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# Computer aided ship design

## Part 2. Ship Motion & Wave Load

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**A**dvanced  
**S**hip  
**D**esign  
**A**utomation  
**L**aboratory

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# Part II. Ship Motion & Wave Load

## : 강의 개요

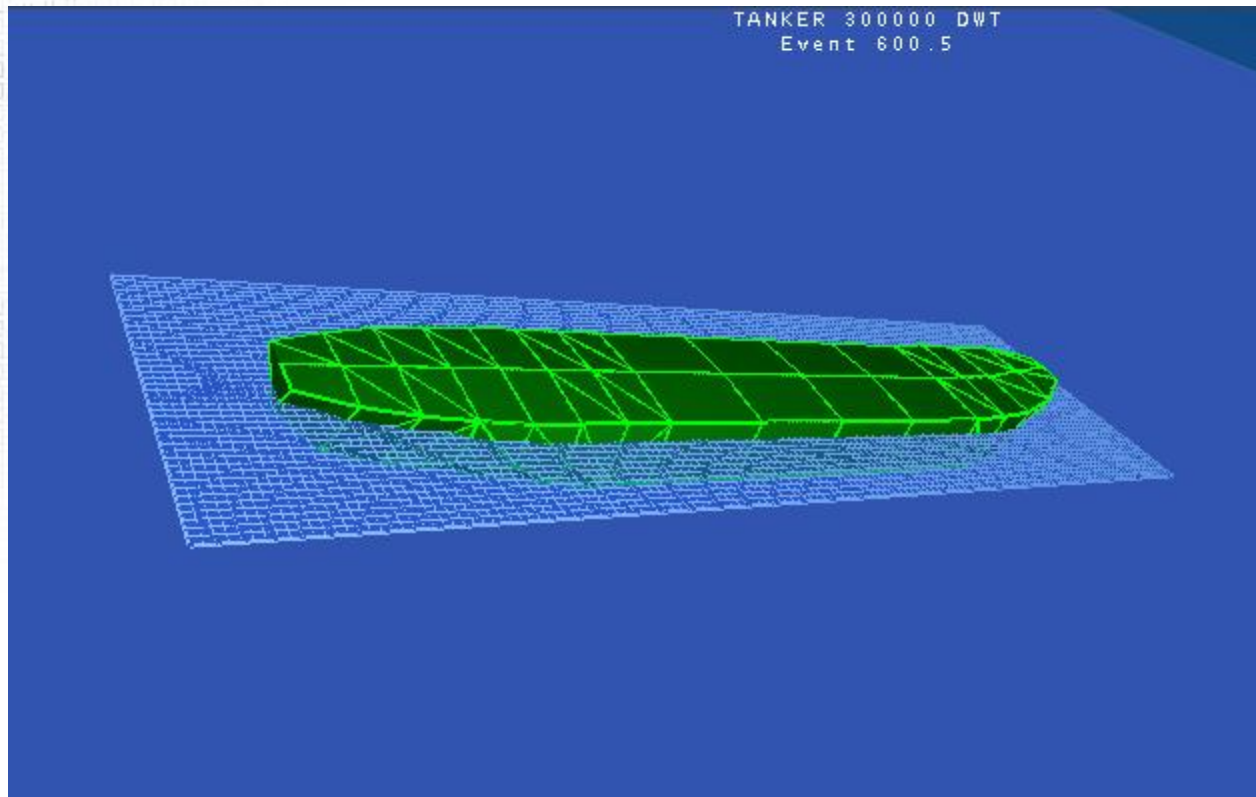
- 유체 역학(학부 2학년), 해양환경정보시스템(학부 3학년), 선박운항 제어론(학부 3학년), 선체구조설계시스템(학부 3학년)의 전공 과목과 공학 수학(학부 2학년), 기초구조정역학(학부 2학년), 동역학(학부 2학년)의 기초 교과목을 토대로 선박의 6자유도 운동 방정식의 유도 과정을 이해한다.  
(Ship Motion)
- 선박의 6자유도 운동 방정식의 계산 결과를 구조 설계에 활용 하는 방안에 대해 이해한다.  
(Wave Load)

# Part II. Ship Motion & Wave Load

## : Term project (I) – 6DOF Ship motion simulation

- 선박의 6자유도 운동 방정식을 수치적으로 계산하여 시간에 따른 선박의 운동을 구하고, 이를 화면에 가시화 한다.

ex) MOSES (Multi-Operational Structural Engineering Simulator)  
: Ultramarine 에서 개발(미국, 휴스턴)



Part I. 곡선, 곡면  
모델링의 결과 사용

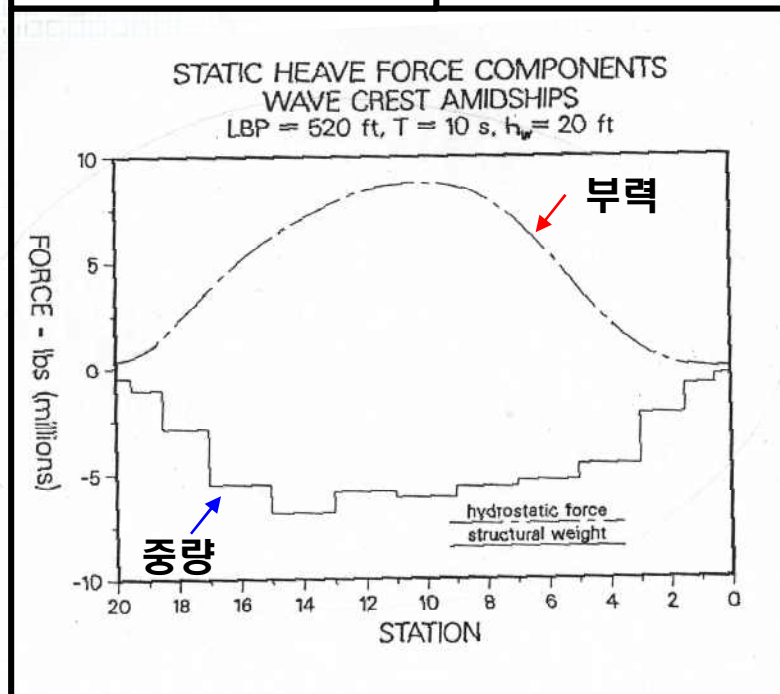
# Part II. Ship Motion & Wave Load

## : Term project (II) – Wave Load에 의한 VWBM<sup>1)</sup> 계산

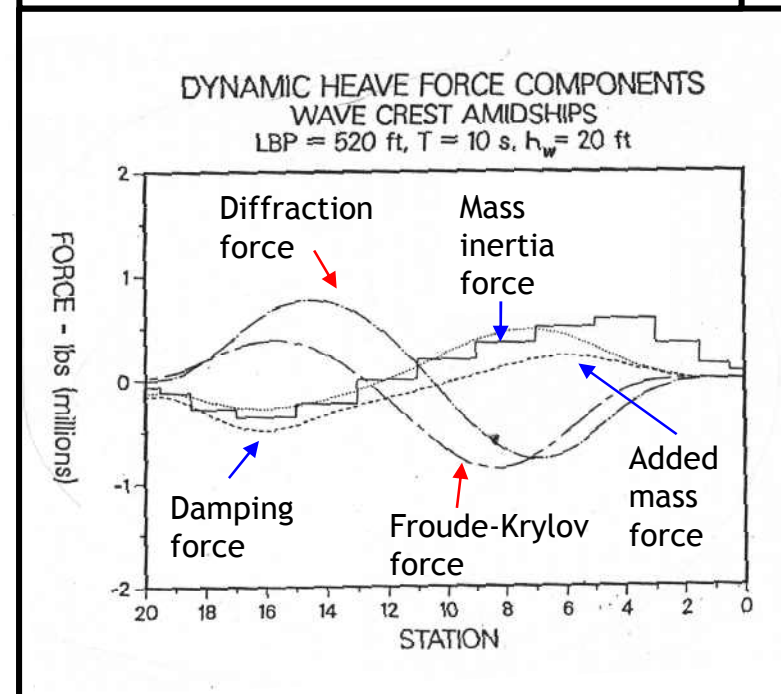
- 1) VWBM : Vertical Wave Bending Moment
- 2) SWBM : Still Water Bending Moment

- 6자유도 운동 방정식의 풀이 결과를 사용하여 선박에 작용하는 Wave Load를 계산하고, Shear force 및 Bending moment를 계산한다.

중량-부력 하중 곡선



Wave와 선박 운동에 의한 하중 곡선





# Part II. Ship Motion & Wave Load

## : 참고 자료

### ■ 참고 서적 및 자료

- 1) 小山健夫 著, 구종도 역, "선박과 해양구조물의 운동학", 연건문화사, 1997
- 2) Bhattacharyya, R. , "Dynamics of Marine Vehicles", John Wiley & Sons, 1978
- 3) Faltinsen, O.M. , "Sea loads on ships and offshore structures", Cambridge Univ. Press, 1998
- 4) Dean, R.G. , "Water wav mechanics for engineers and scientists", Prentice-Hall, Inc , 1984
- 5) Newman, J.N. , "Marine Hydrodynamics", The MIT Press, Cambridge, 1997
- 6) 이승건, "선박운동 조종론", 부산대학교 출판부, 2004
- 7) Journee, J.M.J. , Massie, W.W. , "Offshore Hydrodynamics", Delft University of Technology, 2001 (<http://www.shipmotions.nl/index.html>)
- 8) Journee, J.M.J. , Adegeest, L.J.M. , "Theoretical Manual of Strip Theory program“ Seaway for Windows”", Delft University of Technology, 2003 (<http://www.shipmotions.nl/index.html>)
- 9) Tommy Pedersen, "Wave Load Prediction - a Design Tool", PhD thesis, Department of naval architecture and offshore engineering, 2000
- 10) Cengel & Cimbala, "Fluid Mechanics", Mc Graw Hill, 2005
- 11) Erwin Kreyszig, "Advanced Engineering Mathematics", Wiley, 2005
- 12) Falnes, J. , "Ocean waves and oscillating systems", Cambridge Univ. Press, 2002



# Introduction.

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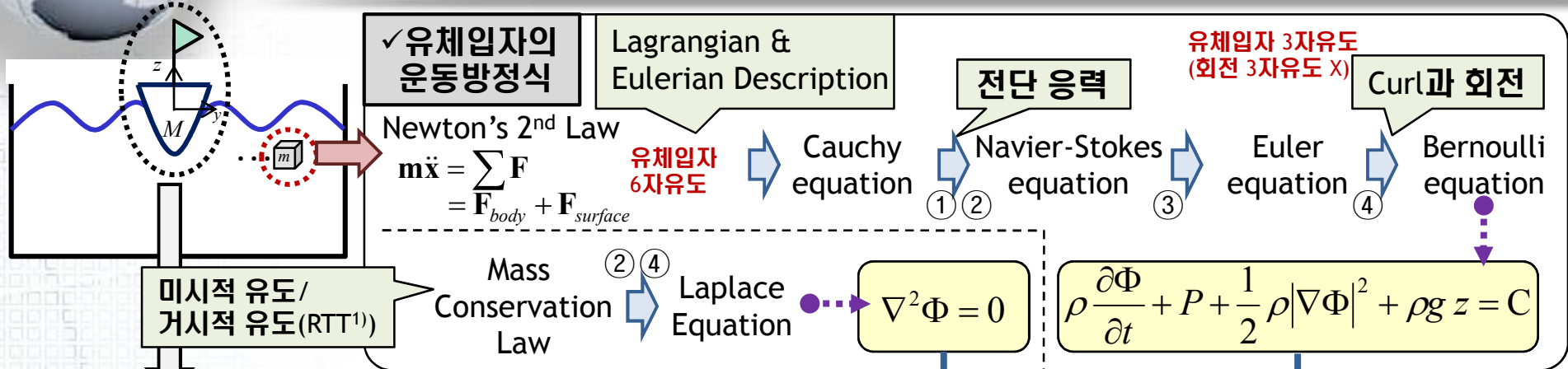
# Introduction

## : 유체 역학 및 선박 운동 방정식 개요

- 1) RTT : Reynold Transport Theorem
- 2) SWBM : Still Water Bending Moment
- 3) VWBM : Vertical Wave Bending Moment

✓ Assumption

- ① 뉴턴 유체 (Newtonian fluid)
- ② 비압축성 유동 (Incompressible flow)
- ③ 비점성 유동 (Inviscid flow)
- ④ 비회전 유동 (Irrotational flow)



### ✓ 선박의 6자유도 운동 방정식

① Coordinate system 정의 (Global & Body-fixed coordinate)

② Newton's 2<sup>nd</sup> Law

$$\begin{aligned}
 \mathbf{M}\ddot{\mathbf{x}} &= \sum \mathbf{F} = \mathbf{F}_{body} + \mathbf{F}_{surface} \\
 &= \mathbf{F}_{gravity} + \mathbf{F}_{Fluid} \\
 &= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{F.K} + \mathbf{F}_D + \mathbf{F}_R \\
 &= \mathbf{F}_{restoring} + \mathbf{F}_{exciting} - \mathbf{A}\ddot{\mathbf{x}} - \mathbf{B}\dot{\mathbf{x}}
 \end{aligned}$$

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} = \mathbf{F}_{restoring} + \mathbf{F}_{exciting}$$

cf) 선형화된 복원력 ( $\mathbf{F}_{restoring} = -\mathbf{C}\mathbf{x}$ )

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$$

z방향 성분

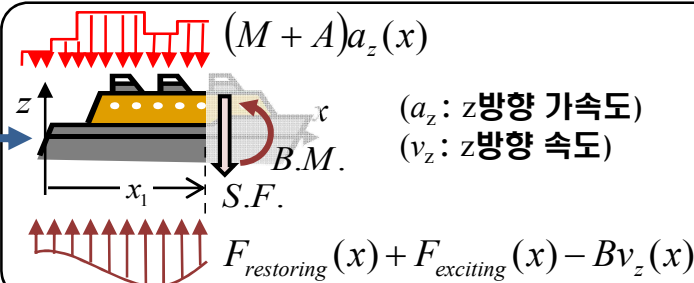
### ✓ 유체력 계산

$$\begin{aligned}
 \Phi_T &= \Phi_I \text{ (Incident wave potential)} \\
 &+ \Phi_D \text{ (Diffraction velocity potential)} \\
 &+ \Phi_R \text{ (Radiation velocity potential)}
 \end{aligned}$$

Linearization

$$P = -\rho g z - \rho \frac{\partial \Phi_T}{\partial t}$$

$$\mathbf{F}_{Fluid} = \iint_{S_B} P n dS = \mathbf{F}_{static} + \mathbf{F}_{F.K} + \mathbf{F}_D + \mathbf{F}_R$$



### ✓ Shear force (S.F.) 및 bending moment (B.M.)

Shear force (S.F.) 계산

↓ (적분) 7

Bending moment (B.M.)

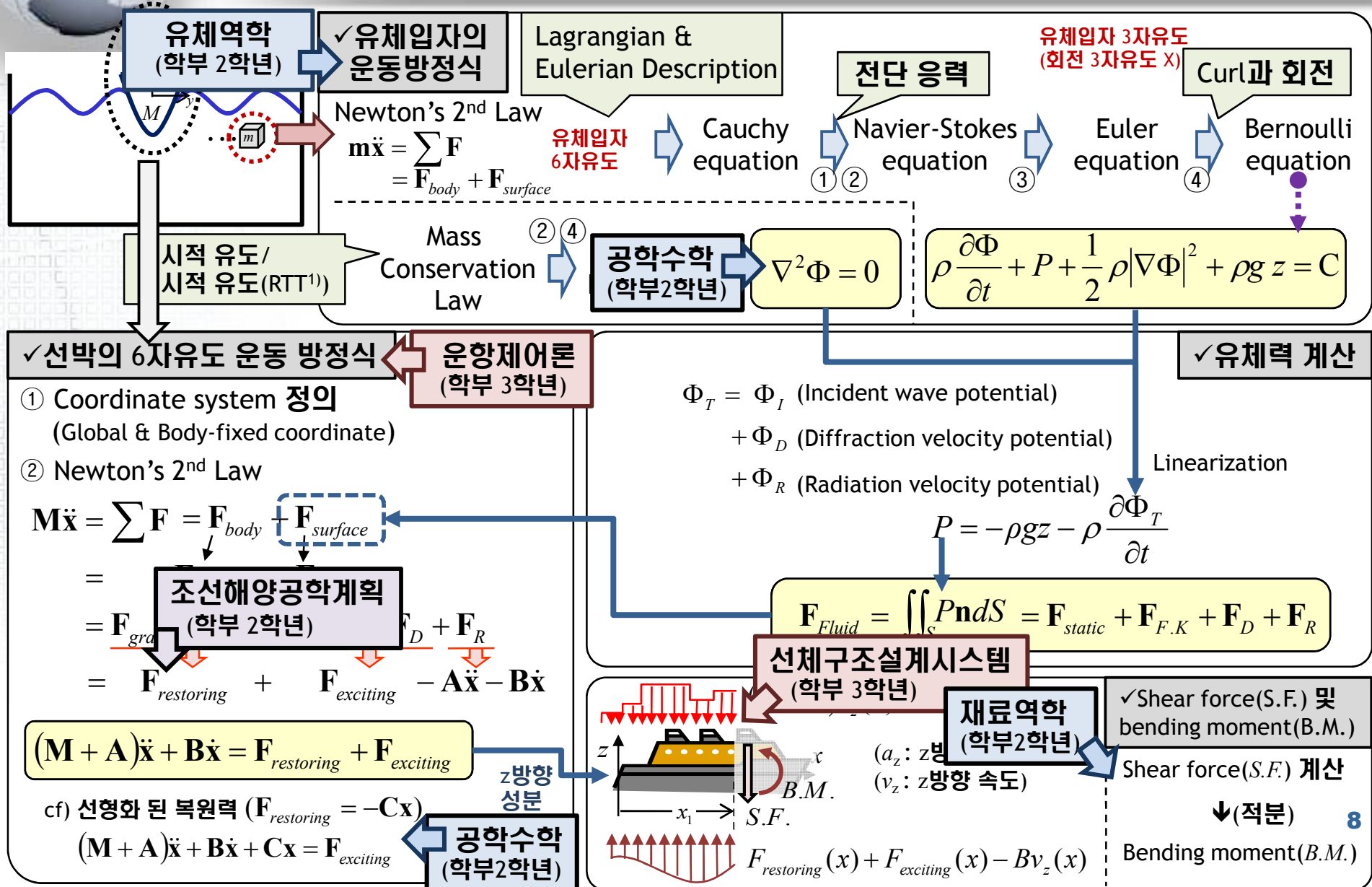
# Introduction

## : 학부 과목과의 연계

- 1) RTT : Reynold Transport Theorem
- 2) SWBM : Still Water Bending Moment
- 3) VWBM : Vertical Wave Bendingdng Moment

✓ Assumption

- ① 뉴턴 유체 (Newtonian fluid)
- ② 비압축성 유동 (Incompressible flow)
- ③ 비점성 유동 (Inviscid flow)
- ④ 비회전 유동 (Irrotational flow)



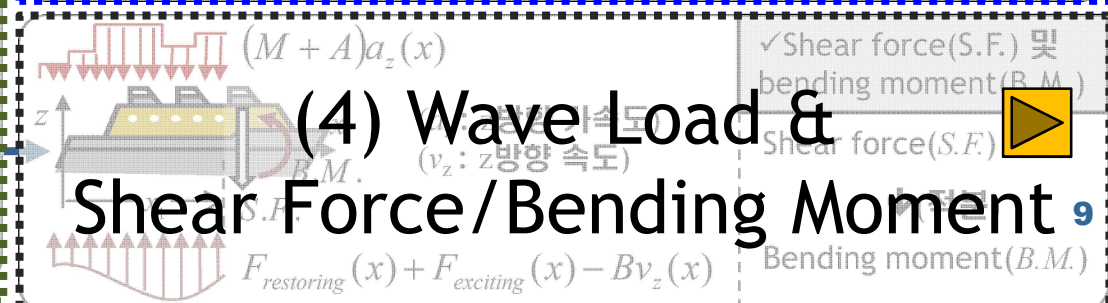
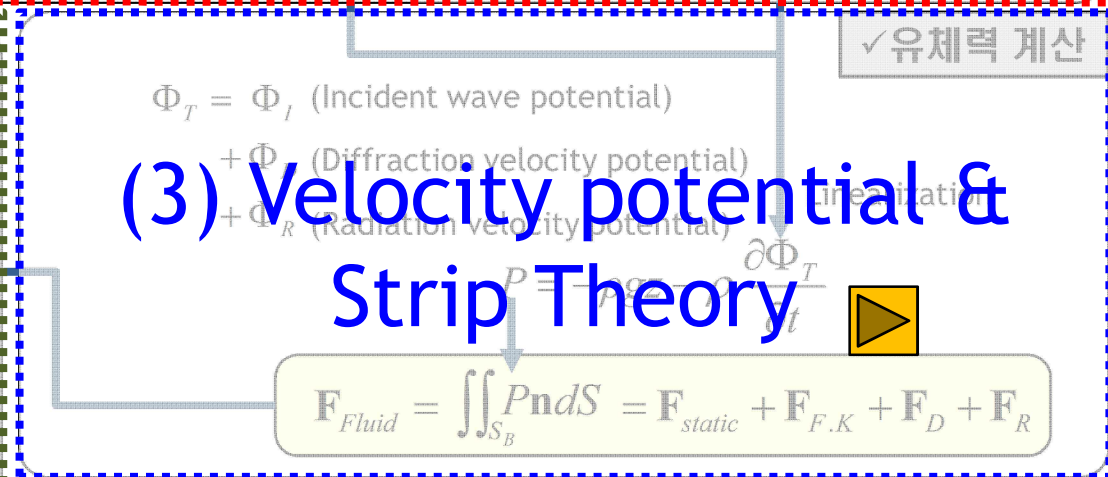
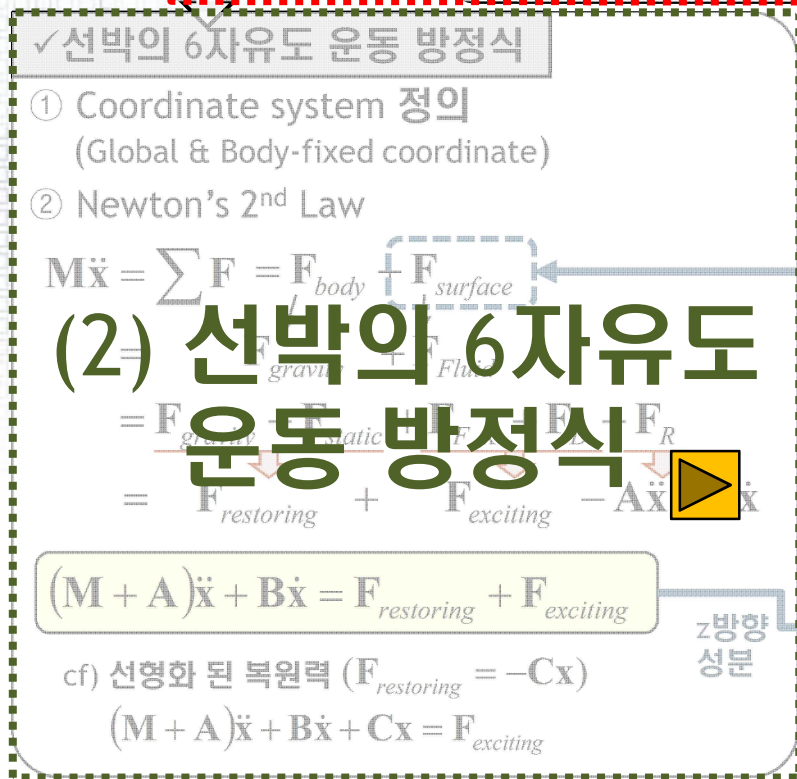
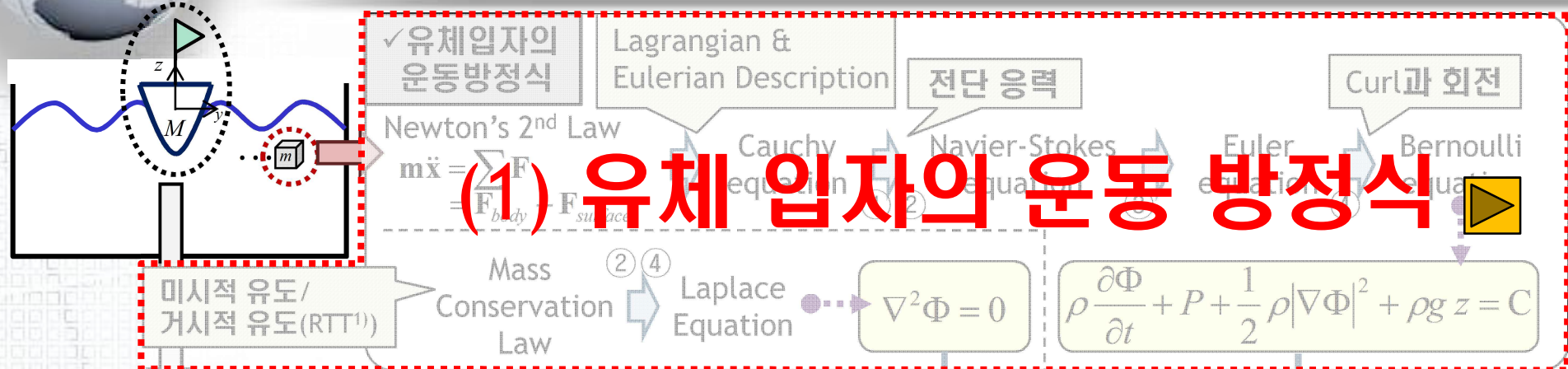


# Introduction

## : 배울 내용

- 1) RTT : Reynold Transport Theorem
- 2) SWBM : Still Water Bending Moment
- 3) VWBM : Vertical Wave Bending Moment

- ✓ Assumption
- ① 뉴턴 유체 (Newtonian fluid)
  - ② 비압축성 유동 (Incompressible flow)
  - ③ 비점성 유동 (Inviscid flow)
  - ④ 비회전 유동 (Irrotational flow)





# 유체 입자의 운동 방정식

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# 유체 입자의 운동 방정식 정리 (Cauchy eq. ~ Bernoulli eq.)

- 1) 전단응력이 전단변형율(의 시간변화율)에 비례하는 유체
- 2) 선형 변형과 등방 팽창에 의한 점성 계수  $\mu, \lambda$ 의 관계식을 정의함

Cauchy Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\sigma}$  ,  $(\mathbf{V} = [u, v, w]^T)$

↓ ① 뉴턴 유체<sup>1)</sup> (Newtonian fluid)  
② Stokes assumption<sup>2)</sup>

Navier-Stokes Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P + \mu \left( \frac{1}{3} \nabla (\nabla \cdot \mathbf{V}) + \nabla^2 \mathbf{V} \right)$   
(in general form)

$(\mu = 0)$  ↓ ③ 비점성 유동 (Inviscid flow)

Euler Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P$

$\rho = \rho(P)$  ↓ ④ barotropic flow

Euler Equation :  $\frac{\partial \mathbf{V}}{\partial t} + \nabla B = \mathbf{V} \times \boldsymbol{\omega}$  ,  $\left( B = \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} , q^2 = u^2 + v^2 + w^2 \right)$   
(Another form)

$\left( \frac{\partial \mathbf{V}}{\partial t} = 0 \right)$  ↓ ⑤ Steady flow

Bernoulli equation (case1)  $B = \text{Constant}$   
 $\left( \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} = C \right)$   
along streamlines and vortex lines

$\left( \mathbf{V} = \nabla \Phi , q^2 = |\nabla \Phi|^2 , \boldsymbol{\omega} = 0 \right)$  ↓ ⑥ Unsteady, irrotational flow

Bernoulli equation (case2)  $\frac{\partial \Phi}{\partial t} + \frac{1}{2} |\nabla \Phi|^2 + gy + \int \frac{dP}{\rho} = F(t)$

$(\rho = \text{constant})$  ↓ ⑦ Incompressible flow

Bernoulli equation (case3)  $\frac{\partial \Phi}{\partial t} + \frac{1}{2} |\nabla \Phi|^2 + gy + \frac{P}{\rho} = F(t)$

Continuity Equation (질량보존)  $\frac{\partial \rho}{\partial t} + \nabla \cdot \rho \mathbf{V} = 0$

↓ ⑦ incompressible flow  
 $\rho = \text{constant} \left( \frac{\partial \rho}{\partial t} = 0 \right)$

$\nabla \cdot \mathbf{V} = 0$

↓ ⑥ irrotational flow  
 $(\mathbf{V} = \nabla \Phi)$

Laplace Equation  $\nabla^2 \Phi = 0$

Newtonian fluid  
Stokes assumption  
Inviscid flow  
Unsteady flow  
Irrotational flow  
Incompressible flow





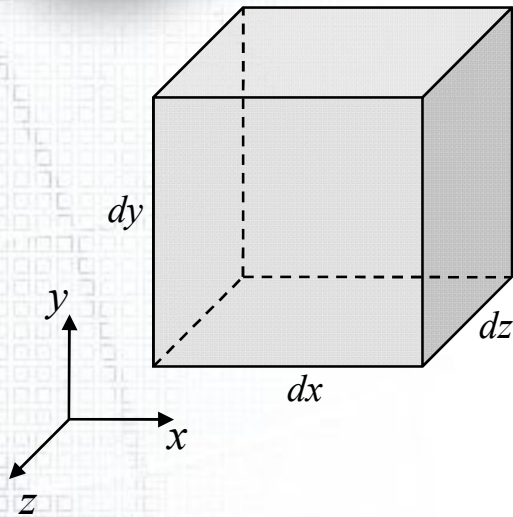
# Cauchy Equation의 유도

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# Cauchy Equation<sup>1),2)</sup> 유도

미소 유체 요소



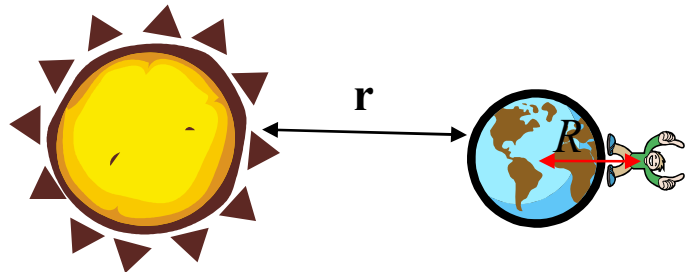
✓ 미소 유체 요소가 받는 힘 (Newton's 2<sup>nd</sup> Law)

질량 X 가속도 → Lagrangian & **Eulerian description**

$$m \frac{dV}{dt} = \sum \mathbf{F} = \mathbf{F}_{Body} + \mathbf{F}_{Surface} \quad (\text{체적력} + \text{표면력})$$

Action at a distance  
without physical Contact<sup>2)</sup>

ex) gravitational force, magnetic, electrostatic,...



Surface forces are exerted on an area element  
by the surroundings through direct contact<sup>2)</sup>



유체의 종류에 따라 다르지만, 저항하는 힘을 느낌

(Q) 공기중에서도 표면력이 있는가?

(A) 있다.

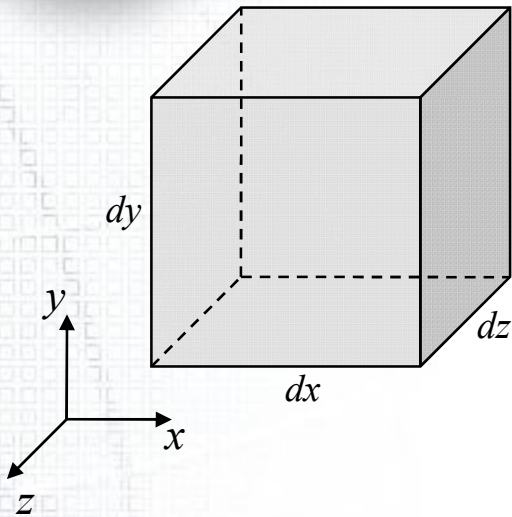
(밀도가 작아서 작용하는 힘을 못 느낄 뿐)

# Cauchy Equation<sup>1),2)</sup> 유도

## - Body Force

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93

미소 유체 요소



✓ Gravitational force is a body force on Earth.

$$\mathbf{F}_{Body} = \mathbf{F}_{Gravity} = m\mathbf{g} = \rho g dx dy dz \quad (m = \rho dV = \rho dx dy dz)$$

$$\begin{cases} F_{x,Body} = F_{x,Gravity} = mg_x = \rho g_x dx dy dz = 0 \\ F_{y,Body} = F_{y,Gravity} = mg_y = \rho g_y dx dy dz = -mg \\ F_{z,Body} = F_{z,Gravity} = mg_z = \rho g_z dx dy dz = 0 \end{cases}$$

지구상에서의 체적력

$$g = 9.81 \quad : \text{중력가속도}$$

(cf) 중력에 의한 포텐셜 에너지 :  $\Pi = gy$

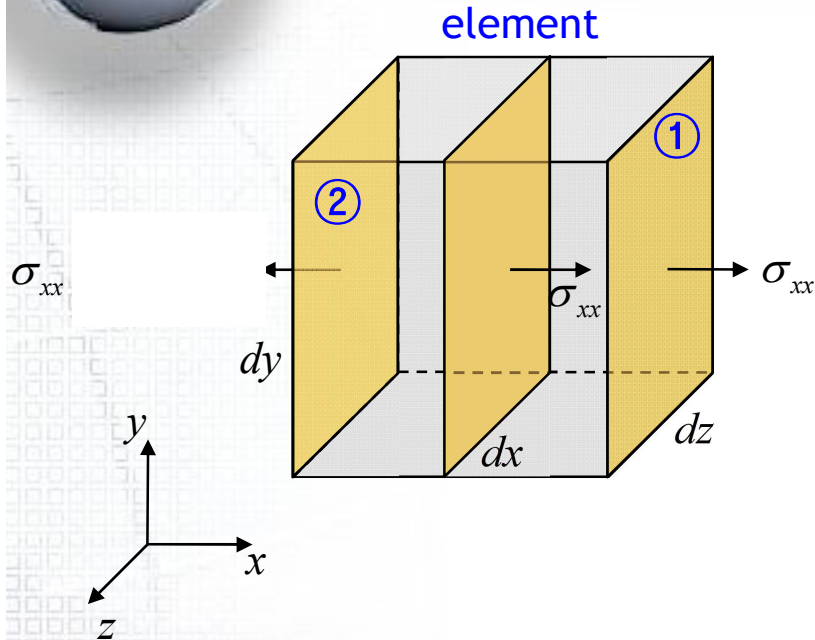
중력력은 보존력(conservative force)이므로,  
 중력에 의한 포텐셜 에너지의 Gradient로부터 구할 수 있음

$$\begin{aligned} \mathbf{g} &= -\nabla \Pi = -\nabla(gy) \\ &= -\left[ \mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y} + \mathbf{k} \frac{\partial}{\partial z} \right] (gy) = -g\mathbf{j} \end{aligned}$$

# Cauchy Equation<sup>1),2)</sup> 유도

## - Stress at a point

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.30-32, 88-93



✓ 한 점에서의 응력 (Stress at a point)

➔ 어느 면에 어느 방향으로 힘이 작용하는지 정의해야 함

Orientation  
of the surface

Direction  
of the force

➔ 2개의 방향 필요

$$\sigma_{xx}$$

➔ point를 element로 간주하여 면에 작용하는 응력을 정의한 후,  $dx, dy, dz$ 가 거의 0에 가깝다 ( $dx, dy, dz \rightarrow 0$ )는 것으로부터 응력을 정의

➔  $dx \rightarrow 0$ 일 때, 두 면이 일치하므로, 응력의 크기는 같고, 방향은 반대로 정의됨

# Cauchy Equation<sup>1),2)</sup> 유도

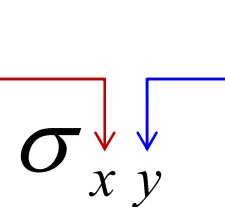
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Orientation of the surface      Direction of the force      ➔ 2개의 방향 필요



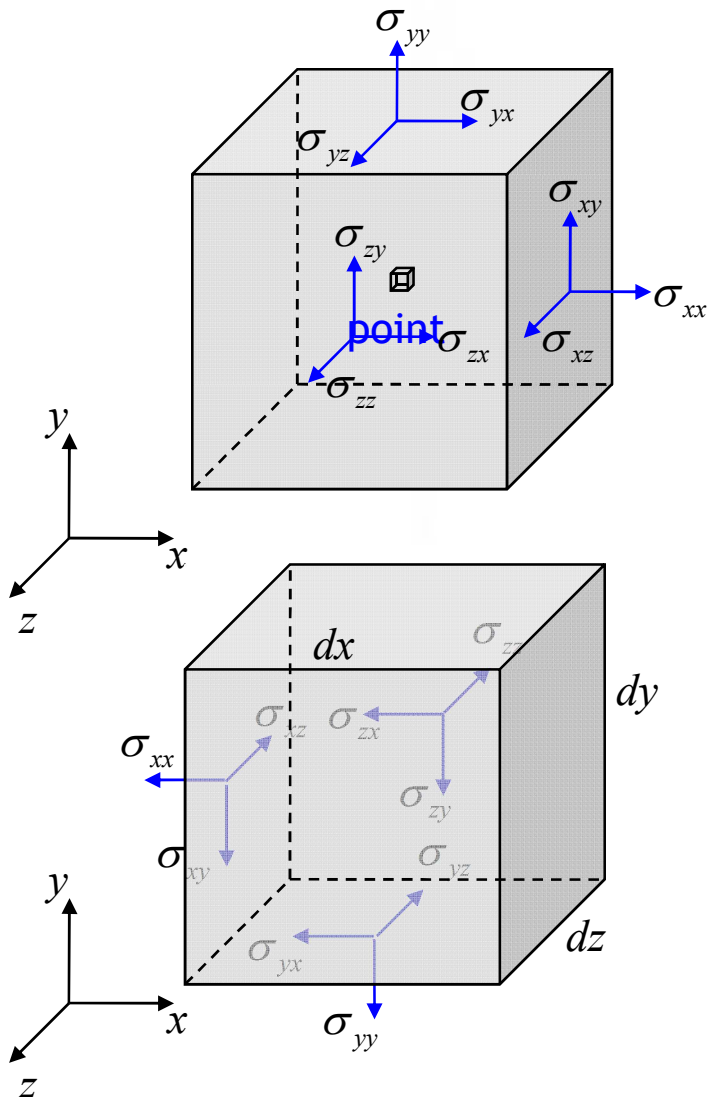
➔ 한 면에 수직인 성분과 2개의 접선 방향 성분으로 나타남 (3개의 응력 성분이 존재함)

➔ 육면체의 세 면에 대해 총 9개의 응력 성분을 정의

$$(\sigma_{xx}, \sigma_{xy}, \sigma_{xz}, \sigma_{yx}, \sigma_{yy}, \sigma_{yz}, \sigma_{zx}, \sigma_{zy}, \sigma_{zz})$$

➔ 3X3 Matrix 형태로 정리함

응력 텐서 : (stress tensor)  $\sigma = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix}$

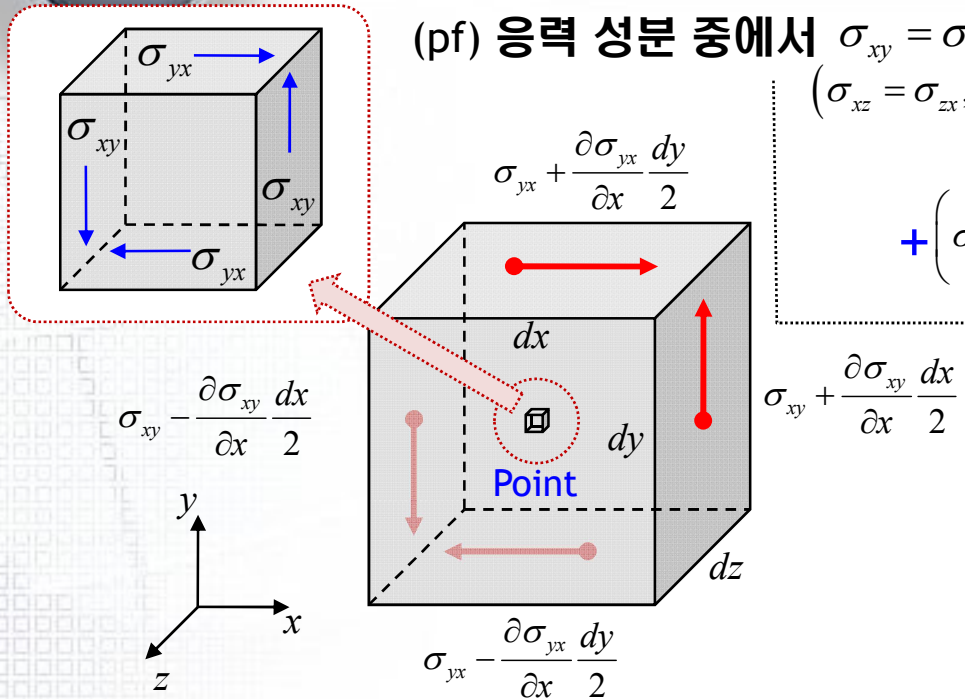




# Cauchy Equation<sup>1),2)</sup> 유도

## - Stress at a point

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 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93



(pf) 응력 성분 중에서  $\sigma_{xy} = \sigma_{yx}$  이다.  
 $(\sigma_{xz} = \sigma_{zx}, \sigma_{yz} = \sigma_{zy})$

$$\begin{aligned}
 & - \left( \sigma_{yx} + \frac{\partial \sigma_{yx}}{\partial x} \frac{dy}{2} \right) dx dz \frac{dy}{2} \\
 & + \left( \sigma_{xy} - \frac{\partial \sigma_{xy}}{\partial x} \frac{dx}{2} \right) dy dz \frac{dx}{2} \\
 & - \left( \sigma_{yx} - \frac{\partial \sigma_{yx}}{\partial x} \frac{dy}{2} \right) dx dz \frac{dy}{2} \\
 & + \left( \sigma_{xy} + \frac{\partial \sigma_{xy}}{\partial x} \frac{dx}{2} \right) dy dz \frac{dx}{2}
 \end{aligned}$$

rotational equilibrium of the element

$$\begin{aligned}
 I_{zz} \dot{\omega} &= M_z \\
 \frac{\rho dx dy dz}{12} (dx^2 + dy^2) \dot{\omega} &= (\sigma_{xy} - \sigma_{yx}) dx dy dz \\
 \frac{\rho}{12} (dx^2 + dy^2) \dot{\omega} &= \sigma_{xy} - \sigma_{yx} = 0
 \end{aligned}$$

point이므로  $dx \rightarrow 0, dy \rightarrow 0$

$$\therefore \sigma_{xy} = \sigma_{yx} \quad (\sigma_{xz} = \sigma_{zx}, \sigma_{yz} = \sigma_{zy})$$

※ Total moment

$$\begin{aligned}
 M_z &= \left( \sigma_{xy} + \frac{\partial \sigma_{xy}}{\partial x} \frac{dx}{2} \right) dy dz \frac{dx}{2} - \left( \sigma_{yx} + \frac{\partial \sigma_{yx}}{\partial x} \frac{dy}{2} \right) dx dz \frac{dy}{2} \\
 &+ \left( \sigma_{xy} - \frac{\partial \sigma_{xy}}{\partial x} \frac{dx}{2} \right) dy dz \frac{dx}{2} - \left( \sigma_{yx} - \frac{\partial \sigma_{yx}}{\partial x} \frac{dy}{2} \right) dx dz \frac{dy}{2} \\
 &= \sigma_{xy} dx dy dz - \sigma_{yx} dx dy dz = (\sigma_{xy} - \sigma_{yx}) dx dy dz
 \end{aligned}$$

※ Mass moment of inertia

$$I_{zz} = \frac{m}{12} (dx^2 + dy^2) = \frac{\rho dx dy dz}{12} (dx^2 + dy^2)$$

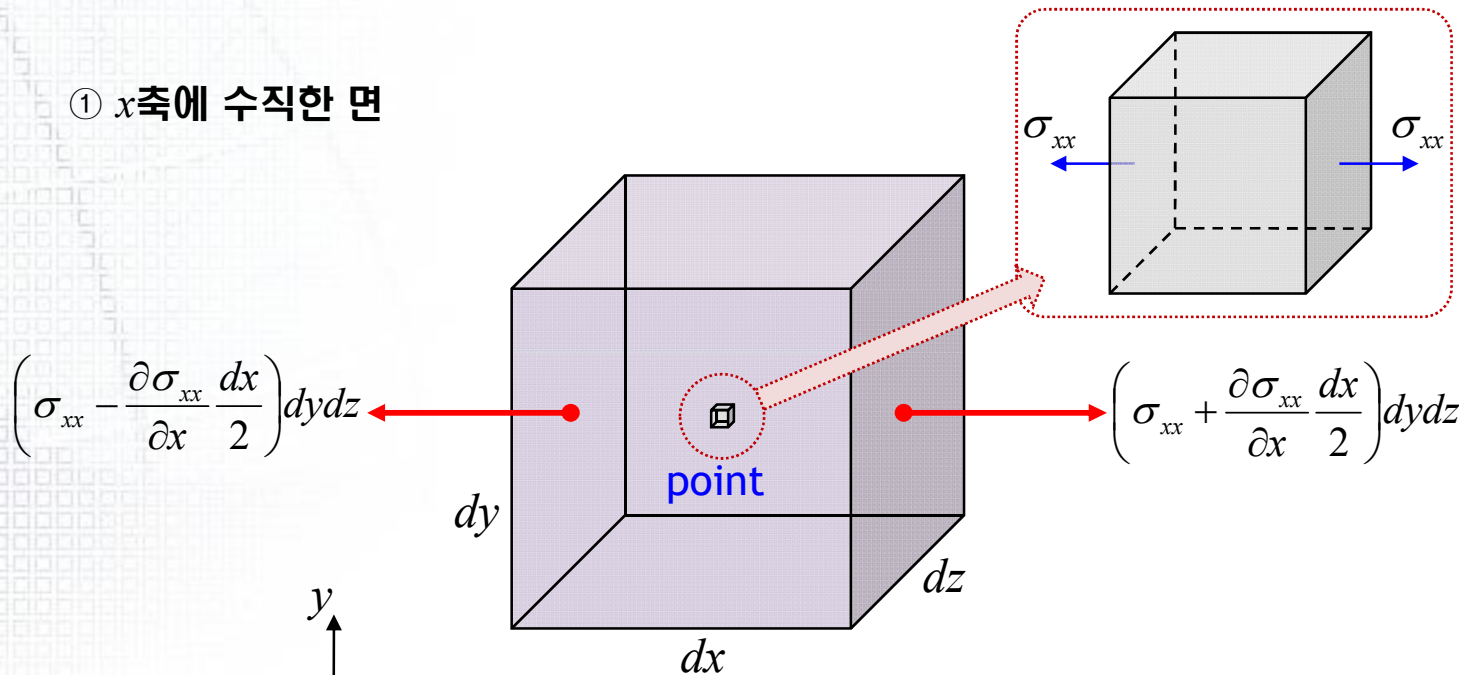
# Cauchy Equation<sup>1),2)</sup> 유도

## - Surface force on Element

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93

표면력 중  $x$ 축 방향 힘 : 중심면으로부터  $\pm \frac{dx}{2}$  만큼 떨어진 곳에서  $x$ 축 방향으로 작용하는 힘

①  $x$ 축에 수직한 면



$$F_{xx} = \left(\sigma_{xx} + \frac{\partial \sigma_{xx}}{\partial x} \frac{dx}{2}\right) dy dz - \left(\sigma_{xx} - \frac{\partial \sigma_{xx}}{\partial x} \frac{dx}{2}\right) dy dz$$

$$= \frac{\partial \sigma_{xx}}{\partial x} dx dy dz$$

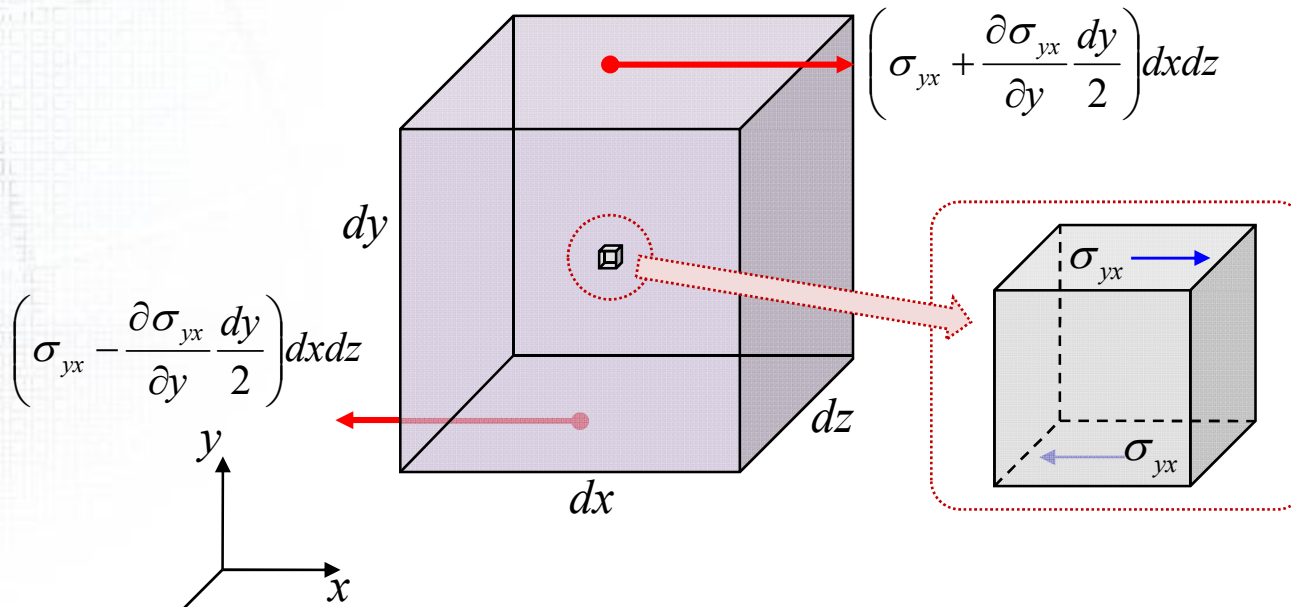
# Cauchy Equation<sup>1),2)</sup> 유도

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- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93

표면력 중  $x$ 축 방향 힘 : 중심면으로부터  $\pm \frac{dy}{2}$  만큼 떨어진 곳에서  $x$ 축 방향으로 작용하는 힘

②  $y$ 축에 수직한 면



$$F_{yx} = \left( \sigma_{yx} + \frac{\partial \sigma_{yx}}{\partial y} \frac{dy}{2} \right) dx dz - \left( \sigma_{yx} - \frac{\partial \sigma_{yx}}{\partial y} \frac{dy}{2} \right) dx dz$$

$$= \frac{\partial \sigma_{yx}}{\partial y} dx dy dz$$

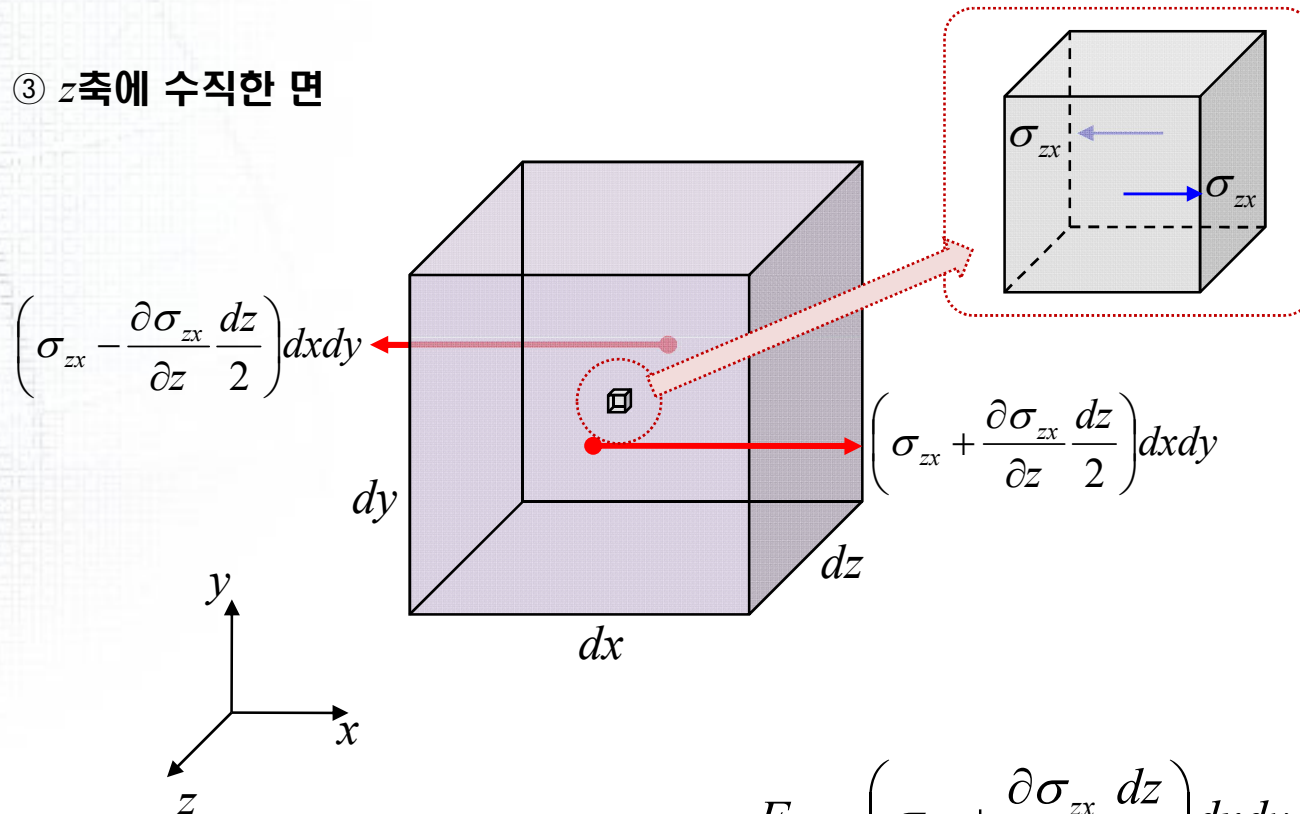
# Cauchy Equation<sup>1),2)</sup> 유도

## - Surface force on Element

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93

표면력 중  $x$ 축 방향 힘 : 중심면으로부터  $\pm \frac{dz}{2}$  만큼 떨어진 곳에서  $x$ 축 방향으로 작용하는 힘

③  $z$ 축에 수직인 면



$$F_{zx} = \left(\sigma_{zx} + \frac{\partial \sigma_{zx}}{\partial z} \frac{dz}{2}\right) dxdy - \left(\sigma_{zx} - \frac{\partial \sigma_{zx}}{\partial z} \frac{dz}{2}\right) dxdy$$

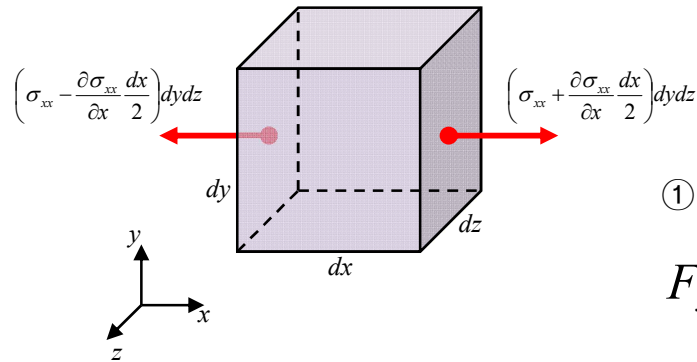
$$= \frac{\partial \sigma_{zx}}{\partial z} dxdydz$$



# Cauchy Equation<sup>1),2)</sup> 유도

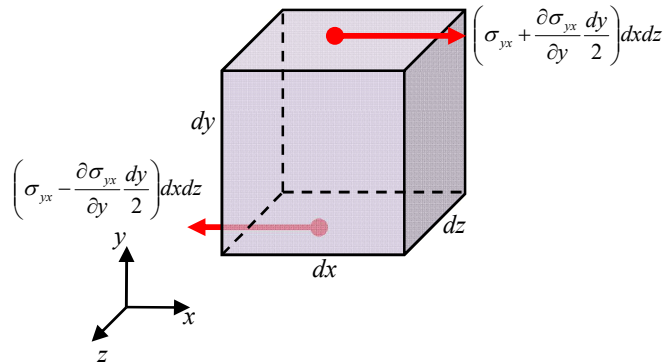
## - Surface force on Element

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93



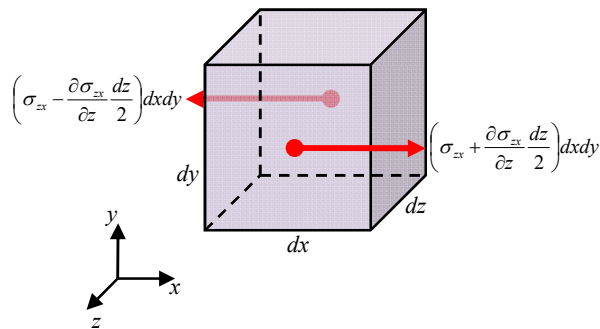
① x축에 수직한 면

$$F_{xx} = \frac{\partial \sigma_{xx}}{\partial x} dx dy dz$$



② y축에 수직한 면

$$F_{yx} = \frac{\partial \sigma_{yx}}{\partial y} dx dy dz$$



③ z축에 수직한 면

$$F_{zx} = \frac{\partial \sigma_{zx}}{\partial z} dx dy dz$$

✓ 표면력 중 x축 방향 힘

$$F_{x,surface} = F_{xx} + F_{yx} + F_{zx}$$

$$= \left( \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} \right) dx dy dz$$



✓ 표면력 중 y축 방향 힘

$$F_{y,surface} = F_{xy} + F_{yy} + F_{zy}$$

$$= \left( \frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z} \right) dx dy dz$$

✓ 표면력 중 z축 방향 힘

$$F_{z,surface} = F_{xz} + F_{yz} + F_{zz}$$

$$= \left( \frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right) dx dy dz$$

# Cauchy Equation<sup>1),2)</sup> 유도

## - Surface force on Element

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.88-93

### (3) 미소 유체 요소에 작용하는 표면력

#### ① 표면력 중 x축 방향 힘

$$F_{x, Surface} = \left( \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} \right) dx dy dz$$

#### ② 표면력 중 y축 방향 힘

$$F_{y, Surface} = \left( \frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z} \right) dx dy dz$$

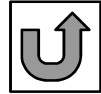
#### ③ 표면력 중 z축 방향 힘

$$F_{z, Surface} = \left( \frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right) dx dy dz$$

↓ 벡터 형태로 나타내면

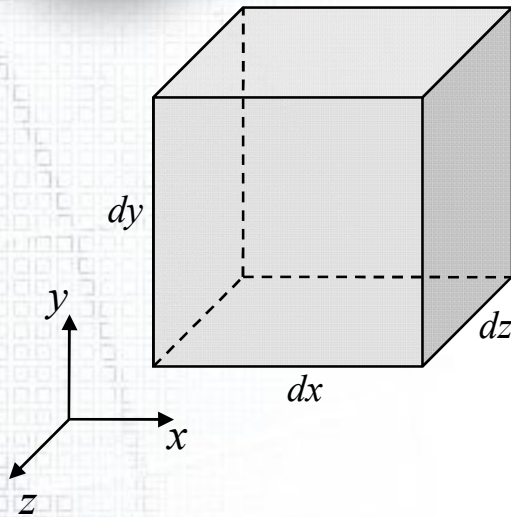
$$\begin{aligned} \mathbf{F}_{Surface} &= \hat{i} F_{x, Surface} + \hat{j} F_{y, Surface} + \hat{k} F_{z, Surface} \\ &= \left[ \left( \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} \right) \hat{i} + \left( \frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z} \right) \hat{j} + \left( \frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z} \right) \hat{k} \right] dx dy dz \\ &= \left[ \frac{\partial}{\partial x} (\sigma_{xx} \hat{i} + \sigma_{xy} \hat{j} + \sigma_{xz} \hat{k}) + \frac{\partial}{\partial y} (\sigma_{yx} \hat{i} + \sigma_{yy} \hat{j} + \sigma_{yz} \hat{k}) + \frac{\partial}{\partial z} (\sigma_{zx} \hat{i} + \sigma_{zy} \hat{j} + \sigma_{zz} \hat{k}) \right] dx dy dz \\ &= \left[ \frac{\partial}{\partial x}, \frac{\partial}{\partial y}, \frac{\partial}{\partial z} \right] \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} \begin{bmatrix} \hat{i} \\ \hat{j} \\ \hat{k} \end{bmatrix} dx dy dz = \underline{[\nabla \cdot \boldsymbol{\sigma}]} dx dy dz \end{aligned}$$

→ 응력 텐서 : Symmetric Matrix ( $\because \sigma_{xy} = \sigma_{yx}, \sigma_{xz} = \sigma_{zx}, \sigma_{yz} = \sigma_{zy}$ ) 22



# Cauchy Equation<sup>1)</sup> 유도

미소 유체 요소



✓ 미소 유체 요소가 받는 힘 (Newton's 2<sup>nd</sup> Law)

$$m \frac{d\mathbf{V}}{dt} = \rho \left( \frac{\partial \mathbf{V}}{\partial t} + u \frac{\partial \mathbf{V}}{\partial x} + v \frac{\partial \mathbf{V}}{\partial y} + w \frac{\partial \mathbf{V}}{\partial z} \right) dx dy dz \quad (m = \rho dx dy dz)$$

$$m \frac{d\mathbf{V}}{dt} = \sum \mathbf{F} = \mathbf{F}_{Body} + \mathbf{F}_{Surface} \quad (\text{체적력} + \text{표면력})$$

$$\mathbf{F}_{Body} = \rho \mathbf{g} dx dy dz$$

$$\mathbf{F}_{Surface} = [\nabla \cdot \boldsymbol{\sigma}] dx dy dz$$

대입

$dx dy dz$  로 나누면

$$\rho \left( \frac{\partial \mathbf{V}}{\partial t} + u \frac{\partial \mathbf{V}}{\partial x} + v \frac{\partial \mathbf{V}}{\partial y} + w \frac{\partial \mathbf{V}}{\partial z} \right) dx dy dz = \rho \mathbf{g} dx dy dz + [\nabla \cdot \boldsymbol{\sigma}] dx dy dz$$

$$\rho \left( \frac{\partial \mathbf{V}}{\partial t} + u \frac{\partial \mathbf{V}}{\partial x} + v \frac{\partial \mathbf{V}}{\partial y} + w \frac{\partial \mathbf{V}}{\partial z} \right) = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\sigma} \Rightarrow \text{Cauchy Equation}$$



# Navier–Stokes Equation

**A**dvanced  
**S**hip  
**D**esign  
**A**utomation  
**L**aboratory

# Navier-Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

✓ Cauchy Equation

$$\rho \frac{d\mathbf{V}}{dt} = \rho \left( \frac{\partial \mathbf{V}}{\partial t} + u \frac{\partial \mathbf{V}}{\partial x} + v \frac{\partial \mathbf{V}}{\partial y} + w \frac{\partial \mathbf{V}}{\partial z} \right) = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\sigma}$$

$$x\text{축 성분} : \rho \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = \rho g_x + \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z}$$

$$y\text{축 성분} : \rho \left( \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} \right) = \rho g_y + \frac{\partial \sigma_{xy}}{\partial x} + \frac{\partial \sigma_{yy}}{\partial y} + \frac{\partial \sigma_{zy}}{\partial z}$$

$$z\text{축 성분} : \rho \left( \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} \right) = \rho g_z + \frac{\partial \sigma_{xz}}{\partial x} + \frac{\partial \sigma_{yz}}{\partial y} + \frac{\partial \sigma_{zz}}{\partial z}$$



(Q) 3개의 식인데, 미지수는 몇 개?

(A) 밀도(1개), 속도 (3개), 응력 텐서(6개) **총 10개의 미지수가 존재**

$$(\rho, u, v, w, \sigma_{xx}, \sigma_{xy}, \sigma_{xz}, \sigma_{yy}, \sigma_{yz}, \sigma_{zz})$$

응력 텐서 9개 성분 중 대칭성으로 인해 6개의 미지수 존재 ( $\sigma_{ij} = \sigma_{ji}$ )

Great Idea !!!

**응력 텐서를 속도로 표현하자!!!**

➔ 유체 입자의 운동과 변형, 뉴턴 유체의 가정으로부터 유도할 수 있음



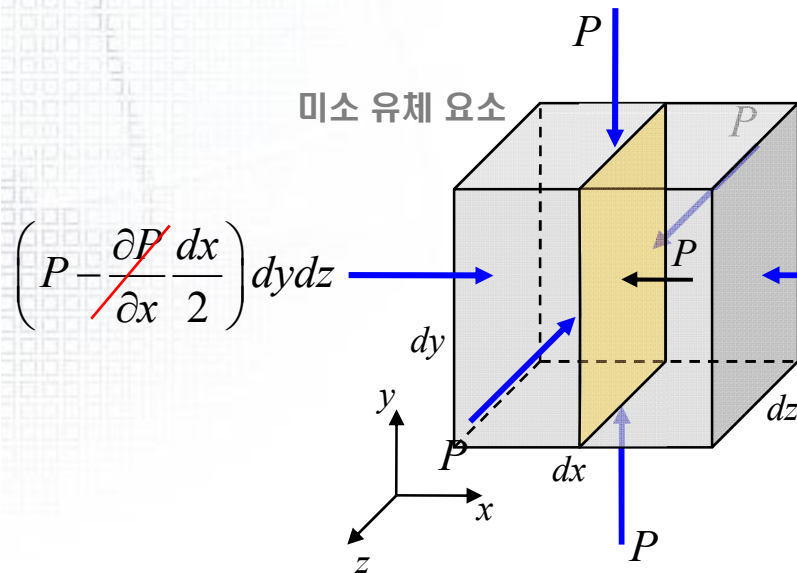
# Navier-Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

✓ 응력 텐서 성분의 분해 : 정지 상태의 유체인 경우

$$\sigma_{ij} = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix}$$

$P$  : 이상기체의 상태방정식과 같은 열역학 방정식으로부터 구해지는 열역학적 압력 (thermodynamic pressure)



$$\begin{aligned} F_{xx} &= -\left(P + \frac{\partial P}{\partial x} \frac{dx}{2}\right) dydz + \left(P - \frac{\partial P}{\partial x} \frac{dx}{2}\right) dydz \\ &= -\frac{\partial P}{\partial x} dx dy dz = m\ddot{x} = 0 \end{aligned}$$

➔ 정지 상태에서의 미소 유체 입자의 양쪽 면에 작용하는 열역학적 압력은 같다.

➔ 동일한 방법으로 다른 두 축에 대한 열역학적 압력도 고려하면, 모든 면에 작용하는 열역학적 압력은 동일함 <sup>36</sup>

# Navier–Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, “공학도를 위한 길잡이 유체공학입문”, 문운당, 2002, pp258-264

✓ 응력 텐서 성분의 분해 : 유체가 운동하는 경우

$$\sigma_{ij} = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix} + \begin{bmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{bmatrix}$$

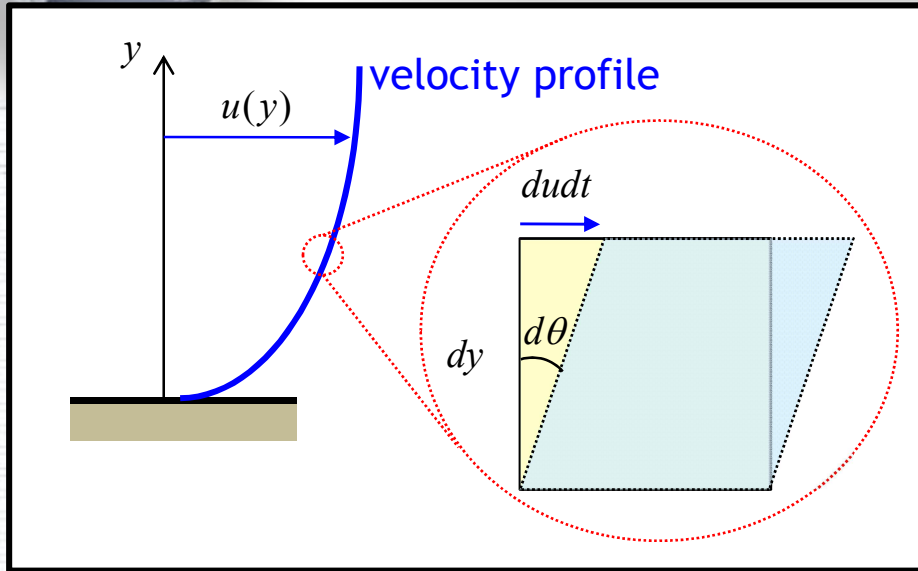
유체 입자가 운동하면, 점성에 의해 각 면에 응력이 발생함



➔ 유체 입자의 운동과 변형, 뉴턴 유체의 가정으로부터 유도할 수 있음

# 뉴턴 유체<sup>1)</sup> (Newtonian Fluid)

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264



✓ 미소 구간에서의 전단변형을

$$d\theta \approx \tan d\theta = \frac{du}{dy} dt$$

✓ 전단변형율의 시간변화율은 속도 구배와 같음

$$\frac{d\theta}{dt} = \frac{du}{dy} \quad \text{--- ①}$$

✓ 뉴턴 유체(Newtonian fluid)

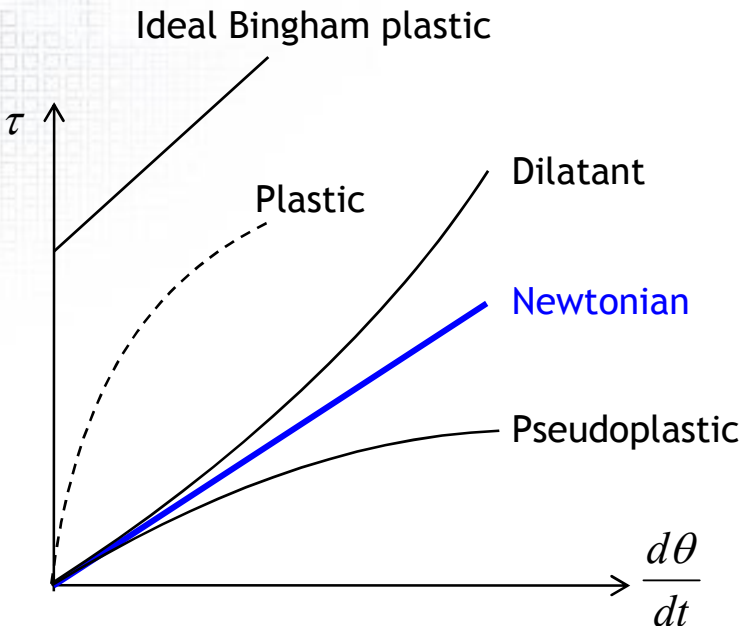
: 전단응력이 전단변형율의 시간변화율에 비례

$$\tau \propto \frac{d\theta}{dt} \quad \text{--- ②}$$

✓ 뉴턴 유체(Newtonian fluid)의 특징

: ①, ②에 의해, 전단응력은 속도구배에 비례함  
(비례 상수  $\mu$ : 점성 계수)

$$\tau \propto \frac{du}{dy} \quad \rightarrow \quad \tau = \mu \frac{du}{dy}$$

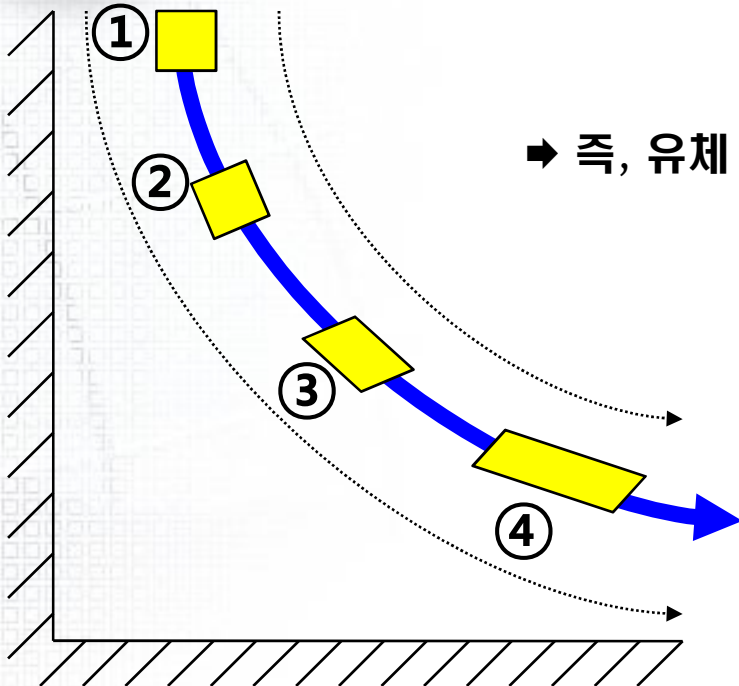


# Navier-Stokes Equation 유도

## - 유체 입자의 운동과 변형

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

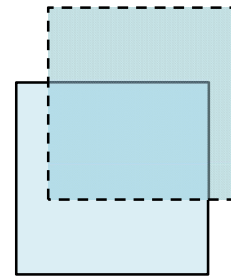
미소 유체 요소



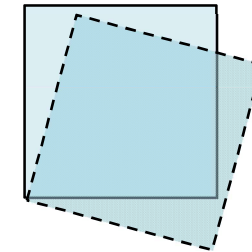
① → ④ : 유체 입자의 이동 경로

■ → ▭ : 유체 입자의 변형

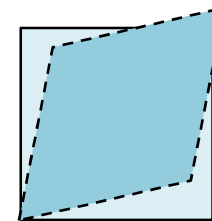
→ 즉, 유체 입자는 운동(병진, 회전)하면서, 변형(전단변형, 선형변형)도 함



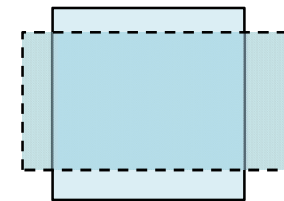
병진 운동 (translation)



회전 (rotation)



전단변형 (angular distortion)



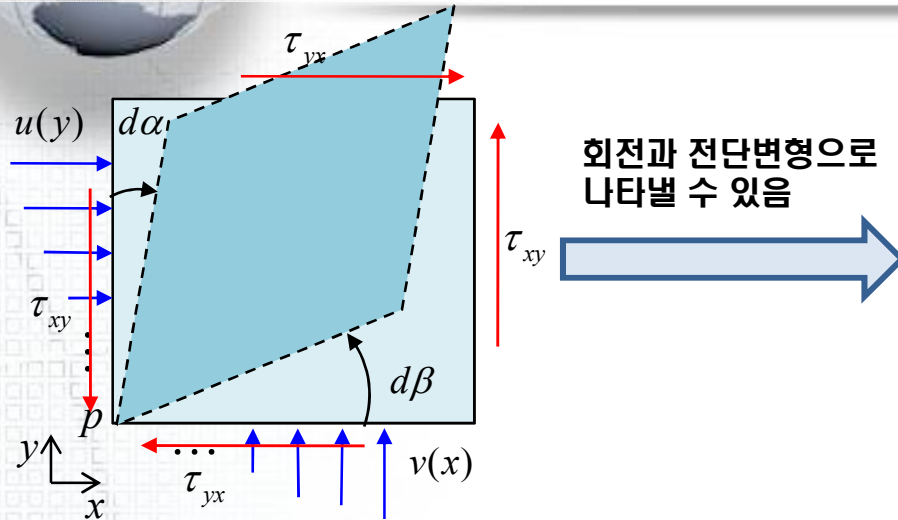
선형변형 (volume distortion)



# Navier-Stokes Equation 유도

## - 전단변형에 의한 전단력

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264



Given :  $d\alpha, d\beta$   
 Find : 전단응력  $\tau_{xy} (= \tau_{yx})$

(1) ②에서 ①을 빼면,  $d\beta + d\alpha = 2d\theta_2$

$$d\theta_2 = \frac{1}{2}(d\alpha + d\beta)$$

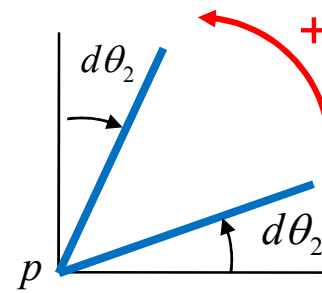
→ 전단변형율

(2)  $dt$ 로 나누면,

$$\frac{d\theta_2}{dt} = \frac{1}{2} \left( \frac{d\alpha}{dt} + \frac{d\beta}{dt} \right)$$

→ 전단변형율(의 시간변화율)

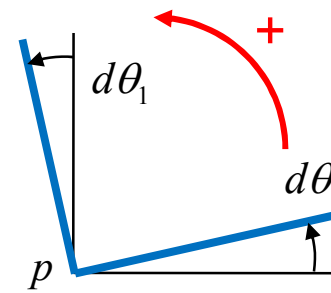
전단변형(shear strain)



$$-d\alpha = -d\theta_2 + d\theta_1 \quad \dots \textcircled{1}$$

$$d\beta = d\theta_2 + d\theta_1 \quad \dots \textcircled{2}$$

회전(rotation)



(3) 뉴턴 유체라 가정하면, 전단응력은 전단변형율의 시간변화율에 비례함 (비례상수  $2\mu$ )

$$\tau = 2\mu \frac{d\theta_2}{dt} = 2\mu \frac{1}{2} \left( \frac{d\alpha}{dt} + \frac{d\beta}{dt} \right) = \mu \left( \frac{d\alpha}{dt} + \frac{d\beta}{dt} \right)$$

(4) 전단변형율의 시간 변화율은 속도 구배와 같음

$$\tau = \mu \left( \frac{du}{dy} + \frac{dv}{dx} \right) = \tau_{xy} = \tau_{yx}$$

(5) 동일한 방법에 의해 다음도 성립함

$$\tau_{xz} = \tau_{zx} = \mu \left( \frac{du}{dz} + \frac{dw}{dx} \right), \quad \tau_{yz} = \tau_{zy} = \mu \left( \frac{dv}{dz} + \frac{dw}{dy} \right) \quad \mathbf{30}$$

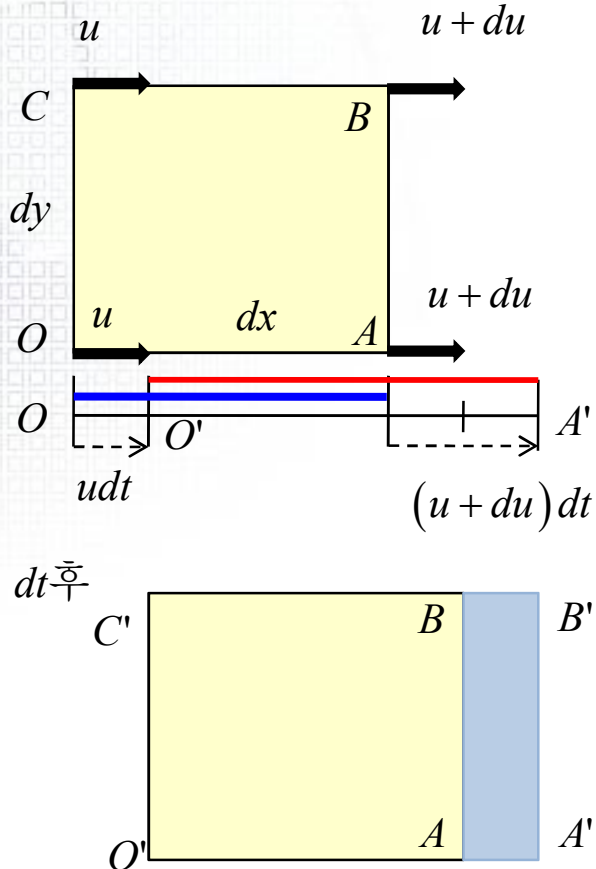
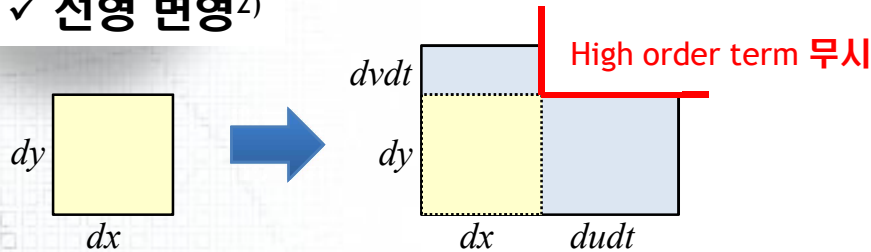
# Navier-Stokes Equation 유도

## - 선형 변형에 의한 수직응력

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

등방 팽창\* : 각 축방향으로 동일한 길이만큼 팽창

✓ 선형 변형<sup>2)</sup>



→ 시간 t에서의 부피 :  $V = dx dy dz$

→ 시간 t+dt에서의 부피 :

$$\begin{aligned}
 V + dV &= (dx + dudt)(dy + dvdt)(dz + dwdt) && \text{High order term 무시} \\
 &= dx dy dz + (dudydz + dvdx dz + dwdxdy) dt + \dots dt^2 + (dudvdw) dt^3 \\
 &\cong dx dy dz + (dudydz + dvdx dz + dwdxdy) dt \\
 &= dx dy dz + dx dy dz \left( \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right) dt = V + V \left( \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right) dt
 \end{aligned}$$

$$\therefore dV = V \Theta dt \quad \left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

$$AA' = (u + du)dt - udt = dudt$$

속도차이에 의해 선형적으로 늘어난 길이

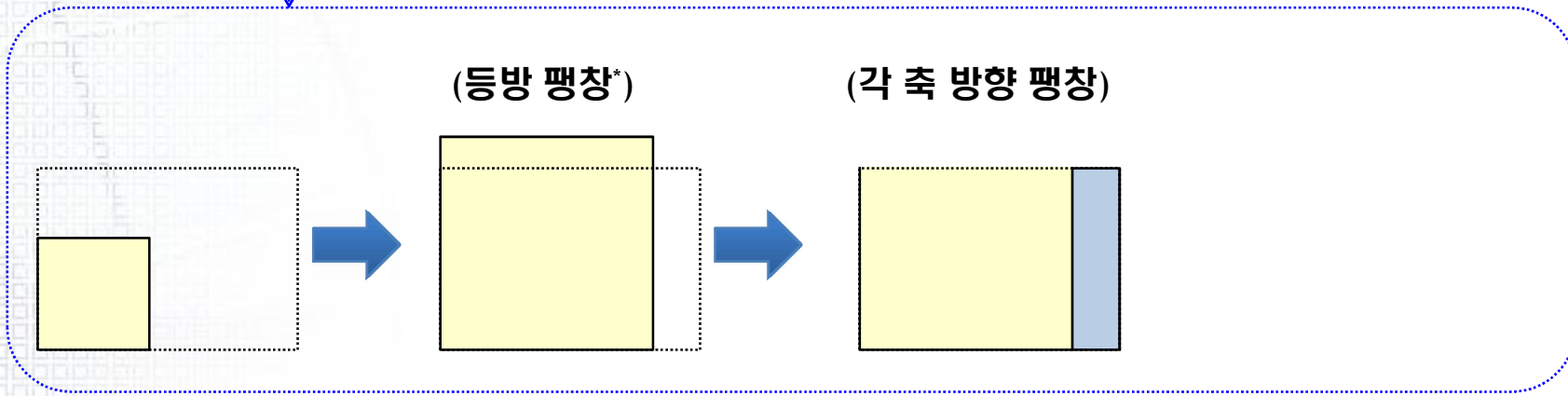
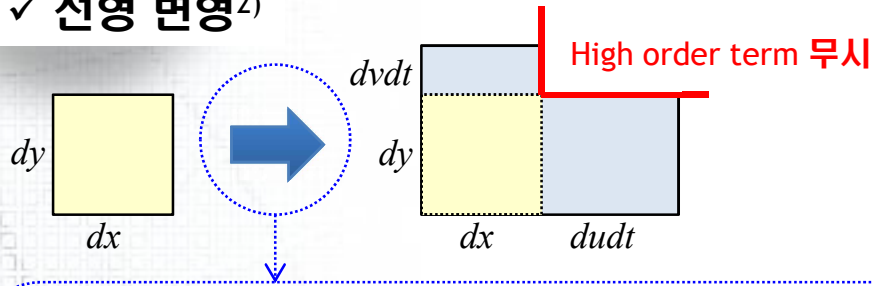
# Navier-Stokes Equation 유도

## - 선형 변형에 의한 수직응력

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

등방 팽창\* : 각 축방향으로 동일한 길이만큼 팽창

✓ 선형 변형<sup>2)</sup>



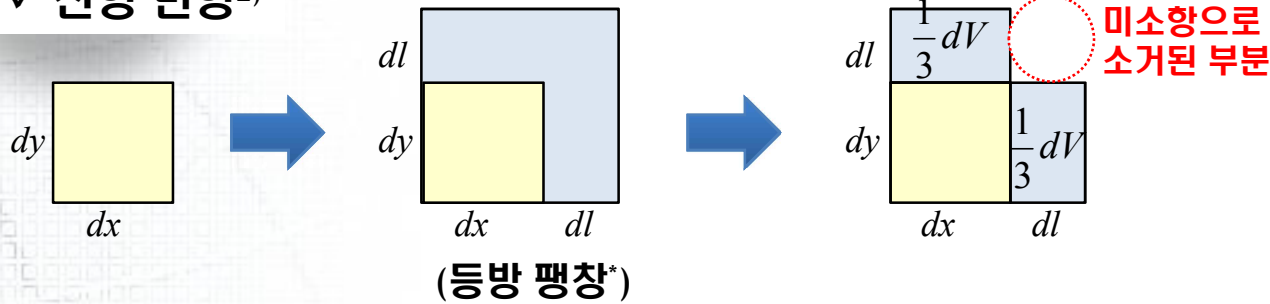
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➔ 등방 팽창이므로,  $t+dt$ 에서 증가한 부피  $dV$  중  $1/3$ 이 각 축방향으로 증가한 부피

$$dV = V\Theta dt \quad \left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

$$\frac{1}{3} dV = \frac{1}{3} V\Theta dt = \underline{dl dy dz}$$

(x축으로 등방 팽창한 부피)

$$dl = \frac{1}{3} \frac{V}{dy dz} \Theta dt = \frac{1}{3} \frac{dx dy dz}{dy dz} \Theta dt = \frac{1}{3} \Theta dx dt$$

$$\frac{1}{dt} \left( \frac{dl}{dx} \right) = \frac{1}{3} \Theta$$

➔ 등방팽창에 의한 변형율

➔ 등방팽창에 의한 변형율의 시간변화율

dl로 정리

dxdt로 나눔

➔ 등방팽창에 의한 수직응력은 등방팽창에 의한 변형율의 시간변화율에 비례함 (비례상수 :  $3\gamma$ )

$$\tau_x^{dilation, isotropic} = 3\gamma \frac{1}{3} \Theta = \gamma \Theta$$

(분수를 소거하기 위해  $3\gamma$ 를 택한 것)



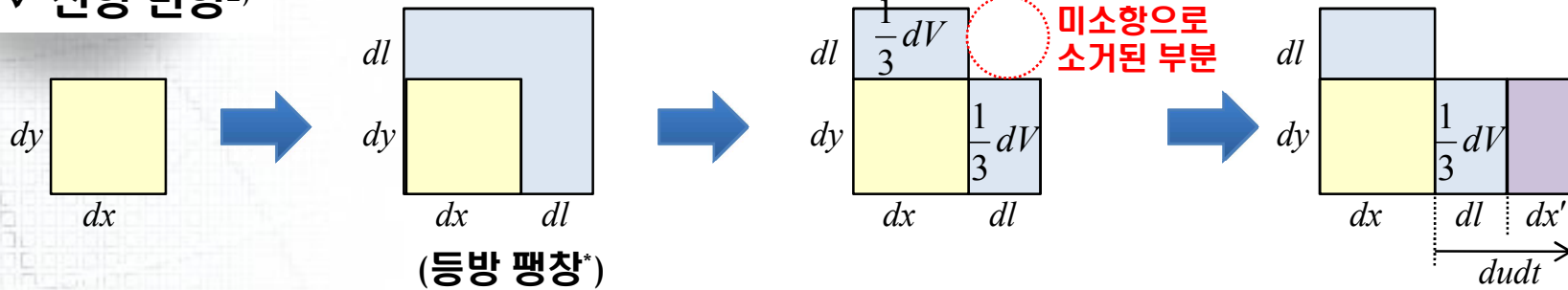
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- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
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$$\Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz}$$

✓ 선형 변형<sup>2)</sup>



➔ 등방 팽창을 제외한  $x$ 축방향의 팽창에 의한 변형율의 시간 변화율

$$-x\text{축방향의 변형율} : \frac{dx'}{dx} = \frac{dudt - dl}{dx} = \frac{dudt}{dx} - \frac{dl}{dx}$$

$$-x\text{축방향의 변형율의 시간변화율} : \frac{1}{dt} \frac{dx'}{dx} = \frac{du}{dx} - \frac{1}{dt} \frac{dl}{dx} = \frac{du}{dx} - \frac{1}{3} \Theta$$

➔ 등방 팽창을 제외한  $x$ 축방향의 팽창에 의한 전단응력 (비례상수 :  $2\mu$ )

$$\tau_x^{dilation,x} = 2\mu \left( \frac{\partial u}{\partial x} - \frac{1}{3} \Theta \right)$$

➔ 등방팽창에 의한 변형율의 시간변화율

$$\frac{1}{dt} \left( \frac{dl}{dx} \right) = \frac{1}{3} \Theta$$

$\lambda, \mu$ 는 물질의 특성에 따라 달라지는 값

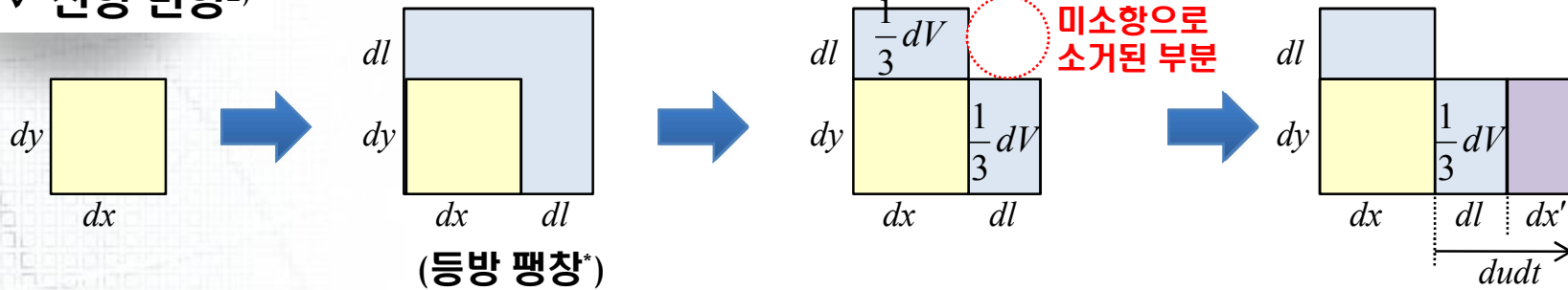
# Navier-Stokes Equation 유도

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$$\Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz}$$

✓ 선형 변형<sup>2)</sup>



→ x축 방향의 팽창에 의한 전단응력 ( $\tau_{xx} = \tau_x^{dilation, isotropic} + \tau_x^{dilation, x}$ )

$$\begin{aligned} \tau_{xx} &= \tau_x^{dilation, isotropic} + \tau_x^{dilation, x} \\ &= \gamma \Theta + 2\mu \left( \frac{\partial u}{\partial x} - \frac{1}{3} \Theta \right) = \left( \gamma - \frac{2\mu}{3} \right) \Theta + 2\mu \frac{\partial u}{\partial x} = \lambda \Theta + 2\mu \frac{\partial u}{\partial x} \\ &= \lambda (\nabla \cdot \mathbf{V}) + 2\mu \frac{\partial u}{\partial x} \end{aligned}$$

$\left( \lambda = \gamma - \frac{2\mu}{3} \right)$

→ 등방팽창을 제외한 각 축방향의 수직 응력  
→ 등방팽창에 의한 수직 응력

→ 등방팽창에 의한 변형율의 시간변화율

$$\frac{1}{dt} \left( \frac{dl}{dx} \right) = \frac{1}{3} \Theta$$

$\lambda, \mu$ 는 물질의 특성에 따라 달라지는 값

→ 동일한 방법에 의해 다음도 성립함 :  $\tau_{yy} = \lambda \Theta + 2\mu \frac{\partial v}{\partial y}$ ,  $\tau_{zz} = \lambda \Theta + 2\mu \frac{\partial w}{\partial z}$

# Navier–Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
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✓ 응력 텐서 성분의 분해 : 유체가 운동하는 경우

$$\sigma_{ij} = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix} + \begin{bmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{bmatrix}$$

$P$  : 열역학적 압력  
(thermodynamic pressure)

유체 입자가 운동하면, 점성에 의해 각 면에 응력이 발생함

➔ 유체 입자의 운동과 변형, 뉴턴 유체의 가정으로부터 유도할 수 있음

✓ 수직응력

$$\tau_{xx} = \lambda\Theta + 2\mu \frac{\partial u}{\partial x}$$

$$\tau_{yy} = \lambda\Theta + 2\mu \frac{\partial v}{\partial y}$$

$$\tau_{zz} = \lambda\Theta + 2\mu \frac{\partial w}{\partial z}$$

$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

✓ 전단응력

$$\tau_{xy} = \tau_{yx} = \mu \left( \frac{du}{dy} + \frac{dv}{dx} \right)$$

$$\tau_{xz} = \tau_{zx} = \mu \left( \frac{du}{dz} + \frac{dw}{dx} \right)$$

$$\tau_{yz} = \tau_{zy} = \mu \left( \frac{dv}{dz} + \frac{dw}{dy} \right)$$

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
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# Navier-Stokes Equation 유도

✓ 응력 텐서 성분의 분해 : 유체가 운동하는 경우 ( $P$  : 열역학적 압력)  
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$$\sigma_{ij} = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix} + \begin{bmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{bmatrix}$$

➔ 응력 텐서의 대각 성분의 합을 구해보면,

$$\sigma_{xx} = -P + \tau_{xx} = -P + \lambda\Theta + 2\mu \frac{\partial u}{\partial x}$$

$$\sigma_{yy} = -P + \tau_{yy} = -P + \lambda\Theta + 2\mu \frac{\partial v}{\partial y}$$

$$+ \sigma_{zz} = -P + \tau_{zz} = -P + \lambda\Theta + 2\mu \frac{\partial w}{\partial z}$$

$$\begin{aligned} \underline{\sigma_{xx} + \sigma_{yy} + \sigma_{zz}} &= -3P + 3\lambda\Theta + 2\mu \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) \\ &= \underline{-3P + (3\lambda + 2\mu)\Theta} \end{aligned}$$

✓ 수직응력

$$\tau_{xx} = \lambda\Theta + 2\mu \frac{\partial u}{\partial x}$$

$$\tau_{yy} = \lambda\Theta + 2\mu \frac{\partial v}{\partial y}$$

$$\tau_{zz} = \lambda\Theta + 2\mu \frac{\partial w}{\partial z}$$

$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right)$$



# Navier-Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264
- 4) 矢川元基, 이형직 역, 유한요소법에 의한 유체역학, 열전도 해석 입문, 피어슨 에듀케이션 코리아, pp8-10

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✓ 수직응력

$$\tau_{xx} = \lambda\Theta + 2\mu \frac{\partial u}{\partial x}$$

$$\tau_{yy} = \lambda\Theta + 2\mu \frac{\partial v}{\partial y}$$

$$\tau_{zz} = \lambda\Theta + 2\mu \frac{\partial w}{\partial z}$$

$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right)$$

➔ 응력 텐서의 대각 성분의 합을 구해보면,

$$\sigma_{xx} + \sigma_{yy} + \sigma_{zz} = -3P + (3\lambda + 2\mu)\Theta$$

➔ \*평균압력 or \*기계적 압력을 정의

$$-\frac{1}{3}(\sigma_{xx} + \sigma_{yy} + \sigma_{zz}) = \bar{P}$$

$$\bar{P} = P - \left( \lambda + \frac{2}{3}\mu \right) \Theta = P - \kappa\Theta$$

➔ 평균압력과 열역학적 압력은 일반적으로 같지 않음

(참고) 비압축성 유동을 가정할 경우,  
열역학적 압력을 결정할 수 있는 상태 방정식이 존재하지 않음<sup>2)</sup>의 p102

( $\rho = const.$  가 상태방정식을 대체함<sup>1)</sup>의 p427)

(참고) 비압축성 유동 (incompressible flow) 일 경우,  $\Theta = 0$  이므로  
비로소 평균압력과 열역학적 압력은 같아짐

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

# Navier-Stokes Equation 유도

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$$\sigma_{ij} = \begin{bmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{bmatrix} = \begin{bmatrix} -P & 0 & 0 \\ 0 & -P & 0 \\ 0 & 0 & -P \end{bmatrix} + \begin{bmatrix} \tau_{xx} & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \tau_{yy} & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \tau_{zz} \end{bmatrix}$$

→ Stokes assumption :  $\kappa = \lambda + \frac{2}{3}\mu = 0 \rightarrow \lambda = -\frac{2}{3}\mu$

$$\left( \bar{P} = P - \left( \lambda + \frac{2}{3}\mu \right) \Theta = P - \kappa \Theta \right)$$

$$\sigma_{xx} = -P + \tau_{xx} = -P + \lambda \Theta + 2\mu \frac{\partial u}{\partial x} = -P - \frac{2}{3}\mu \Theta + 2\mu \frac{\partial u}{\partial x}$$

$$\sigma_{yx} = \tau_{yx} = \mu \left( \frac{du}{dy} + \frac{dv}{dx} \right)$$

$$\sigma_{zx} = \tau_{zx} = \mu \left( \frac{du}{dz} + \frac{dw}{dx} \right)$$

✓ 수직응력

$$\tau_{xx} = \lambda \Theta + 2\mu \frac{\partial u}{\partial x}$$

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$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

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# Navier-Stokes Equation 유도

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
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$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

$$\begin{aligned} \sigma_{xx} &= -P - \frac{2}{3} \mu \Theta + 2\mu \frac{\partial u}{\partial x} \\ \sigma_{yx} &= \mu \left( \frac{du}{dy} + \frac{dv}{dx} \right) \\ \sigma_{zx} &= \mu \left( \frac{du}{dz} + \frac{dw}{dx} \right) \end{aligned}$$

➔ Navier-Stokes Equation의 유도 (x축 성분)

$$\begin{aligned} \rho \frac{du}{dt} &= \rho g_x + \frac{\partial \sigma_{xx}}{\partial x} + \frac{\partial \sigma_{yx}}{\partial y} + \frac{\partial \sigma_{zx}}{\partial z} \\ &= \rho g_x + \frac{\partial}{\partial x} \left( -P - \frac{2}{3} \mu \Theta + 2\mu \frac{\partial u}{\partial x} \right) + \mu \frac{\partial}{\partial y} \left( \frac{du}{dy} + \frac{dv}{dx} \right) + \mu \frac{\partial}{\partial z} \left( \frac{du}{dz} + \frac{dw}{dx} \right) \\ &= \rho g_x - \frac{\partial P}{\partial x} - \frac{2}{3} \mu \frac{\partial}{\partial x} (\Theta) + \mu \left( 2 \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 v}{\partial y \partial x} + \frac{\partial^2 u}{\partial z^2} + \frac{\partial^2 w}{\partial z \partial x} \right) \\ &= \rho g_x - \frac{\partial P}{\partial x} - \frac{2}{3} \mu \frac{\partial}{\partial x} (\Theta) + \mu \left( \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 v}{\partial y \partial x} + \frac{\partial^2 w}{\partial z \partial x} + \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right) \\ &= \rho g_x - \frac{\partial P}{\partial x} - \frac{2}{3} \mu \frac{\partial}{\partial x} (\Theta) + \mu \left\{ \frac{\partial}{\partial x} \left( \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} \right) + \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2} \right\} \\ &= \rho g_x - \frac{\partial P}{\partial x} - \frac{2}{3} \mu \frac{\partial}{\partial x} (\Theta) + \mu \left\{ \frac{\partial}{\partial x} \Theta + \nabla^2 u \right\} = \rho g_x - \frac{\partial P}{\partial x} + \frac{\mu}{3} \frac{\partial}{\partial x} (\Theta) + \mu \nabla^2 u \end{aligned}$$

# Navier-Stokes Equation 유도

(Continue)

$$\left( \Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz} \right)$$

- x축 성분 :

$$\rho \frac{du}{dt} = \rho g_x - \frac{\partial P}{\partial x} + \frac{\mu}{3} \frac{\partial}{\partial x} (\Theta) + \mu \nabla^2 u$$

- y축 성분

$$\rho \frac{dv}{dt} = \rho g_y - \frac{\partial P}{\partial y} + \frac{\mu}{3} \frac{\partial}{\partial y} (\Theta) + \mu \nabla^2 v$$

- z축 성분

$$\rho \frac{dw}{dt} = \rho g_z - \frac{\partial P}{\partial z} + \frac{\mu}{3} \frac{\partial}{\partial z} (\Theta) + \mu \nabla^2 w$$

벡터 형태로 정리하면 ( $\mathbf{v} = [u, v, w]^T$ )

$$\begin{aligned} \rho \frac{d\mathbf{V}}{dt} &= \rho \mathbf{g} - \nabla P + \frac{\mu}{3} \nabla \Theta + \mu \nabla^2 \mathbf{V} \\ &= \rho \mathbf{g} - \nabla P + \mu \left( \frac{1}{3} \nabla \Theta + \nabla^2 \mathbf{V} \right) \end{aligned}$$

=> Navier-Stokes' Equation  
(General form)



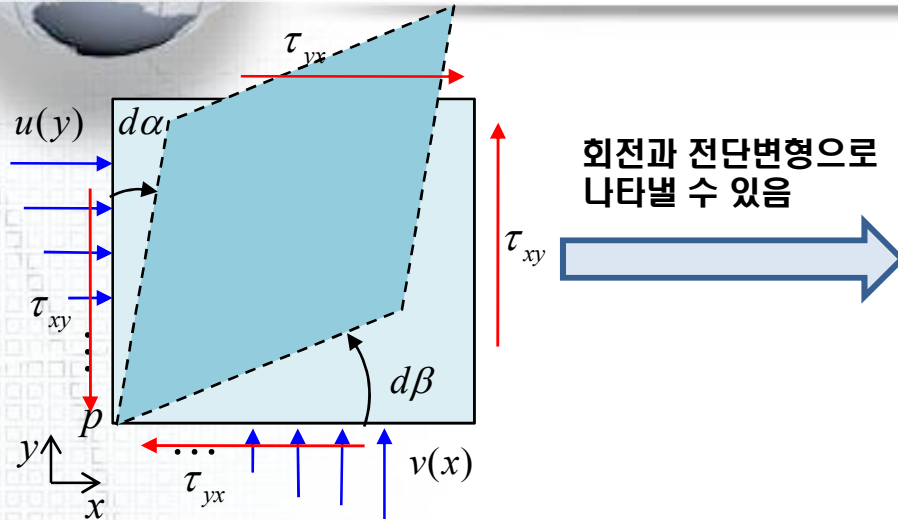
# 유체 입자의 회전

**A**dvanced  
**S**hip  
**D**esign  
**A**utomation  
**L**aboratory



# 유체 입자의 회전과 Curl

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264



Given :  $d\alpha, d\beta$   
Find : 회전각속도  $\omega_z$

(1) ②와 ①을 더하면,  $d\beta - d\alpha = 2d\theta_1$

$$d\theta_1 = \frac{1}{2}(d\beta - d\alpha)$$

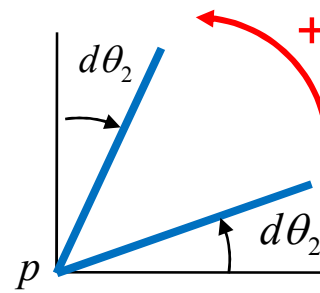
→ z축에 대한 회전 각도

(2)  $dt$ 로 나누면,

$$\frac{d\theta_1}{dt} = \frac{1}{2} \left( \frac{d\beta}{dt} - \frac{d\alpha}{dt} \right)$$

→ z축에 대한 회전 각속도

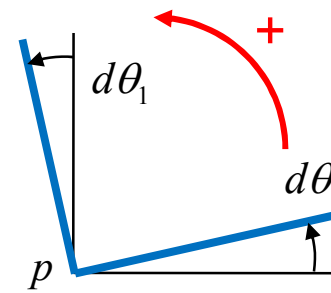
전단변형 (shear strain)



$$-d\alpha = -d\theta_2 + d\theta_1 \quad \dots \textcircled{1}$$

$$d\beta = d\theta_2 + d\theta_1 \quad \dots \textcircled{2}$$

회전 (rotation)



(3) 전단변형율의 시간 변화율은 속도 구배와 같음

$$\omega_z = \frac{1}{2} \left( \frac{dv}{dx} - \frac{du}{dy} \right)$$

(5) 동일한 방법에 의해 다음도 성립함

→ y축에 대한 회전각속도 :  $\omega_y = \frac{1}{2} \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$

→ x축에 대한 회전각속도 :  $\omega_x = \frac{1}{2} \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right)$

# 유체 입자의 회전과 Curl

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

## ✓ curl의 정의

$$\text{curl } \mathbf{V} = \nabla \times \mathbf{V} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ u & v & w \end{vmatrix}$$

$$= \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right) \mathbf{i} + \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right) \mathbf{j} + \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) \mathbf{k}$$

$$= 2\omega_x \mathbf{i} + 2\omega_y \mathbf{j} + 2\omega_z \mathbf{k}$$

=> curl  $\mathbf{V}$  는 유체 입자의 (회전 각속도/2)에 해당하는 값

## ✓ x축에 대한 각속도

$$\omega_x = \frac{1}{2} \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right)$$

## ✓ y축에 대한 각속도

$$\omega_y = \frac{1}{2} \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

## ✓ z축에 대한 각속도

$$\omega_z = \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

✓ Irrotational Flow :  $\nabla \times \mathbf{v} = 0 \left( \frac{\partial w}{\partial y} = \frac{\partial v}{\partial z}, \frac{\partial u}{\partial z} = \frac{\partial w}{\partial x}, \frac{\partial v}{\partial x} = \frac{\partial u}{\partial y} \right)$

✓  $\nabla \times \mathbf{v} = 0$  일 때,  $\mathbf{v} = \nabla \Phi$  인 Scalar Function  $\Phi$  가 존재함



# Euler Equation & Bernoulli Equation<sup>1),2)</sup> 유도

# Euler Equation<sup>1),2)</sup> 유도

✓ Navier-Stokes Equation (general form)

$$\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P + \mu \left( \frac{1}{3} \nabla \Theta + \nabla^2 \mathbf{V} \right)$$



비점성 유동 (inviscid flow) 라 가정하면  
 $(\mu = 0)$

$$\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P$$



$$\frac{d\mathbf{V}}{dt} = \mathbf{g} - \frac{\nabla P}{\rho}$$

=> Euler Equation

$$\mathbf{V} = [u, v, w]^T$$

$$\Theta = \nabla \cdot \mathbf{V} = \frac{du}{dx} + \frac{dv}{dy} + \frac{dw}{dz}$$

$$\left( \begin{array}{l} x\text{축 성분:} \\ y\text{축 성분:} \\ z\text{축 성분:} \end{array} \right. \begin{array}{l} \frac{du}{dt} = \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = g_x - \frac{1}{\rho} \frac{\partial P}{\partial x} \\ \frac{dv}{dt} = \frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} + w \frac{\partial v}{\partial z} = g_y - \frac{1}{\rho} \frac{\partial P}{\partial y} \\ \frac{dw}{dt} = \frac{\partial w}{\partial t} + u \frac{\partial w}{\partial x} + v \frac{\partial w}{\partial y} + w \frac{\partial w}{\partial z} = g_z - \frac{1}{\rho} \frac{\partial P}{\partial z} \end{array} \right)$$

# Bernoulli Equation<sup>1),2)</sup>의 유도 - Euler equation의 변경

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.118-122

$$\frac{d\mathbf{V}}{dt} = \mathbf{g} - \frac{\nabla P}{\rho} \quad \xrightarrow{x\text{축 성분}} \quad \frac{\partial u}{\partial t} + \boxed{u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z}} = g_x - \frac{1}{\rho} \frac{\partial P}{\partial x}$$

$$\boxed{u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z}} = u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} + \left( v \frac{\partial v}{\partial x} - v \frac{\partial v}{\partial x} \right) + \left( w \frac{\partial w}{\partial x} - w \frac{\partial w}{\partial x} \right)$$

$$= \left( u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial x} + w \frac{\partial w}{\partial x} \right) + \left( v \frac{\partial u}{\partial y} - v \frac{\partial v}{\partial x} \right) + \left( w \frac{\partial u}{\partial z} - w \frac{\partial w}{\partial x} \right)$$

$$= \left[ \frac{\partial}{\partial x} \left( \frac{1}{2} u^2 \right) + \frac{\partial}{\partial x} \left( \frac{1}{2} v^2 \right) + \frac{\partial}{\partial x} \left( \frac{1}{2} w^2 \right) \right] + v \left( \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} \right) + w \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right)$$

$$= \frac{\partial}{\partial x} \left( \frac{1}{2} (u^2 + v^2 + w^2) \right) + v \left( \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} \right) + w \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right) = \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 \right) - v\omega_z + w\omega_y$$

$$\text{curl } \mathbf{v} = \nabla \times \mathbf{v} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ u & v & w \end{vmatrix} = \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right) \mathbf{i} + \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right) \mathbf{j} + \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) \mathbf{k} = \omega_x \mathbf{i} + \omega_y \mathbf{j} + \omega_z \mathbf{k}$$



# Bernoulli Equation<sup>1),2)</sup>의 유도 - Euler equation의 변경

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, p17, pp.118-122

\*등엔트로피 : 유체의 엔트로피가 변하지 않음, 열의 공급이 일어나지 않고(adiabatic), 마찰이 없음(frictionless)

$$\frac{dV}{dt} = \mathbf{g} - \frac{\nabla P}{\rho}$$

**x축 성분** → 
$$\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} = g_x - \frac{1}{\rho} \frac{\partial P}{\partial x}$$

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 \right) - v\omega_z + w\omega_y = g_x - \frac{1}{\rho} \frac{\partial P}{\partial x}$$

Body force를 포텐셜 에너지의 gradient로 표현 ( $\mathbf{g} = -\nabla(gy)$ )

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 \right) + \frac{\partial}{\partial x} (gy) + \frac{1}{\rho} \frac{\partial P}{\partial x} = v\omega_z - w\omega_y$$

① barotropic flow

(유체의 밀도가 압력만의 함수로 표현)

$$\rho = \rho(P)$$

①에 의해 다음이 성립함 → 
$$\frac{1}{\rho(P)} \frac{\partial P}{\partial x} = \frac{\partial}{\partial x} \int \frac{dP}{\rho(P)}$$

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 \right) + \frac{\partial}{\partial x} (gy) + \frac{\partial}{\partial x} \int \frac{dP}{\rho} = v\omega_z - w\omega_y$$

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} \right) = v\omega_z - w\omega_y$$

# Bernoulli Equation<sup>1),2)</sup>의 유도 - Euler equation의 변경

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.118-122

✓ Euler Equation

$$\frac{d\mathbf{V}}{dt} = \mathbf{g} - \frac{\nabla P}{\rho}$$

x축 성분

$$\frac{\partial u}{\partial t} + \frac{\partial}{\partial x} \left( \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} \right) = v\omega_z - w\omega_y$$

y축 성분

$$\frac{\partial v}{\partial t} + \frac{\partial}{\partial y} \left( \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} \right) = w\omega_x - u\omega_z$$

z축 성분

$$\frac{\partial w}{\partial t} + \frac{\partial}{\partial z} \left( \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} \right) = u\omega_y - v\omega_x$$



$$\frac{\partial \mathbf{V}}{\partial t} + \nabla B = \mathbf{V} \times \boldsymbol{\omega} \quad \text{: Another form of Euler equation}$$

$$\mathbf{V} = [u, v, w]^T$$

$$\boldsymbol{\omega} = [\omega_x, \omega_y, \omega_z]^T$$

$$B = \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho}$$

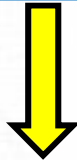
$$\mathbf{V} \times \boldsymbol{\omega} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ u & v & w \\ \omega_x & \omega_y & \omega_z \end{vmatrix} = \mathbf{i}(v\omega_z - w\omega_y) + \mathbf{j}(w\omega_x - u\omega_z) + \mathbf{k}(u\omega_y - v\omega_x)$$

: Bernoulli function  
(단위 질량당 운동에너지, 포텐셜에너지, 유동에너지의 합)

# Bernoulli Equation<sup>1),2)</sup>의 유도 - 현재까지의 가정 정리

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182  
2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.118-122

Cauchy Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} + \nabla \cdot \boldsymbol{\sigma}$



- ① 뉴턴 유체 (Newtonian fluid)
- ② Stokes assumption

Navier-Stokes Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P + \mu \left( \frac{1}{3} \nabla \Theta + \nabla^2 \mathbf{V} \right)$   
(in general form)



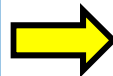
- ③ 비점성 유동 (Inviscid flow)

Euler Equation :  $\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P$



- ③ barotropic flow

Euler Equation :  $\frac{\partial \mathbf{V}}{\partial t} + \nabla B = \mathbf{V} \times \boldsymbol{\omega}$   
(Another form)



다음에 대해 Bernoulli equation을 유도

- 1) Steady flow / rotational or irrotational flow
- 2) Unsteady flow / irrotational flow

※ Euler equation을 유도하기 까지 비압축성 유동(Incompressible flow)은 가정하지 않았음

# Bernoulli Equation<sup>1),2)</sup>의 유도

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.118-122

## 1) Steady flow에서의 Bernoulli equation

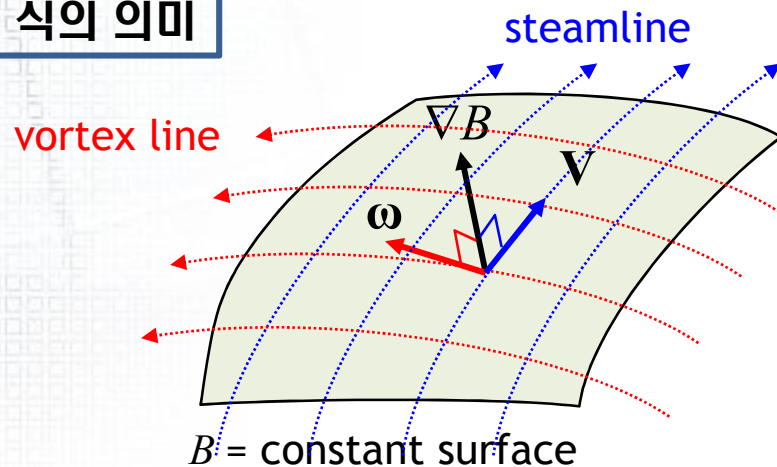
Euler Equation :  $\frac{\partial \mathbf{V}}{\partial t} + \nabla B = \mathbf{V} \times \boldsymbol{\omega}$

$$B = \frac{1}{2}q^2 + gy + \int \frac{dP}{\rho}$$

↓ steady flow  
 (유체 입자 운동이 시간과 무관)  $\left( \frac{\partial \mathbf{V}}{\partial t} = 0 \right)$

$$\nabla B = \mathbf{V} \times \boldsymbol{\omega}$$

식의 의미



$B = \text{constant}$  : 곡면(surface)

$\nabla B$  : 곡면(surface)에 수직인 벡터

$\mathbf{V}$  : 유체입자의 속도

$\boldsymbol{\omega}$  : 유체입자의 회전 (회전축 벡터)

$\nabla B$  는  $\mathbf{V}$  와  $\boldsymbol{\omega}$  에 수직

➔  $B = \text{constant}$  인 곡면 위에 stream line과 vortex line을 포함하고 있음

➔ 즉, stream line과 vortex line 이 이루는 곡면에서  $B = \text{constant}$ 를 만족함

$$\therefore B = \frac{1}{2}q^2 + gy + \int \frac{dP}{\rho} = \text{constant along streamlines and vortex lines}$$

➔ Inviscid, steady, barotropic flow

# Bernoulli Equation<sup>1),2)</sup>의 유도

1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, p179-182

2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.118-122

## 2) Unsteady irrotational flow 에서의 Bernoulli equation

✓ irrotational flow의 경우 다음을 만족함

$$\rightarrow \mathbf{V} = \nabla \phi$$

$$\rightarrow \text{curl } \mathbf{V} = \boldsymbol{\omega} = 0$$

✓ Euler equation에 대입함

$$\frac{\partial \mathbf{V}}{\partial t} + \nabla B = \mathbf{V} \times \boldsymbol{\omega} \quad \rightarrow \quad \frac{\partial}{\partial t} (\nabla \phi) + \nabla B = 0 \quad \rightarrow \quad \nabla \frac{\partial \phi}{\partial t} + \nabla B = 0$$

$$\rightarrow \nabla \left( \frac{\partial \phi}{\partial t} + B \right) = \nabla \left( \frac{\partial \phi}{\partial t} + \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} \right) = 0$$

$$\rightarrow \frac{\partial \phi}{\partial t} + \frac{1}{2} q^2 + gy + \int \frac{dP}{\rho} = F(t)$$

➔ Inviscid, unsteady, barotropic, irrotational flow

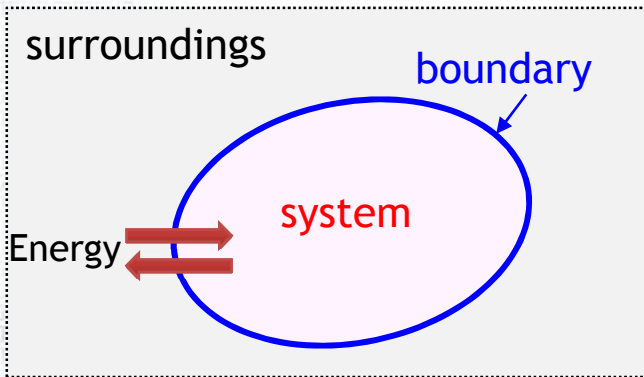




# 에너지 방정식과 Bernoulli equation의 비교<sup>1)</sup>

# 에너지 방정식

✓ 시스템(system)의 정의 : 연구의 대상이 되는 공간 내의 영역이나 영역 내에 포함된 물질의 양

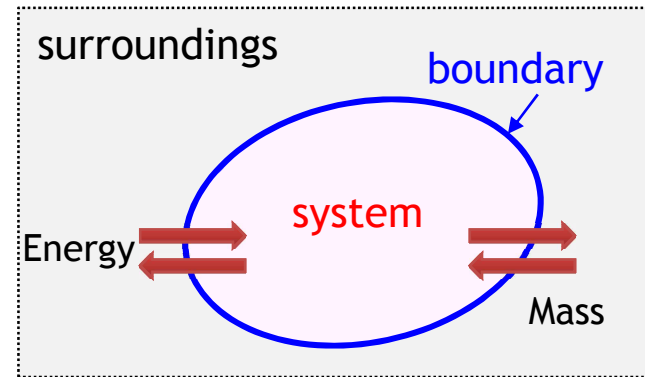


➔ 밀폐 시스템(closed system)  
: 에너지(열, 일)의 출입은 가능하나  
질량의 통과가 없음

➔ 시스템의 단위 질량당 에너지

$$e = u + \frac{V^2}{2} + gz$$

내부E      운동E      포텐셜E  
 (E:Energy)



➔ 개방 시스템(open system)  
: 에너지(열, 일)의 출입 및  
질량의 유출입도 가능

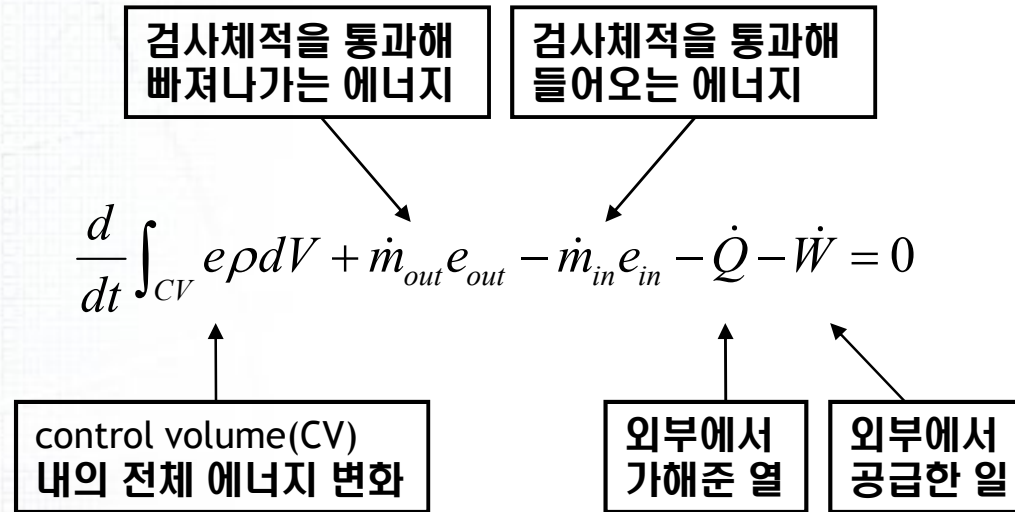
➔ 시스템의 단위 질량당 에너지

$$e = u + \frac{P}{\rho} + \frac{V^2}{2} + gz$$

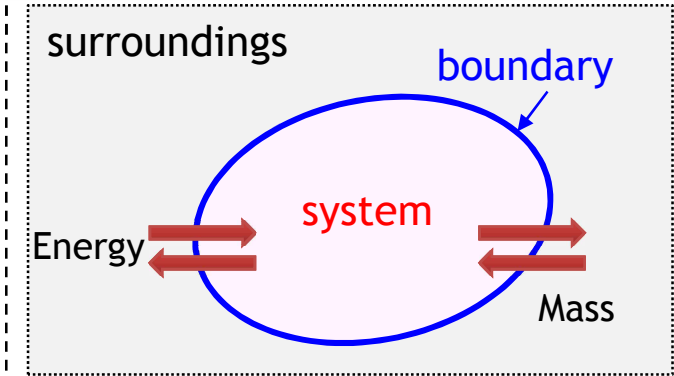
유동E  
 : 질량의 유출입이 가능한 경우,  
 유동을 유지시키기 위한 에너지

# 에너지 방정식

✓ 에너지 방정식 :



$$\frac{d}{dt} \int_{CV} e \rho dV + \dot{m}_{out} \left( u_{out} + \left( \frac{P}{\rho} \right)_{out} + \frac{\mathbf{V}_{out}^2}{2} + gz_{out} \right) - \dot{m}_{in} \left( u_{in} + \left( \frac{P}{\rho} \right)_{in} + \frac{\mathbf{V}_{in}^2}{2} + gz_{in} \right) = \dot{Q} + \dot{W}$$



➔ 시스템의 단위 질량당 에너지

$$e = u + \frac{P}{\rho} + \frac{\mathbf{V}^2}{2} + gz$$

# 에너지 방정식

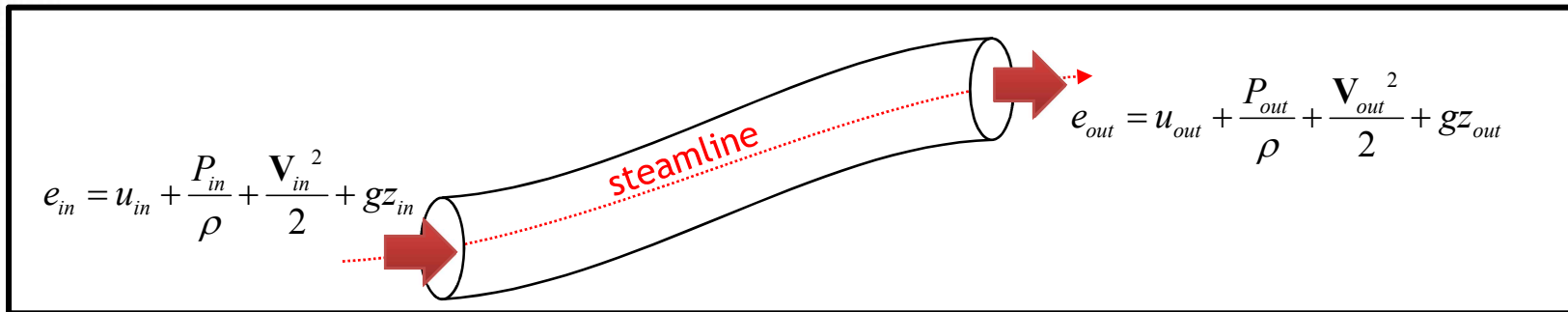
- ✓ 정상 상태 (steady state)라고 가정하면, CV 내부의 에너지 변화 없음

$$\frac{d}{dt} \int_{CV} e \rho dV + \dot{m}_{out} \left( u_{out} + \left( \frac{P}{\rho} \right)_{out} + \frac{V_{out}^2}{2} + gz_{out} \right) - \dot{m}_{in} \left( u_{in} + \left( \frac{P}{\rho} \right)_{in} + \frac{V_{in}^2}{2} + gz_{in} \right) = \dot{Q} + \dot{W}$$

- ✓ 유선을 따라 이동하는 경우, 유선을 포함하는 관을 CV으로 선택하면, 질량 유량이 동일함 ( $\dot{m}_{out} = \dot{m}_{in} = \dot{m}$ )

$$\dot{m} \left( u_{out} + \left( \frac{P}{\rho} \right)_{out} + \frac{V_{out}^2}{2} + gz_{out} \right) - \dot{m} \left( u_{in} + \left( \frac{P}{\rho} \right)_{in} + \frac{V_{in}^2}{2} + gz_{in} \right) = \dot{Q} + \dot{W}$$

$$\dot{m} \left( u_{out} - u_{in} + \frac{P_{out}}{\rho} - \frac{P_{in}}{\rho} + \frac{V_{out}^2 - V_{in}^2}{2} + g(z_{out} - z_{in}) \right) = \dot{Q} + \dot{W}$$



# 에너지 방정식과 Bernoulli equation 비교

✓ 에너지 방정식에서 외부로부터의 열의 공급이나 일이 없다고 하면,

$$\dot{m} \left( u_{out} - u_{in} + \frac{P_{out}}{\rho} - \frac{P_{in}}{\rho} + \frac{V_{out}^2 - V_{in}^2}{2} + g(z_{out} - z_{in}) \right) = \dot{Q} + \dot{W} = 0$$

$$u_{out} - u_{in} + \frac{P_{out}}{\rho} - \frac{P_{in}}{\rho} + \frac{V_{out}^2 - V_{in}^2}{2} + g(z_{out} - z_{in}) = 0$$

$$\frac{P_{out}}{\rho} + \frac{V_{out}^2}{2} + gz_{out} = \frac{P_{in}}{\rho} + \frac{V_{in}^2}{2} + gz_{in} - (u_{out} - u_{in}) \rightarrow \text{내부에너지 변화(마찰에 의한 에너지 손실)}$$

➔ steady, incompressible Bernoulli equation

➔ 마찰에 의한 에너지 손실이 없다고 가정하면,

$$\frac{P_{out}}{\rho} + \frac{V_{out}^2}{2} + gz_{out} = \frac{P_{in}}{\rho} + \frac{V_{in}^2}{2} + gz_{in}$$





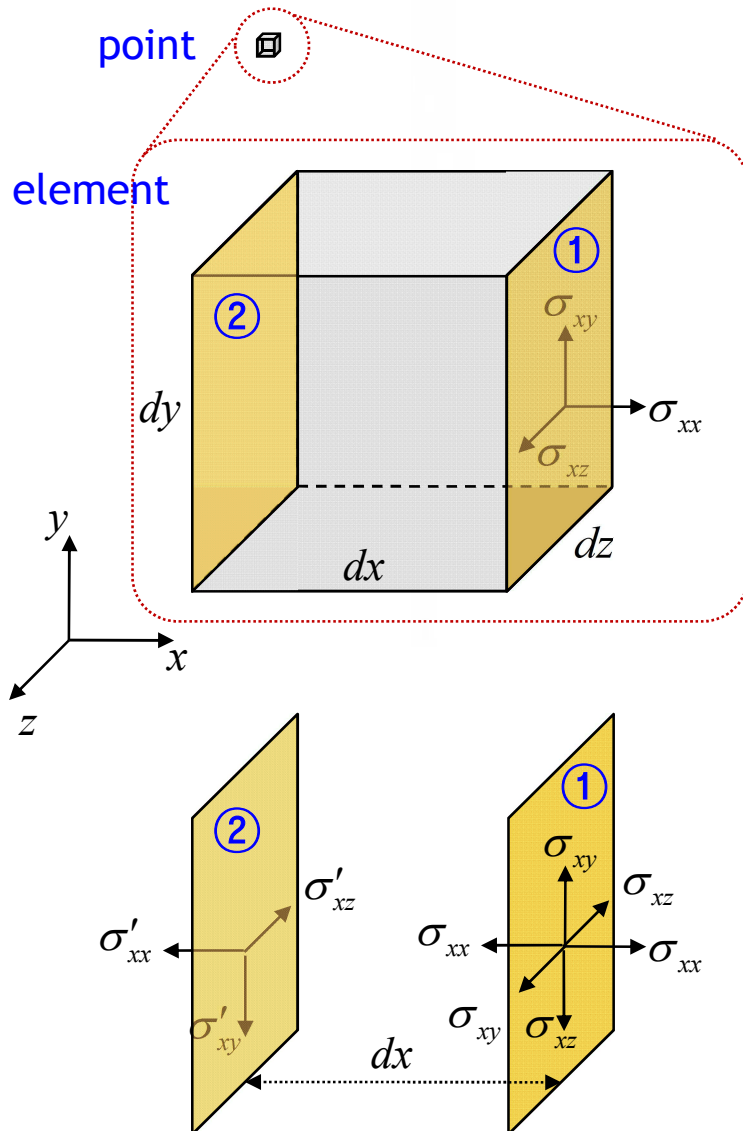
# 보충 슬라이드

**A**dvanced  
**S**hip  
**D**esign  
**A**utomation  
**L**aboratory

# Cauchy Equation<sup>1),2)</sup> 유도

## - Stress at a point

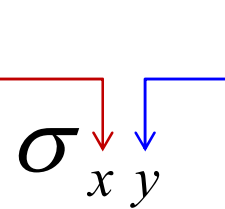
1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp.396-401  
 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.30-32, 88-93



✓ 한 점에서의 응력 (Stress at a point)

➔ 어느 면에 어느 방향으로 힘이 작용하는지 정의해야 함

Orientation of the surface      Direction of the force      ➔ 2개의 방향 필요



➔ 한 면에 수직인 성분과 2개의 접선 방향 성분으로 나타남 (3개의 응력 성분이 존재함)

➔ 면 ②에 작용하는 응력 성분은 ?

면 ②에서는 Normal vector의 방향이 면 ①과 반대이므로, 응력의 방향을 반대로 정의함 (sign convention)

면 ①과 ②는  $dx$ 만큼 떨어져 있기 때문에 일반적으로 응력이 다름

point이므로  $dx \rightarrow 0 \rightarrow$  즉, 두 면은 서로 일치함.

따라서, Point에서 두 면의 응력의 크기는 같고 방향은 반대

# Bernoulli Equation 유도

$$\text{curl } \mathbf{V} = \left( \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} \right) \mathbf{i} + \left( \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \right) \mathbf{j} + \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) \mathbf{k} = 0$$

$$\frac{\partial w}{\partial y} = \frac{\partial v}{\partial z}, \quad \frac{\partial u}{\partial z} = \frac{\partial w}{\partial x}, \quad \frac{\partial v}{\partial x} = \frac{\partial u}{\partial y}$$

$$(\mathbf{V} = \nabla \Phi)$$

✓ Euler Equation

$$\rho \frac{d\mathbf{V}}{dt} = \rho \mathbf{g} - \nabla P \quad \xrightarrow{\text{x축 성분}} \quad \rho \left( \frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} + w \frac{\partial u}{\partial z} \right) = \rho g_x - \frac{\partial P}{\partial x}$$

$$\downarrow \quad u = \frac{\partial \Phi}{\partial x}, \quad \frac{\partial u}{\partial z} = \frac{\partial w}{\partial x}, \quad \frac{\partial u}{\partial y} = \frac{\partial v}{\partial x} \quad \text{대입}$$

$$\rho \left( \frac{\partial}{\partial t} \left( \frac{\partial \Phi}{\partial x} \right) + u \frac{\partial u}{\partial x} + v \frac{\partial v}{\partial x} + w \frac{\partial w}{\partial x} \right) = \rho g_x - \frac{\partial P}{\partial x}$$

$$\downarrow \quad \rho \left( \frac{\partial}{\partial x} \left( \frac{\partial \Phi}{\partial t} \right) + \frac{1}{2} \frac{\partial}{\partial x} (u^2) + \frac{1}{2} \frac{\partial}{\partial x} (v^2) + \frac{1}{2} \frac{\partial}{\partial x} (w^2) \right) = \rho g_x - \frac{\partial P}{\partial x}$$

$$\downarrow \quad \rho \left( \frac{\partial}{\partial x} \left( \frac{\partial \Phi}{\partial t} \right) + \frac{1}{2} \frac{\partial}{\partial x} (u^2 + v^2 + w^2) \right) = \rho g_x - \frac{\partial P}{\partial x}$$

# Bernoulli Equation<sup>1)</sup> 유도

$$\text{(continue)} \quad \rho \left( \frac{\partial}{\partial x} \left( \frac{\partial \Phi}{\partial t} \right) + \frac{1}{2} \frac{\partial}{\partial x} (u^2 + v^2 + w^2) \right) = \rho g_x - \frac{\partial P}{\partial x}$$

양변을 x로 적분해 주면

$$\int \rho \left[ \frac{\partial}{\partial x} \left( \frac{\partial \Phi}{\partial t} \right) + \frac{1}{2} \frac{\partial}{\partial x} (u^2 + v^2 + w^2) \right] dx = \int \left[ \rho g_x - \frac{\partial P}{\partial x} \right] dx$$

$$\rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) = \rho g_x x - P + f_1(y, z, t)$$

$$\rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) + P = \rho g_x x + f_1(y, z, t)$$

# Bernoulli Equation<sup>1)</sup> 유도

$$\begin{aligned}
 x\text{축 성분} &: \rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) + P = \rho g_x x + f_1(y, z, t) \\
 y\text{축 성분} &: \rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) + P = \rho g_y y + f_2(x, z, t) \\
 z\text{축 성분} &: \rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) + P = \rho g_z z + f_3(x, y, t)
 \end{aligned}$$

① 좌변이 모두 같으므로, 우변도 같다.

$$\rho g_x x + f_1(y, z, t) = \rho g_y y + f_2(x, z, t) = \rho g_z z + f_3(x, y, t)$$

② 중력장만을 고려한다면,  $g_x = g_y = 0, g_z = -g$  라 할 수 있다.

$$f_1(y, z, t) = f_2(x, z, t) = -\rho g z + f_3(x, y, t)$$

③  $f_1$ 과  $f_2$ 를 비교하면,  $x, y$ 의 함수가 아님을 알 수 있다.  $f_3$ 은  $t$ 의 함수가 된다.

$$f_1(z, t) = f_2(z, t) = -\rho g z + f(t)$$

④ 원래 식에 대입하면,

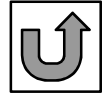
$$\rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho (u^2 + v^2 + w^2) + P + \rho g z = f(t) \longrightarrow$$

⑤ 속도항을 속도 포텐셜로 표현하면,

$$\rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho |\nabla \Phi|^2 + P + \rho g z = f(t)$$

=> Bernoulli Equation

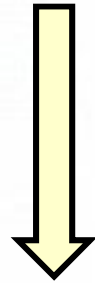




# Bernoulli Equation<sup>1)</sup>

$$\rho \frac{\partial \Phi}{\partial t} + \frac{1}{2} \rho |\nabla \Phi|^2 + P + \rho g z = f(t)$$

=> Bernoulli Equation



정상 유동 (Steady flow)라 하면,

$$\frac{\partial \Phi}{\partial t} = 0, \quad f(t) = C \quad (C: \text{const})$$

$$\frac{1}{2} \rho |\nabla \Phi|^2 + P + \rho g z = C$$

Kinetic Energy + Pressure Energy + Potential Energy = Constant

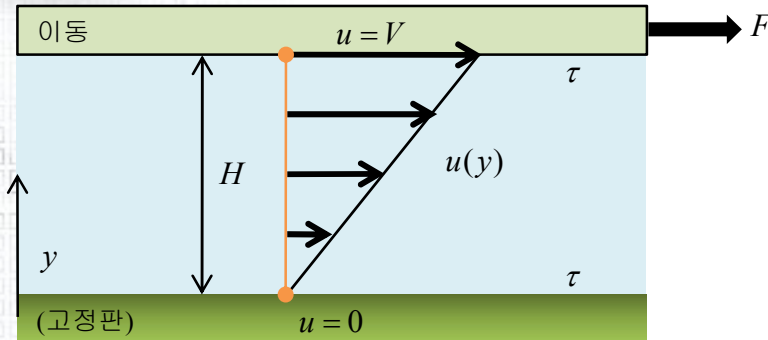
=> Mechanical Energy conservation

# 뉴턴 유체<sup>1)</sup> (Newtonian Fluid)

- 1) Cengel & Cimbala, Fluid Mechanics Mc Graw Hill, 2005, pp46-50, pp426-431
- 2) Kundu, P.K., Fluid Mechanics, Academic Press, 2008, pp.100-105
- 3) 김찬중, "공학도를 위한 길잡이 유체공학입문", 문운당, 2002, pp258-264

## ✓ 뉴턴 유체<sup>1)</sup> (Newtonian Fluid)

A: 접촉면적



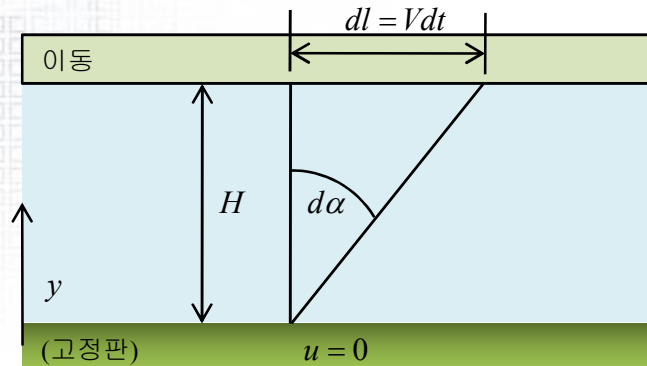
No-Slip condition  
정상상태

$$\text{전단응력: } \tau = \frac{F}{A}$$

$$\text{유체의 속도분포: } u(y) = \frac{V}{H} y$$

$$\text{유체의 속도변화: } \frac{du}{dy} = \frac{V}{H}$$

시간이 무한이 지나면,  
높이에 따라 속도 분포가  
선형적으로 변함



(assump.  $d\alpha \ll 1$ )

$$\text{전단변형율 (shear strain rate): } d\alpha \approx \tan d\alpha = \frac{dl}{H} = \frac{V}{H} dt = \frac{du}{dy} dt$$

$$\text{전단변형율 (의 시간변화율)*: } \frac{d\alpha}{dt} = \frac{du}{dy}$$

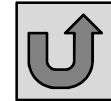
$$\text{전단응력은 전단변형율에 비례함: } \tau \propto \frac{d\alpha}{dt}$$

(실험적으로 증명 가능함 - 구체적으로)

$$\text{즉, 전단응력은 속도 구배에 비례함: } \tau \propto \frac{du}{dy} \rightarrow \tau = \mu \frac{du}{dy}$$

(Newtonian Fluid)

\*[김찬중, 유체공학입문, 문운당, 2004]에서는 이것을 전단변형율의 시간변화율이라고 명기  
[Cengel, Fluid Mechanics, Mc Graw Hill, 2006], [Whit, Fluid Mechanics 6th edition]은 전단변형율이라는 용어사용



# Bernoulli equation의 상수 'C' 의 의미와 계기 압력<sup>1)</sup>

## Bernoulli Equation

$$\rho \frac{\partial \Phi}{\partial t} + P + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = C$$

↪ 항등식

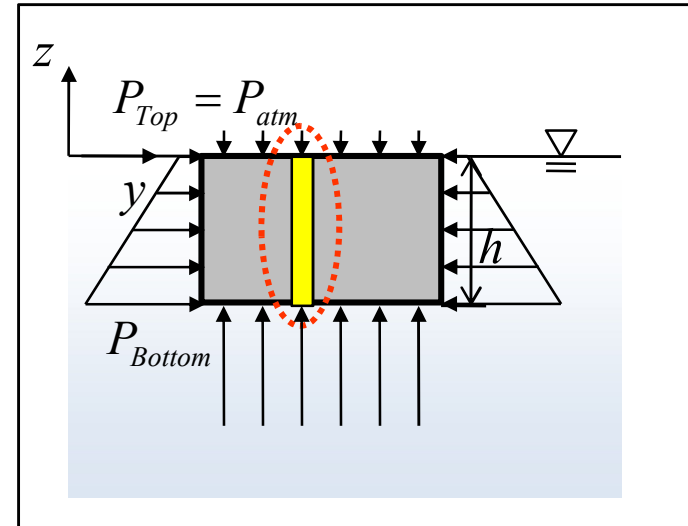
✓ 정적인 상태 :  $\Phi = 0 \left( \frac{\partial \Phi}{\partial t} = 0, \nabla \Phi = 0 \right)$

$$P + \rho g z = C$$

항등식 이므로,  $z = 0$  일 때도 성립

$$\frac{P_{atm}}{z=0 \text{ 일 때 압력}} = C = \text{대기압}(P_{atm})$$

$$\therefore \rho \frac{\partial \Phi}{\partial t} + P + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = P_{atm}$$



✓ 물체 바닥에서의 압력은?

$$\rho \frac{\partial \Phi}{\partial t} + P_{Bottom} + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = P_{atm}$$

$$\rho \frac{\partial \Phi}{\partial t} + (P_{atm} + P_{Fluid}) + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = P_{atm}$$

$$\therefore \rho \frac{\partial \Phi}{\partial t} + P_{Fluid} + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = 0$$

‘계기 압력’

※ Bernoulli equation에서 우변의 상수 C=0으로 표기한 경우, 압력 P는 대기압을 뺀 유체만의 압력임을 의미함

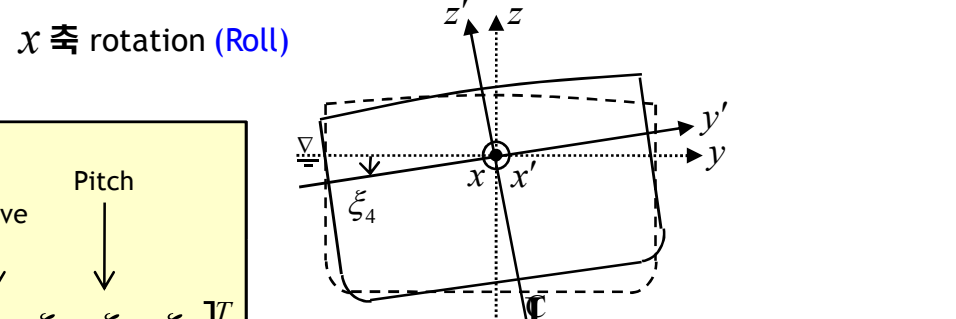
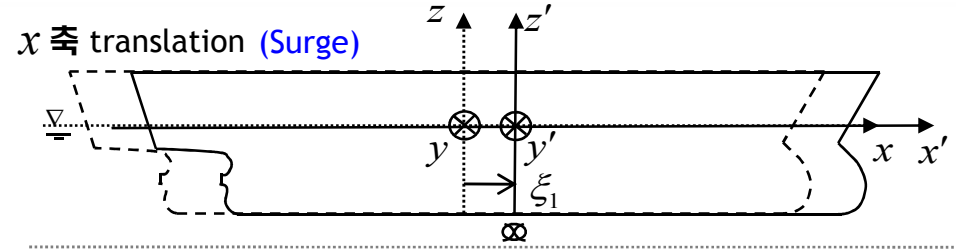
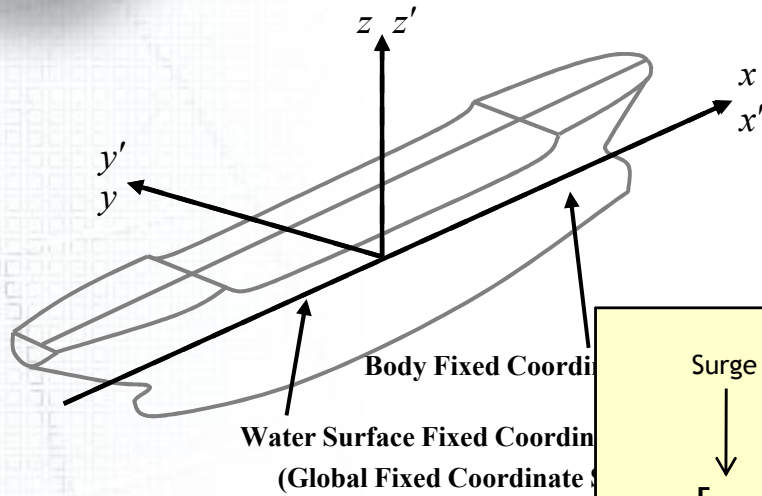


# 선박의 6자유도 운동 방정식 (6DOF Equations of Ship Motion)

# Coordinate System

$x'$  축 - 원점: Midship, (+): 선수  
 $y'$  축 - 원점: Centerline, (+): 좌현  
 $z'$  축 - 원점: 수선면, (+): 선박의 위

$x$  축 - 원점: Midship, (+):  $x'$  축을 포함하고 수선면과 직교인 평면과 수선면 사이의 교선  
 $y$  축 - 원점: Centerline, (+):  $z$  축과  $x$  축의 외적 방향  
 $z$  축 - 원점: 수면, (+): 수선면에 수직한 위 방향



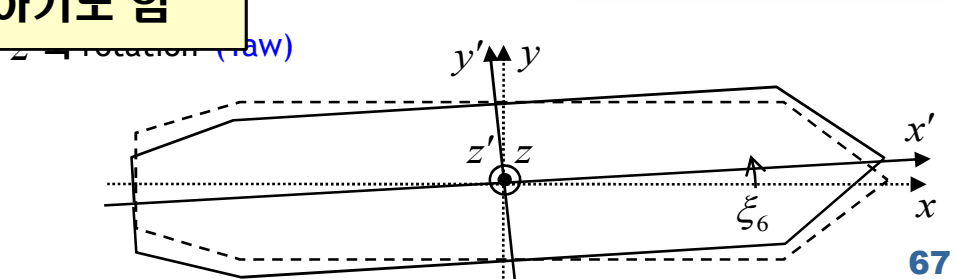
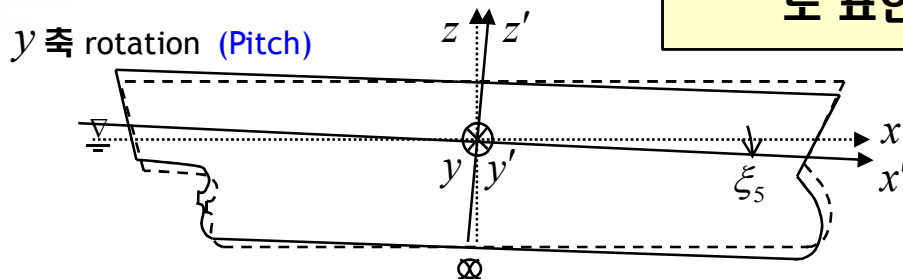
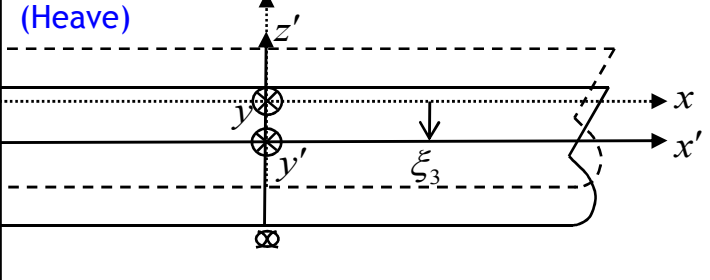
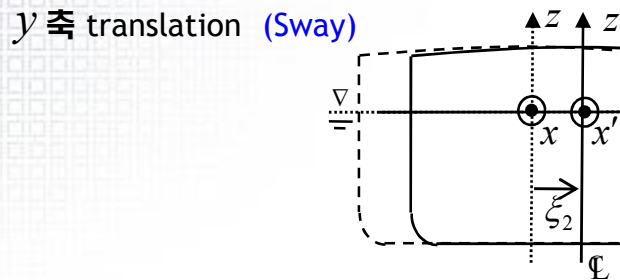
Surge	Heave	Pitch
↓	↓	↓
$\xi_1$	$\xi_2$	$\xi_3$
↑ Sway	↑ Roll	↑ Yaw

$$\mathbf{X} = [\xi_1, \xi_2, \xi_3, \xi_4, \xi_5, \xi_6]^T$$


---


$$\mathbf{x} = [x, y, z, \phi, \theta, \psi]^T$$

**로 표현하기도 함**





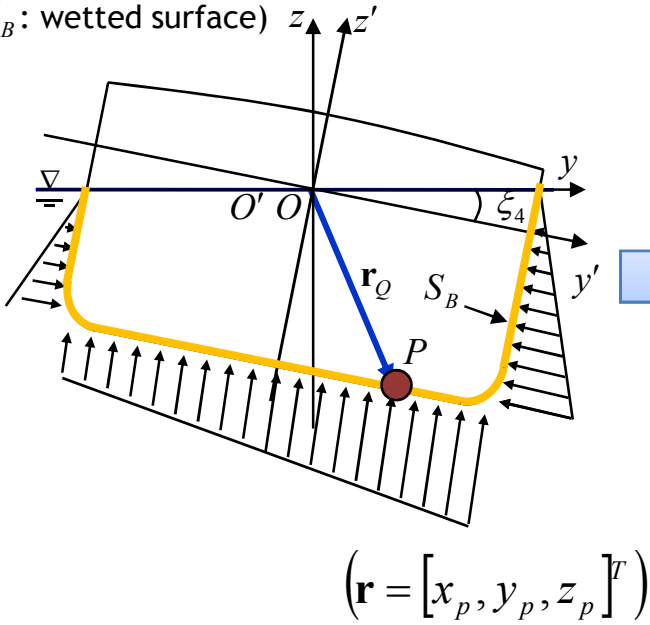
# 선박의 6자유도 운동 방정식

: Force & moment acting on the surface

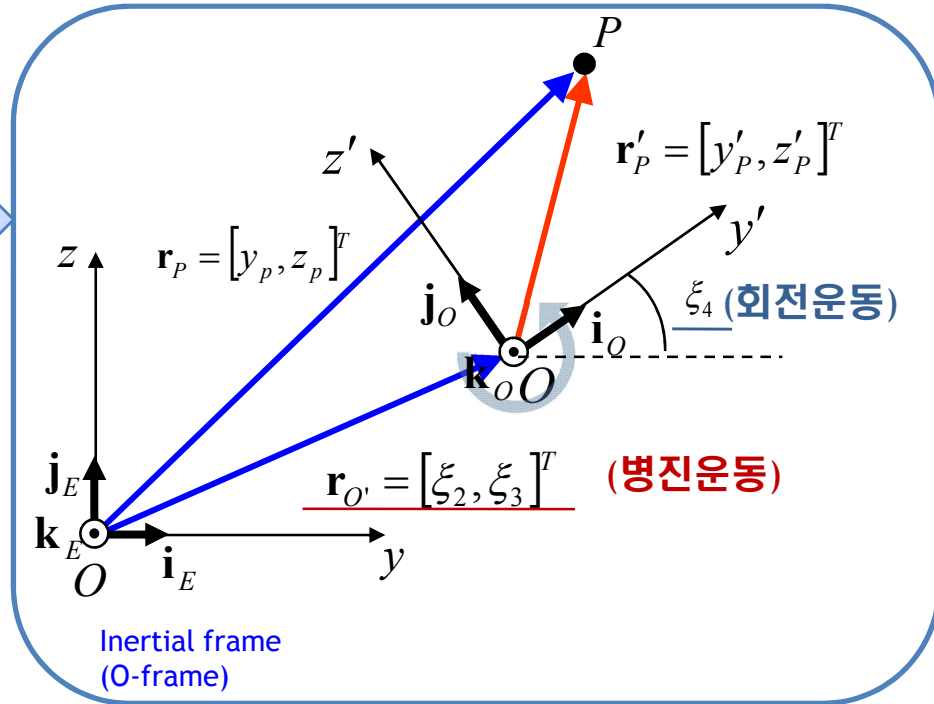
( $S_B$ : wetted surface)

좌현으로 기울어진 상태  
(선박을 정면에서 바라봄)

( $S_B$ : wetted surface)



✓ 지구 고정 좌표계  $O-xyz$  에서 표현된 선박 표면의 점 P와  
선박 운동 변위  $\mathbf{x} = [\xi_1, \xi_2, \xi_3, \xi_4, \xi_5, \xi_6]^T$  의 관계

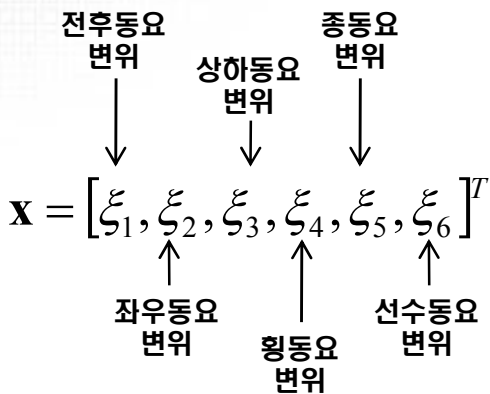


$$\mathbf{r}_p = \mathbf{r}_{O'} + Rot(\xi_4)\mathbf{r}'_p$$

$$\begin{bmatrix} y_p \\ z_p \end{bmatrix} = \begin{bmatrix} \xi_2 \\ \xi_3 \end{bmatrix} + \begin{bmatrix} \cos \xi_4 & -\sin \xi_4 \\ \sin \xi_4 & \cos \xi_4 \end{bmatrix} \begin{bmatrix} y'_p \\ z'_p \end{bmatrix}$$

(병진운동) (회전운동)

지구고정 좌표계  
 $O-xyz$  에서  
선박 표면의 점 P는  
운동 변위로 표현됨



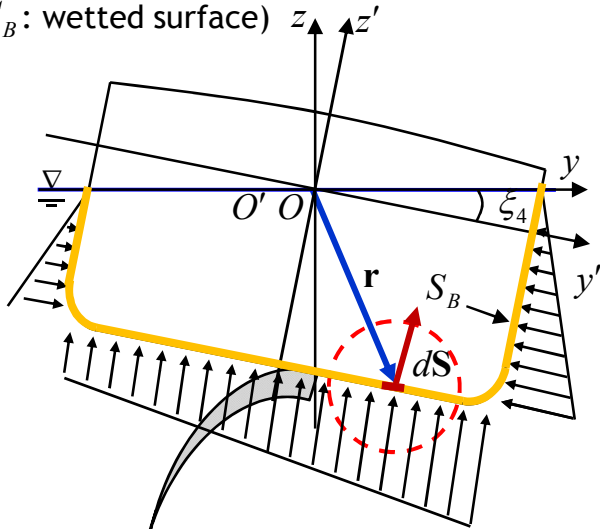
# 선박의 6자유도 운동 방정식

## : Force & moment acting on the surface

( $S_B$ : wetted surface)

좌현으로 기울어진 상태  
(선박을 정면에서 바라봄)

( $S_B$ : wetted surface)



(미소 면적에 작용하는 힘)

$$d\mathbf{F} = P d\mathbf{S} = P \mathbf{n} dS$$

$$(\mathbf{n} = [n_1, n_2, n_3]^T)$$

$$d\mathbf{M} = \mathbf{r}_p \times d\mathbf{F}$$

(미소 면적에 작용하는 모멘트)

$$(\mathbf{r} = [x, y, z]^T)$$

Force : 표면에 작용하는 모든 힘을 적분하여 구함

✓ 미소 면적에 작용하는 단위 길이당 힘 :

$$d\mathbf{F} = P \cdot d\mathbf{S} = P \cdot \mathbf{n} dS = -\rho g z \cdot \mathbf{n} dS$$

✓ Total force

$$\mathbf{F} = \iint_{S_B} P \mathbf{n} dS$$

Moment : (모멘트) = (거리) X (힘)

✓ 미소 면적에 작용하는 단위 길이당 모멘트 :

$$d\mathbf{M} = \mathbf{r} \times d\mathbf{F} = \mathbf{r} \times P \mathbf{n} dS = (\mathbf{r} \times \mathbf{n}) P dS$$

✓ Total moment

$$\mathbf{M} = \iint_{S_B} P (\mathbf{r} \times \mathbf{n}) dS$$

왜 r이 먼저 오는가? (좌표축에서 양의 방향을 고려함)

# 선박의 6자유도 운동 방정식

## :Notation

$$\mathbf{r} \times \mathbf{n} = \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ x & y & z \\ n_1 & n_2 & n_3 \end{vmatrix} = \mathbf{i}(yn_3 - zn_2) + \mathbf{j}(zn_1 - xn_3) + \mathbf{k}(xn_2 - yn_1)$$

✓ Fluid force acting on the surface

$$\mathbf{F} = \iint_{S_B} P \mathbf{n} dS$$

$(\mathbf{n} = [n_1, n_2, n_3]^T)$

성분별로 나눠쓰면,

$$\begin{cases} F_1 = \iint_{S_B} P n_1 dS \\ F_2 = \iint_{S_B} P n_2 dS \\ F_3 = \iint_{S_B} P n_3 dS \end{cases}$$

$$F_j = \iint_{S_B} P n_j dS$$

$(j = 1, \dots, 6)$

✓ Fluid moment acting on the surface

$$\mathbf{M} = \iint_{S_B} P (\mathbf{r} \times \mathbf{n}) dS$$

$(\mathbf{n} = [n_1, n_2, n_3]^T)$   
 $(\mathbf{r} = [x, y, z]^T)$

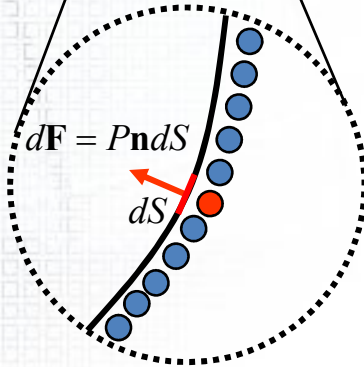
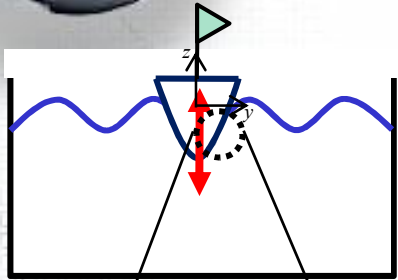
성분별로 나눠쓰면,

$$\begin{cases} M_1 = \iint_{S_B} P (yn_3 - zn_2) dS, & \left\{ F_4 = \iint_{S_B} P n_4 dS \right. \\ M_2 = \iint_{S_B} P (zn_1 - xn_3) dS, & \left\{ F_5 = \iint_{S_B} P n_5 dS \right. \\ M_3 = \iint_{S_B} P (xn_2 - yn_1) dS, & \left\{ F_6 = \iint_{S_B} P n_6 dS \right. \end{cases}$$

# 선박의 6자유도 운동 방정식 유도

$F_{F.K}$ : Froude- krylov force  
 $F_D$ : Diffraction force  
 $F_R$ : Radiation force

$\Phi_I$  : Incident wave velocity potential  
 $\Phi_D$  : Diffraction velocity potential  
 $\Phi_R$  : Radiation velocity potential



$d\mathbf{F}$  : 하나의 유체 입자가  
선박 표면에 가하는 힘

$dS$  : 미소 면적

$\mathbf{n}$  : 미소 면적의 Normal 벡터

✓ Bernoulli Equation

$$\rho \frac{\partial \Phi}{\partial t} + P + \frac{1}{2} \rho |\nabla \Phi|^2 + \rho g z = 0$$

Linearization

✓ Laplace Equation

$$\nabla^2 \Phi = 0$$

$$\Phi = \Phi_I + \Phi_D + \Phi_R$$

Linear combination  
of basic solutions

Basic solutions

$$P_{Fluid} = -\rho g z - \rho \frac{\partial \Phi}{\partial t} = \boxed{-\rho g z} - \rho \left( \frac{\partial \Phi_I}{\partial t} + \frac{\partial \Phi_D}{\partial t} + \frac{\partial \Phi_R}{\partial t} \right)$$

유체입자가 표면에  
작용하는 압력

$$= \boxed{P_{static}} + \underbrace{P_{F.K} + P_D + P_R}_{P_{dynamic}}$$

$$\mathbf{F}_{Fluid} = \iint_{S_B} P_{Fluid} \mathbf{n} dS = \mathbf{F}_{static} + \mathbf{F}_{F.K} + \mathbf{F}_D + \mathbf{F}_R$$

(  $S_B$ : wetted surface)

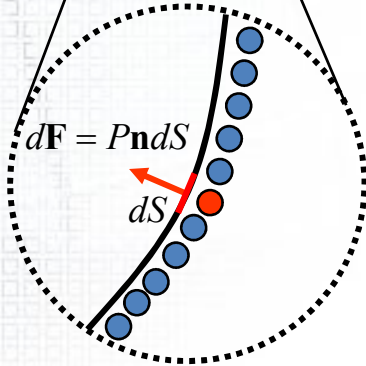
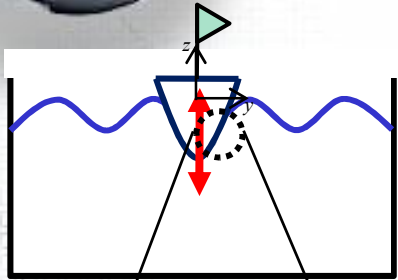
선박의 침수 표면 전체에 대하여 적분 (표면력)  
 (유체입자가 선박에 작용하는 힘과 모멘트)

✓ 유체입자 하나에 작용하는 Body force 와 Surface force로부터 구한 압력을 선박의 침수 표면 전체에 대해 적분하여 선박에 작용하는 유체력을 계산함

# 선박의 6자유도 운동 방정식 유도

$F_{F.K}$ : Froude- krylov force  
 $F_D$ : Diffraction force  
 $F_R$ : Radiation force

$\Phi_I$ : Incident wave velocity potential  
 $\Phi_D$ : Diffraction velocity potential  
 $\Phi_R$ : Radiation velocity potential



$dF$ : 하나의 유체 입자가 선박 표면에 가하는 힘

$dS$ : 미소 면적

$n$ : 미소 면적의 Normal 벡터

$$\mathbf{x} = [\xi_1, \xi_2, \xi_3, \xi_4, \xi_5, \xi_6]^T$$

$\xi_1$ : surge     $\xi_3$ : roll

$\xi_2$ : sway     $\xi_4$ : pitch

$\xi_5$ : heave     $\xi_6$ : yaw

$M_4$ : 6×6 added mass matrix

$B$ : 6×6 damping coeff. matrix

$C$ : 6×6 restoring coeff. matrix

## ✓ 선박에 작용하는 유체력

$$\mathbf{F}_{Fluid} = \iint_{S_B} P \mathbf{n} dS = \mathbf{F}_{static} + \mathbf{F}_{F.K} + \mathbf{F}_D + \mathbf{F}_R$$

: 유체입자 하나에 작용하는 (Body force) 와 (Surface force)로부터 구한 압력을 선박의 침수 표면 전체에 대해 적분하여 선박에 작용하는 유체력을 계산함

## ✓ 선박의 6자유도 운동방정식

### Newton's 2<sup>nd</sup> Law

선박의 Surface force로 작용

$$M\ddot{\mathbf{x}} = \sum \mathbf{F} = (\text{Body Force}) + (\text{Surface Force})$$

$$= \mathbf{F}_{gravity} + \mathbf{F}_{Fluid} + \mathbf{F}_{external}$$

Body force    Surface force

기타 외부에서 작용하는 외력

$$M\ddot{\mathbf{x}} = \underbrace{\mathbf{F}_{Gravity}}_{\mathbf{F}_{Restoring}} + \underbrace{\mathbf{F}_{static} + \mathbf{F}_{F.K} + \mathbf{F}_D + \mathbf{F}_R}_{\mathbf{F}_{wave\ exciting}} + \mathbf{F}_{external, dynamic} + \mathbf{F}_{external, static}$$

$$\mathbf{F}_R = -\mathbf{A}\ddot{\mathbf{x}} - \mathbf{B}\dot{\mathbf{x}}$$

added mass
Damping Coefficient

$$M\ddot{\mathbf{x}} = (\mathbf{F}_{gravity} + \mathbf{F}_{static}) + \mathbf{F}_{wave\ exciting} - \mathbf{A}\ddot{\mathbf{x}} - \mathbf{B}\dot{\mathbf{x}} + \mathbf{F}_{external, dynamic} + \mathbf{F}_{external, static}$$

Linearization,  $(\mathbf{F}_{restoring} = (\mathbf{F}_{gravity} + \mathbf{F}_{static}) \approx -\mathbf{C}\mathbf{x})$

복원력 계수 C는 어떻게 구할까?

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{wave\ exciting} + \mathbf{F}_{external, dynamic} + \mathbf{F}_{external, static}$$

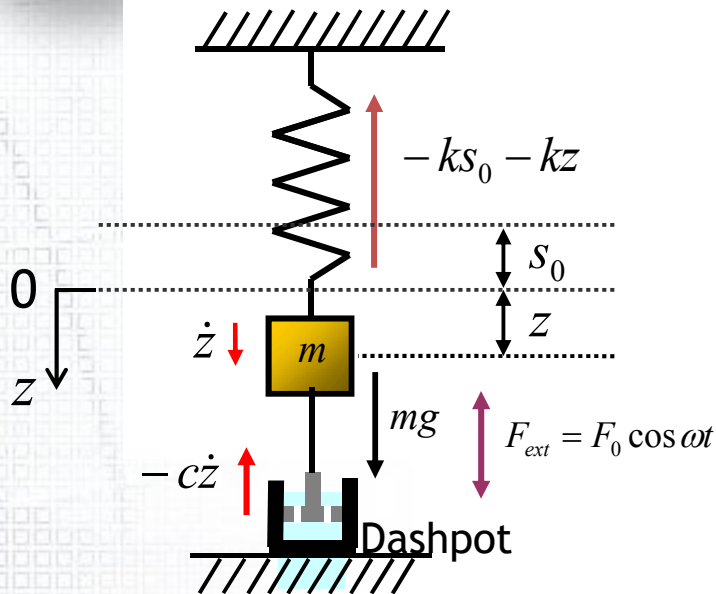


# 선박의 6자유도 운동 방정식

## : Heave 운동에서의 복원력 계수 C구하기

$$\mathbf{F}_{restoring} = (\mathbf{F}_{gravity} + \mathbf{F}_{static}) \approx -\mathbf{C}\mathbf{x}$$

✓ Mass-Spring-Damper system

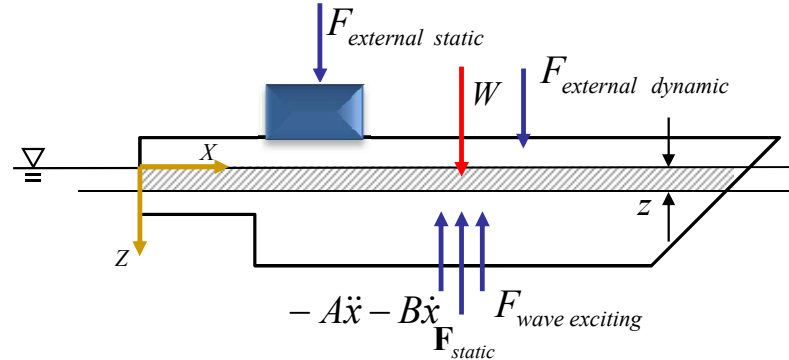


$$\begin{aligned} m\ddot{z} &= F \\ &= mg - ks_0 - kz - c\dot{z} + F_0 \cos \omega t \\ &= -kz - c\dot{z} + F_0 \cos \omega t \end{aligned}$$

$$m\ddot{z} + c\dot{z} + kz = F_0 \cos \omega t$$

$\rho$ : density of sea water  
 $A_{WP}$ : Waterplane Area

✓ 선박의 heave 운동방정식 예



$$\mathbf{M}\ddot{z} = \sum \mathbf{F} = \underbrace{(\mathbf{F}_{Gravity} + \mathbf{F}_{static})}_{\mathbf{F}_{restoring}} + \mathbf{F}_{wave\ exciting} - \mathbf{A}\ddot{z} - \mathbf{b}\dot{z} + \mathbf{F}_{external,dynamic} + \mathbf{F}_{external,static}$$

✓ 선박의 heave 운동에서의 복원력의 예

$$\mathbf{F}_{restoring} = (\mathbf{F}_{Gravity} + \mathbf{F}_{static})$$

변위에 대해 반대방향으로 부력이 작용함을 의미

$$\mathbf{F}_{gravity} = Mg\mathbf{k}$$

M: 선박 중량

$$\mathbf{F}_{static} = (-\rho g V_0 - \rho g A_{WP} \cdot z)\mathbf{k}$$

: heave 변위 z에 의한 추가 부력

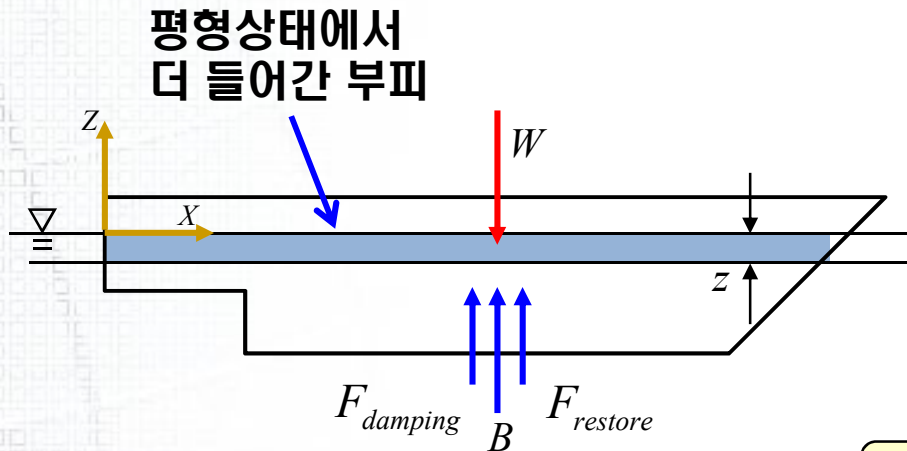
$$\mathbf{F}_{restoring} = (Mg - \rho g V_0 - \rho g A_{WP} \cdot z)\mathbf{k}$$

$$= -(\rho g A_{WP}) \cdot z\mathbf{k} = -C z\mathbf{k} \quad , (C = \rho g A_{WP})$$

# 선박의 6자유도 운동 방정식

## : Uncoupled equation of motion

✓ Heave 운동 방정식 유도



$$\begin{aligned}
 M\ddot{\xi}_3 &= \sum F = F_{Body} + F_{Surface} \\
 &= F_{gravity} + F_{Fluid} \\
 &= \underbrace{F_{gravity}}_{-Mg} + \underbrace{F_{static}}_{\rho g V_0 - \rho g A_{wp} \cdot \xi_3} + \underbrace{F_{F.K} + F_D}_{F_{exciting,3}} + \underbrace{F_R}_{-A_{33}\ddot{\xi}_3 - B_{33}\dot{\xi}_3} \\
 &= -Mg + (\rho g V_0 - \rho g A_{wp} \xi_3) + F_{exciting,3} - A_{33}\ddot{\xi}_3 - B_{33}\dot{\xi}_3
 \end{aligned}$$

$$\therefore (M + A_{33})\ddot{\xi}_3 + B_{33}\dot{\xi}_3 + \rho g A_{wp} \xi_3 = F_{exciting,3}$$

✓ Surge 운동 방정식 유도  
(Heave에서 복원력 성분만 제외)

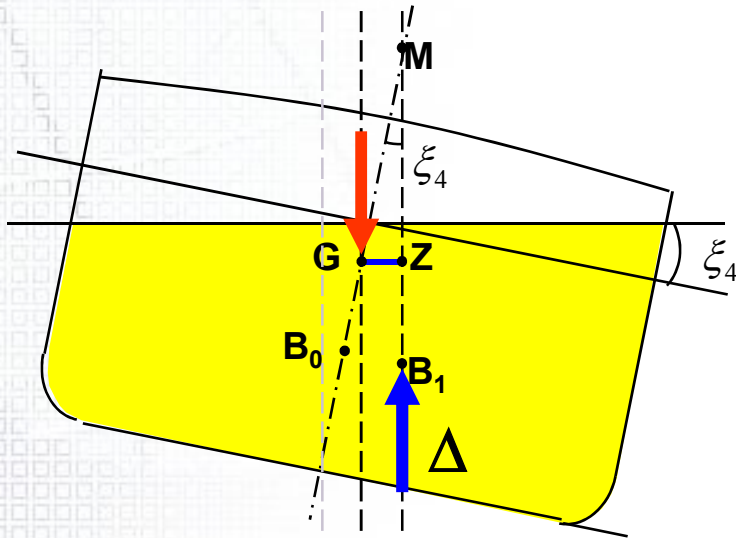
$$\therefore (M + A_{11})\ddot{\xi}_1 + B_{11}\dot{\xi}_1 = F_{exciting,1}$$

✓ Sway 운동 방정식 유도  
(Heave에서 복원력 성분만 제외)

$$\therefore (M + A_{22})\ddot{\xi}_2 + B_{22}\dot{\xi}_2 = F_{exciting,2}$$

# 선박의 6자유도 운동 방정식 : Uncoupled equation of motion

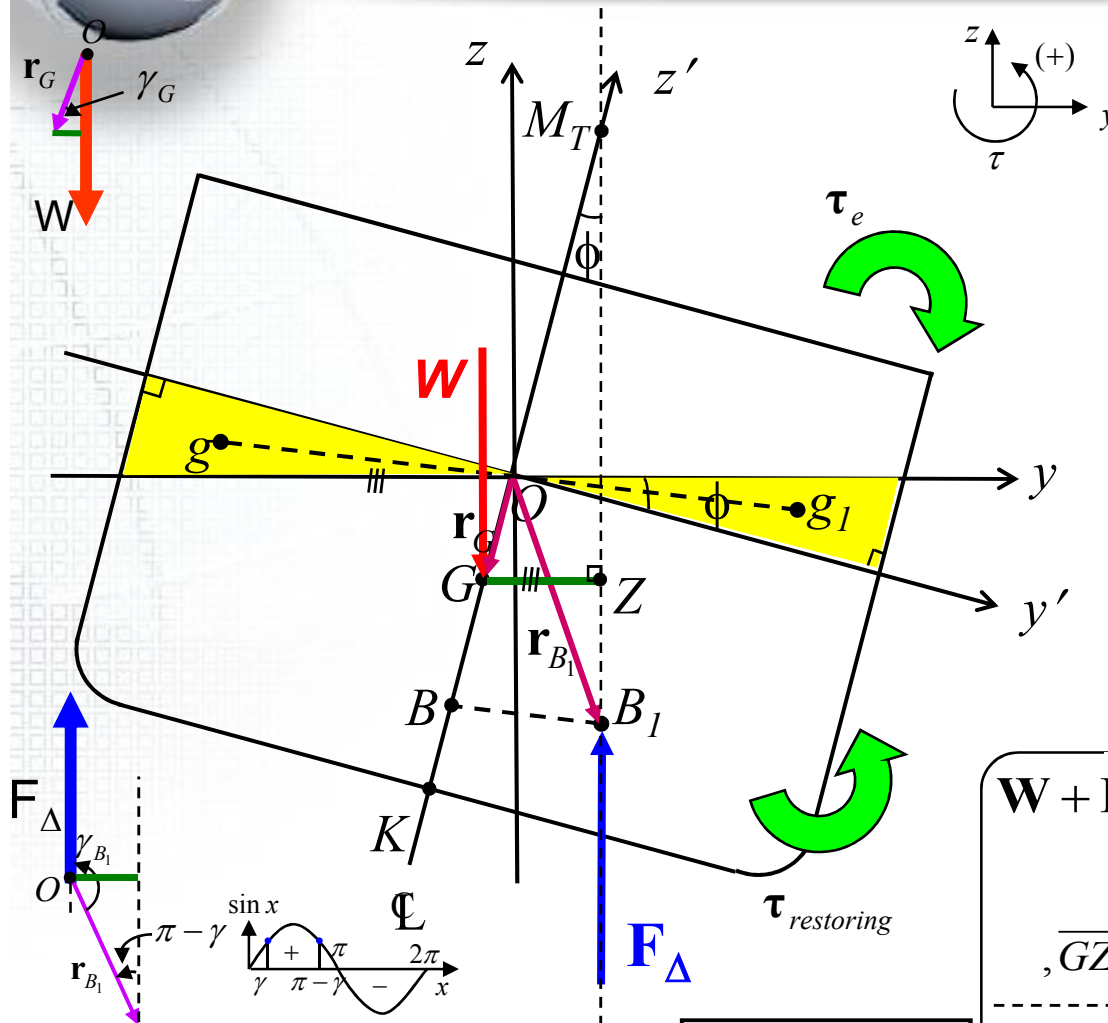
✓ Roll 운동 방정식 유도



$$\begin{aligned}
 I_{xx} \ddot{\xi}_4 &= \sum M = M_{Body} + M_{Surface} \\
 &= M_{gravity} + M_{Fluid} \\
 &= \underbrace{M_{gravity}}_{\mathbf{r}_G \times W\mathbf{k}} + \underbrace{M_{static}}_{\mathbf{r}_{B_1} \times \Delta\mathbf{k}} + \underbrace{M_{F.K} + M_D}_{M_{exciting,4}} + \underbrace{M_R}_{-A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4} \\
 &= \Delta \cdot \overline{GZ} + M_{exciting,4} - A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4
 \end{aligned}$$

# 5. 횡 복원력

## - 선박의 횡 복원 안정성(Transv. Stability) - 안정 상태 (3)



$$\sin(\pi - \gamma) = \sin \gamma$$

$oy'z'$  : Body fixed coordinate

$oyz$  : Global coordinate

$M_T$ :  $B_1$ 을 지나는 부력 작용선과 선체 중심선과의 교점

- G: 수직방향 무게중심
- B: 수직방향 부력 중심
- W: 선박 무게
- $F_{\Delta}$ : 부력

⑥-2) 회전축  $O$ 를 지나고,  $y-z$  평면에 수직인 축에 대한 중력과 부력에 의한 모멘트를 구해보면,

$$\tau = \tau_G + \tau_{\Delta}$$

$$= \mathbf{r}_G \times \mathbf{W} + \mathbf{r}_{B_1} \times \mathbf{F}_{\Delta}$$

$$= \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ 0 & y_G & z_G \\ 0 & 0 & W_z \end{vmatrix} + \begin{vmatrix} \mathbf{i} & \mathbf{j} & \mathbf{k} \\ 0 & y_{B_1} & z_{B_1} \\ 0 & 0 & F_{\Delta,z} \end{vmatrix}$$

$$= y_G \cdot W_z \mathbf{i} + y_{B_1} \cdot F_{\Delta,z} \mathbf{i}$$

$$\mathbf{W} + \mathbf{F}_{\Delta} = 0 \Rightarrow W_z + F_{\Delta,z} = 0 \Rightarrow W_z = -F_{\Delta,z}$$

$$\tau = (-y_G + y_{B_1}) \cdot F_{\Delta,z} \mathbf{i}$$

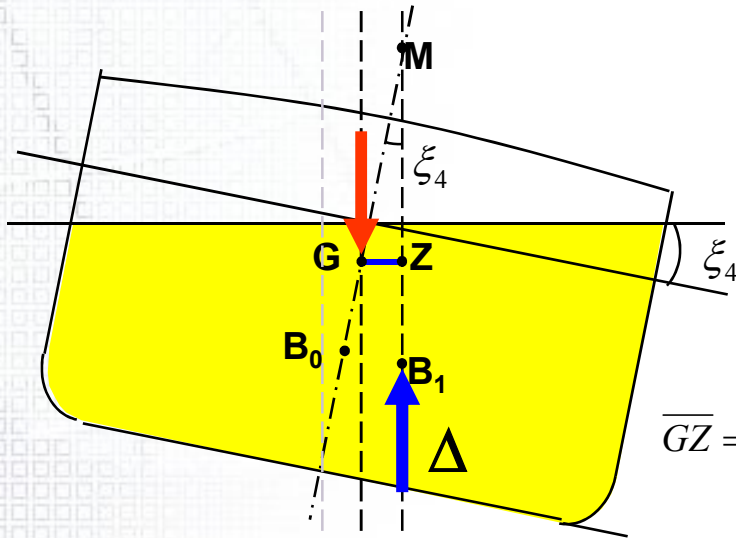
$$, \overline{GZ} = |-y_G + y_{B_1}|$$

$Z$ :  $B_1$ 을 통한 부력작용선과  $G$ 를 지나고  $y$ 축과 평행한 선이 만나는 점

왼쪽 예제의 경우,  $\tau$ 가 (+)이면(무게중심이 부력 작용선에 대해 왼쪽에 위치)  $\tau$ 는 복원 모멘트  $\tau_{restoring}$  로 작용함  $\tau = \tau_{restoring}$

# 선박의 6자유도 운동 방정식 : Uncoupled equation of motion

✓ Roll 운동 방정식 유도



$$\overline{GZ} = -\overline{GM} \sin \xi_4$$

$$\sin \xi_4 \approx \xi_4$$

$$\begin{aligned}
 I_{xx} \ddot{\xi}_4 &= \sum M = M_{Body} + M_{Surface} \\
 &= M_{gravity} + M_{Fluid} \\
 &= \underbrace{M_{gravity}}_{\mathbf{r}_G \times W\mathbf{k}} + \underbrace{M_{static}}_{\mathbf{r}_{B_1} \times \Delta\mathbf{k}} + \underbrace{M_{F.K} + M_D}_{M_{exciting,4}} + \underbrace{M_R}_{-A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4} \\
 &= \Delta \cdot \overline{GZ} + M_{exciting,4} - A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4 \\
 &= -\Delta \cdot \overline{GM} \sin \xi_4 + M_{exciting,4} - A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4 \\
 &\approx -\Delta \cdot \overline{GM} \xi_4 + M_{exciting,4} - A_{44}\ddot{\xi}_4 - B_{44}\dot{\xi}_4
 \end{aligned}$$

$$\therefore (I_{xx} + A_{44})\ddot{\xi}_4 + B_{44}\dot{\xi}_4 + \Delta \cdot \overline{GM}_T \xi_4 = M_{exciting,4}$$

✓ Pitch 운동 방정식 유도  
(Roll 운동 방정식과 동일)

$$\therefore (I_{yy} + A_{55})\ddot{\xi}_5 + B_{55}\dot{\xi}_5 + \Delta \cdot \overline{GM}_L \xi_5 = M_{exciting,5}$$

✓ Yaw 운동 방정식 유도  
(RollHeave에서 복원력 성분만 제외)

$$\therefore (I_{zz} + A_{66})\ddot{\xi}_6 + B_{66}\dot{\xi}_6 = M_{exciting,6}$$



1) Journee, J.M.J. , Adegeest, L.J.M. , Theoretical Manual of Strip Theory program " Seaway for Windows", Delft University of Technology, 2003, pp38-42  
 2) Newman, J.N. , Marine Hydrodynamics, The MIT Press, Cambridge, 1997, pp 307-311  
 3) Journee, J.M.J. , Massie, W.W. , Offshore Hydrodynamics, Delft University of Technology, 2001, pp8-1-4

# 선박의 6자유도 운동 방정식

(변위 :  $\mathbf{x} = [\xi_1, \dots, \xi_6]^T$ )  
 ( $\mathbf{M}, \mathbf{A}, \mathbf{B}, \mathbf{C}$  :  $6 \times 6$  Matrix)

✓ 6DOF Equations of Ship Motion : 6 coupled equation

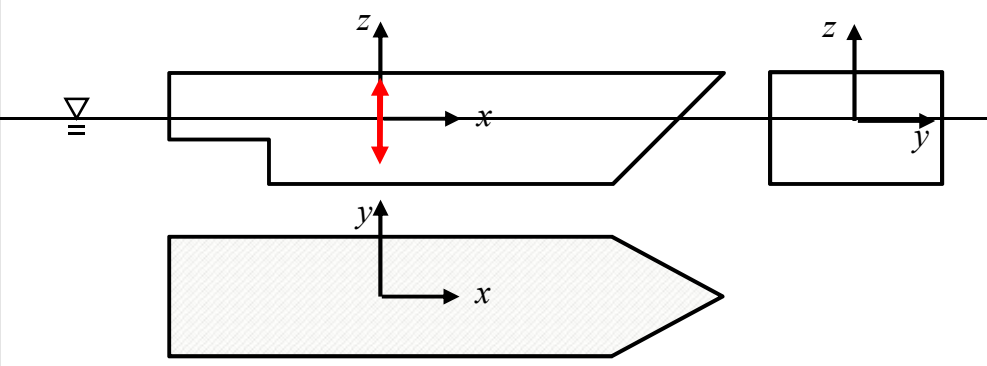
Given  $(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$  Find  $\mathbf{x}, \dot{\mathbf{x}}, \ddot{\mathbf{x}}$

(단, 계수 A,B,C 및 외력  $\mathbf{F}_{exciting}$  은 주어져 있다고 가정)

✓ Assumption

1. Small amplitude water wave (파장에 비해 파고가 작음)
2. small amplitude motion (선박의 운동이 작음)
3. Slender body (선박의 길이에 비해 폭이 작음, Strip theory에서 자세히 설명)  
 → 물체의 전후 동요(Surge)는 독립적으로 취급 (Coupling 고려 안함)
4. Lateral symmetry (symmetric about  $xz$ -plane)  
 → 물체 운동이 종운동(Longitudinal motion) 과 횡운동(Transverse motion)으로 나뉨  
 surge, heave, pitch ↔ sway, roll, yaw

서로 영향을 주지 않음



(ex) 좌우 대칭 선박

Heave운동에 의해 영향을 받는 운동은?

surge , pitch

~~sway~~ , ~~roll~~ , ~~yaw~~

- 1) Journée, J.M.J. , Adegeest, L.J.M. , Theoretical Manual of Strip Theory program“ Seaway for Windows”, Delft University of Technology, 2003, pp38-42
- 2) Newman, J.N. , Marine Hydrodynamics, The MIT Press, Cambridge, 1997, pp 307-311
- 3) Journée, J.M.J. , Massie, W.W. , Offshore Hydrodynamics, Delft University of Technology, 2001, pp8-1-4

# 선박의 6자유도 운동 방정식

✓ 6DOF Equations of Ship Motion

: 3 kinds of coupled motions (surge / heave-pitch / sway-roll-yaw)

$$\mathbf{M} = \begin{bmatrix} M & 0 & 0 & 0 & Mz_C & -My_C \\ 0 & M & 0 & -Mz_C & 0 & Mx_C \\ 0 & 0 & M & My_C & -Mx_C & 0 \\ 0 & -Mz_C & My_C & I_{xx} & 0 & -I_{xz} \\ Mz_C & 0 & -Mx_C & 0 & I_{yy} & 0 \\ -My_C & Mx_C & 0 & -I_{zx} & 0 & I_{zz} \end{bmatrix} \quad \mathbf{B} = \begin{bmatrix} B_{11} & 0 & B_{13} & 0 & B_{15} & 0 \\ 0 & B_{22} & 0 & B_{24} & 0 & B_{26} \\ B_{31} & 0 & B_{33} & 0 & B_{35} & 0 \\ 0 & B_{42} & 0 & B_{44} & 0 & B_{46} \\ B_{51} & 0 & B_{53} & 0 & B_{55} & 0 \\ 0 & B_{62} & 0 & B_{64} & 0 & B_{66} \end{bmatrix} \quad \mathbf{x} = \begin{bmatrix} \xi_1 \\ \xi_2 \\ \xi_3 \\ \xi_4 \\ \xi_5 \\ \xi_6 \end{bmatrix}$$

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$$

$$\mathbf{F}_{exciting} = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \end{bmatrix}$$

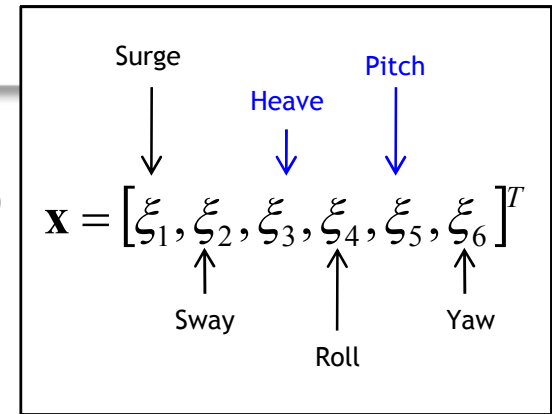
$$\mathbf{A} = \begin{bmatrix} A_{11} & 0 & A_{13} & 0 & A_{15} & 0 \\ 0 & A_{22} & 0 & A_{24} & 0 & A_{26} \\ A_{31} & 0 & A_{33} & 0 & A_{35} & 0 \\ 0 & A_{42} & 0 & A_{44} & 0 & A_{46} \\ A_{51} & 0 & A_{53} & 0 & A_{55} & 0 \\ 0 & A_{62} & 0 & A_{64} & 0 & A_{66} \end{bmatrix} \quad \mathbf{C} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & C_{33} & 0 & C_{35} & 0 \\ 0 & 0 & 0 & C_{44} & 0 & 0 \\ 0 & 0 & C_{53} & 0 & C_{55} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

- 1) Journée, J.M.J. , Adegeest, L.J.M. , Theoretical Manual of Strip Theory program " Seaway for Windows", Delft University of Technology, 2003, pp38-42
- 2) Newman, J.N. , Marine Hydrodynamics, The MIT Press, Cambridge, 1997, pp 307-311
- 3) Journée, J.M.J. , Massie, W.W. , Offshore Hydrodynamics, Delft University of Technology, 2001, pp8-1-4

# 선박의 6자유도 운동 방정식

✓ 6DOF Equations of Ship Motion

: 3 kinds of coupled motions (surge / heave-pitch / sway-roll-yaw)



$$\mathbf{M} = \begin{bmatrix} M & 0 & 0 & 0 & Mz_c & -My_c \\ 0 & M & 0 & -Mz_c & 0 & Mx_c \\ 0 & 0 & M & My_c & -Mx_c & 0 \\ 0 & -Mz_c & My_c & I_{xx} & 0 & -I_{xz} \\ Mz_c & 0 & -Mx_c & 0 & I_{yy} & 0 \\ -My_c & Mx_c & 0 & -I_{xz} & 0 & I_{zz} \end{bmatrix} \quad \mathbf{B} = \begin{bmatrix} B_{11} & 0 & B_{13} & 0 & B_{15} & 0 \\ 0 & B_{22} & 0 & B_{24} & 0 & B_{26} \\ B_{31} & 0 & B_{33} & 0 & B_{35} & 0 \\ 0 & B_{42} & 0 & B_{44} & 0 & B_{46} \\ B_{51} & 0 & B_{53} & 0 & B_{55} & 0 \\ 0 & B_{62} & 0 & B_{64} & 0 & B_{66} \end{bmatrix}$$

운동 방정식 :  $(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$

$$\mathbf{F}_{exciting} = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \end{bmatrix}$$

$$\mathbf{A} = \begin{bmatrix} A_{11} & 0 & A_{13} & 0 & A_{15} & 0 \\ 0 & A_{22} & 0 & A_{24} & 0 & A_{26} \\ A_{31} & 0 & A_{33} & 0 & A_{35} & 0 \\ 0 & A_{42} & 0 & A_{44} & 0 & A_{46} \\ A_{51} & 0 & A_{53} & 0 & A_{55} & 0 \\ 0 & A_{62} & 0 & A_{64} & 0 & A_{66} \end{bmatrix} \quad \mathbf{C} = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & C_{33} & 0 & C_{35} & 0 \\ 0 & 0 & 0 & C_{44} & 0 & 0 \\ 0 & 0 & C_{53} & 0 & C_{55} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

➔ heave-pitch

motion of equation :

$$\begin{bmatrix} M + A_{33} & -Mx_c + A_{35} \\ -Mx_c + A_{53} & A_{55} + I_{yy} \end{bmatrix} \begin{bmatrix} \ddot{\xi}_3 \\ \ddot{\xi}_5 \end{bmatrix} + \begin{bmatrix} B_{33} & B_{35} \\ B_{53} & B_{55} \end{bmatrix} \begin{bmatrix} \dot{\xi}_3 \\ \dot{\xi}_5 \end{bmatrix} + \begin{bmatrix} C_{33} & C_{35} \\ C_{53} & C_{55} \end{bmatrix} \begin{bmatrix} \xi_3 \\ \xi_5 \end{bmatrix} = \begin{bmatrix} F_3 \\ F_5 \end{bmatrix}$$

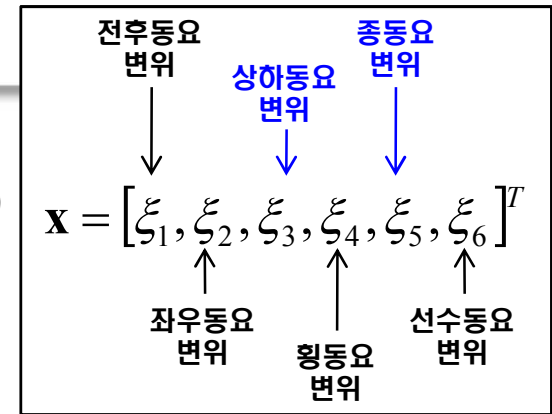
( $y_c = 0$  으로 가정)

- 1) Journée, J.M.J. , Adegeest, L.J.M. , Theoretical Manual of Strip Theory program " Seaway for Windows", Delft University of Technology, 2003, pp38-42
- 2) Newman, J.N. , Marine Hydrodynamics, The MIT Press, Cambridge, 1997, pp 307-311
- 3) Journée, J.M.J. , Massie, W.W. , Offshore Hydrodynamics, Delft University of Technology, 2001, pp8-1-4

# 선박의 6자유도 운동 방정식

## ✓ 6DOF Equations of Ship Motion

: 3 kinds of coupled motions (surge / heave-pitch / sway-roll-yaw)



$$M = \begin{bmatrix} M & 0 & 0 & 0 & Mz_c & -My_c \\ 0 & M & 0 & -Mz_c & 0 & Mx_c \\ 0 & 0 & M & My_c & -Mx_c & 0 \\ 0 & -Mz_c & My_c & I_{xx} & 0 & -I_{xz} \\ Mz_c & 0 & -Mx_c & 0 & I_{yy} & 0 \\ -My_c & Mx_c & 0 & -I_{zx} & 0 & I_{zz} \end{bmatrix}$$

$$B = \begin{bmatrix} B_{11} & 0 & B_{13} & 0 & B_{15} & 0 \\ 0 & B_{22} & 0 & B_{24} & 0 & B_{26} \\ B_{31} & 0 & B_{33} & 0 & B_{35} & 0 \\ 0 & B_{42} & 0 & B_{44} & 0 & B_{46} \\ B_{51} & 0 & B_{53} & 0 & B_{55} & 0 \\ 0 & B_{62} & 0 & B_{64} & 0 & B_{66} \end{bmatrix}$$

$$(M + A)\ddot{x} + B\dot{x} + Cx = F_{exciting}$$

$$F_{exciting} = \begin{bmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \end{bmatrix}$$

$$A = \begin{bmatrix} A_{11} & 0 & A_{13} & 0 & A_{15} & 0 \\ 0 & A_{22} & 0 & A_{24} & 0 & A_{26} \\ A_{31} & 0 & A_{33} & 0 & A_{35} & 0 \\ 0 & A_{42} & 0 & A_{44} & 0 & A_{46} \\ A_{51} & 0 & A_{53} & 0 & A_{55} & 0 \\ 0 & A_{62} & 0 & A_{64} & 0 & A_{66} \end{bmatrix}$$

$$C = \begin{bmatrix} 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & C_{33} & 0 & C_{35} & 0 \\ 0 & 0 & 0 & C_{44} & 0 & 0 \\ 0 & 0 & C_{53} & 0 & C_{55} & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 \end{bmatrix}$$

→ sway-roll-yaw :

$$\begin{bmatrix} m + A_{22} & -mz_c + A_{24} & mx_c + A_{26} \\ -mz_c + A_{42} & I_{xx} + A_{44} & -I_{xz} + A_{46} \\ mx_c + A_{62} & -I_{zx} + A_{64} & I_{zz} + A_{66} \end{bmatrix} \begin{bmatrix} \ddot{\xi}_2 \\ \ddot{\xi}_4 \\ \ddot{\xi}_6 \end{bmatrix} + \begin{bmatrix} B_{22} & B_{24} & B_{26} \\ B_{42} & B_{44} & B_{46} \\ B_{62} & B_{64} & B_{66} \end{bmatrix} \begin{bmatrix} \dot{\xi}_2 \\ \dot{\xi}_4 \\ \dot{\xi}_6 \end{bmatrix} + \begin{bmatrix} 0 & 0 & 0 \\ 0 & C_{44} & 0 \\ 0 & 0 & 0 \end{bmatrix} \begin{bmatrix} \xi_2 \\ \xi_4 \\ \xi_6 \end{bmatrix} = \begin{bmatrix} F_2 \\ F_4 \\ F_6 \end{bmatrix}$$

# 선박의 6자유도 운동 방정식

## : Heave 운동 방정식 풀이 (Assuption : Harmonic motion)

### ✓ 선박의 6자유도 운동 방정식

$$(M + A)\ddot{x} + B\dot{x} + Cx = F_{exciting}$$

ex) Heave  
↓

$$(M + A_{33})\ddot{\xi}_3 + B_{33}\dot{\xi}_3 + C_{33}\xi_3 = F_{exciting,3}$$

시간이 충분히 흘러 외력의 주기와 같은 운동을 함  
(Harmonic motion)  
=> 초기 Transient motion은 고려하지 않음

$\xi_3(t) = \xi_3^A e^{i\omega t}$	$F_{exciting,3} = F_3^A e^{i\omega t} = \eta_0 f_3^A e^{i\omega t}$
$\dot{\xi}_3(t) = i\omega \xi_3^A e^{i\omega t}$	
$\ddot{\xi}_3(t) = -\omega^2 \xi_3^A e^{i\omega t}$	
	( $\eta_0$ : Wave Amplitude, Real)
	( $f_3^A$ : Wave exciting force Amplitude/1m wave amplitude, Complex)

ex) If  $\xi_3^A$  is not a complex (real)

Let  $\xi_3^A = a$

$$\xi_3 = \xi_3^A e^{i\omega t} \quad \downarrow \text{(Euler 공식)}$$

$$= a(\cos \omega t + i \sin \omega t)$$

$$= a \cos \omega t + ia \sin \omega t$$

$$\text{Re}\{\xi_3\} = a \cos \omega t$$

If  $\xi_3^A$  is a complex

Let  $\xi_3^A = a + ib$

$$\xi_3 = \xi_3^A e^{i\omega t} \quad \downarrow \text{(Euler 공식)}$$

$$= (a + ib)(\cos \omega t + i \sin \omega t)$$

$$= (a \cos \omega t - b \sin \omega t) + i(b \cos \omega t + a \sin \omega t)$$

$$\text{Re}\{\xi_3\} = a \cos \omega t - b \sin \omega t = c \cos(\omega t - \varepsilon)$$

Phase가 나타남



# 선박의 6자유도 운동 방정식

## : Heave 운동 방정식 풀이 (Assuption : Harmonic motion)

### ✓ 선박의 6자유도 운동 방정식

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$$

ex) Heave  
↓

$$(M + A_{33})\ddot{\xi}_3 + B_{33}\dot{\xi}_3 + C_{33}\xi_3 = F_{exciting,3}$$

↓

$\xi_3(t) = \xi_3^A e^{i\omega t}$	$F_{exciting,3} = F_3^A e^{i\omega t} = \eta_0 f_3^A e^{i\omega t}$ <p>( <math>\eta_0</math> : Wave Amplitude, Real )                  ( <math>f_3^A</math> : Wave exciting force Amplitude, Complex )</p>
$\dot{\xi}_3(t) = i\omega \xi_3^A e^{i\omega t}$	
$\ddot{\xi}_3(t) = -\omega^2 \xi_3^A e^{i\omega t}$	

$$(M + A_{33})(-\omega^2 \xi_3^A e^{i\omega t}) + B_{33}(i\omega \xi_3^A e^{i\omega t}) + C_{33}(\xi_3^A e^{i\omega t}) = \eta_0 f_3^A e^{i\omega t}$$

↓

$$\{-\omega^2(M + A_{33}) + i\omega B_{33} + C_{33}\} \xi_3^A e^{i\omega t} = \eta_0 f_3^A e^{i\omega t}$$

↓

$$\underline{\{-\omega^2(M + A_{33}) + i\omega B_{33} + C_{33}\} \xi_3^A} = \eta_0 f_3^A \implies \xi_3^A = \eta_0 f_3^A D^{-1} \implies \frac{\xi_3^A}{\eta_0} = f_3^A \mathbf{D}^{-1}$$

$\mathbf{D} \Rightarrow \text{Complex}$

✓ **RAO** (Response Amplitude Operator)  
 : 1m wave Amplitude 를 가지는  
 주파수  $\omega$  인 wave 에 대한  
 선박의 6자유도 운동 변위

↑



# 선박의 6자유도 운동 방정식

## : Heave-pitch 연성 운동 방정식 풀이

(continue)

① 운동 방정식에  
변위, 속도, 가속도 대입

$$\begin{bmatrix} M + A_{33} & -Mx_C + A_{35} \\ -Mx_C + A_{53} & A_{55} + I_{yy} \end{bmatrix} \begin{bmatrix} -\omega^2 \xi_3^A e^{i\omega t} \\ -\omega^2 \xi_5^A e^{i\omega t} \end{bmatrix} + \begin{bmatrix} B_{33} & B_{35} \\ B_{53} & B_{55} \end{bmatrix} \begin{bmatrix} i\omega \xi_3^A e^{i\omega t} \\ i\omega \xi_5^A e^{i\omega t} \end{bmatrix} + \begin{bmatrix} C_{33} & C_{35} \\ C_{53} & C_{55} \end{bmatrix} \begin{bmatrix} \xi_3^A e^{i\omega t} \\ \xi_5^A e^{i\omega t} \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A e^{i\omega t} \\ \eta_0 F_5^A e^{i\omega t} \end{bmatrix}$$



② 양변을  $e^{i\omega t}$  로 나눔

$$-\omega^2 \begin{bmatrix} M + A_{33} & -Mx_C + A_{35} \\ -Mx_C + A_{53} & A_{55} + I_{yy} \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} + i\omega \begin{bmatrix} B_{33} & B_{35} \\ B_{53} & B_{55} \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} + \begin{bmatrix} C_{33} & C_{35} \\ C_{53} & C_{55} \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix}$$



③  $\begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix}$  로 묶어서 정리

$$\left( -\omega^2 \begin{bmatrix} M + A_{33} & -Mx_C + A_{35} \\ -Mx_C + A_{53} & A_{55} + I_{yy} \end{bmatrix} + i\omega \begin{bmatrix} B_{33} & B_{35} \\ B_{53} & B_{55} \end{bmatrix} + \begin{bmatrix} C_{33} & C_{35} \\ C_{53} & C_{55} \end{bmatrix} \right) \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix}$$



$$\begin{bmatrix} -\omega^2(M + A_{33}) + i\omega B_{33} + C_{33} & -\omega^2(-Mx_C + A_{35}) + i\omega B_{35} + C_{35} \\ -\omega^2(-Mx_C + A_{53}) + i\omega B_{53} + C_{53} & -\omega^2(A_{55} + I_{yy}) + i\omega B_{55} + C_{55} \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix}$$



$$\begin{bmatrix} P & Q \\ R & S \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix} \quad \left( \begin{array}{l} P = -\omega^2(M + A_{33}) + i\omega B_{33} + C_{33} \\ Q = -\omega^2(-Mx_C + A_{35}) + i\omega B_{35} + C_{35} \\ R = -\omega^2(-Mx_C + A_{53}) + i\omega B_{53} + C_{53} \\ S = -\omega^2(I_{yy} + A_{55}) + i\omega B_{55} + C_{55} \end{array} \right)$$

# 선박의 6자유도 운동 방정식

## : Heave-pitch 연성 운동 방정식 풀이

(continue)

③  $\begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix}$  로 묶어서 정리

$$\begin{bmatrix} P & Q \\ R & S \end{bmatrix} \begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix} \quad \left( \begin{array}{l} P = -\omega^2 (M + A_{33}) + i\omega B_{33} + C_{33} \\ Q = -\omega^2 (-Mx_C + A_{35}) + i\omega B_{35} + C_{35} \\ R = -\omega^2 (-Mx_C + A_{53}) + i\omega B_{53} + C_{53} \\ S = -\omega^2 (I_{yy} + A_{55}) + i\omega B_{55} + C_{55} \end{array} \right)$$



④ 역행렬을 곱하여  
변위를  $\xi_3^A, \xi_5^A$  구함

$$\begin{bmatrix} \xi_3^A \\ \xi_5^A \end{bmatrix} = \begin{bmatrix} P & Q \\ R & S \end{bmatrix}^{-1} \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix} = \frac{1}{PS - QR} \begin{bmatrix} S & -Q \\ -R & P \end{bmatrix} \begin{bmatrix} \eta_0 F_3^A \\ \eta_0 F_5^A \end{bmatrix}$$

$$= \frac{1}{PS - QR} \begin{bmatrix} \eta_0 (F_3^A S - F_5^A Q) \\ \eta_0 (-F_3^A R + F_5^A P) \end{bmatrix} = \begin{bmatrix} \eta_0 \frac{F_3^A S - F_5^A Q}{PS - QR} \\ \eta_0 \frac{F_5^A P - F_3^A R}{PS - QR} \end{bmatrix}$$



⑤ 1m 파고에 대한  
운동 변위 (\*RAO)

$$\therefore \frac{\xi_3^A}{\eta_0} = \frac{F_3^A S - F_5^A Q}{PS - QR}$$

$$\frac{\xi_5^A}{\eta_0} = \frac{F_5^A P - F_3^A R}{PS - QR}$$

\*RAO(Response Amplitude Operator) : 1m v에 대한 선박의 운동 응답

# 선박의 6자유도 운동 방정식

## : 주파수 영역(Frequency domain)에서의 6자유도 운동 방정식 풀이

### 선박의 6자유도 운동 방정식

$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$$

General Case

$$\mathbf{x} = \begin{bmatrix} \xi_1 \\ \xi_2 \\ \xi_3 \\ \xi_4 \\ \xi_5 \\ \xi_6 \end{bmatrix} = \begin{bmatrix} \xi_1^A \\ \xi_2^A \\ \xi_3^A \\ \xi_4^A \\ \xi_5^A \\ \xi_6^A \end{bmatrix}$$

→ Complex amplitude

$$e^{i\omega t} = \mathbf{x}^A e^{i\omega t}, \quad \dot{\mathbf{x}} = i\omega \mathbf{x}^A e^{i\omega t}, \quad \ddot{\mathbf{x}} = -\omega^2 \mathbf{x}^A e^{i\omega t}, \quad \mathbf{F}_{exciting} = \eta_0 \begin{bmatrix} f_1^A \\ f_2^A \\ f_3^A \\ f_4^A \\ f_5^A \\ f_6^A \end{bmatrix} e^{i\omega t} = \eta_0 \mathbf{f}^A e^{i\omega t}$$

$$(\mathbf{M} + \mathbf{A})(-\omega^2 \mathbf{x}^A e^{i\omega t}) + \mathbf{B}(i\omega \mathbf{x}^A e^{i\omega t}) + \mathbf{C}(\mathbf{x}^A e^{i\omega t}) = \eta_0 \mathbf{f}^A e^{i\omega t}$$

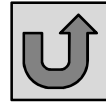
$$\{-\omega^2(\mathbf{M} + \mathbf{A}) + i\omega\mathbf{B} + \mathbf{C}\} \mathbf{x}^A e^{i\omega t} = \eta_0 \mathbf{f}^A e^{i\omega t}$$

$$\underbrace{\{-\omega^2(\mathbf{M} + \mathbf{A}) + i\omega\mathbf{B} + \mathbf{C}\}}_{= \mathbf{D}} \mathbf{x}^A = \eta_0 \mathbf{f}^A \quad \Rightarrow \quad \mathbf{x}^A = \eta_0 \mathbf{D}^{-1} \mathbf{f}^A \quad \Rightarrow \quad \frac{\mathbf{x}^A}{\eta_0} = \mathbf{D}^{-1} \mathbf{f}^A$$

✓ **RAO** (Response Amplitude Operator)  
: 1m wave Amplitude 를 가지는  
주파수  $\omega$  인 wave 에 대한  
선박의 6자유도 운동 변위

# 선박 6자유도 운동 방정식

: 시간 영역(Time domain)에서의 운동 방정식 풀이



$$(\mathbf{M} + \mathbf{A})\ddot{\mathbf{x}} + \mathbf{B}\dot{\mathbf{x}} + \mathbf{C}\mathbf{x} = \mathbf{F}_{exciting}$$

6자유도 선박의 운동 방정식

↓ 변경

$$\ddot{\mathbf{x}} = (\mathbf{M} + \mathbf{A})^{-1}(-\mathbf{B}\dot{\mathbf{x}} - \mathbf{C}\mathbf{x} + \mathbf{F}_{exciting})$$

↓ 초기 위치 및 자세와 초기 속도 대입 ( $\dot{\mathbf{x}}_0, \mathbf{x}_0$ )

$$\ddot{\mathbf{x}}_1 = (\mathbf{M} + \mathbf{A})^{-1}(-\mathbf{B}\dot{\mathbf{x}}_0 - \mathbf{C}\mathbf{x}_0 + \mathbf{F}_{exciting})$$

↓ 수치 적분<sup>1)</sup>을 통해 가속도로부터 속도, 변위를 구함

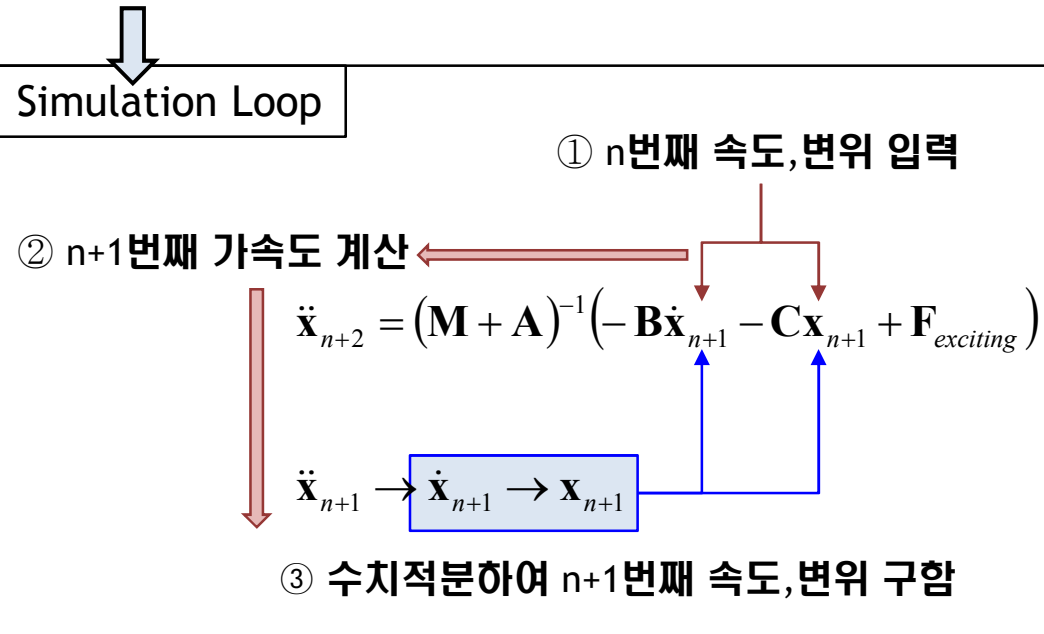
$$\ddot{\mathbf{x}}_1 \rightarrow \dot{\mathbf{x}}_1 \rightarrow \mathbf{x}_1$$

↓ 속도, 변위 대입

$$\ddot{\mathbf{x}}_2 = (\mathbf{M} + \mathbf{A})^{-1}(-\mathbf{B}\dot{\mathbf{x}}_1 - \mathbf{C}\mathbf{x}_1 + \mathbf{F}_{exciting})$$

↓ 수치 적분

$$\ddot{\mathbf{x}}_2 \rightarrow \dot{\mathbf{x}}_2 \rightarrow \mathbf{x}_2$$



Term Project#4 Ship motion simulation  
프로그램에서는 시간 영역에서의 풀이 방법을 사용하여 매시간 선박의 운동을 계산함



## Cf) Strip theory를 사용하여 Added mass, Damping coefficient, Wave exciting force 계산

○ 안의 값은 Library File에서 구할 수 있음

$$\left( A_{jk}^0 = \int_L a_{jk} dx, B_{jk}^0 = \int_L b_{jk} dx \right)$$

$$A_{33} = \int_L a_{33} dx - \frac{U}{\omega^2} b_{33}^A$$

$$B_{33} = \int_L b_{33} dx + U a_{33}^A$$

$$A_{35} = - \int_L x a_{33} dx - \frac{U}{\omega^2} B_{33}^0 + \frac{U}{\omega^2} x_A b_{33}^A - \frac{U^2}{\omega^2} a_{33}^A$$

$$B_{35} = - \int_L x b_{33} dx + U A_{33}^0 - U x_A a_{33}^A - \frac{U^2}{\omega^2} b_{33}^A$$

$$A_{53} = - \int_L x a_{33} dx + \frac{U}{\omega^2} B_{33}^0 + \frac{U}{\omega^2} x_A b_{33}^A$$

$$B_{53} = - \int_L x b_{33} dx - U A_{33}^0 - U x_A a_{33}^A$$

$$A_{55} = \int_L x^2 a_{33} dx + \frac{U^2}{\omega^2} A_{33}^0 - \frac{U}{\omega^2} x_A b_{33}^A + \frac{U^2}{\omega^2} x_A a_{33}^A$$

$$B_{55} = \int_L x^2 b_{33} dx + \frac{U^2}{\omega^2} B_{33}^0 + U x_A^2 a_{33}^A + \frac{U^2}{\omega^2} x_A b_{33}^A$$

$$F_3 = \rho \alpha \int_L (f_3 + h_3) dx + \rho \alpha \frac{U}{i \omega} h_3^A$$

$$F_5 = - \rho \alpha \int_L \left[ x (f_3 + h_3) + \rho \alpha \frac{U}{i \omega} h_3 \right] dx - \rho \alpha \frac{U}{i \omega} x_A h_3^A$$

$U$ : 선박의 전진 속도

$\rho$ : 유체의 밀도

$\alpha$ : Wave amplitude

$f_j$ : Sectional Froude Krylov force ( $j^{\text{th}}$  mode)

$h_j$ : Sectional Diffraction force ( $j^{\text{th}}$  mode)

$\omega$ : Encounter wave frequency

$x_A, a_{jk}^A, b_{jk}^A$ : Values at the aftermost section

## Cf) Strip theory를 사용하여 Added mass, Damping coefficient, Wave exciting force 계산

$$\left( A_{jk}^0 = \int_L a_{jk} dx, B_{jk}^0 = \int_L b_{jk} dx \right)$$

$$A_{22} = \int_L a_{22} dx - \frac{U}{\omega^2} b_{22}^A$$

$$B_{22} = \int_L b_{22} dx + Ua_{22}^A$$

$$A_{24} = A_{42} = \int_L a_{24} dx - \frac{U}{\omega^2} b_{24}^A$$

$$B_{24} = B_{42} = \int_L b_{24} dx + Ua_{24}^A$$

$$A_{26} = \int_L xa_{22} dx + \frac{U}{\omega^2} B_{22}^0 - \frac{U}{\omega^2} x_A b_{22}^A + \frac{U^2}{\omega^2} a_{22}^A$$

$$B_{26} = \int_L xb_{22} dx - UA_{22}^0 + Ux_A a_{22}^A + \frac{U^2}{\omega^2} b_{22}^A$$

$$A_{44} = \int_L a_{44} dx - \frac{U}{\omega^2} b_{44}^A$$

$$B_{44} = \int_L b_{44} dx + Ua_{44}^A + B_{44}^*$$

$$A_{46} = \int_L xa_{24} dx + \frac{U}{\omega^2} B_{24}^0 - \frac{U}{\omega^2} x_A b_{24}^A + \frac{U^2}{\omega^2} a_{24}^A$$

$$B_{46} = \int_L xb_{24} dx - UA_{24}^0 + Ux_A a_{24}^A + \frac{U^2}{\omega^2} b_{24}^A$$

$$A_{62} = \int_L xa_{22} dx - \frac{U}{\omega^2} B_{22}^0 - \frac{U}{\omega^2} x_A b_{22}^A$$

$$B_{62} = \int_L xb_{22} dx + UA_{22}^0 + Ux_A a_{22}^A$$

$$A_{64} = \int_L xa_{24} dx - \frac{U}{\omega^2} B_{24}^0 - \frac{U}{\omega^2} x_A b_{24}^A$$

$$B_{64} = \int_L xb_{24} dx + UA_{24}^0 + Ux_A a_{24}^A$$

$$A_{66} = \int_L x^2 a_{22} dx + \frac{U^2}{\omega^2} A_{22}^0 - \frac{U}{\omega^2} x_A^2 b_{22}^A + \frac{U^2}{\omega^2} x_A a_{22}^A$$

$$B_{66} = \int_L x^2 b_{22} dx + \frac{U^2}{\omega^2} B_{22}^0 + Ux_A^2 a_{22}^A + \frac{U^2}{\omega^2} x_A b_{22}^A$$

$U$ : 선박의 전진 속도

$\rho$ : 유체의 밀도

$\alpha$ : Wave amplitude

$f_j$ : Sectional Froude Krylov force ( $j^{\text{th}}$  mode)

$h_j$ : Sectional Diffraction force ( $j^{\text{th}}$  mode)

$\omega$ : Encounter wave frequency

$x_A, a_{jk}^A, b_{jk}^A$ : Values at the aftermost section

$B_4^*$ : Roll Damping

Cf) Strip theory를 사용하여 Added mass, Damping coefficient, Wave exciting force 계산

$$\left( A_{jk}^0 = \int_L a_{jk} dx, B_{jk}^0 = \int_L b_{jk} dx \right)$$

$$F_2 = \rho\alpha \int_L (f_2 + h_2) dx + \rho\alpha \frac{U}{i\omega} h_2^A$$

$$F_4 = \rho\alpha \int_L (f_4 + h_4) dx + \rho\alpha \frac{U}{i\omega} h_4^A$$

$$F_6 = \rho\alpha \int_L \left[ x(f_2 + h_2) + \rho\alpha \frac{U}{i\omega} h_2 \right] dx + \rho\alpha \frac{U}{i\omega} x_A h_2^A$$

$U$ : 선박의 전진 속도

$\rho$ : 유체의 밀도

$\alpha$ : Wave amplitude

$f_j$ : Sectional Froude Krylov force ( $j^{\text{th}}$  mode)

$h_j$ : Sectional Diffraction force ( $j^{\text{th}}$  mode)

$\omega$ : Encounter wave frequency

$x_A, a_{jk}^A, b_{jk}^A$ : Values at the aftermost section

$B_4^*$ : Roll Damping

## (cf) Fourth-Order Runge-Kutta Method

The most commonly used set of values for the parameters

$$y_{n+1} = y_n + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$

$$k_1 = f(x_n, y_n)$$

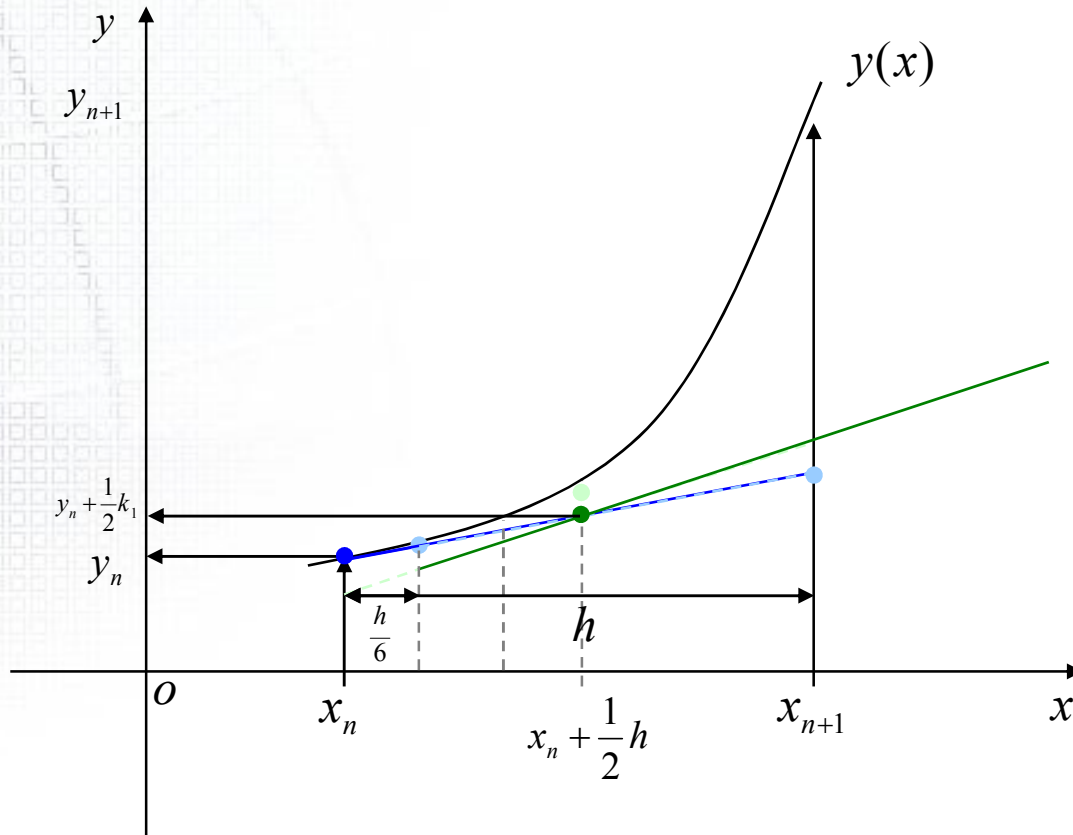
$$k_2 = f\left(x_n + \frac{1}{2}h, y_n + \frac{1}{2}hk_1\right)$$

$$k_3 = f\left(x_n + \frac{1}{2}h, y_n + \frac{1}{2}hk_2\right)$$

$$k_4 = f(x_n + h, y_n + \beta_6 hk_3)$$

# (cf) Fourth-Order Runge-Kutta Method

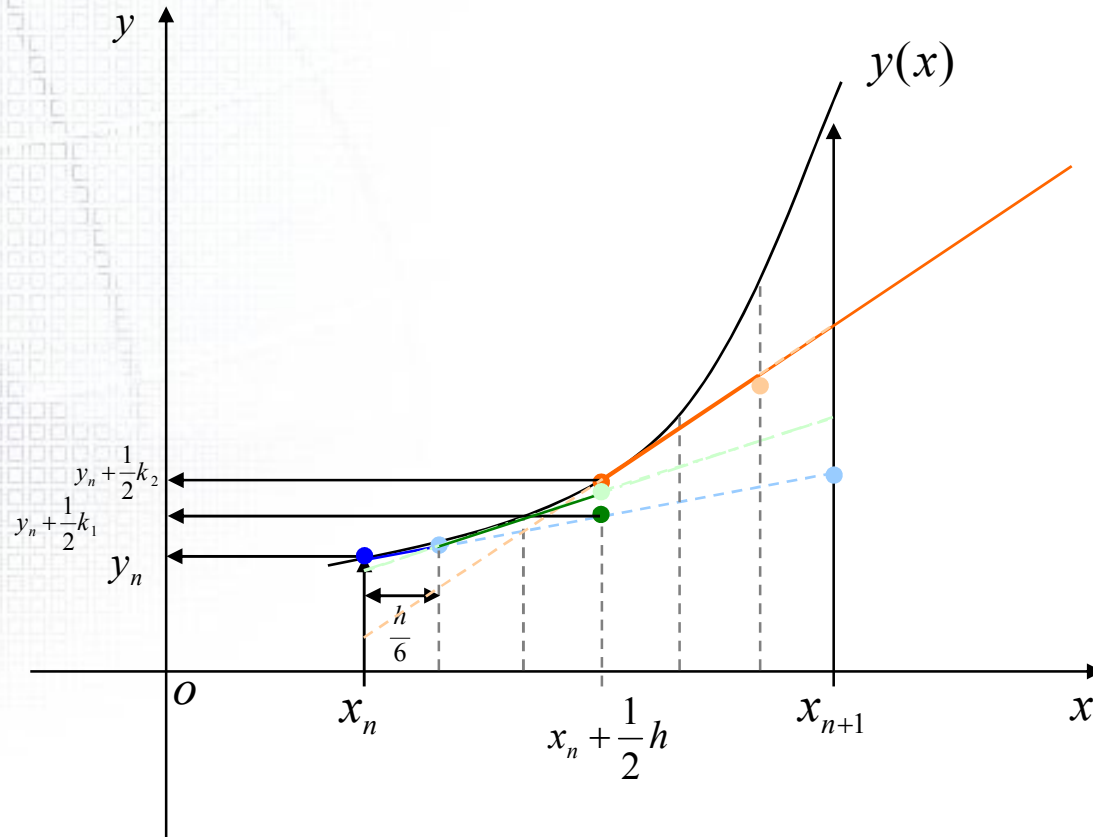
$x_n$ 과  $x_{n+1} = x_n + h$  사이를 6등분 하여 Improved Euler method와 비슷한 방법으로  $y(x)$ 를 구하는 방법.



- ①  $x_n, y_n, f(x_n, y_n)$
- ②  $h \cdot f(x_n, y_n) = k_1$
- ③  $\frac{1}{6} k_1, y_{n+1} = y_n + \frac{1}{6} k_1$
- ④  $h \cdot f(x_n + \frac{1}{2} h, y_n + \frac{1}{2} k_1) = k_2$
- ⑤  $\frac{2}{6} k_2 = \frac{1}{3} k_2, y_{n+1} = y_n + \frac{1}{6} k_1 + \frac{1}{3} k_2$

# (cf) Fourth-Order Runge-Kutta Method

$x_n$ 과  $x_{n+1} = x_n + h$  사이를 6등분 하여 Improved Euler method와 비슷한 방법으로  $y(x)$ 를 구하는 방법.

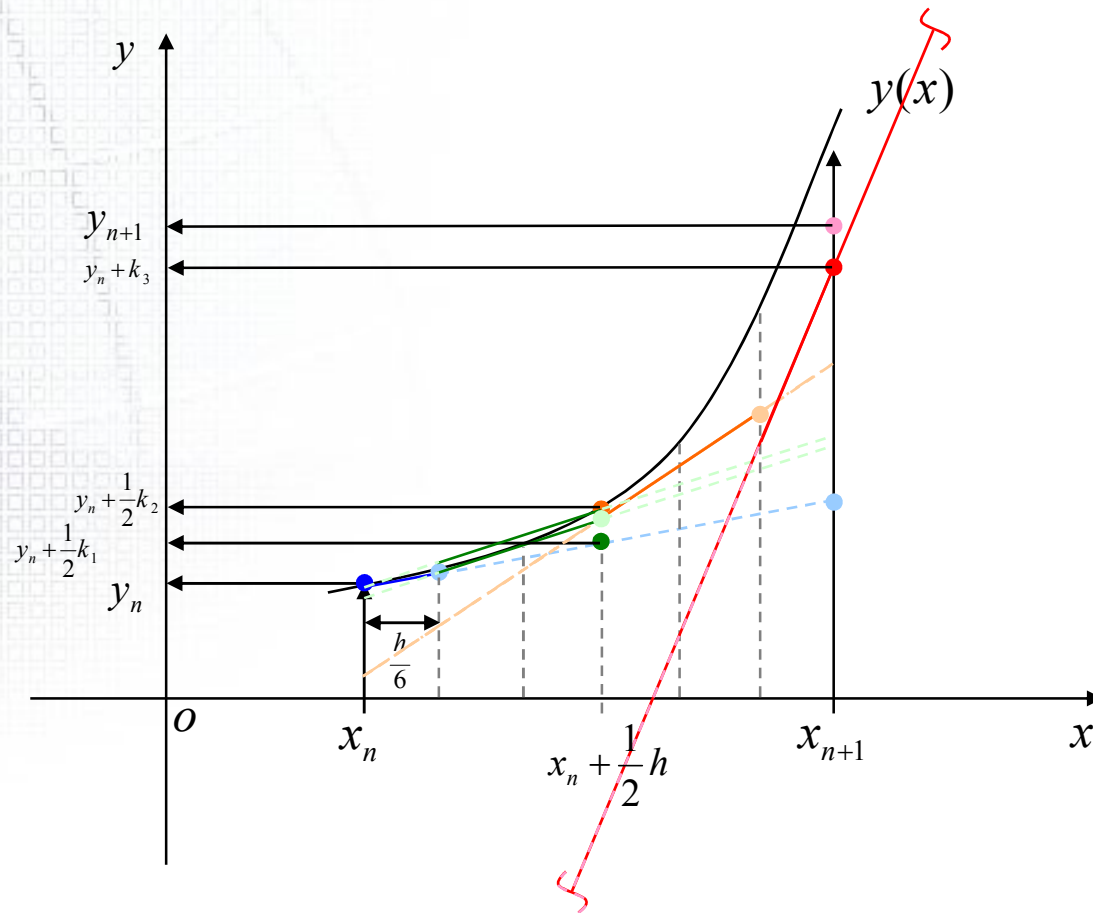


- ①  $x_n, y_n, f(x_n, y_n)$
- ②  $h \cdot f(x_n, y_n) = k_1$
- ③  $\frac{1}{6}k_1, y_{n+1} = y_n + \frac{1}{6}k_1$
- ④  $h \cdot f(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1) = k_2$
- ⑤  $\frac{2}{6}k_2 = \frac{1}{3}k_2, y_{n+1} = y_n + \frac{1}{6}k_1 + \frac{1}{3}k_2$
- ⑥  $y_n + \frac{1}{2}k_2$
- ⑦  $h \cdot f(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_2) = k_3$
- ⑧  $\frac{2}{6}k_3 = \frac{1}{3}k_3, y_{n+1} = y_n + \frac{1}{6}k_1 + \frac{1}{3}k_2 + \frac{1}{3}k_3$



# (cf) Fourth-Order Runge-Kutta Method

$x_n$ 과  $x_{n+1} = x_n + h$  사이를 6등분 하여 Improved Euler method와 비슷한 방법으로  $y(x)$ 를 구하는 방법.



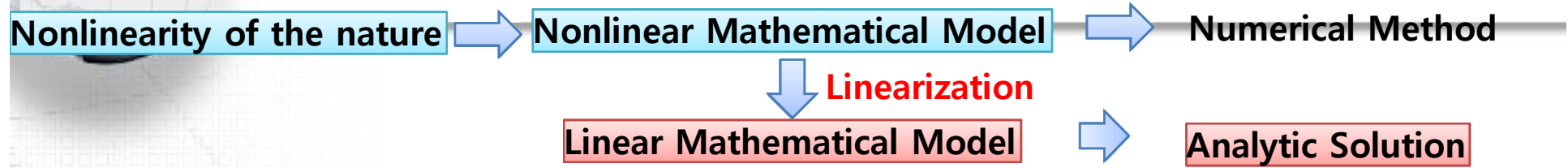
- ①  $x_n, y_n, f(x_n, y_n)$
- ②  $h \cdot f(x_n, y_n) = k_1$
- ③  $\frac{1}{6}k_1, y_{n+1} = y_n + \frac{1}{6}k_1$
- ④  $h \cdot f(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_1) = k_2$
- ⑤  $\frac{2}{6}k_2 = \frac{1}{3}k_2, y_{n+1} = y_n + \frac{1}{6}k_1 + \frac{1}{3}k_2$
- ⑥  $y_n + \frac{1}{2}k_2$
- ⑦  $h \cdot f(x_n + \frac{1}{2}h, y_n + \frac{1}{2}k_2) = k_3$
- ⑧  $\frac{2}{6}k_3 = \frac{1}{3}k_3, y_{n+1} = y_n + \frac{1}{6}k_1 + \frac{1}{3}k_2 + \frac{1}{3}k_3$
- ⑨  $y_n + k_3$
- ⑩  $h \cdot f(x_n + h, y_n + k_3) = k_4$
- ⑪  $\frac{1}{6}k_4$

$$\textcircled{12} \quad y_{n+1} = y_n + \frac{1}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$

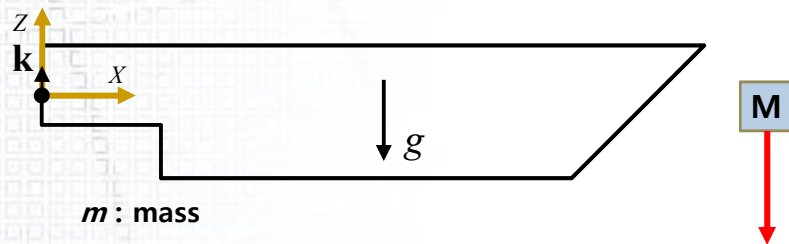


## [참고] 선박의 Heave 운동과 Spring-mass-damping system의 비교

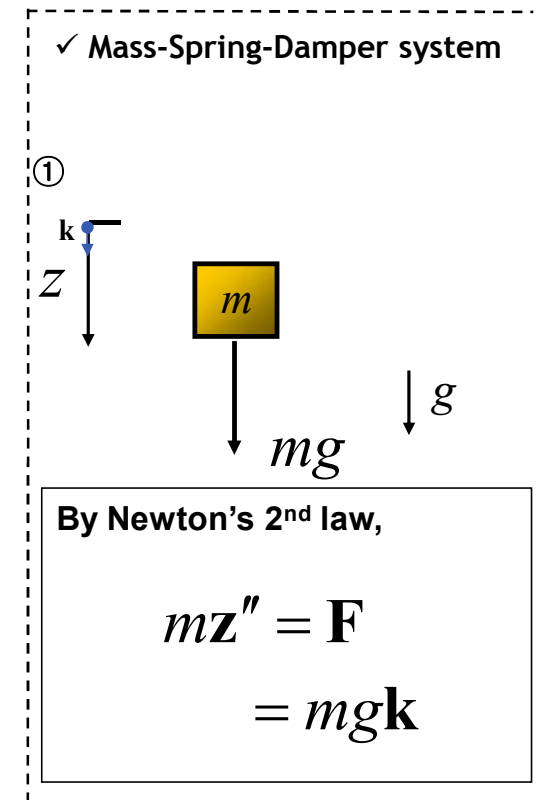
# Nonlinearity



## Ex) Heave Motion of a Ship – step 1



$$\begin{aligned}
 m\ddot{\mathbf{z}} &= \mathbf{F} \\
 &= \mathbf{F}_{gravity} \\
 &= -mg\mathbf{k}
 \end{aligned}$$



# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

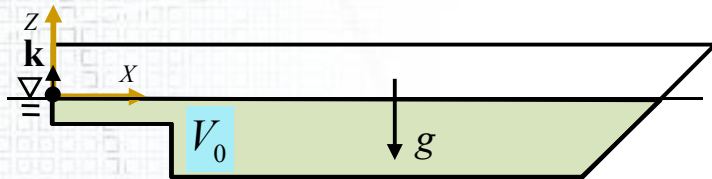
Numerical Method

Linearization

Linear Mathematical Model

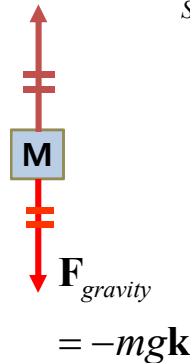
Analytic Solution

## Ex) Heave Motion of a Ship – step 2



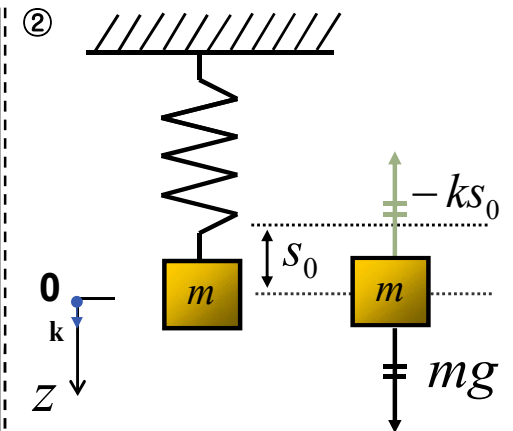
$m$  : mass  
 $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area

$$\mathbf{F}_{static} = \iint_{S_B} P_{static} \mathbf{n} dS = \rho g V_0 \mathbf{k}$$

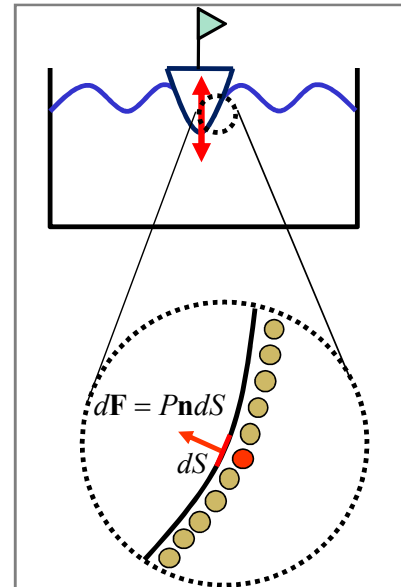


✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned} m\ddot{\mathbf{z}} &= \mathbf{F} \\ &= \mathbf{F}_{gravity} + \mathbf{F}_{static} \\ &= -mg\mathbf{k} + \rho g V_0 \mathbf{k} \\ &= 0 \quad (\because \ddot{\mathbf{z}} = 0) \quad : \text{static equilibrium} \end{aligned}$$



$dS$  : infinitesimal submerged surface area  
 $d\mathbf{F}$  : force exerted by the infinitesimal fluid element on  $dS$   
 $\mathbf{n}$  : normal vector of  $dS$

$$\begin{aligned} m\mathbf{z}'' &= \mathbf{F} \\ &= mg\mathbf{k} - ks_0\mathbf{k} \\ &= 0 \quad (\because \mathbf{z}'' = 0) \end{aligned}$$

: static equilibrium

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# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

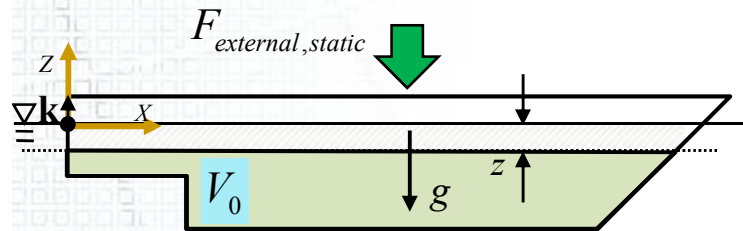
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 3



$m$  : mass  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area  
 $\rho$  : density of sea water

$$\mathbf{F}_{static} = \iint_{S_B} P_{static} \mathbf{n} dS$$

$$= \rho g V_0 \mathbf{k} + \mathbf{F}_{additional\ bouyancy}$$

$$\mathbf{F}_{external,static}$$

$$\mathbf{F}_{gravity}$$

$$= -mg\mathbf{k}$$

additional bouyancy caused by additional displacement  $z$

if,  $z$  is small

$$\mathbf{F}_{additional\ bouyancy} = -\rho g A_{WP} z$$

$$= -kz$$

$$, k = \rho g A_{WP}$$

$$m\ddot{\mathbf{z}} = \mathbf{F}$$

$$= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{external,static}$$

$$= -mg\mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} z + \mathbf{F}_{external,static}$$

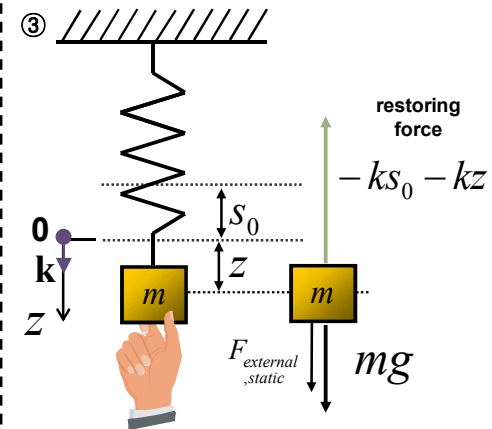
$$= -\rho g A_{wp} z + \mathbf{F}_{external,static}$$

$$= -kz + \mathbf{F}_{external,static}$$

$$= 0 \quad (\because \ddot{z} = 0)$$

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$mz'' = \mathbf{F}$$

$$= mg\mathbf{k} - ks_0\mathbf{k} - kz\mathbf{k} + \mathbf{F}_{external,static}$$

$$= -kz\mathbf{k} + \mathbf{F}_{external,static}$$

$$= 0 \quad (\because z'' = 0)$$



# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

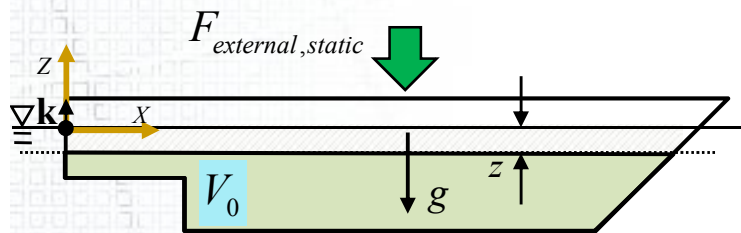
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 4



$m$  : mass       $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

$$\begin{aligned}
 m\ddot{z} &= \mathbf{F} \\
 &= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{external,static} \\
 &= -mg\mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} + \mathbf{F}_{external,static} \\
 &= -\rho g A_{wp} \mathbf{z} + \mathbf{F}_{external,static} \\
 &= -k\mathbf{z} + \mathbf{F}_{external,static}, \quad k = \rho g A_{wp} \\
 &= 0 \quad (\because \ddot{z} = 0)
 \end{aligned}$$

$$\begin{aligned}
 \mathbf{F}_{static} &= \iint_{S_B} P_{static} \mathbf{n} dS \\
 &= \rho g V_0 \mathbf{k} + \mathbf{F}_{additional\ bouyancy} \\
 \mathbf{F}_{external,static} &= -mg\mathbf{k}
 \end{aligned}$$

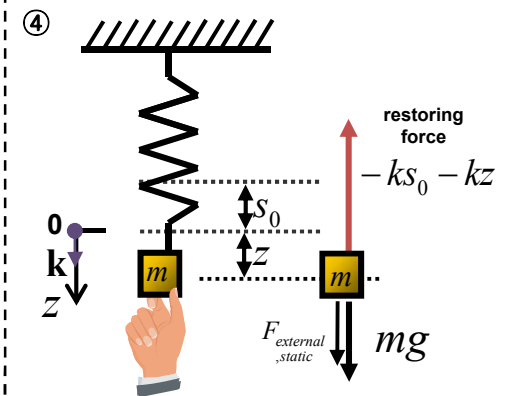
additional bouyancy caused by additional displacement  $z$

$$\begin{aligned}
 \text{if, } z \text{ is small} \\
 \mathbf{F}_{additional\ bouyancy} &= -\rho g A_{wp} \mathbf{z} \\
 &= -k\mathbf{z} \\
 &, k = \rho g A_{wp}
 \end{aligned}$$

Linearized Restoring Force

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned}
 m\mathbf{z}'' &= \mathbf{F} \\
 &= mg\mathbf{k} - ks_0\mathbf{k} - kz\mathbf{k} + \mathbf{F}_{external,static} \\
 &= -kz\mathbf{k} + \mathbf{F}_{external,static}
 \end{aligned}$$

$$m\mathbf{z}'' + k\mathbf{z} = 0 \quad \text{Oscillation by the restoring force}$$



# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

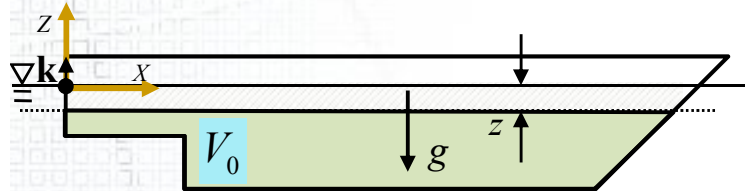
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 4



$m$  : mass       $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

$$\begin{aligned} \mathbf{F}_{static} &= \iint_{S_B} P_{static} \mathbf{n} dS \\ &= \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} \\ &= \rho g V_0 \mathbf{k} - k \mathbf{z} \end{aligned}$$

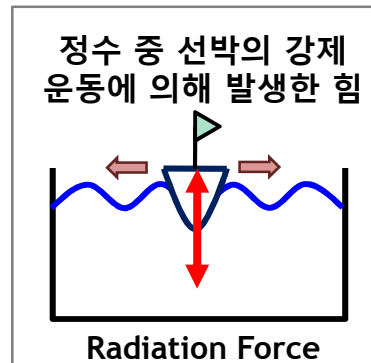
$$\mathbf{F}_{gravity} = -mg \mathbf{k}$$

$$\begin{aligned} m\ddot{\mathbf{z}} &= \mathbf{F} \\ &= \mathbf{F}_{gravity} + \mathbf{F}_{static} \\ &= -mg \mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} \\ &= -\rho g A_{wp} \mathbf{z} \\ &= -k \mathbf{z} \end{aligned}$$



Ship will oscillate forever?

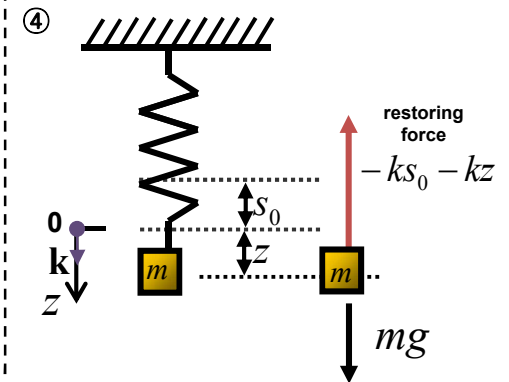
Energy is dissipated by radiation wave



$$\mathbf{F}_{radiation} = \iint_{S_B} P_{radiation} \mathbf{n} dS$$

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned} m\mathbf{z}'' &= \mathbf{F} \\ &= mg \mathbf{k} - ks_0 \mathbf{k} - k \mathbf{z} \\ &= -k \mathbf{z} \end{aligned}$$

$$m\mathbf{z}'' + k \mathbf{z} = 0 \quad \text{Oscillation by the restoring force}$$

# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

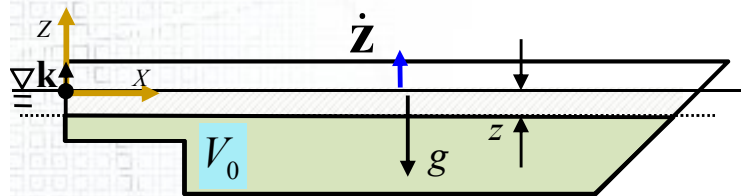
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 5



$m$  : mass       $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

$$\begin{aligned}
 m\ddot{z} &= F \\
 &= F_{gravity} + F_{static} + F_{radiation} \\
 &= -mg\mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} z - c\dot{z} \\
 &= -\rho g A_{wp} z - c\dot{z} \\
 &= -kz - c\dot{z}
 \end{aligned}$$

$$\begin{aligned}
 F_{static} &= \iint_{S_B} P_{static} n dS \\
 &= \rho g V_0 \mathbf{k} - \rho g A_{wp} z \\
 &= \rho g V_0 \mathbf{k} - kz \\
 F_{radiation} &= -c\dot{z} \\
 F_{gravity} &= -mg\mathbf{k}
 \end{aligned}$$

opposite to velocity

정수 중 선박의 강제 운동에 의해 발생한 힘

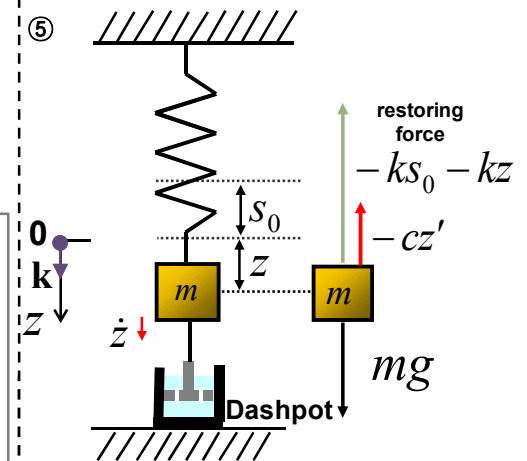
Radiation Force

$$\begin{aligned}
 F_{radiation} &= \iint_{S_B} P_{radiation} n dS \\
 &= -c\dot{z}
 \end{aligned}$$

$c$  : damping coefficient

✓ Archimedes' Principle  $F_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned}
 mz'' &= F \\
 &= mg\mathbf{k} - ks_0\mathbf{k} - kz\mathbf{k} - cz'\mathbf{k} \\
 &= -kz\mathbf{k} - cz'\mathbf{k}
 \end{aligned}$$

# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

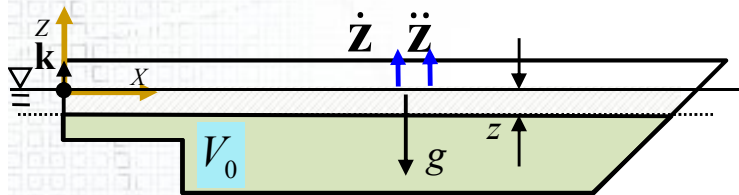
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 5

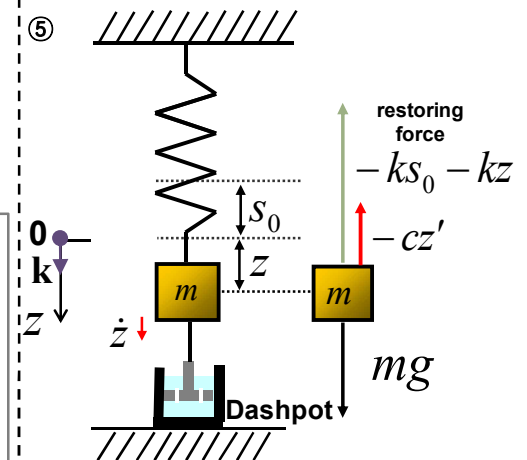


$m$  : mass       $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

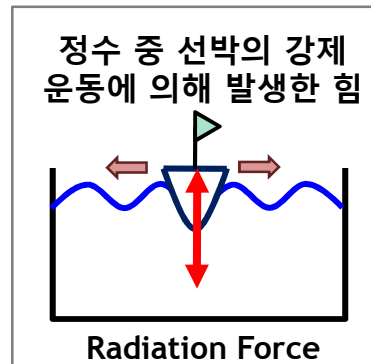
$$\begin{aligned} \mathbf{F}_{static} &= \iint_{S_B} P_{static} \mathbf{n} dS \\ &= \rho g V_0 \mathbf{k} - \rho g A_{wp} z \\ &= \rho g V_0 \mathbf{k} - kz \\ \mathbf{F}_{radiation} &= -c\dot{z} - m_a \ddot{z} \\ \mathbf{F}_{gravity} &= -mg\mathbf{k} \end{aligned}$$

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned} m\ddot{z} &= \mathbf{F} \\ &= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{radiation} \\ &= -mg\mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} z - c\dot{z} - m_a \ddot{z} \\ &= -\rho g A_{wp} z - c\dot{z} - m_a \ddot{z} \\ &= -kz - c\dot{z} - m_a \ddot{z} \end{aligned}$$



$$\begin{aligned} \mathbf{F}_{radiation} &= \iint_{S_B} P_{radiation} \mathbf{n} dS \\ &= -c\dot{z} - m_a \ddot{z} \end{aligned}$$

$c$  : damping coefficient  
 $m_a$  : added mass

$$\begin{aligned} m\mathbf{z}'' &= \mathbf{F} \\ &= mg\mathbf{k} - ks_0\mathbf{k} - kz\mathbf{k} - cz'\mathbf{k} \\ &= -kz\mathbf{k} - cz'\mathbf{k} \end{aligned}$$

# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

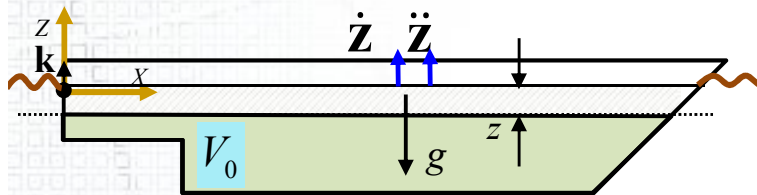
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 6

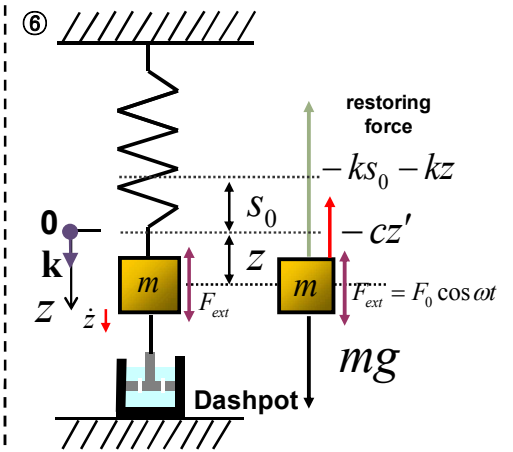


$m$  : mass       $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

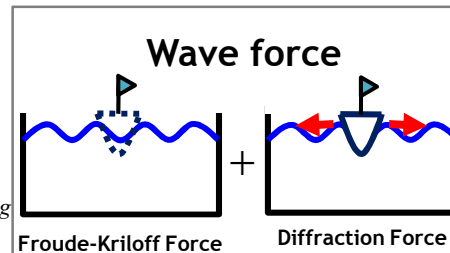
$$\begin{aligned} \mathbf{F}_{static} &= \iint_{S_B} P_{static} \mathbf{n} dS \\ &= \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} \\ &= \rho g V_0 \mathbf{k} - k \mathbf{z} \\ \mathbf{F}_{gravity} &= -m g \mathbf{k} \\ \mathbf{F}_{radiation} &= -c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} \\ \mathbf{F}_{exciting} &= F_{exciting} \end{aligned}$$

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned} m \ddot{\mathbf{z}} &= \mathbf{F} \\ &= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{radiation} + \mathbf{F}_{exciting} \\ &= -m g \mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + \mathbf{F}_{exciting} \\ &= -\rho g A_{wp} \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + \mathbf{F}_{exciting} \\ &= -k \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + \mathbf{F}_{exciting} \end{aligned}$$



$$\begin{aligned} \mathbf{F}_{wave\ exciting} &= \iint_{S_B} P_{wave\ exciting} \mathbf{n} dS \\ &= \mathbf{F}_{exciting} \end{aligned}$$

$c$  : damping coefficient  
 $m_a$  : added mass

$$\begin{aligned} m \mathbf{z}'' &= \mathbf{F} \\ &= m g \mathbf{k} - k s_0 \mathbf{k} - k \mathbf{z} - c \dot{\mathbf{z}} + F_0 \cos \omega t \\ &= -k \mathbf{z} - c \dot{\mathbf{z}} + F_0 \cos \omega t \end{aligned}$$

# Nonlinearity

Nonlinearity of the nature

Nonlinear Mathematical Model

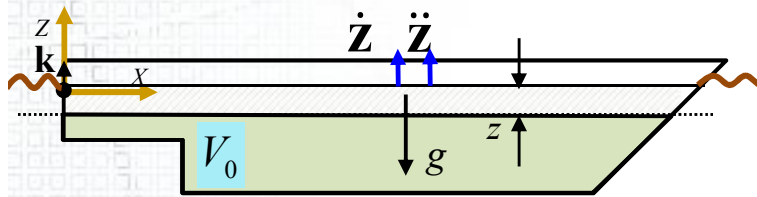
Numerical Method

Linearization

Linear Mathematical Model

Analytic Solution

## Ex) Heave Motion of a Ship – step 6



$m$  : mass  
 $\rho$  : density of sea water  
 $V_0$  : submerged volume  
 $S_B$  : submerged surface area  
 $A_{wp}$  : waterplane area

$$\begin{aligned} \mathbf{F}_{static} &= \iint_{S_B} P_{static} \mathbf{n} dS \\ &= \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} \\ &= \rho g V_0 \mathbf{k} - k \mathbf{z} \\ \mathbf{F}_{exciting} &= F_0 \cos \omega t \\ \mathbf{F}_{radiation} &= -c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} \\ \mathbf{F}_{gravity} &= -m g \mathbf{k} \end{aligned}$$

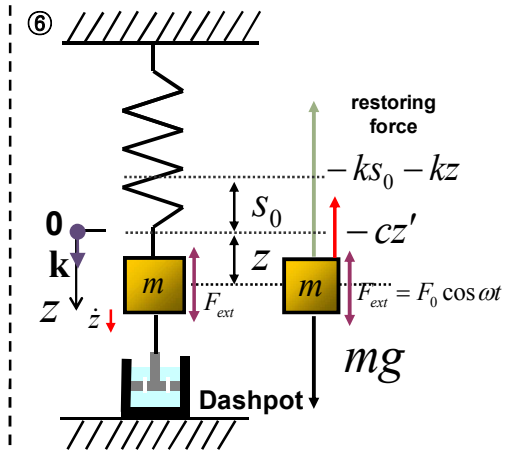
$$\begin{aligned} m \ddot{\mathbf{z}} &= \mathbf{F} \\ &= \mathbf{F}_{gravity} + \mathbf{F}_{static} + \mathbf{F}_{radiation} + \mathbf{F}_{exciting} \\ &= -m g \mathbf{k} + \rho g V_0 \mathbf{k} - \rho g A_{wp} \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + F_0 \cos \omega t \\ &= -\rho g A_{wp} \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + F_0 \cos \omega t \\ &= -k \mathbf{z} - c \dot{\mathbf{z}} - m_a \ddot{\mathbf{z}} + F_0 \cos \omega t \end{aligned}$$

$$(m + m_a) \ddot{\mathbf{z}} + c \dot{\mathbf{z}} + k \mathbf{z} = F_0 \cos \omega t$$

$c$  : damping coefficient  
 $m_a$  : added mass

✓ Archimedes' Principle  $\mathbf{F}_{static} = \rho g V_0 \mathbf{k}$

✓ Mass-Spring-Damper system



$$\begin{aligned} m \ddot{\mathbf{z}} &= \mathbf{F} \\ &= m g \mathbf{k} - k S_0 \mathbf{k} - k z \mathbf{k} - c \dot{\mathbf{z}} \mathbf{k} + F_0 \cos \omega t \\ &= -k z \mathbf{k} - c \dot{\mathbf{z}} \mathbf{k} + F_0 \cos \omega t \end{aligned}$$

$$m \ddot{\mathbf{z}} + c \dot{\mathbf{z}} + k \mathbf{z} = F_0 \cos \omega t$$





## Linked slide

: Homogeneous / Particular solution &  
Zero Input / Zero State



# Comparison : example

$$y(t) = y_h(t) + y_p(t)$$

example\*

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

solution

## 1) Homogeneous Solution

Homogeneous

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

Try :  $y_h(t) = e^{mt}$

$$(\ddot{m} + 3\dot{m} + 2)e^{mt} = 0$$

$$(m+1)(m+2) = 0$$

$$\therefore m = -1, m = -2$$

$$y_h(t) = e^{mt}$$

$e^{-t}$  and  $e^{-2t}$  : linearly independent

$$\therefore y_h(t) = c_1 e^{-t} + c_2 e^{-2t}$$

! Notation changed in the example

$$z(t) \longrightarrow y(t)$$

$$z'(t), z''(t) \longrightarrow y'(t), y''(t)$$

$$Z_{\text{transient}}, Z_{\text{steady}} \longrightarrow y_h(t), y_p(t)$$

$$Z_{\text{zero-input}}, Z_{\text{zero-state}} \longrightarrow y_0(t), y_1(t)$$

# Comparison : example

$$y(t) = y_h(t) + y_p(t)$$

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

Solution

## 2) Particular Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$\text{Try : } y_p(t) = A \sin 2t + B \cos 2t$$

$$\ddot{y}_p = -4A \sin 2t - 4B \cos 2t$$

$$\dot{y}_p = 2A \cos 2t - 2B \sin 2t$$

$$\text{L.H.S.: } -4(A \sin 2t + B \cos 2t) + 6(A \cos 2t - B \sin 2t) + 2(A \sin 2t + B \cos 2t) = (-2A - 6B) \sin 2t + (6A - 2B) \cos 2t$$

$$\text{R.H.S.: } \sin 2t$$

$$\text{L.H.S.} = \text{R.H.S.}$$

$$\begin{aligned} -2A - 6B &= 1 \\ 6A - 2B &= 0 \end{aligned} \quad \Rightarrow \quad A = -\frac{1}{20}, B = -\frac{3}{20}$$

$$\therefore y_p = -\frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

# Comparison : example

$$y(t) = y_h(t) + y_p(t)$$

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

solution

### 3) General Solution

$$y_h(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$y_p = -\frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

$$y(t) = c_1 e^{-t} + c_2 e^{-2t} - \frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

Initial condition :  $y(0) = 1, \dot{y}(0) = 5$

$$y(0) : c_1 + c_2 - \frac{3}{20} = 1$$

$$\dot{y}(0) : -c_1 - 2c_2 - \frac{2}{20} = 5$$



$$c_1 = \frac{37}{5}$$

$$c_2 = -\frac{25}{4}$$

$$y_h(t) = \frac{37}{5} e^{-t} - \frac{25}{4} e^{-2t}, \quad y_p(t) = -\frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

$$y(t) = \frac{37}{5} e^{-t} - \frac{25}{4} e^{-2t} - \frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

# Comparison : example

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

1)  
Homogeneous  
Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y_{0\_h}(t) = e^{mt}$$

$$(m+1)(m+2)e^{mt} = 0$$

$$y_{0\_h}(t) = c_1e^{-t} + c_2e^{-2t}$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = 0$$

$$y_{1\_h}(t) = e^{mt}$$

$$(m+1)(m+2)e^{mt} = 0$$

$$y_{1\_h}(t) = c_1e^{-t} + c_2e^{-2t}$$

# Comparison : example

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

1) Homogeneous Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = 0$$

$$y_{0\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$y_{1\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

2) Particular Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y_{1\_p} = A \sin 2t + B \cos 2t$$

$$\text{L.H.S.: } (-2A - 6B) \sin 2t + (6A - 2B) \cos 2t$$

$$\text{R.H.S.: } \sin 2t$$

$$\text{L.H.S.} = \text{R.H.S.} \quad A = -\frac{1}{20}, B = -\frac{3}{20}$$

$$y_{0\_p}(t) = 0$$

$$y_{1\_p} = -\frac{1}{20} \sin 2t - \frac{3}{20} \cos 2t$$

# Comparison : example

예제

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

1) Homogeneous Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y_{0\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = 0$$

$$y_{1\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

2) Particular Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y_{0\_p}(t) = 0$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y_{1\_p} = -\frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

3) General Solution

$$y_{0\_g}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$y(0) = 1, \dot{y}(0) = 5$$

$$\begin{aligned} c_1 + c_2 &= 1 \\ -c_1 - 2c_2 &= 5 \end{aligned} \quad \rightarrow \quad \begin{aligned} c_1 &= 7 \\ c_2 &= -6 \end{aligned}$$

$$y_{1\_g}(t) = c_1 e^{-t} + c_2 e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

$$\begin{aligned} y(0) : c_1 + c_2 - \frac{3}{20} &= 1 \\ \dot{y}(0) : -c_1 - 2c_2 - \frac{2}{20} &= 5 \end{aligned} \quad \rightarrow \quad \begin{aligned} c_1 &= \frac{2}{5} \\ c_2 &= -\frac{1}{4} \end{aligned}$$



# Comparison : example

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

1) Homogeneous Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = 0$$

$$y_{0\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$y_{1\_h}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

2) Particular Solution

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$\ddot{y}(t) - 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y_{0\_p}(t) = 0$$

$$y_{1\_p} = -\frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

3) General Solution

$$y_{0\_g}(t) = c_1 e^{-t} + c_2 e^{-2t}$$

$$y_{1\_g}(t) = c_1 e^{-t} + c_2 e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Initial condition

$$y(0) = 1, \dot{y}(0) = 5$$

$$y(0) = 0, \dot{y}(0) = 0$$

General Solution

$$y_{0\_g}(t) = 7e^{-t} - 6e^{-2t}$$

$$y_{1\_g}(t) = \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

# Comparison : example

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

- 1) Homogeneous Solution
- 2) Particular Solution
- 3) General Solution

$$y_{0\_h}(t) = 7e^{-t} - 6e^{-2t}$$

$$y_{0\_p}(t) = 0$$

$$y_{0\_g}(t) = y_0(t) = 7e^{-t} - 6e^{-2t}$$

$$y_{1\_h}(t) = \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t}$$

$$y_{1\_p} = -\frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

$$y_{1\_g}(t) = y_1(t) = \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

# Comparison : example

example

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t \quad y(0) = 1, \dot{y}(0) = 5$$

$$y(t) = \left(\frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t}\right) - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

Zero Input

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = 0$$

$$y(0) = 1, \dot{y}(0) = 5$$

Zero State

$$\ddot{y}(t) + 3\dot{y}(t) + 2y(t) = \sin 2t$$

$$y(0) = 0, \dot{y}(0) = 0$$

- 1) Homogeneous Solution
- 2) Particular Solution

$$y_{0\_h}(t) = 7e^{-t} - 6e^{-2t}$$

$$y_{1\_h}(t) = \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t}$$

$$y_{0\_p}(t) = 0$$

$$y_{1\_p}(t) = -\frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

- 3) General Solution

$$y_{0\_g}(t) = y_0(t) = (7e^{-t} - 6e^{-2t}) \quad y_{1\_g}(t) = y_1(t) = \left(\frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t}\right) - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

$$\begin{array}{ccccccc}
 y(t) & = & & + & & & \\
 || & & || & & || & & \\
 y_0(t) & = & & + & & & \\
 + & & + & & + & & \\
 y_1(t) & = & & + & & & 
 \end{array}$$



# Comparison : example-proof

$$y(t) = y_0(t) + y_1(t)$$

$$m\ddot{y}(t) + c\dot{y}(t) + ky(t) = u(t), u(t) \neq 0 \quad \dots(1)$$

$$y(0) = a, \dot{y}(0) = b, a \neq 0, b \neq 0 \quad \dots(2)$$

Zero Input solution :  $y_0(t)$

$$m\ddot{y}_0(t) + c\dot{y}_0(t) + ky_0(t) = 0$$

$$y_0(0) = a, \dot{y}_0(0) = b$$

Zero state solution :  $y_1(t)$

$$m\ddot{y}_1(t) + c\dot{y}_1(t) + ky_1(t) = u(t)$$

$$y_1(0) = 0, \dot{y}_1(0) = 0$$

assum. :  $y(t) = y_0(t) + y_1(t)$

$$\therefore y(t) = y_0(t) + y_1(t)$$

$y_t \rightarrow (1) :$

L.H.S.:

$$m\ddot{y}(t) + c\dot{y}(t) + ky(t)$$

$$= m(\ddot{y}_0(t) + \ddot{y}_1(t)) + c(\dot{y}_0(t) + \dot{y}_1(t)) + k(y_0(t) + y_1(t))$$

$$= [m\ddot{y}_0(t) + c\dot{y}_0(t) + ky_0(t)] + [m\ddot{y}_1(t) + c\dot{y}_1(t) + ky_1(t)]$$

$$= 0 + u(t)$$

R.H.S.:  $u(t)$

$\therefore$  L.H.S.=R.H.S “(1) satisfied”

$y_t \rightarrow (2) :$

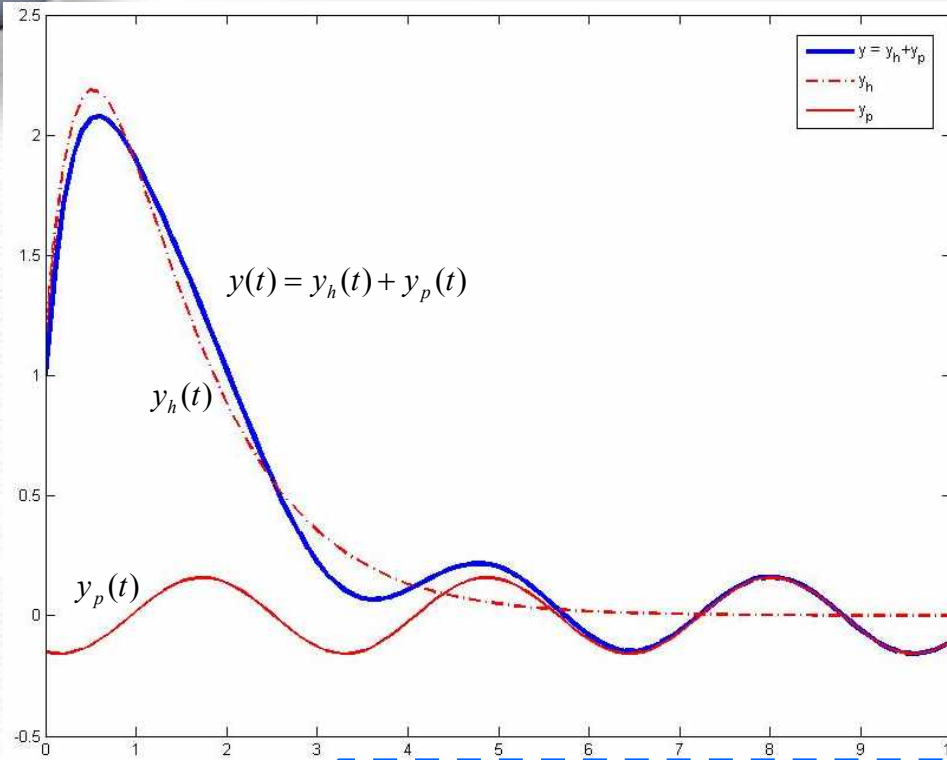
$$y(0) = y_0(0) + y_1(0) = a + 0$$

$$\dot{y}(0) = \dot{y}_0(0) + \dot{y}_1(0) = b + 0$$

$\therefore y(0) = a$  “(2) satisfied”

$$\dot{y}(0) = b$$

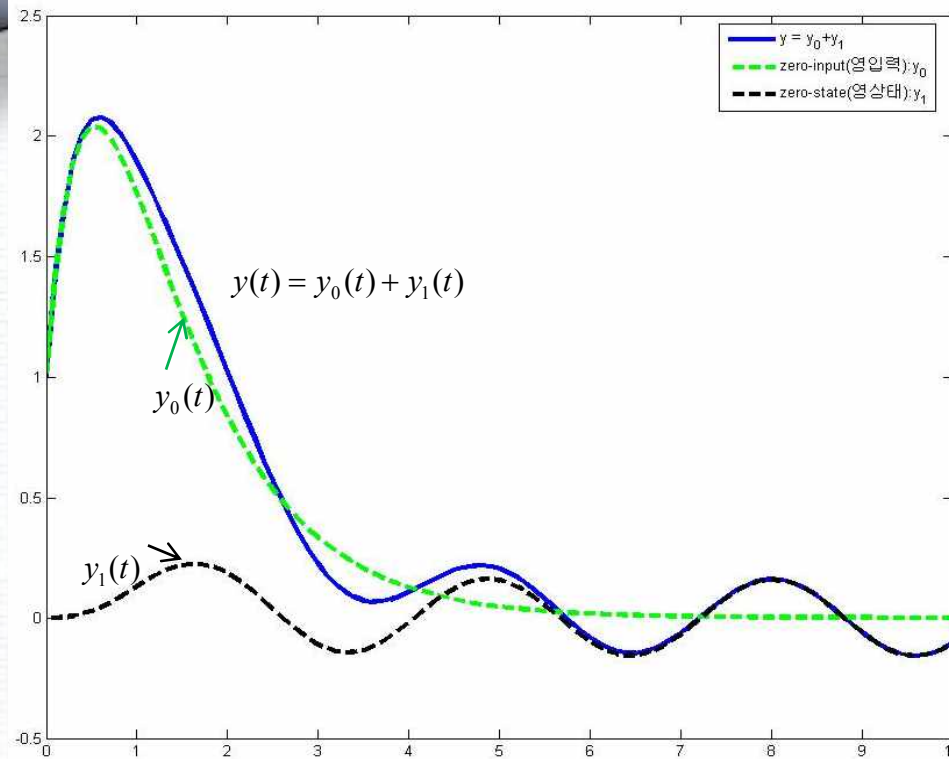
# Comparison : graph



$$y(t) = y_h(t) + y_p(t)$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

# Comparison : graph



Zero input solution  
 $u(t)=0$

Zero state solution  
 $y(0)=0, \dot{y}(0)=0$

$$y(t) = y_h(t) + y_p(t)$$

$$\parallel$$

$$y_0(t)$$

$$+$$

$$y_1(t)$$

$$y(t) = \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$

$$\parallel$$

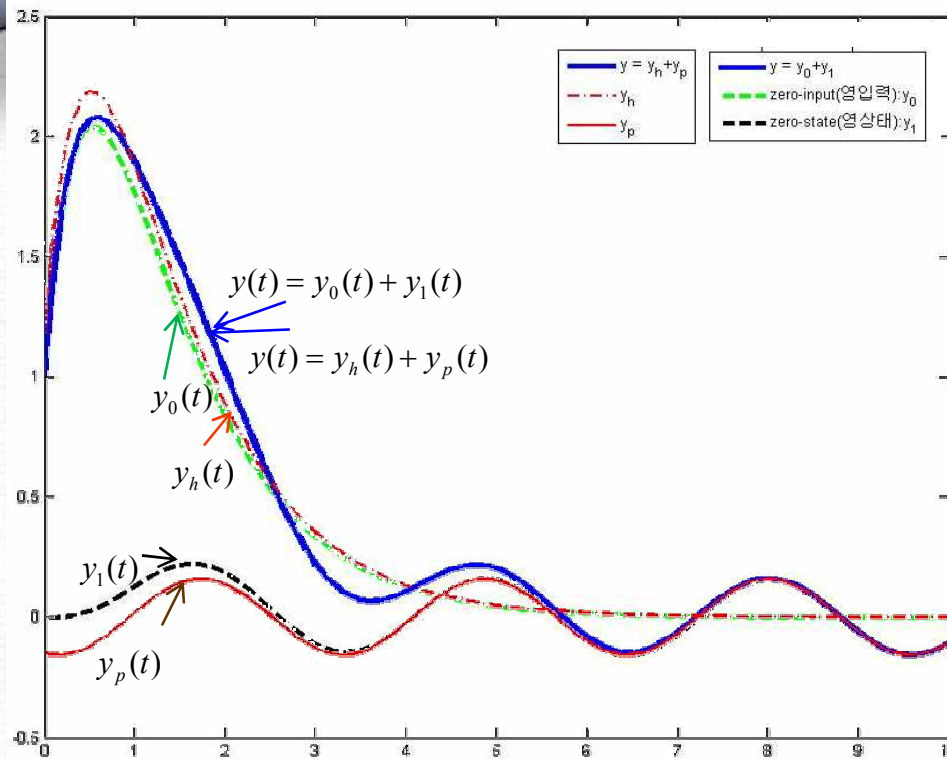
$$y_0(t) = 7e^{-t} - 6e^{-2t}$$

$$+$$

$$y_1(t) = \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t$$



# Comparison : graph



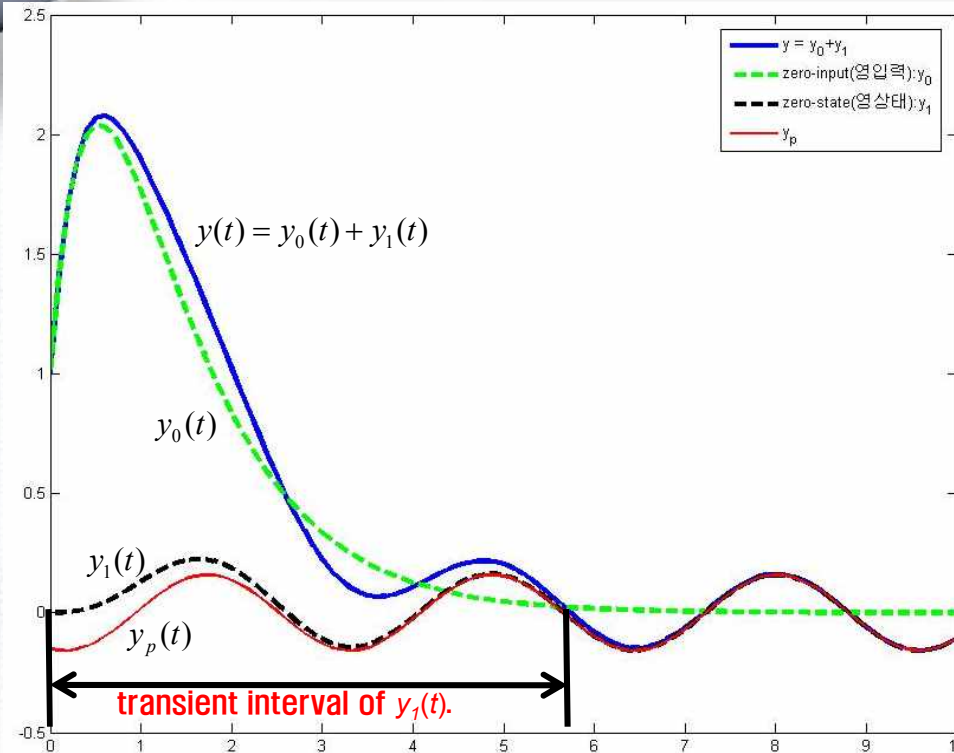
Zero input solution  
 $u(t)=0$

Zero state solution  
 $y(0)=0, \dot{y}(0)=0$

$$\begin{aligned}
 & y(t) = y_h(t) + y_p(t) \\
 & \parallel \\
 & y_0(t) \\
 & + \\
 & y_1(t)
 \end{aligned}$$

$$\begin{aligned}
 y(t) &= \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t \\
 &\parallel \\
 y_0(t) &= 7e^{-t} - 6e^{-2t} \\
 + \\
 y_1(t) &= \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t
 \end{aligned}$$

# Comparison : graph



Zero Input solution  
 $u(t)=0$

Zero state solution  
 $y(0)=0, \dot{y}(0)=0$

$$\begin{aligned}
 y(t) &= y_h(t) + y_p(t) \\
 &\parallel \quad \parallel \quad \parallel \\
 y_0(t) &= y_{0\_h}(t) + y_{0\_p}(t) \\
 &+ \quad + \\
 y_1(t) &= \boxed{y_{1\_h}(t)} + y_{1\_p}(t)
 \end{aligned}$$

$$\begin{aligned}
 y(t) &= \frac{37}{5}e^{-t} - \frac{25}{4}e^{-2t} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t \\
 y_0(t) &= 7e^{-t} - 6e^{-2t} \\
 y_1(t) &= \boxed{\left[ \frac{2}{5}e^{-t} - \frac{1}{4}e^{-2t} \right]} - \frac{1}{20}\sin 2t - \frac{3}{20}\cos 2t
 \end{aligned}$$