

재료상변태

Phase Transformation of Materials

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박은수

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Contents for previous class

Substitutional Diffusion

- 1. Self diffusion in pure material**
- 2. Vacancy diffusion**
- 3. Diffusion in substitutional alloys**

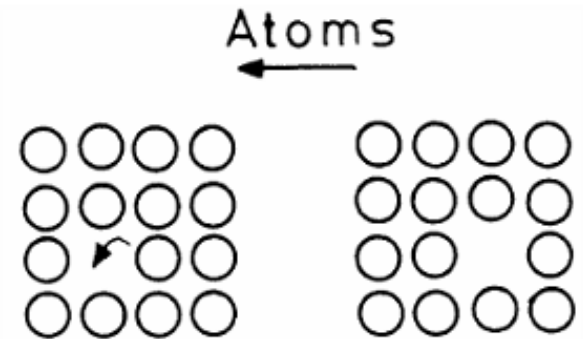
Contents for today's class

- **Diffusion in substitutional alloys**
 - 1) Interdiffusion
 - 2) Velocity of the movement of the lattice plane
 - 3) Fick's 2nd law for substitutional alloys
 - 4) Diffusion in dilute substitutional alloys
- **Atomic mobility:** Relationship between M_B and D_B

Substitutional diffusion

1. Self diffusion in pure material

Probability of vacancy x probability of jump



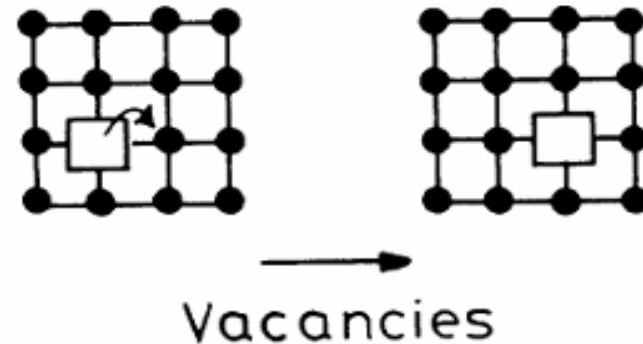
$$D_A = \frac{1}{6} \alpha^2 z v \exp\left(\frac{\Delta S_m + \Delta S_v}{R}\right) \exp\left(-\frac{(\Delta H_m + \Delta H_v)}{RT}\right)$$

$$= D_0 \exp\left(-\frac{Q_{SD}}{RT}\right)$$

2. Vacancy diffusion

$$D_v = \frac{1}{6} \alpha^2 \Gamma_v$$

$$= \frac{1}{6} \alpha^2 z v \exp(\Delta S_m / R) \exp(-\Delta H_m / RT)$$



Comparing D_v with the self-diffusion coefficient of A, D_A ,

$$D_v = D_A / X_v^e$$

3. Diffusion in substitutional alloys

1) Interdiffusion

Fluxes of A and B atoms across a given lattice plane

$$J_A = -D_A \frac{\partial C_A}{\partial x}, \quad J_B = -D_B \frac{\partial C_B}{\partial x}$$

Assume that

$$C_0 = C_A + C_B$$

C_0 : total number of atoms (A, B) per unit volume

$$\frac{\partial C_A}{\partial x} = -\frac{\partial C_B}{\partial x}$$

$$J_A = -D_A \frac{\partial C_A}{\partial x}$$

$$J_B = -D_B \frac{\partial C_B}{\partial x} = D_B \frac{\partial C_A}{\partial x}$$

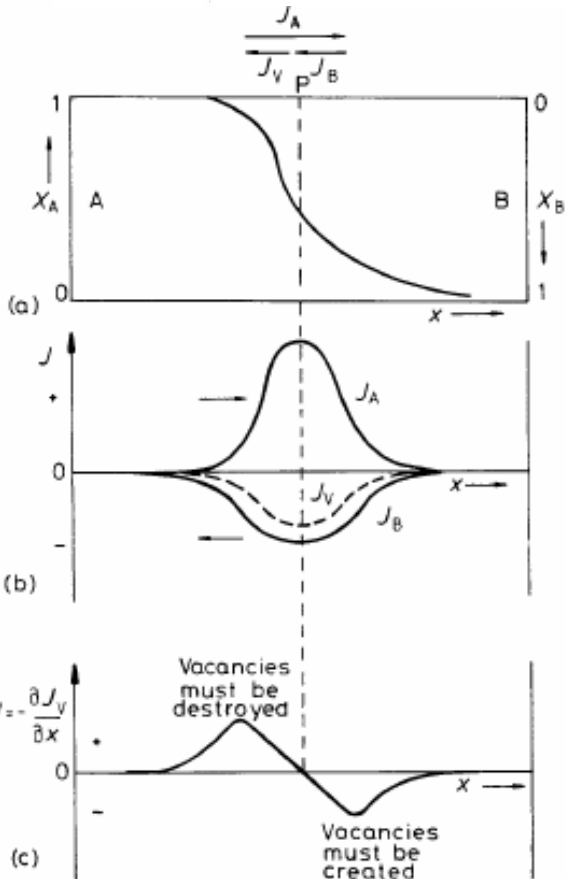
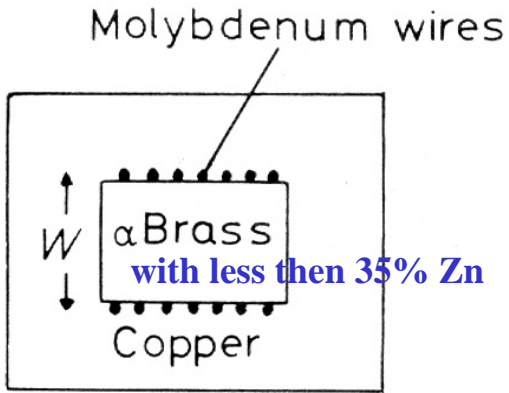
$$J_V = -J_A - J_B \quad (\text{a net flux of vacancies})$$

$$= (D_A - D_B) \frac{\partial C_A}{\partial x}$$

$$\because J_B = -D_B \frac{\partial C_B}{\partial x} = D_B \frac{\partial C_A}{\partial x}$$

$$\frac{\partial C_V}{\partial t} = -\frac{\partial J_V}{\partial x}$$

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$



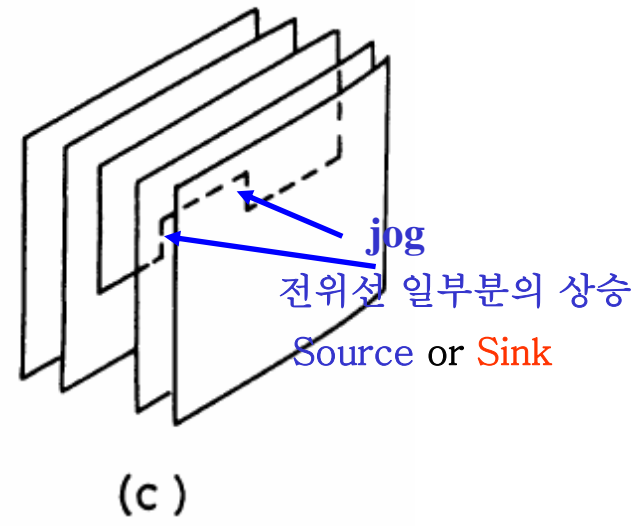
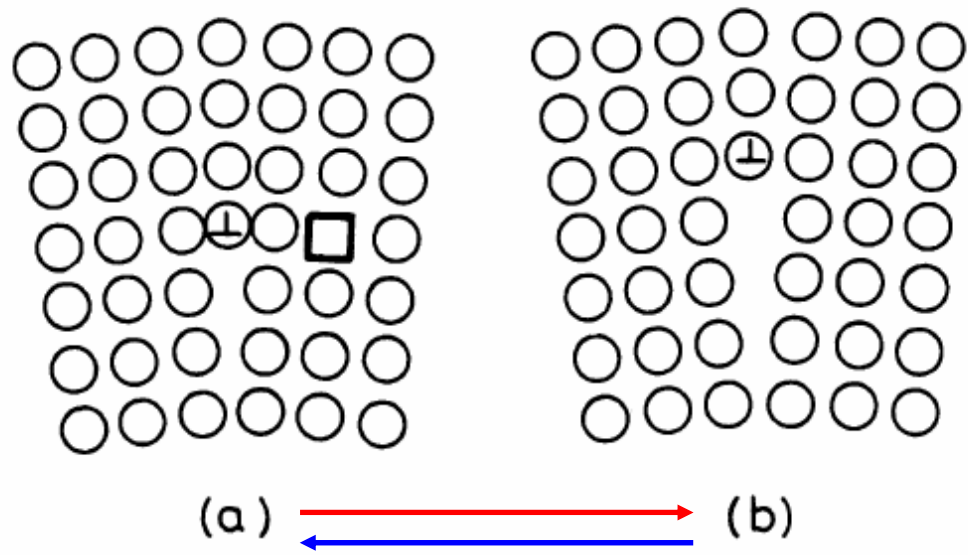
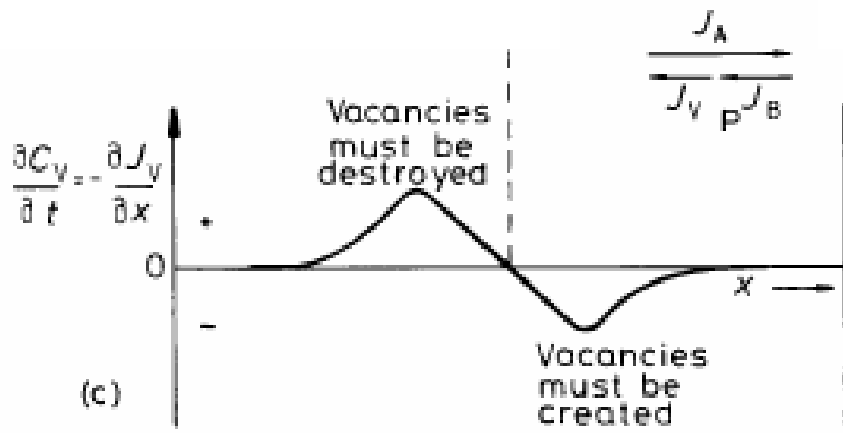


Fig. 2.17 (a) before, (b) after: a vacancy is absorbed at a jog on an edge dislocation [positive climb]. (b) before, (a) after: a vacancy is created by [negative limb] of an edge dislocation. (c) Perspective drawing of a jogged edge dislocation.

If dislocation climbs continue to occur, what would happen?

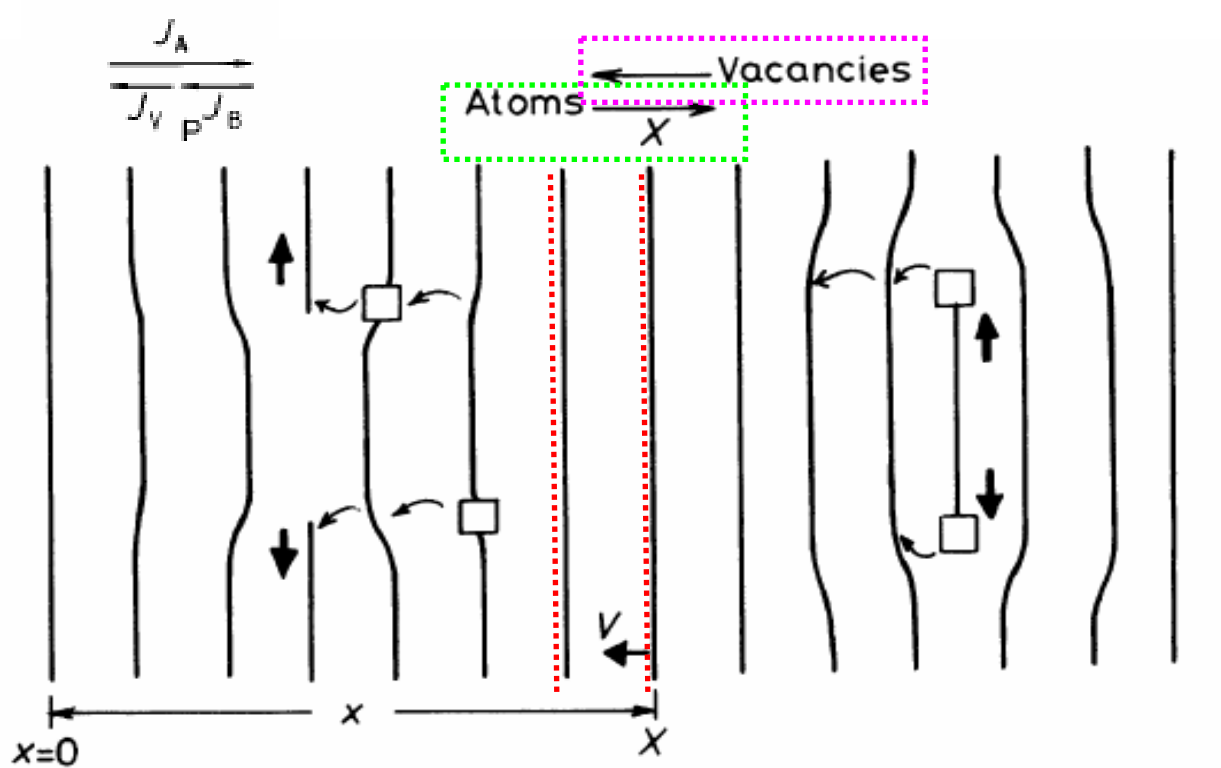


Fig. 2.18 A flux of vacancies causes the atomic planes to move through the specimen.

2) velocity of the movement of the lattice plane is related to the net flux of vacancies across the middle of the diffusion couple, such that

$$Av \cdot \delta t \cdot C_0 = J_v A \cdot \delta t \longrightarrow J_v = C_0 v$$

$$v = \frac{1}{C_0} (D_A - D_B) \frac{\partial C_A}{\partial x} = (D_A - D_B) \frac{\partial X_A}{\partial x}$$

X_A : mole fraction of A atoms, v : marker velocity

In practice, internal movement of lattice planes are usually not directly of interest. More practical questions concern **how long** homogenization of an alloy takes, or **how rapidly** the composition will change at a fixed position relative to the ends of a specimen.

➔ **Derivation of the Fick's 2nd law for substitutional alloys**

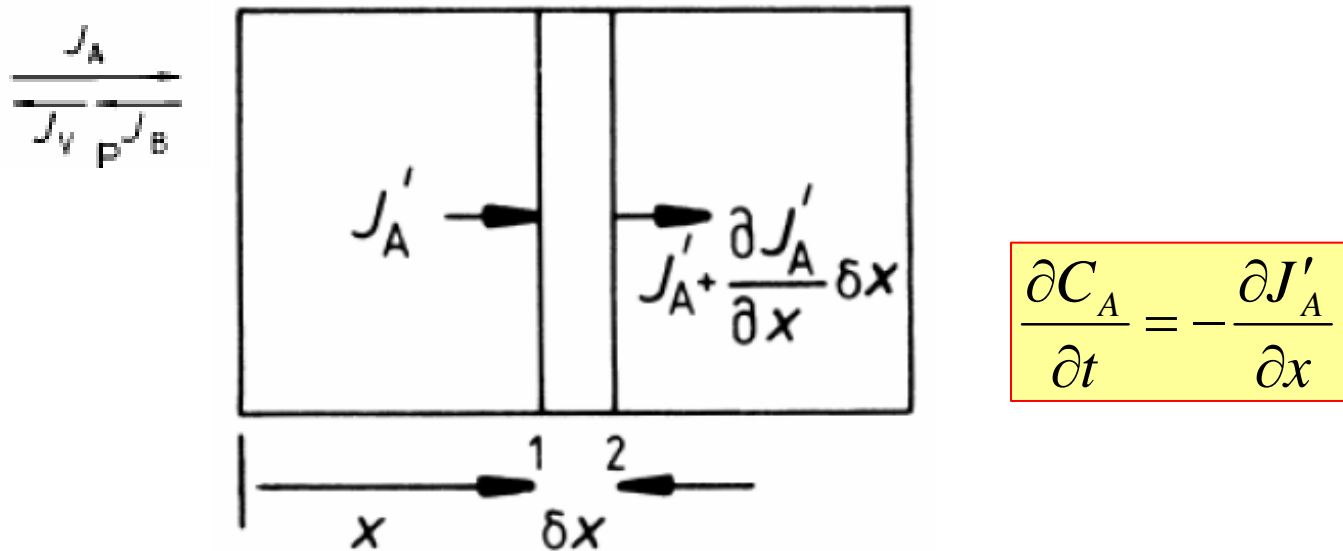


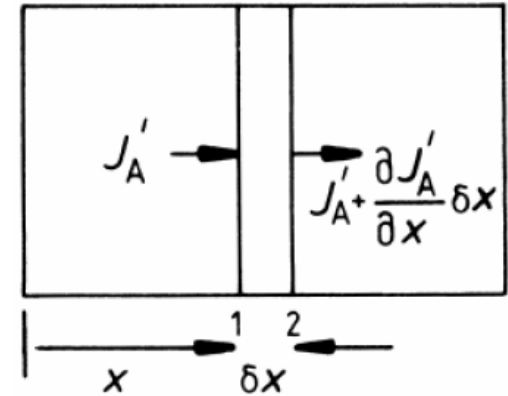
Fig. 2.19 Derivation of Fick's second law for interdiffusion. (See text for details.)

J'_A : total flux of A atoms across a stationary plane with respect to the specimen

3) Derivation of the Fick's 2nd law for substitutional alloys

J'_A : total flux of A atoms across a stationary plane with respect to the specimen

$$\frac{\partial C_A}{\partial t} = -\frac{\partial J'_A}{\partial x}$$



A Diffusive flux due to diffusion relative to the lattice \times
A flux due to the velocity of the lattice

$$J'_A = J_A + vC_A$$

$$= -D_A \frac{\partial C_A}{\partial x} + vC_A \quad v = (D_A - D_B) \frac{\partial X_A}{\partial x}$$

$$= (X_B D_A + X_A D_B) \frac{\partial C_A}{\partial x}$$

$$= -\tilde{D} \frac{\partial C_A}{\partial x}$$

$$\tilde{D} = X_B D_A + X_A D_B$$

$$J'_B = -\tilde{D} \frac{\partial C_B}{\partial x} = -\tilde{D} \frac{\partial C_A}{\partial x}$$

(Darken's equation, **interdiffusion coefficient**)

$$\frac{\partial C_A}{\partial t} = \frac{\partial}{\partial x} \left(\tilde{D} \frac{\partial C_A}{\partial x} \right)$$

Fick's 2nd law for substitutional alloy

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

$$\therefore J'_B = -J'_A$$

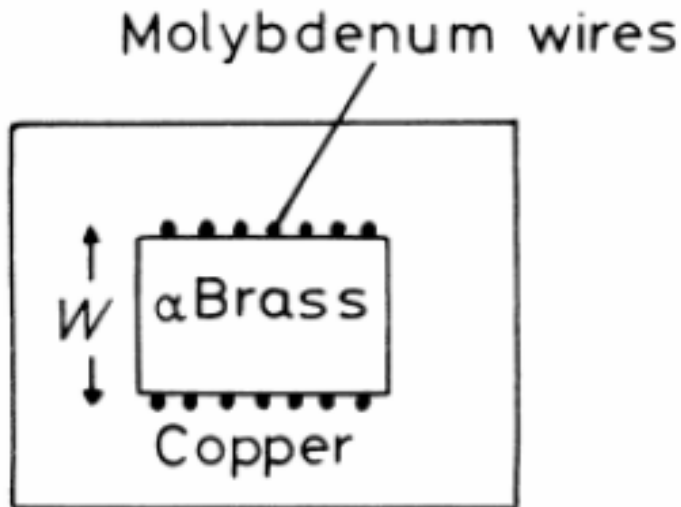
The interdiffusion coefficient (\tilde{D}) can be **experimentally measured** by determining the variation of X_A and X_B after annealing a diffusion couple.

How can we determine D_A and D_B ?

$$v = (D_A - D_B) \frac{\partial X_A}{\partial x}$$

When the velocity of a lattice (v) and interdiffusion coefficient (\tilde{D}) are known, D_A and D_B for the composition at the markers can be calculated.

How can we determine the velocity of a lattice (v)?



The displacement of wires during diffusion was first observed by Smigelskas and Kirkendall in 1947 and is usually known as the **Kirkendall effect**.

→ **Creation/destruction of vacancies is accomplished by dislocation climb.**

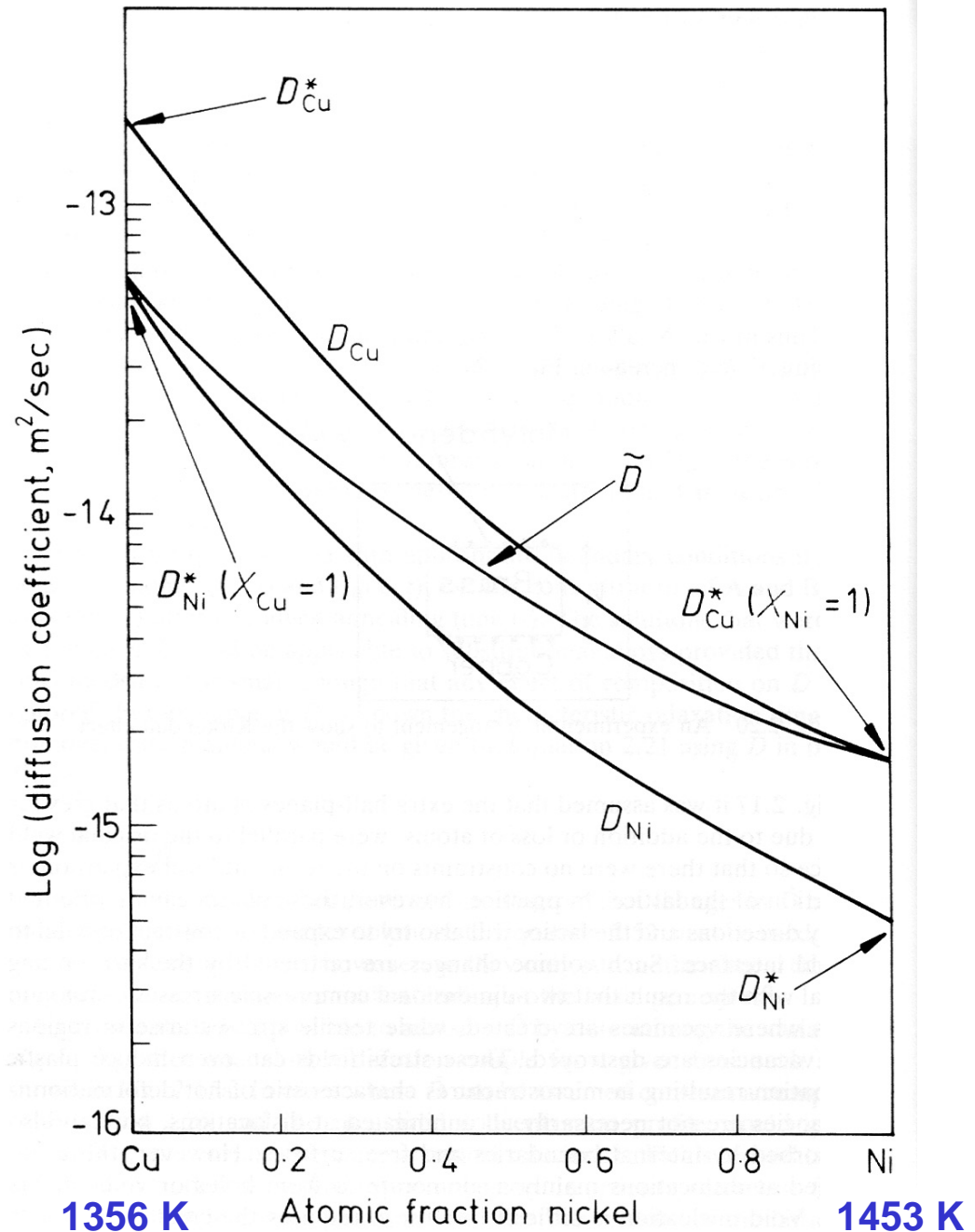
$$D_{Zn} > D_{Cu}$$

After annealing at a high temperature, it was found that the separation of the markers (w) had decreased.

The relationship between the various diffusion coefficients in the Cu-Ni system at 1273 K.

Atoms with the lower melting point possess a higher D .

D_{Cu} , D_{Ni} , (\tilde{D}) are all composition dependent, increasing as X_{Cu} increases.



- In concentrated alloys, the experimentally determined values of \tilde{D} , D_A and D_B are also found to show the same form of temperature dependence.

$$\tilde{D} = \tilde{D}_0 \exp(-Q / RT)$$

Variation of \tilde{D} with composition:

- *For a given crystal structure, \tilde{D} at T_m is roughly constant. Therefore if adding B to A decreases T_m , \tilde{D} will increase at a given temperature, and vice versa.*
- *For both interstitial and substitutional alloys, diffusion is more rapid in bcc (0.68) than in a close-packed lattice (FCC~0.74).*

Ferrite

Austenite

Ex) diffusion of carbon in Fe at 1183 K, $D_C^\alpha / D_C^\gamma \sim 100$

Self-diffusion coefficients for Fe at 1133 K, $D_{Fe}^\alpha / D_{Fe}^\gamma \sim 100$

 **More open and less lattice distortion**

4) Diffusion in dilute substitutional alloy

$$\tilde{D} = X_B D_A + X_A D_B$$

(interdiffusion coefficient)

For Dilute Substitutional Alloys

$$\text{if } X_A \approx 1, \quad \tilde{D} = D_B$$

In this case, D_B is called **'impurity diffusion coefficient'**.

- The reason for this is that the solute atoms can attract vacancies so that there is more than a random probability of finding a vacancy next to a solute atom with the result that they can diffuse faster than solvent.
- This is caused by the larger size or higher valency of the B atoms compared to those of A atoms. If the binding energy is very large, the vacancy will be unable to escape' from the solute atom.

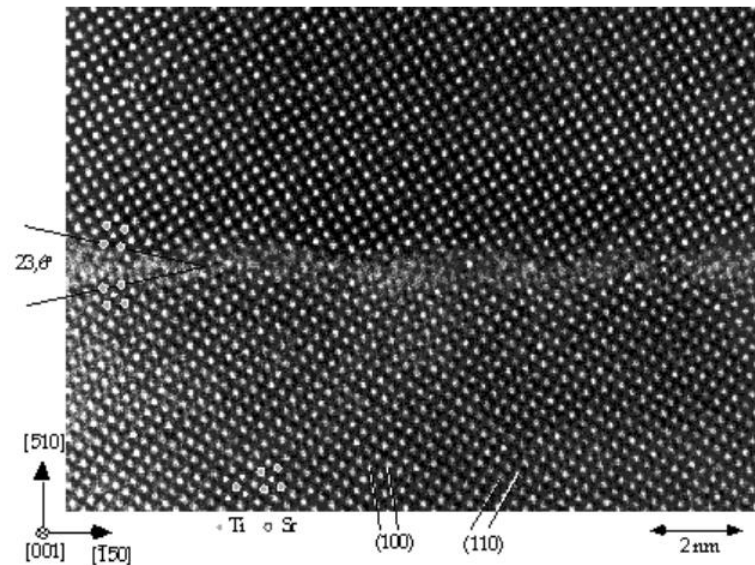
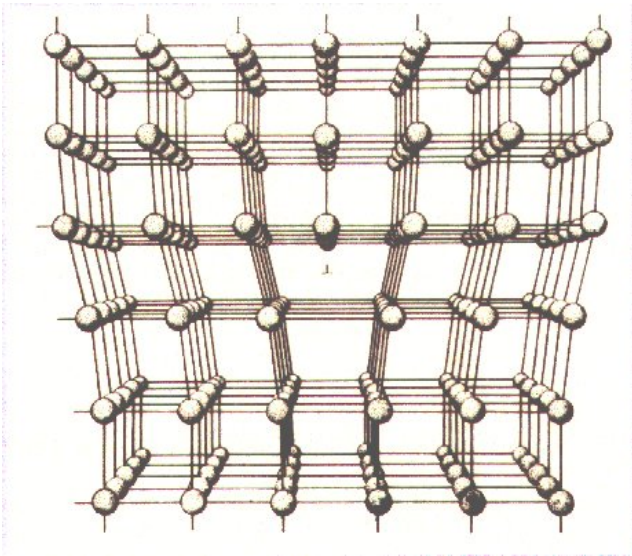
Atomic mobility

Ideal case, Diffusion eventually stops, that is equilibrium is reached.

➔ Fick's first law

Real materials contain lattice defects.

= more open structure = fast diffusion path.



➔ Some atoms can lower their free energies if they migrate to such defects and at 'equilibrium' their concentrations will be higher in the vicinity of the defect than in the matrix.

➔ Fick's law alone is insufficient to describe how the concentration will vary with distance and time.

Atomic mobility

- **Segregation** of atoms occur at crystal defects where the strain energy can be reduced. Segregation causes problems like **temper embrittlement** and **dynamic strain aging**. Fundamental aspects of phase transformation are also affected by segregation.
- The problem of atom migration can be solved by considering the **thermodynamic condition for equilibrium**; namely that the chemical potential of an atom must be the same everywhere. In general the flux of atoms at any point in the lattice is proportional to the **chemical potential gradient**: diffusion occurs down the slope of the chemical potential.

$$J_B = v_B c_B$$

A **diffusion flux** is a combined **quantity of a drift velocity** and **random jumping motion**.

$$v_B = -M_B \frac{\partial \mu_B}{\partial x}$$

$-\frac{\partial \mu_B}{\partial x}$: chemical force causing atom to migrate

M_B : mobility of B atoms



$$J_B = -M_B c_B \frac{\partial \mu_B}{\partial x}$$

How the mobility of an atom is related to its diffusion coefficient?

Relationship between M_B and D_B

$$J_B = -M_B C_B \frac{\partial \mu_B}{\partial x}$$

$$\begin{aligned} \therefore J_B &= -M_B \frac{X_B}{V_m} \frac{RT}{X_B} \left(1 + \frac{\partial \ln \gamma_B}{\partial \ln X_B}\right) \frac{\partial X_B}{\partial x} \\ &= -M_B RT F \frac{\partial C_B}{\partial x} = -D_B \frac{\partial C_B}{\partial x} \end{aligned}$$

$$D_B = M_B RT F$$

For ideal or dilute solutions,
near $X_B \approx 0$, $\gamma_B = \text{const.}$ with respect to X_B

$$\therefore F = 1 \quad D_B = M_B RT$$

$$\mu_B = G_B + RT \ln a_B = G_B + RT \ln \gamma_B X_B$$

$$\begin{aligned} \frac{\partial \mu_B}{\partial x} &= \frac{\partial}{\partial x} (G_B^0 + RT \ln \gamma_B X_B) \\ &= RT \left(\frac{\partial \ln \gamma_B}{\partial x} + \frac{\partial \ln X_B}{\partial x} \right) \\ &= RT \left(1 + \frac{\partial \ln \gamma_B}{\partial \ln X_B} \right) \frac{\partial \ln X_B}{\partial x} \end{aligned}$$

$$C_B = \frac{n_B}{V} = \frac{n_B}{(n_A + n_B)V_m} = \frac{X_B}{V_m}$$

$$F = \left(1 + \frac{d \ln \gamma_B}{d \ln X_B} \right)$$

For non-ideal concentrated solutions,
thermodynamic factor must be included.

→ Related to the curvature of the molar
free energy-composition curve.

The diffusive flux is also affected by the gradient of strain energy, $\partial E/\partial x$.

Ex) Effect of E depends on the position (x) relative to a dislocation.

$$\mu_B = G_B + RT \ln \gamma_B X_B + E \quad \rightarrow \quad J_B = -D_B \cdot \frac{\partial C_B}{\partial x} - \frac{D_B C_B}{RT} \cdot \frac{\partial E}{\partial x}$$

Atoms diffusing towards regions of high concentration can be found

- when diffusion occurs in the presence of an electric field or a temperature gradient.
- when the free energy curve has a negative curvature, which is known as spinodal decomposition.

2008 년 10 월

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5	6	7	8	9	10	11
12	13	14	15	16	17	18
19	20	21 연습시간	22	23 휴강	24	25 중간고사
26	27	28	29	30	31	