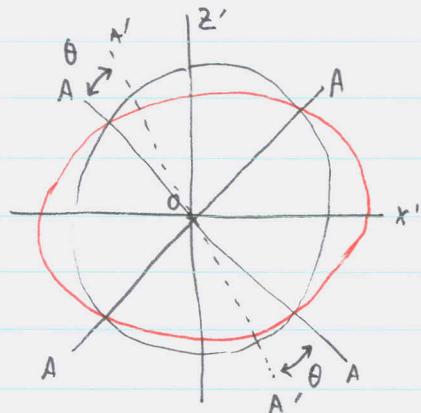
$$[\bar{1}\bar{1}0]_\gamma \rightarrow [100]_{\alpha'} \quad (\text{N-W relation})$$

$$[11\bar{2}]_\gamma \rightarrow [01\bar{1}]_{\alpha'}$$



- Invariant Plane in Bain Distortion. : not moving atoms

just distance change



Sphere - represent the fcc.

ellipsoid. - " the struc. after Bain D. (B.D)

$x'$  &  $y'$  axis 12% expansion

$z'$  axis 20% contraction

after B.D., the only vectors not shortened =  $OA$  or  $OA'$  (line)

How about plane? the vector to be a plane (undistorted)

there shouldn't be distortion in  $OY'$  axis, but it is not so.  
(cylinder)

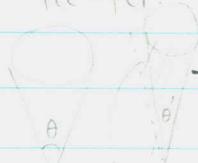
- the theory of Mart. transf.

requires an additional distortion. (postulation).

that reduce the actual expansion along  $OY'$  to zero.

- in the form of  $\pm$  slip or twinning

In Th - an internally twinned martensite plate  
fcc-fct.



- the habit plane of the M: a plane w/ no net distortion in  $\gamma$ .  
distortion avg. over many twins  $\approx 0$ .

Slip required  
.. strain energy w/  $\alpha'/\gamma$  interface of the twins exists  
if the M plate is thin (a few atomic spacings),  $\Sigma \rightarrow$  small

Slip/twin on  $\langle 111 \rangle [11\bar{2}]_{\alpha'}$ , equivalent to  $\langle \bar{1}10 \rangle [\bar{1}10]_{\gamma}$   
common slip plane in bcc slip or twinning.

6.2.2. Comparison of X-raylographic Theory w/ Exp. Results.

- Alloying addition have a great effect on the habit plane.
- As C content  $\uparrow$  in Fe-C habit plane  $\{111\} \rightarrow \{225\} \rightarrow \{259\}$ , twinning plates, lens morpholgy associated w/ a high  $t_p$ , lath morphology & bundles of needles on  $\{111\}$ ,
- Morphology quite irregular Twinning predominant at high C or Ni stainless steel  $\{259\}$ ,  $\{112\}$
- difficult to predict habit plane based on lattice para., dilatation.
- Theory only phenomenological. : no kinetic. use.

### 6.3 Theories of M. Nucleation.

- growth speed :  $10^{-5}$  to  $10^7$ /s to grow its full size (800-1000 m/s measured using resistivity).
- nucleation important in determining final g.s. of M related to toughness and strength # of nuclei  $\uparrow \rightarrow$  g.struc. finer !!

- the initial M. nucleus is coherent w/ the parent  $\delta$ .

#### 6.3.1. Formation of Coherent Nuclei of M.

$$(6.5) \quad \Delta G = A\gamma + V\Delta G_s - VAG_v. \quad V: \text{vol. of nucleus}$$

No consideration of the energy due to thermal stress during cooling applied stress etc.

$\rightarrow$  strain energy important 부교재 학과 공과대학 공기 재료 과학과  
Department of Inorganic Materials Engineering the shear comp. of Bain strain  $\approx 0.32$

the interfacial (surface) energy, relatively small  $\therefore$  coh. nature.

- Nucl. of a thin ellipsoidal nucleus (radius:  $a$  thickness:  $c$  vol.:  $V$ )

Assumptions: nucl. in the matrix, homogeneous nucl.

the nucleus forms by a simple shear,  $s$ ,  $\parallel$  to the plane of the disc. (complete coherency at the interface)

$$(6.6) \text{ from (6.5)} \Delta G = 2\pi a^2 \gamma + V \cdot 2\mu \left(\frac{s}{2}\right)^2 \frac{2(2-\nu)}{\delta(1-\nu)} \pi \frac{c}{a} - \frac{4}{3} \pi a^2 c \cdot \Delta G_v$$

$\mu$  = shear modulus

$\nu = \frac{1}{3}$  for metal

Fig. 6.3(a)

(6.7)

$$\Delta G = 2\pi a^2 \gamma + \frac{16\pi}{3} a c^2 \left(\frac{s}{2}\right)^2 - \frac{4}{3} \pi a^2 c \cdot \Delta G_v$$

only shear w/o  $\epsilon$  due to the dilatation.  $\perp$  to the disc.

→ the most favorable nucl. path if habit plane = invariant plane (exact)

- min. free energy barrier to nucl.

differentiate wrt  $a$  &  $c$ .

$$(6.8) \Delta G^* = \frac{\pi l^2}{3} \left(\frac{\gamma^3}{\Delta G_v}\right)^4 \cdot \left(\frac{s}{2}\right)^4 \frac{1}{l^2} \text{ J/nucleus.}$$

↳ sensitive to  $\gamma$ ,  $\Delta G_v$ ,  $s$

$$(6.9)(6.10) c^* = \frac{2\gamma}{\Delta G_v} \quad a^* = \frac{16\gamma\mu \left(\frac{s}{2}\right)^2}{(\Delta G_v)^2}$$

$$\Delta G_v = 174 \text{ MJ/m}^3, \gamma = \sim 20 \text{ mJ. } c^*/a^* \approx 1/40, \text{ then } \Delta G^* = 20 \text{ eV}$$

→ high for thermal fluctuation to overcome (at 700K,  $kT = 0.06$  eV)

∴ M. nucl.: a heterogeneous process

exp. w/ small droplets (submicron to millimeter)

evidences ① Not all particles transformed even at 40°C w/  $\Delta T = 300^\circ\text{C}$

② # of nuclei less than that expected from homo.

( $\sim 10^4/\text{mm}^3$ )

Department of Inorganic Materials Engineering ③ # of nuclei ↑ as  $\Delta T \uparrow$  # of nuclei indep of  $\Delta T$  ④ simple.

⑤ no surface, not preferred nucl. sites. Dislocation is it! poly crystalline