

2009 fall

# Phase Transformation of Materials

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# Contents for previous class

- **Effect of Temperature on Solid Solubility**

$$X_B^e = A \exp\left\{-\frac{Q}{RT}\right\}$$

- **Equilibrium Vacancy Concentration**

$$X_V^e = \exp\frac{-\Delta G_V}{RT}$$

- **Influence of Interfaces on Equilibrium**

$$\Delta G = \frac{2\gamma V_m}{r} \quad \text{Gibbs-Thomson effect}$$

- **Gibbs-Duhem Equation:** Be able to calculate the change in chemical potential that result from a change in alloy composition.

$$X_A X_B \frac{d^2 G}{dX^2} = RT \left\{ 1 + \frac{d \ln \gamma_A}{d \ln X_A} \right\} = RT \left\{ 1 + \frac{d \ln \gamma_B}{d \ln X_B} \right\}$$

- **Ternary Equilibrium: Ternary Phase Diagram**

# Contents in Phase Transformation

상변태를  
이해하는데  
필요한 배경

(Ch1) 열역학과 상태도: **Thermodynamics**

(Ch2) 확산론: **Kinetics**

(Ch3) 결정계면과 미세조직

대표적인 상변태

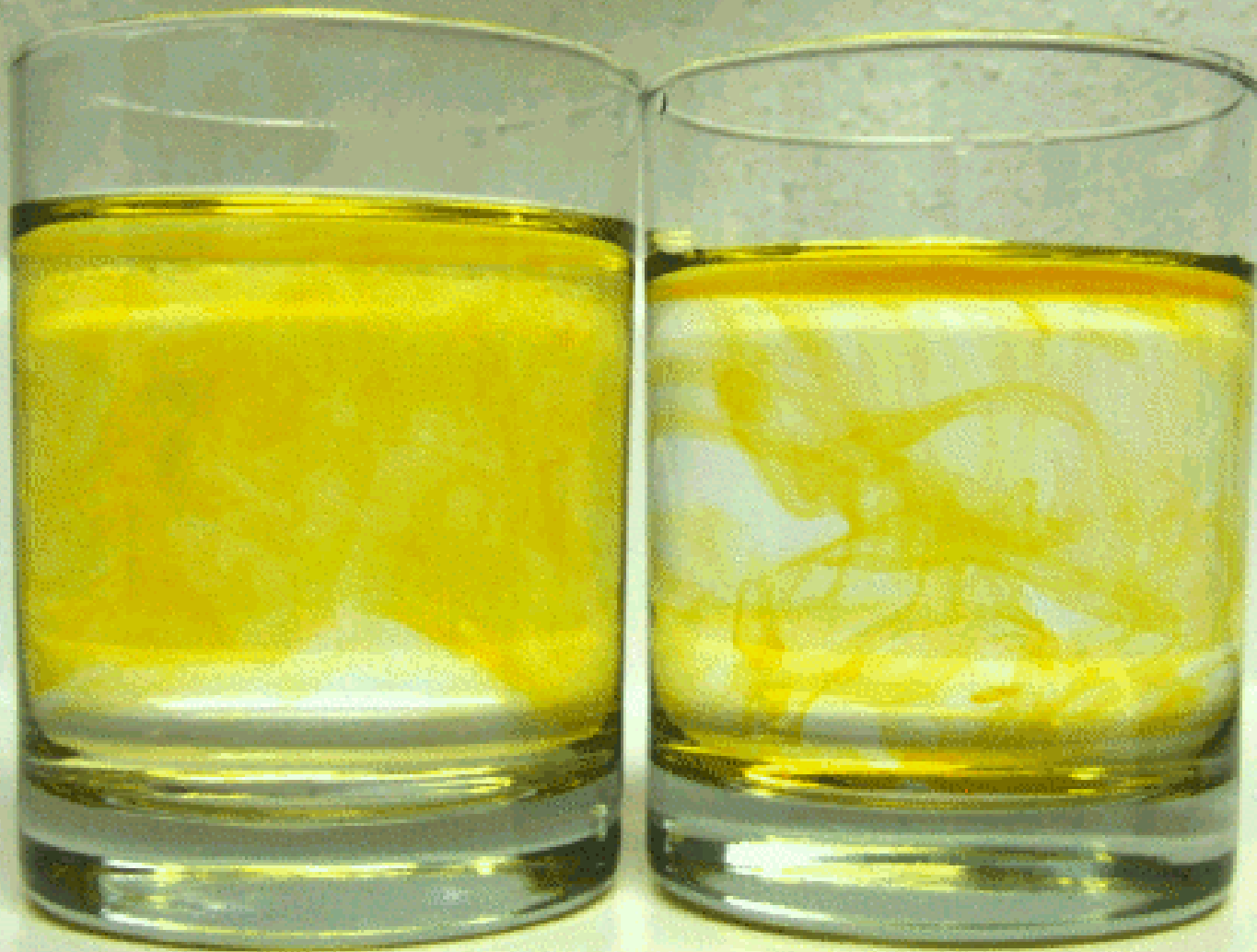
(Ch4) 응고: **Liquid → Solid**

(Ch5) 고체에서의 확산 변태: **Solid → Solid (Diffusional)**

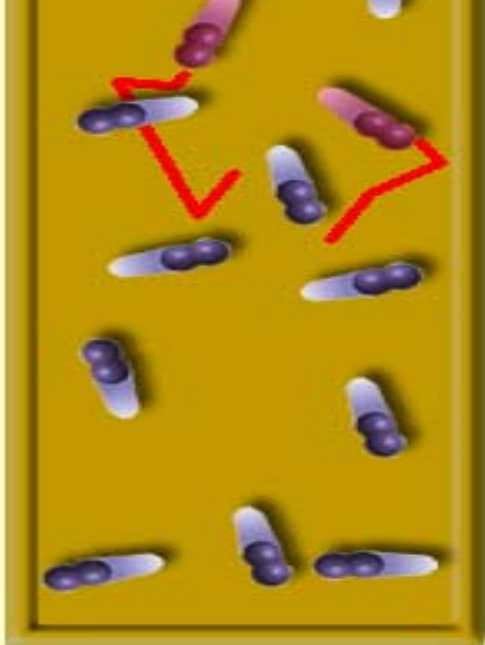
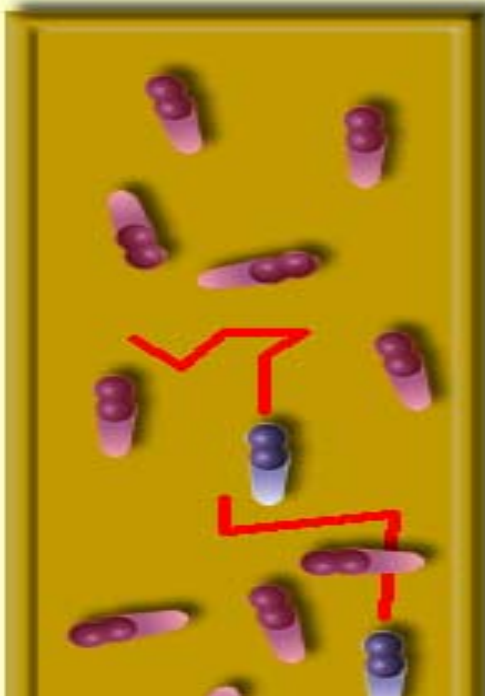
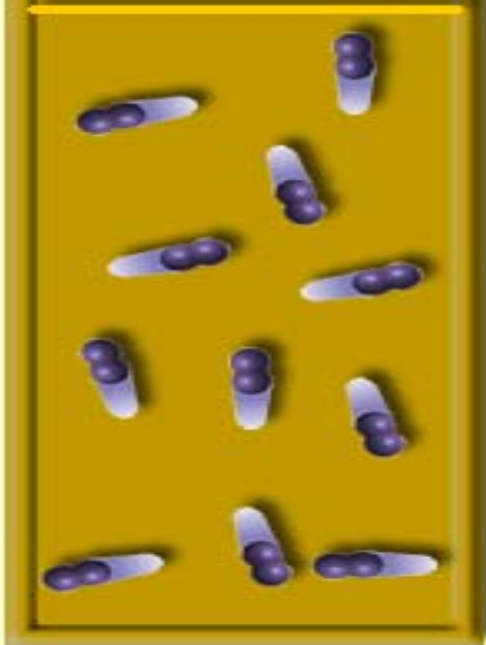
(Ch6) 고체에서의 무확산 변태: **Solid → Solid (Diffusionless)**

# Contents for today's class

- **Diffusion**
- **Interstitial Diffusion – Fick's First Law**
- **Effect of Temperature on Diffusivity**
- **Nonsteady-state diffusion – Fick's Second Law**

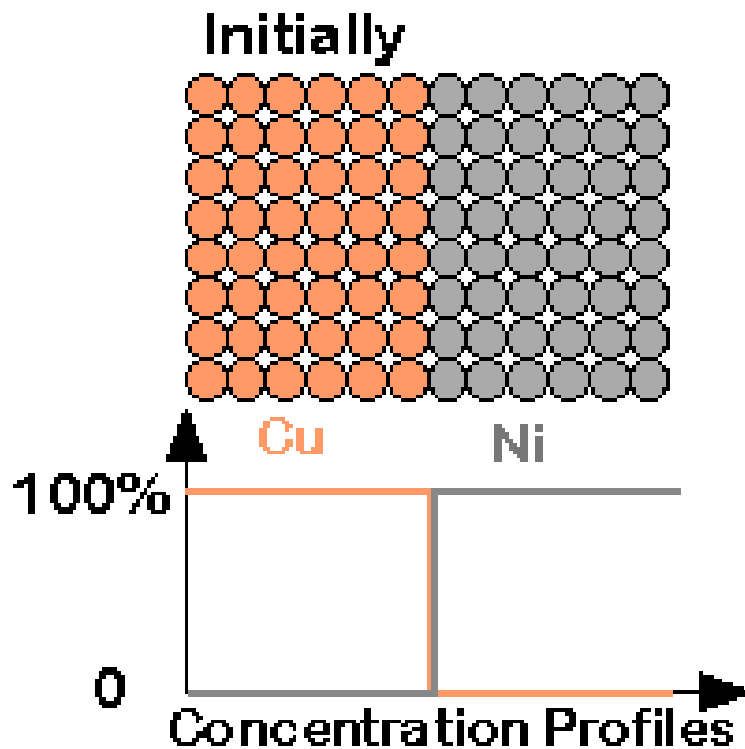


**Diffusion**

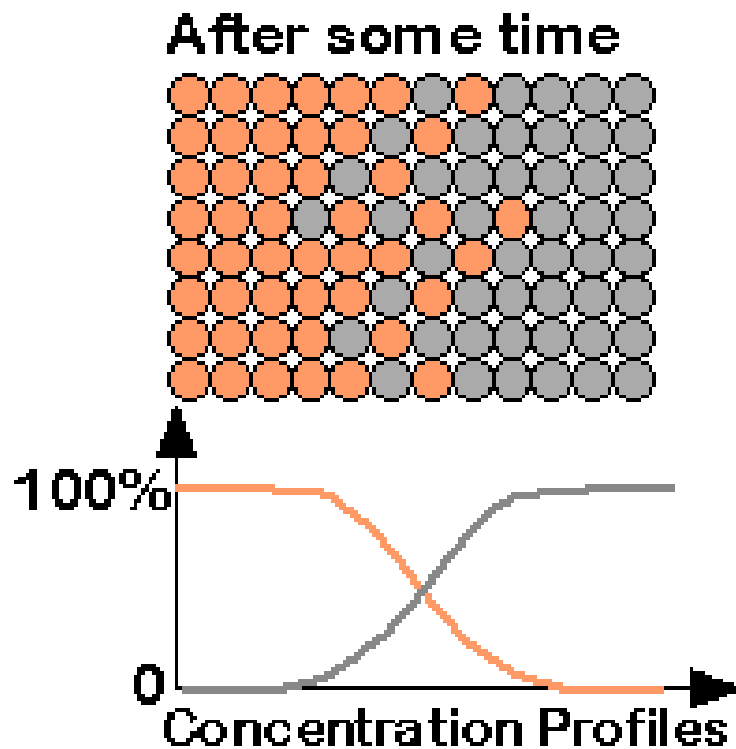


# Diffusion: THE PHENOMENON

- **Interdiffusion:** in a solid with more than one type of element (an alloy), atoms tend to migrate from regions of large concentration.

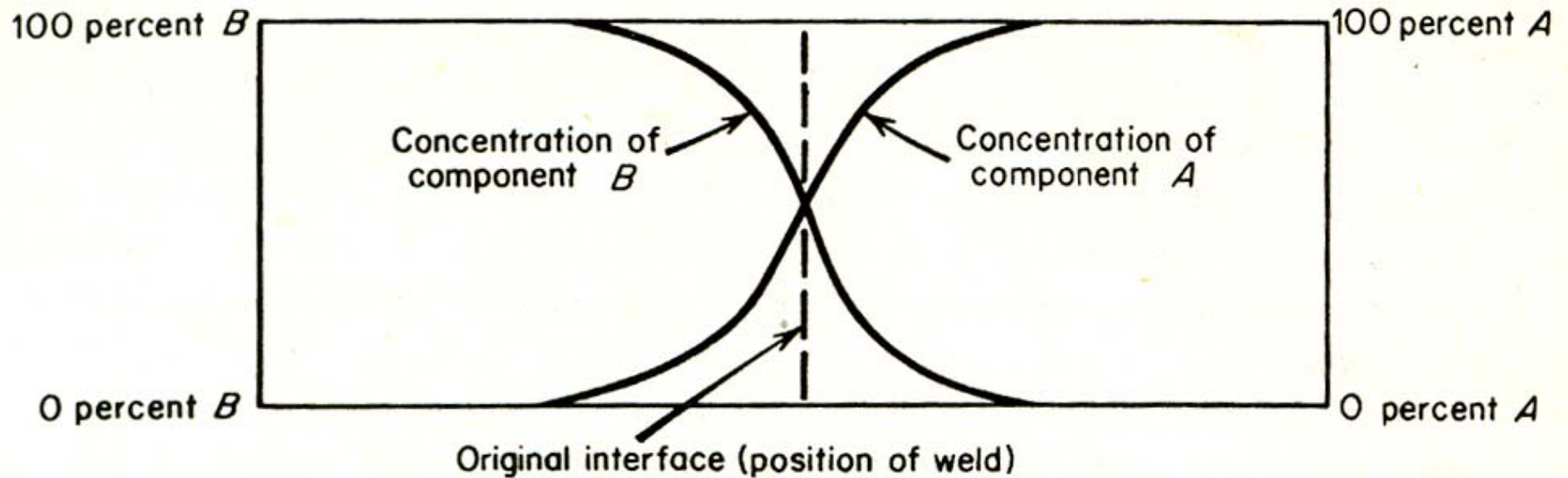


**Substitutional diffusion**



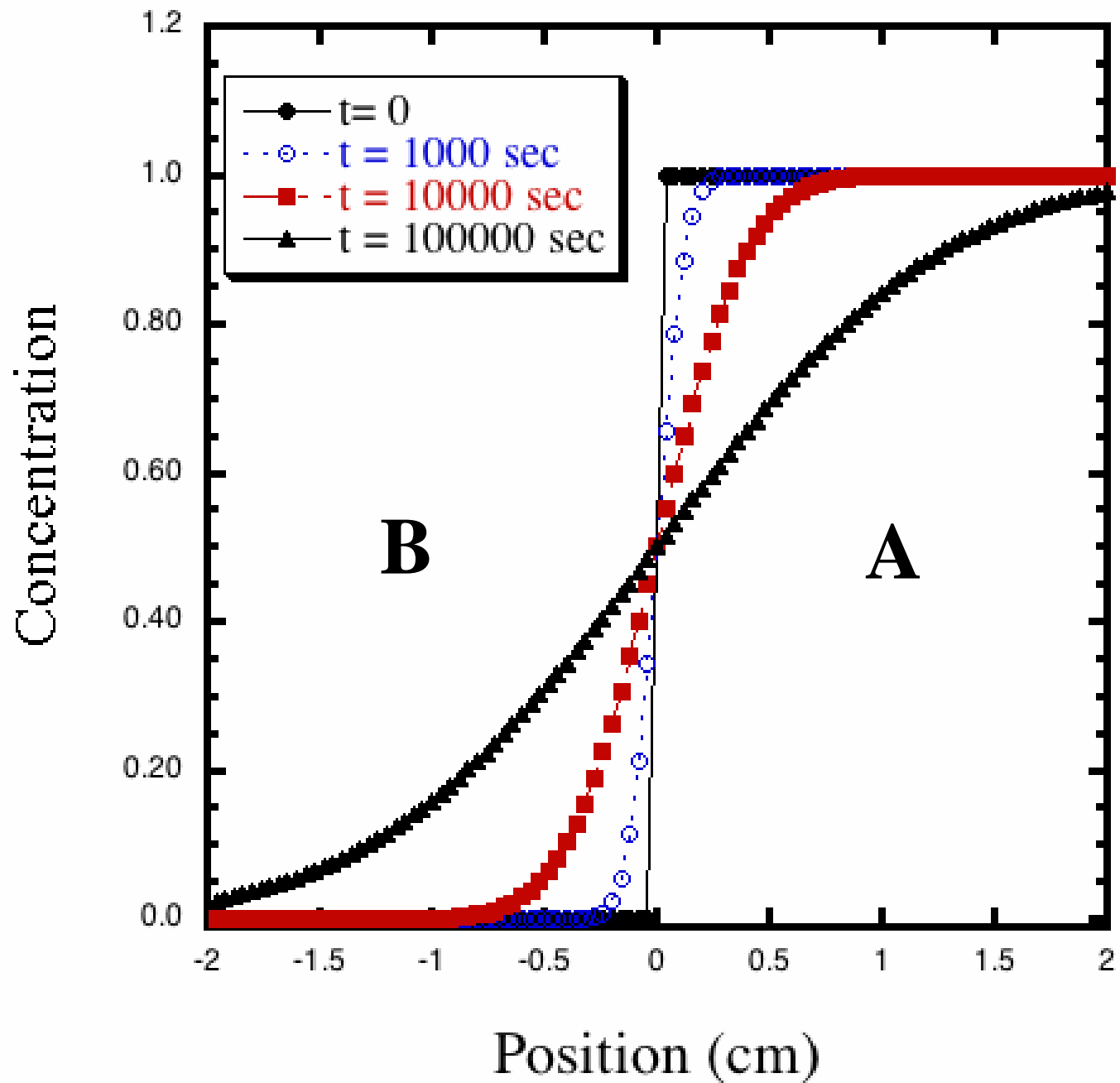
**Interstitial diffusion**

# Mechanism of Solid-State Diffusion



**Fig. 10.6** Curves showing concentration as a function of distance along a diffusion couple. Curves of this type are usually called *penetration curves*.





# Diffusion

Movement of atoms to reduce its chemical potential  $\mu$ .

**Driving force: Reduction of G**

## Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.

## Up-hill diffusion

movement of atoms from a low concentration region to high concentration region

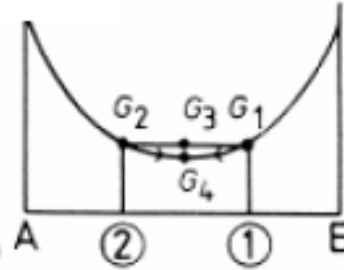
$$\begin{aligned} \mu_A^1 &< \mu_A^2 \\ \mu_B^1 &> \mu_B^2 \end{aligned}$$

'down-hill' diffusion,



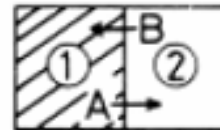
B-rich A-rich

(a)



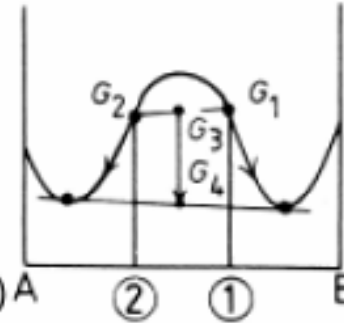
(b)

'up hill' diffusion.

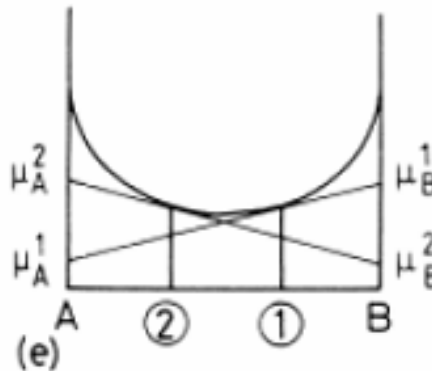


B-rich A-rich

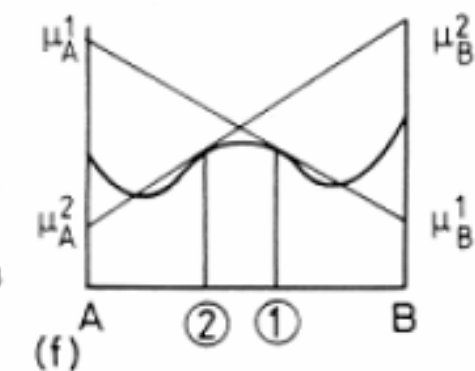
(c)



(d)



(e)

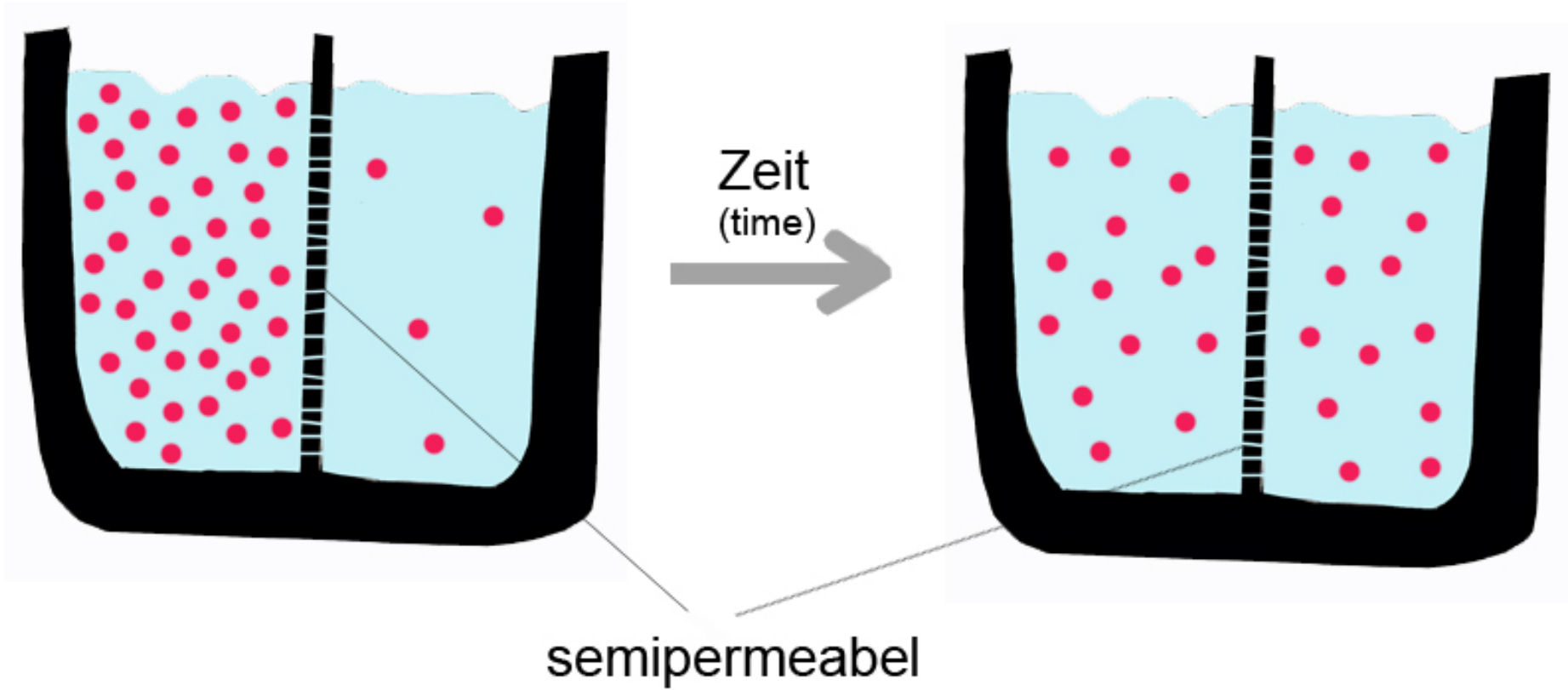


(f)

$$\begin{aligned} \mu_A^1 &> \mu_A^2 \\ \mu_B^1 &< \mu_B^2 \end{aligned}$$

## Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.



# Diffusion

Movement of atoms to reduce its chemical potential  $\mu$ .

Driving force: Reduction of G

## Down-hill diffusion

movement of atoms from a high concentration region to low concentration region.

## Up-hill diffusion

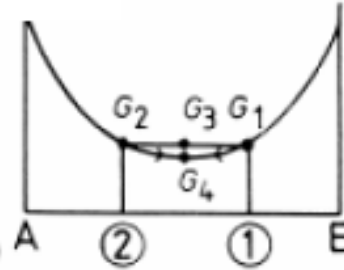
movement of atoms from a low concentration region to high concentration region

'down-hill' diffusion,



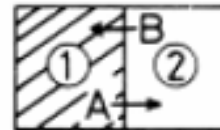
B-rich A-rich

(a)



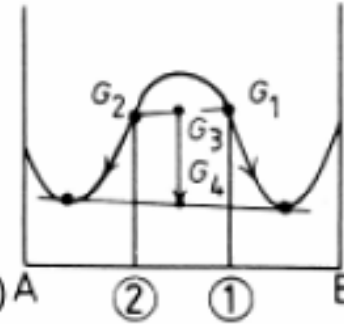
(b) A ② ① B

'up hill' diffusion.



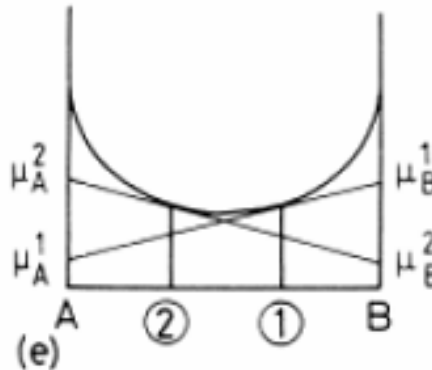
B-rich A-rich

(c)

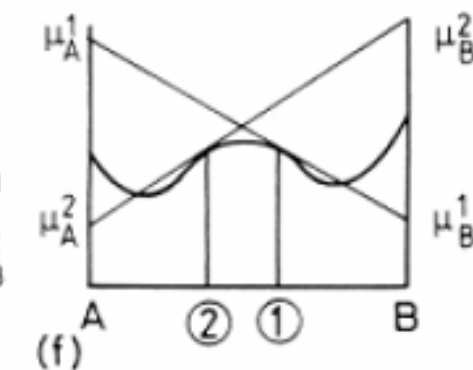


(d) A ② ① B

$$\begin{aligned} \mu_A^1 &< \mu_A^2 \\ \mu_B^1 &> \mu_B^2 \end{aligned}$$



(e)



(f)

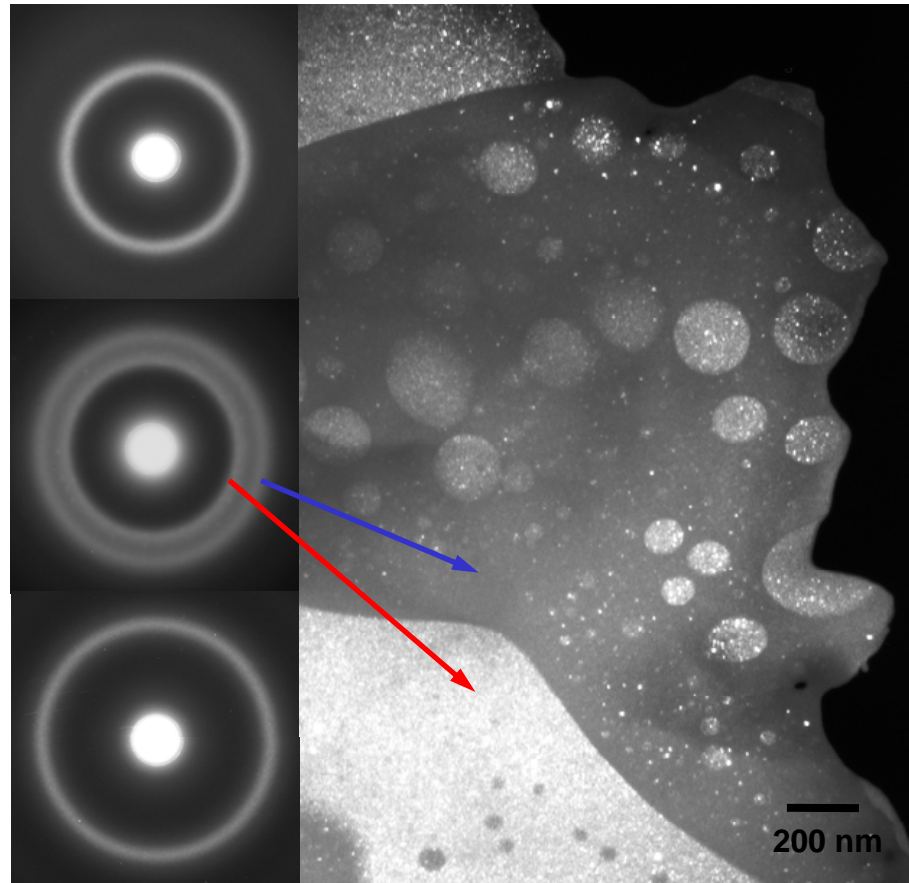
$$\begin{aligned} \mu_A^1 &> \mu_A^2 \\ \mu_B^1 &< \mu_B^2 \end{aligned}$$

# Phase separation in metallic glasses

$\text{Nd}_{60}\text{Al}_{10}\text{Co}_{30}$   
2.91 Å

$\text{Nd}_{30}\text{Zr}_{30}\text{Al}_{10}\text{Co}_{30}$   
2.37 Å, 2.99 Å

$\text{Zr}_{60}\text{Al}_{10}\text{Co}_{30}$   
2.40 Å



TEM results for  $\text{Nd}_{30}\text{Zr}_{30}\text{Al}_{10}\text{Co}_{30}$  alloy

# Diffusion

Diffusion : Mechanism by which matter transported through matter

What is the driving force for diffusion?

⇒ a concentration gradient (x)

⇒ a chemical potential (o)

But this chapter will explain with

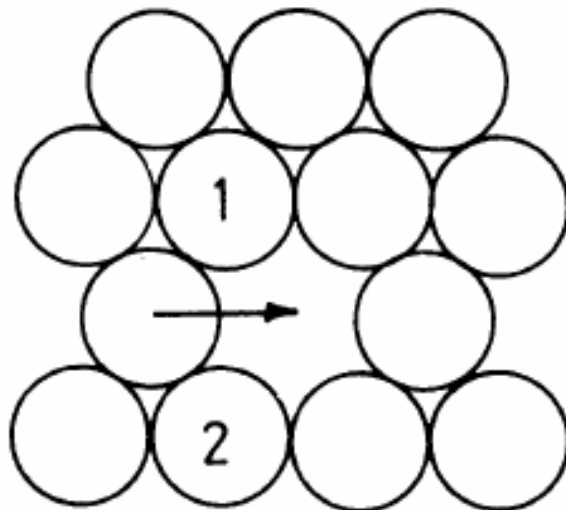
“concentration gradients for a convenience”.

# Atomic mechanisms of diffusion

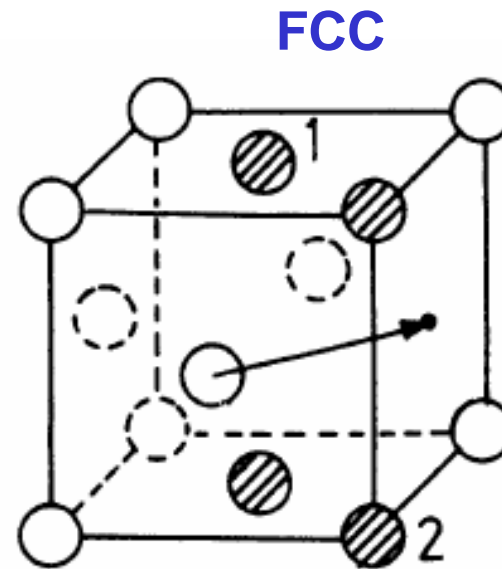
## Substitutional vs. Interstitial diffusion

The mean vibrational energy is  $3kT$ , therefore the amplitude of thermal oscillation of an atom increases markedly with temperature. If a neighboring atom site is vacant, then the oscillation atom can change site with a certain probability (atom jumping). The probability consists of creation of the vacant site and movement of a neighboring atom into the site.

## Substitutional diffusion



(a)



(b)

# Interstitial diffusion

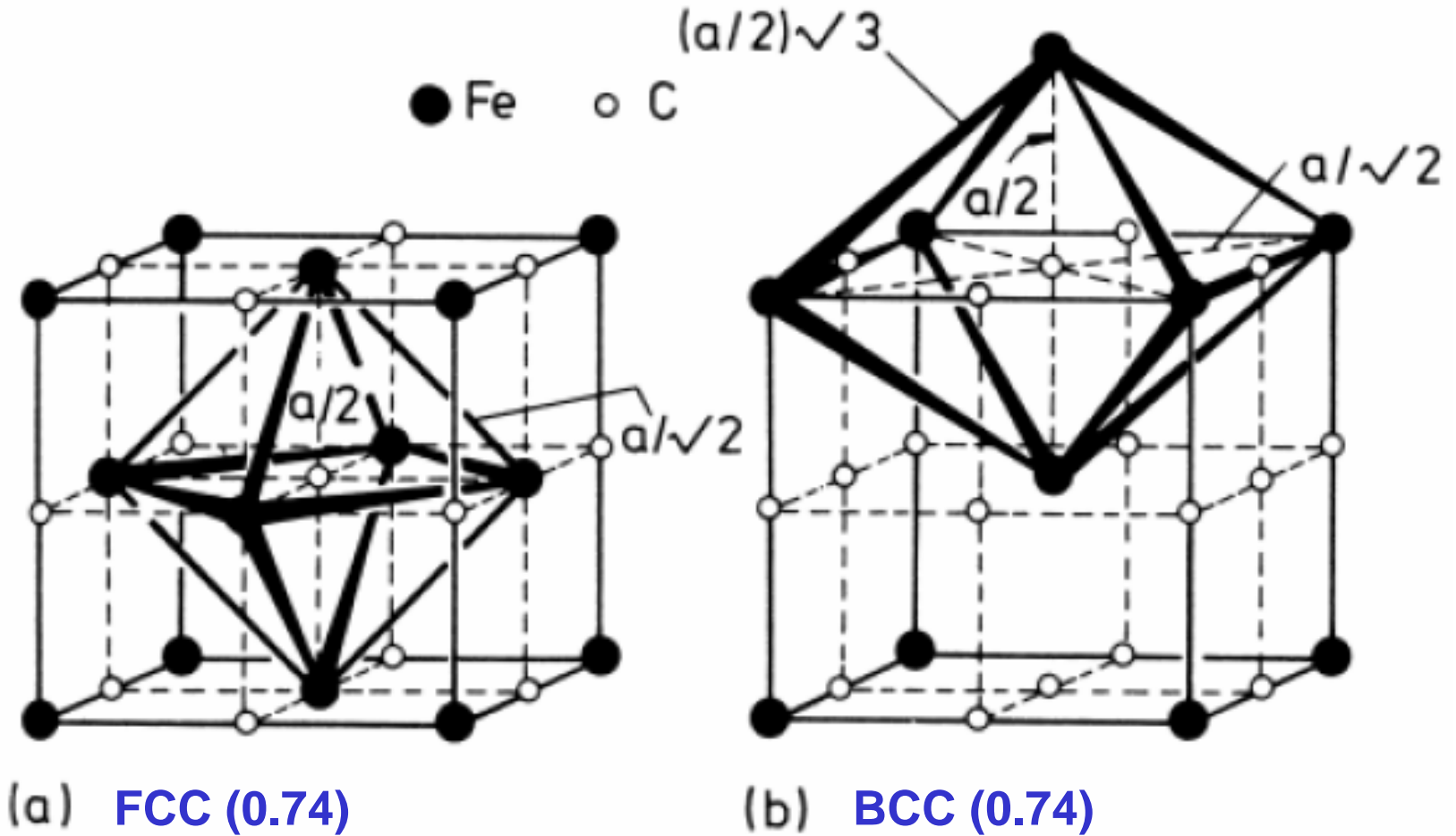
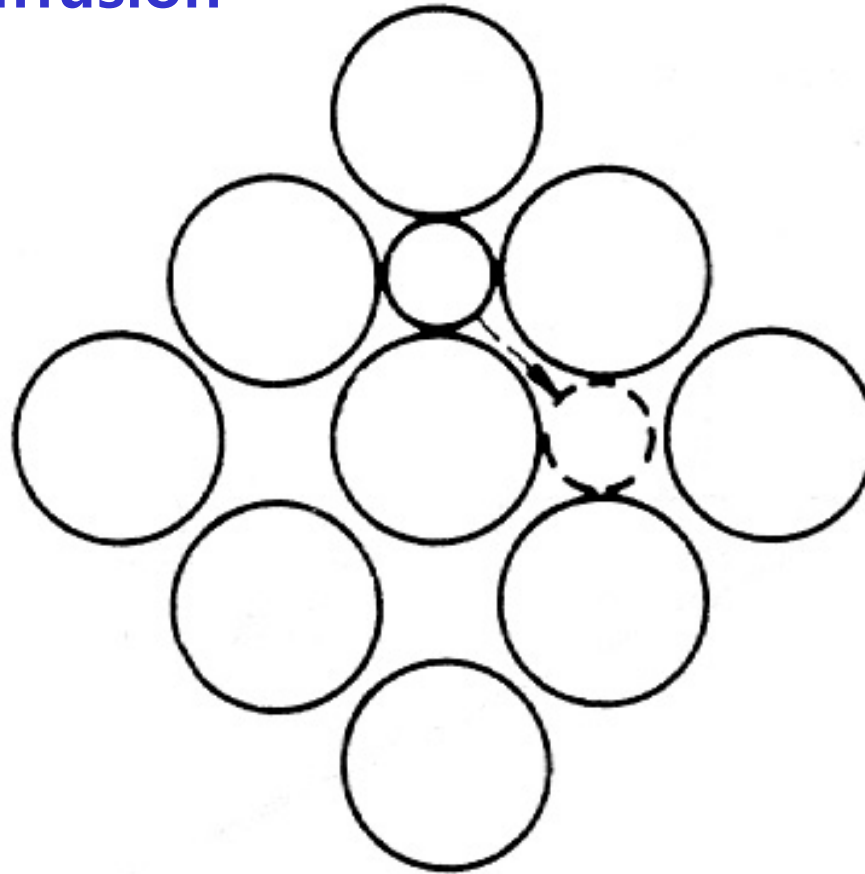


Fig. 2.3 (a) Octahedral interstices (0) in an fcc crystal. (b) Octahedral interstices in a bcc crystal. (After P. Haasen, *Physical Metallurgy*, Cambridge University Press, Cambridge, 1978.)



## Interstitial diffusion



**How interstitial diffusion differs from substitutional diffusion?**

# Interstitial diffusion

Random jump of solute B atoms in a dilute solid solution of a simple cubic lattice

Assume that there is **no lattice distortion** and also that there are **always six vacant sites around the diffusion atom**.

$$J_B = \frac{1}{6} \Gamma_B n_1 - \frac{1}{6} \Gamma_B n_2$$

$J_B$  : Flux of B atom

$\Gamma_B$  : Average jump rate of B atoms

$n_1$  : # of atoms per unit area of plane 1

$n_2$  : # of atoms per unit area of plane 2

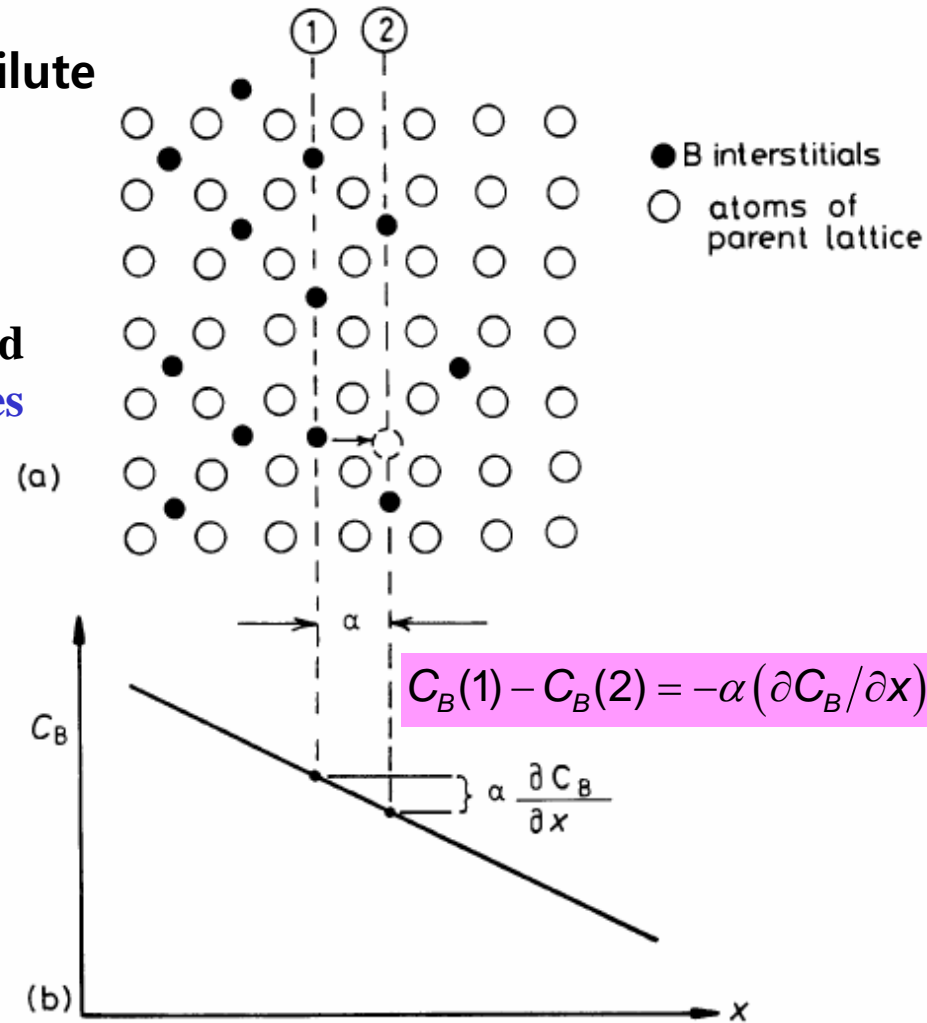


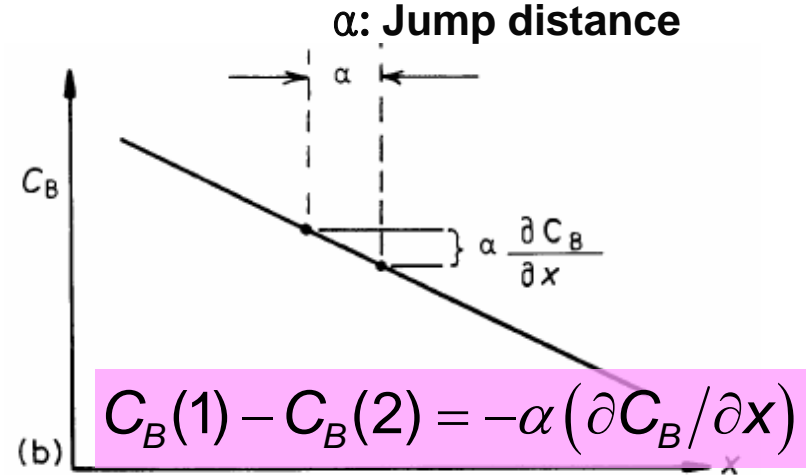
Fig. 2.5 Interstitial diffusion by random jumps in a concentration gradient.

# Fick's First Law of Diffusion

$$J_B = \frac{1}{6} \Gamma_B n_1 - \frac{1}{6} \Gamma_B n_2 = \frac{1}{6} \Gamma_B (n_1 - n_2)$$

$$C_B(1) = n_1/\alpha, \quad C_B(2) = n_2/\alpha$$

$$(n_1 - n_2) = \alpha (C_B(1) - C_B(2))$$



$$J_B = \frac{1}{6} \Gamma_B (n_1 - n_2) = - \left( \frac{1}{6} \Gamma_B \alpha^2 \right) \frac{\partial C_B}{\partial x} = -D_B \frac{\partial C_B}{\partial x}$$

➔ Concentration varies with position. (atoms m<sup>-2</sup> s<sup>-1</sup>)

**D<sub>B</sub>**: Intrinsic diffusivity or  
**Diffusion coefficient of B**  
 ⇒ depends on  
 microstructure of materials

**Magnitude of D in various media**

Gas :  $D \approx 10^{-1} \text{ cm}^2/\text{s}$

Liquid :  $D \approx 10^{-4} \sim 10^{-5} \text{ cm}^2/\text{s}$

Solid : Materials near melting temp.  $D \approx 10^{-8} \text{ cm}^2/\text{s}$

Elemental semiconductor (Si, Ge)  $D \approx 10^{-12} \text{ cm}^2/\text{s}$

Estimate the jump frequency of a carbon atom in  $\gamma$ -Fe (FCC) at 1000 °C.

$$D_B = \frac{1}{6} \Gamma_B \alpha^2$$

lattice parameter of  $\gamma$ -Fe :  $\sim 0.37$  nm

$$\alpha = 0.37 / \sqrt{2} = 0.26 \text{ nm} \quad D_C = 2.5 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}$$

$$\Gamma = 2 \times 10^9 \text{ jumps s}^{-1}$$

the vibration frequency of carbon :  $\sim 10^{13}$

If the crystal structure is not cubic, then the probability of jump is **anisotropic**. For example the **probability of jumps** along the basal direction and the axial direction of hcp crystal are **different**.

$D_B$ , in general, is concentration-dependent.

For example, in  $\gamma$ -Fe at 1000 °C,

$$D_c = 2.5 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}, \text{ when } C_c = 0.15 \text{ wt\%}$$

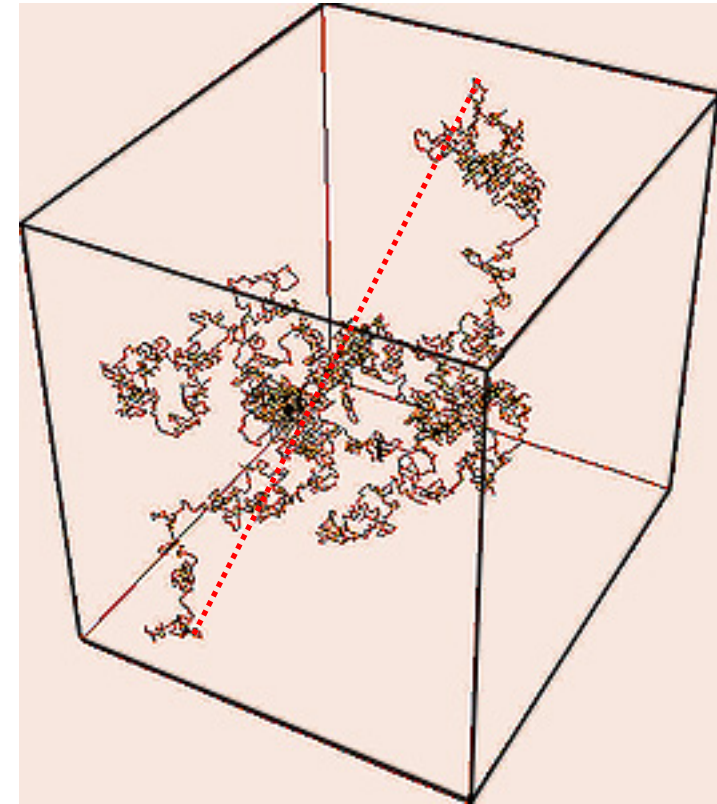
$$D_c = 7.7 \times 10^{-11} \text{ m}^2 \text{ s}^{-1}, \text{ when } C_c = 1.4 \text{ wt\%}$$

$$C_c \uparrow \rightarrow D_c \uparrow$$

For random walk in 3 dimensions,  
after n steps of length  $\alpha$

$$\rightarrow \alpha\sqrt{n}$$

During random walk,  
an atom will move in time (t)  
a radial distance (r) from the origin



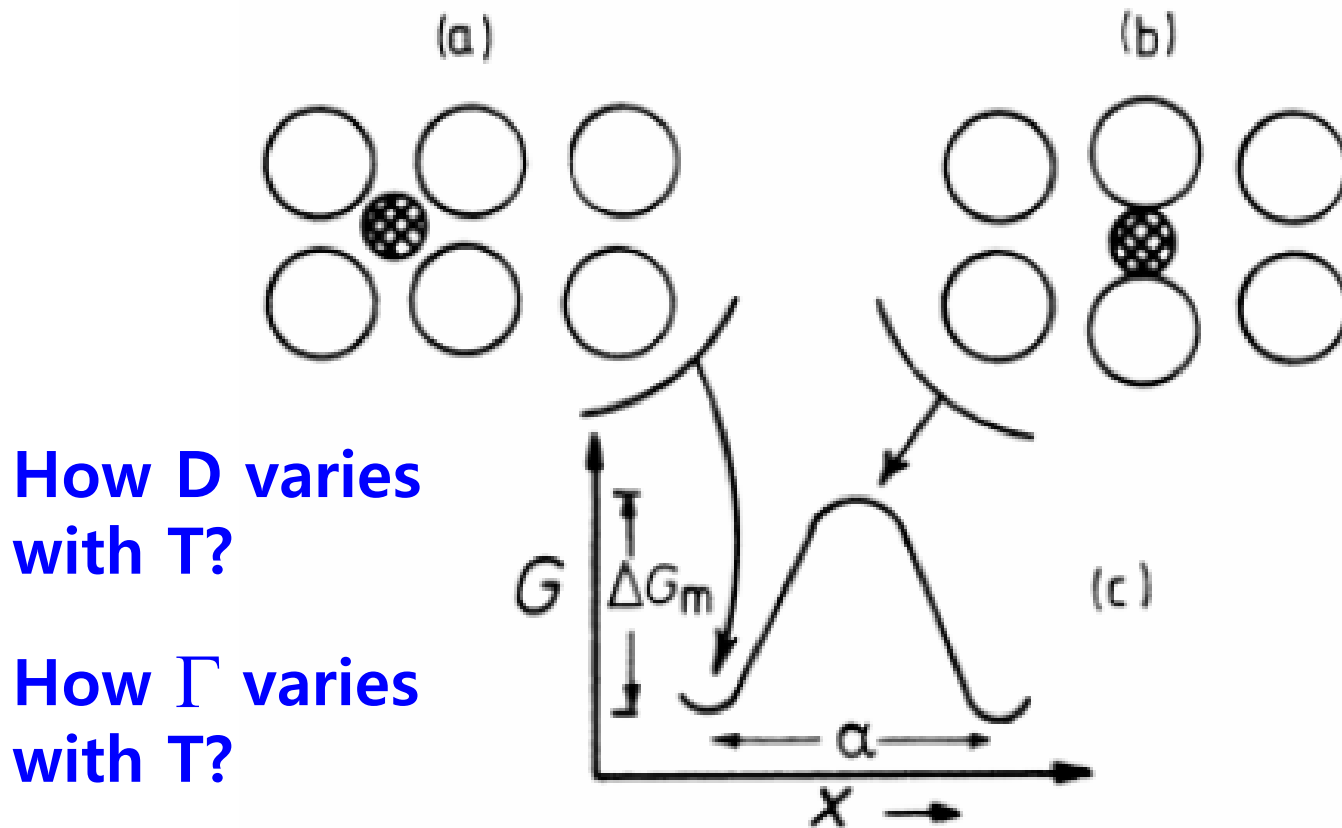
$$r = \alpha\sqrt{(\Gamma t)} = f(D) = 2.4\sqrt{Dt} \quad \because \Gamma = \frac{6D}{\alpha^2} \rightarrow r = \sqrt{6}\sqrt{Dt}$$

$$D_B = \frac{1}{6}\Gamma_B\alpha^2$$

In 1 s, each carbon atom will move a total distance of **~ 0.5 m**  
→ a net displacement : **~ 10  $\mu\text{m}$ .**

# EFFECT OF TEMPERATURE on Diffusivity

## Thermal Activation



How  $D$  varies  
with  $T$ ?

How  $\Gamma$  varies  
with  $T$ ?

Fig. 2.6 Interstitial atom, (a) in equilibrium position, (b) at the position of maximum lattice distortion. (c) Variation of the free energy of the lattice as a function of the position of interstitial. (After P.G. Shewmon, in *physical Metallurgy*, 2<sup>nd</sup> edn., R.W. Cahn (Ed.), North-Holland, Amsterdam, 1974.)

## Thermally activated process *jump frequency* $\Gamma_B$ ?

$$\Gamma_B = Z \nu \exp(-\Delta G_m / RT)$$

$Z$  : nearest neighbor sites

$\nu$  : vibration frequency

$\Delta G_m$  : activation energy for moving

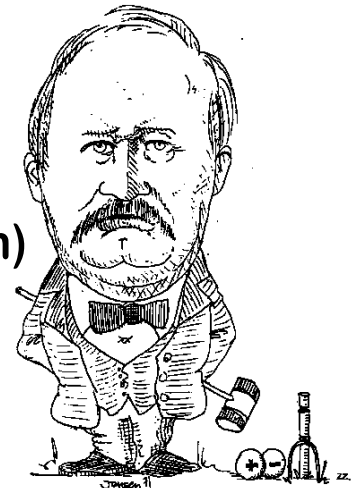
$$\left\{ \Delta G_m = \Delta H_m - T \Delta S_m, D_B = \frac{1}{6} \Gamma_B \alpha^2 \right\}$$

$$D_B = \left[ \frac{1}{6} \alpha^2 Z \nu \exp(\Delta S_m / R) \right] \exp(-\Delta H_m / RT)$$

$$\Delta H_m \equiv Q_{ID}$$

$$D_B = D_{B0} \exp \frac{-Q_{ID}}{RT}$$

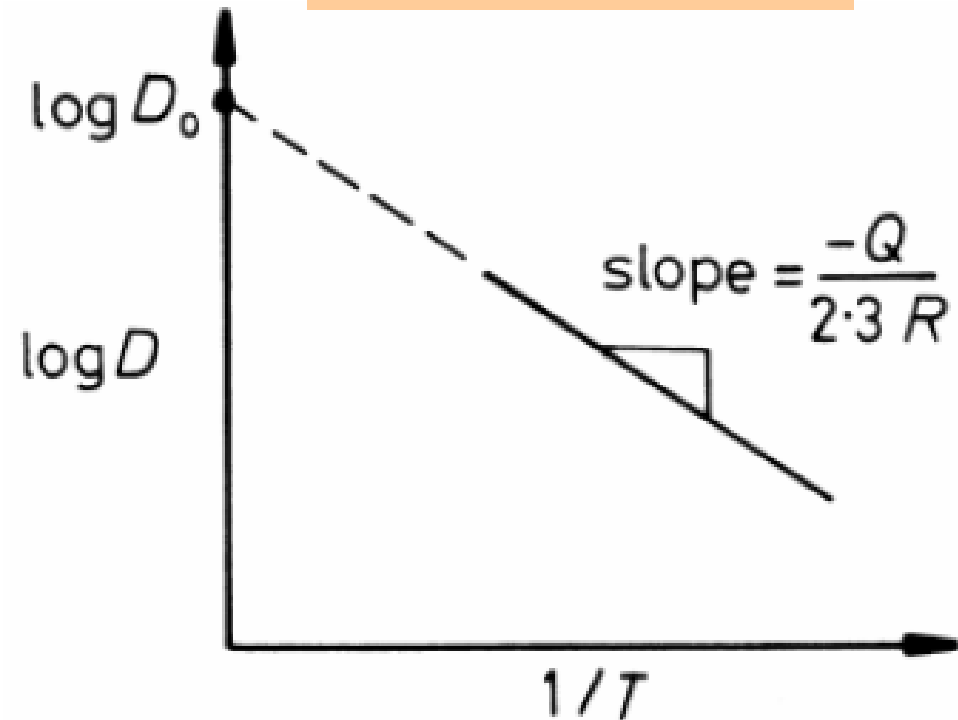
(Arrhenius-type equation)



# Temperature Dependence of Diffusion

$$D_B = D_{B0} \exp \frac{-Q_{ID}}{RT}$$

How to determine  $Q_{ID}$  experimentally?



$$\log D = \log D_0 - \frac{Q}{2.3R} \left( \frac{1}{T} \right)$$

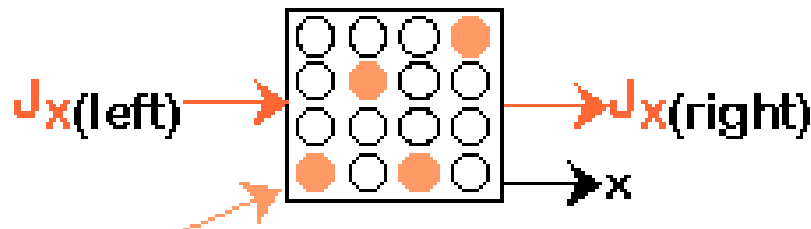
➔ Therefore, from the slope of the  $D$ -curve in an  $\log D$  vs  $1/T$  coordinate, the activation energy may be found.

Fig. 2.7 The slope of  $\log D$  v.  $1/T$  gives the activation energy for diffusion  $Q$ .



# Steady-state diffusion

The simplest type of diffusion to deal with is when the concentration at every point does not change with time.



**Steady State:**  
 $J_x(\text{left}) = J_x(\text{right})$

Concentration,  $C$ , in the box

- Apply Fick's First Law:

$$J_x = -D \frac{dC}{dx}$$

If  $J_x(\text{left}) = J_x(\text{right})$ , then

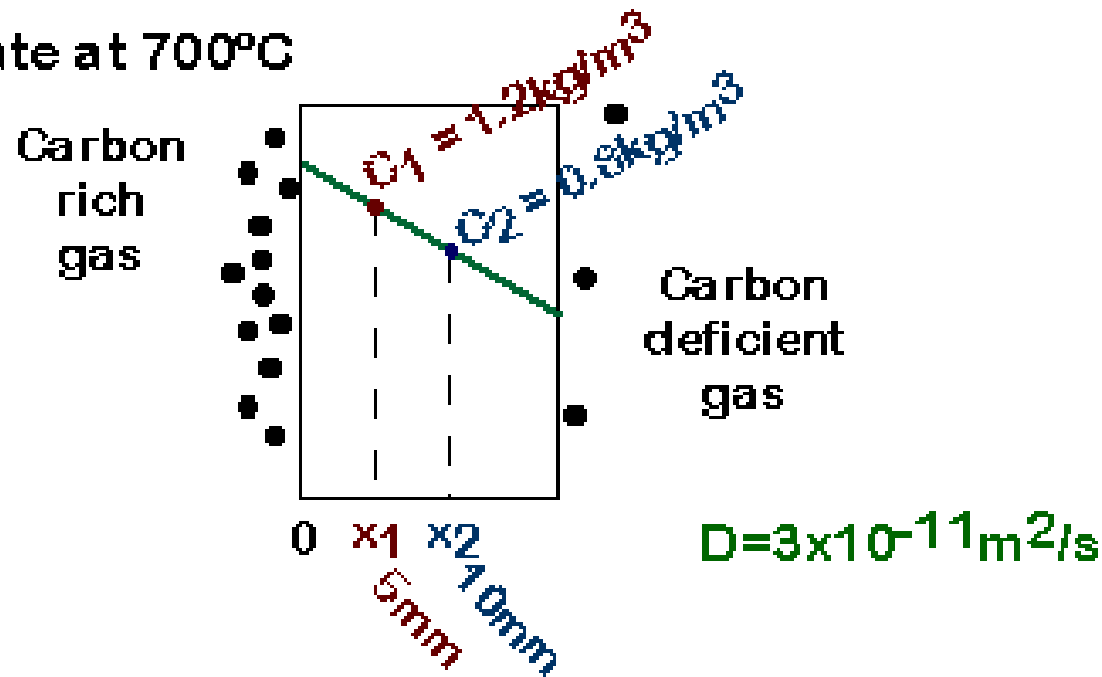
$$\frac{dC}{dx}(\text{left}) = \frac{dC}{dx}(\text{right})$$

$\Rightarrow$  the slope,  $dC/dx$ , is constant  
(does not vary with position)!

# Steady-state diffusion

The simplest type of diffusion to deal with is when the concentration at every point does not change with time.

- Steel plate at 700°C



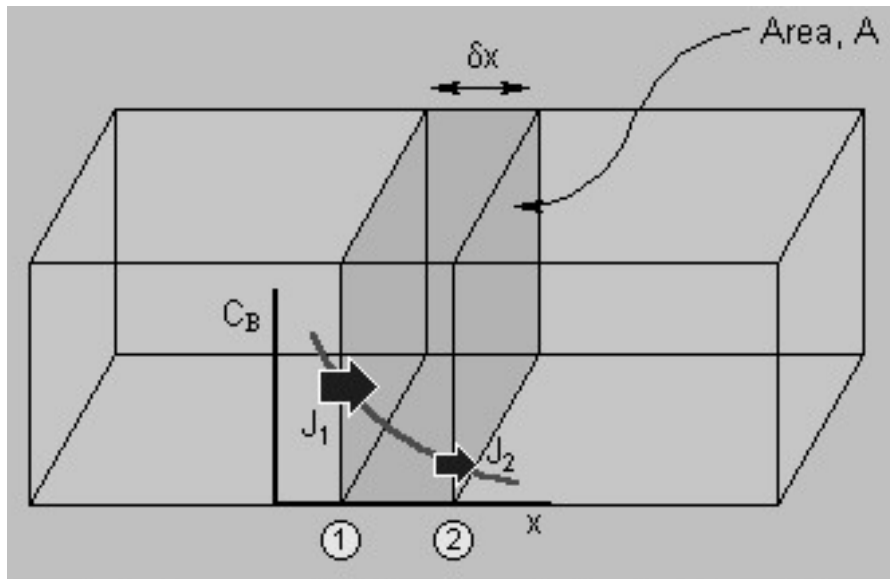
Q: How much carbon is transferring from the rich to deficient side?

$$J = -D \frac{C_2 - C_1}{x_2 - x_1} = 2.4 \times 10^{-9} \frac{\text{kg}}{\text{m}^2 \text{ s}}$$

# Nonsteady-state diffusion

In most practical situations steady-state conditions are not established, i.e. concentration varies with both distance and time, and Fick's 1st law can no longer be used.

How do we know the variation of  $C_B$  with time? → Fick's 2nd law



The number of interstitial B atoms that diffuse into the slice across plane (1) in a small time interval  $dt$  :

$$\rightarrow J_1 A dt$$

$$\text{Likewise : } J_2 A dt$$

Since  $J_2 < J_1$ , the concentration of B within the slice will have increased by

Due to mass conservation

$$(J_1 - J_2) A \delta t = \delta C_B A \delta x \quad \delta C_B = \frac{(J_1 - J_2) A \delta t}{A \delta x}$$

# Nonsteady-state diffusion

$$J_2 = J_1 + \frac{\partial J}{\partial x} \delta x \quad (\delta x \text{ is small})$$

as  $\delta t \Rightarrow 0$

$$\frac{\partial C_B}{\partial t} = - \frac{\partial J_B}{\partial x}$$

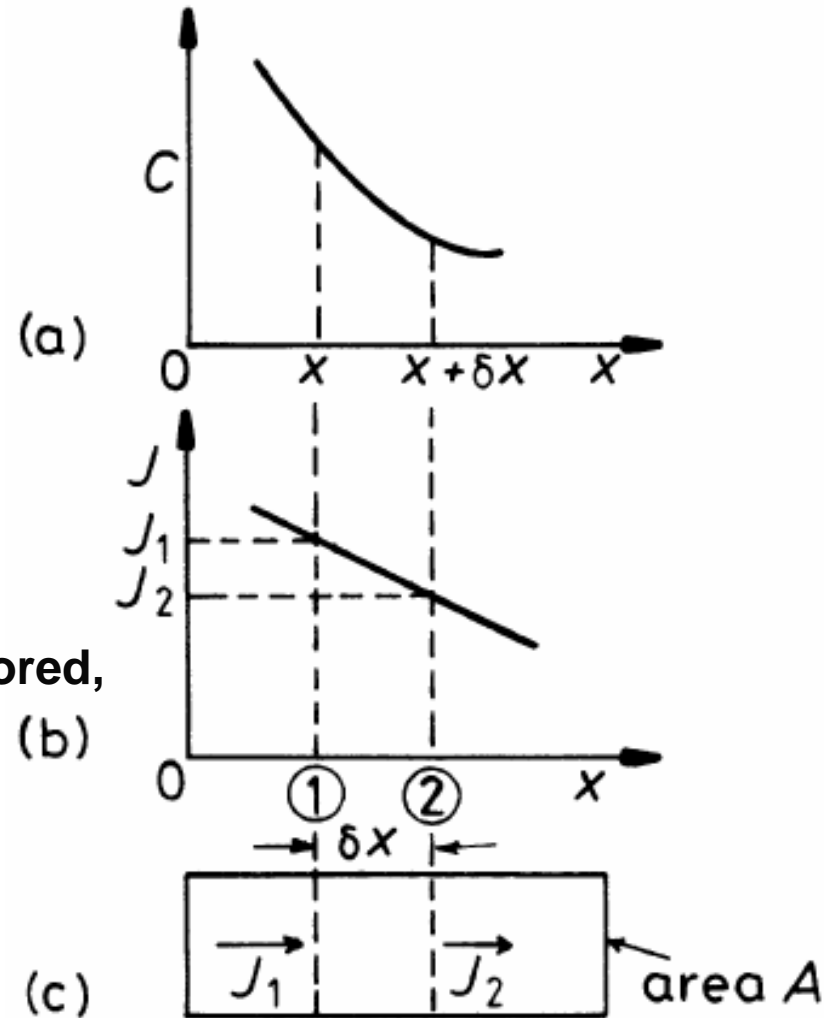
substituting Fick's 1st law gives

$$J_B = -D \frac{\partial C_B}{\partial x} \quad \frac{\partial C_B}{\partial t} = \frac{\partial}{\partial x} \left( D_B \frac{\partial C_B}{\partial x} \right)$$

If variation of D with concentration can be ignored,

$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

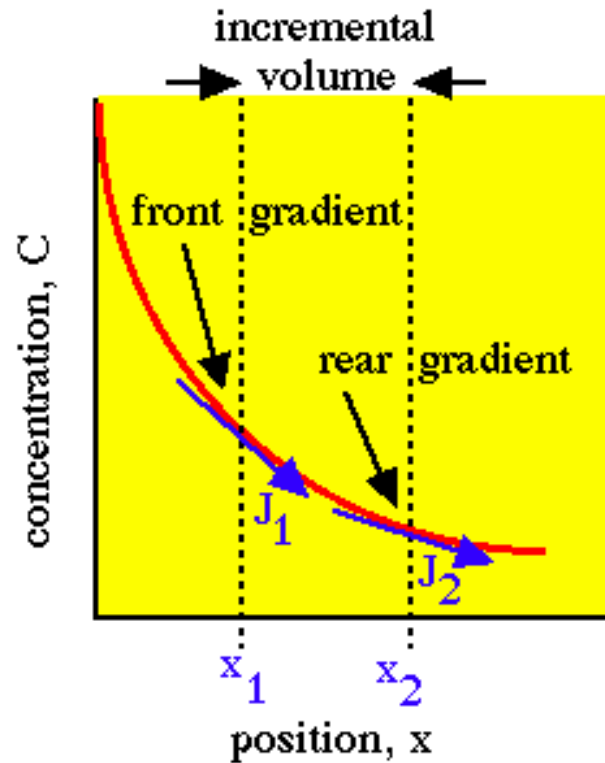
**Fick's Second Law**



# Fick's Second Law

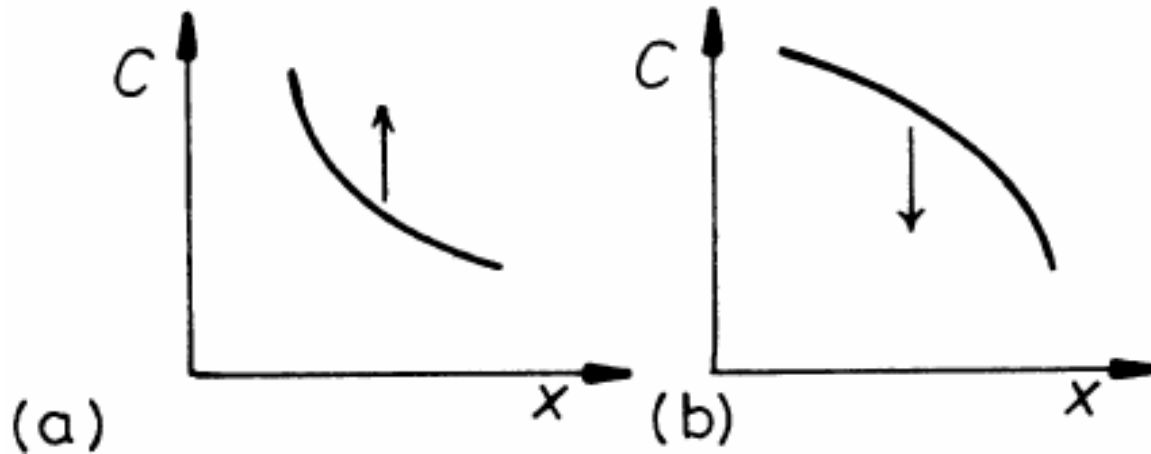
$$\frac{\partial C_B}{\partial t} = D_B \frac{\partial^2 C_B}{\partial x^2}$$

➔ Concentration varies with time and position.



$$\frac{\partial^2 C_B}{\partial x^2} = ?$$

Note that  $\frac{\partial^2 C_B}{\partial x^2}$  is the curvature.



# Diffusion

