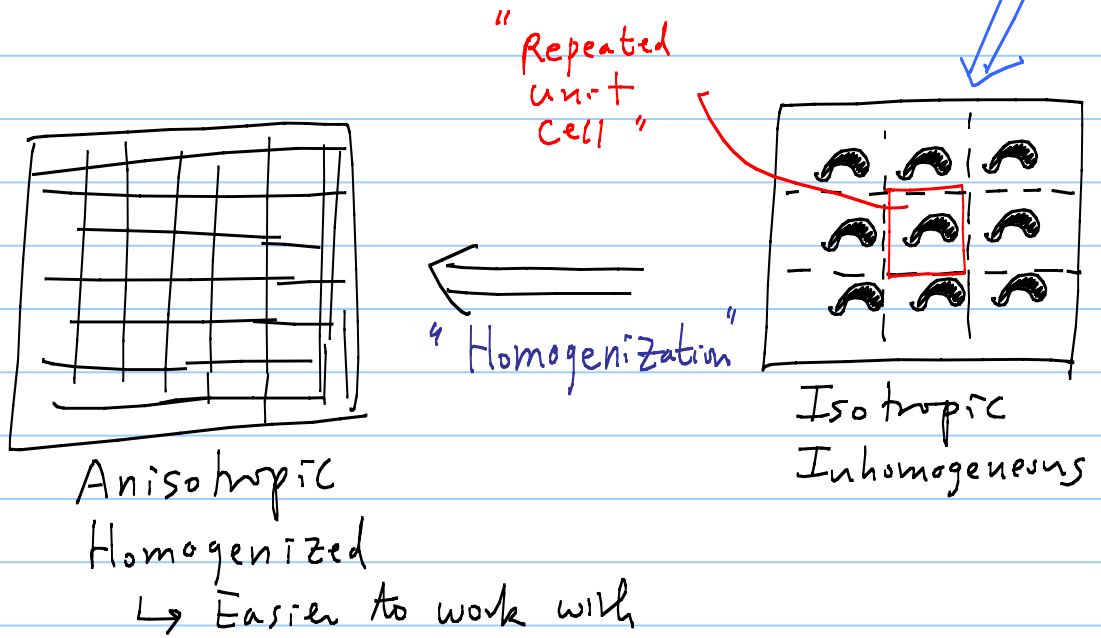
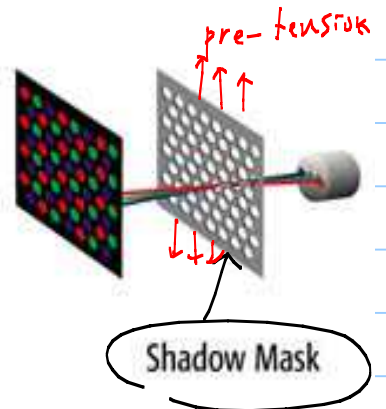
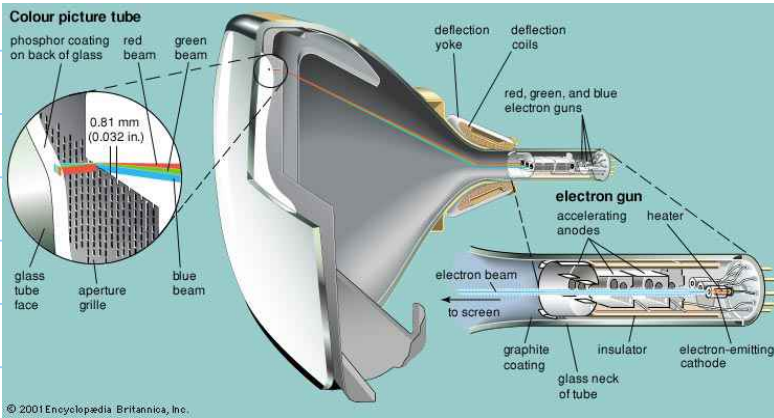


Lect 1-6

Homogenization Method

노트 제목

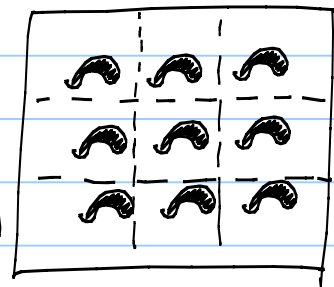
< Motivation 1: For Easier Analysis >



< Motivation 2: Material Microstructure Design >

We want
a certain
Material
property!

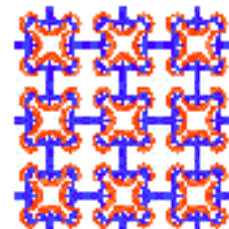
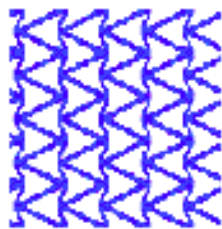
→
Inverse
Homogenization



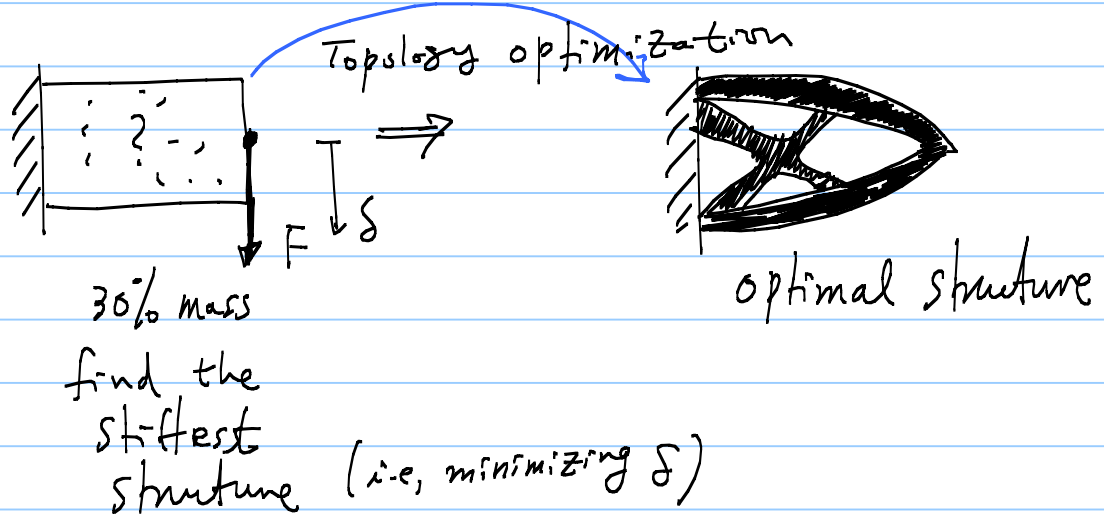
Example: Microstructure Design for

① Negative Poisson's
ratio

② zero thermal
expansion

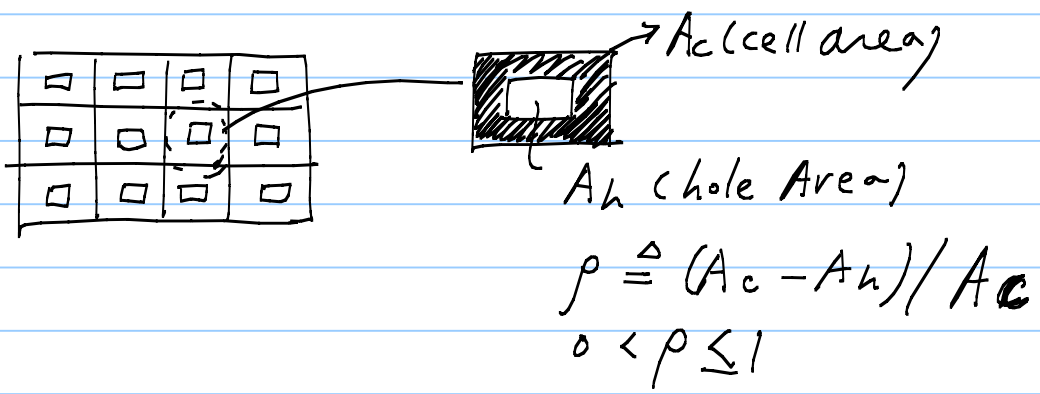


< Motivation 3: Topology optimization >




Theoretical Background

- Discretization into small cells (or elements) and consider a hole of variable size in each cell




* limiting case

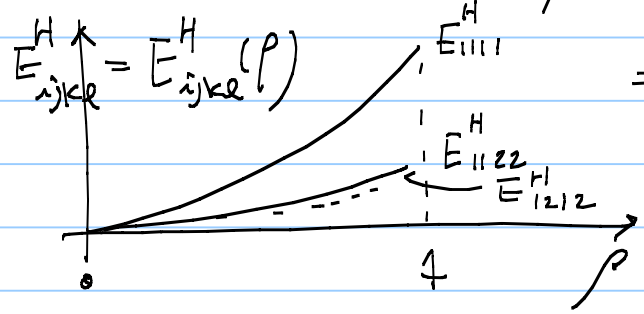
as $\rho \rightarrow 1$

 \rightarrow solid
(existence of material)

as $\rho \rightarrow 0$

 \rightarrow void
(no material in the cell)

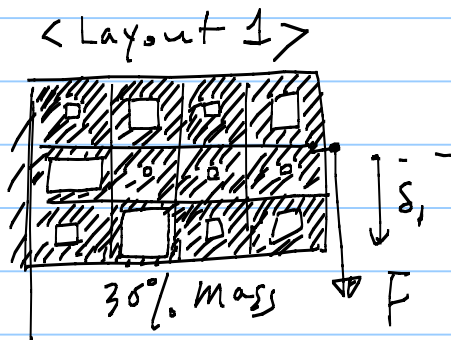
2. Homogenization of holey cell



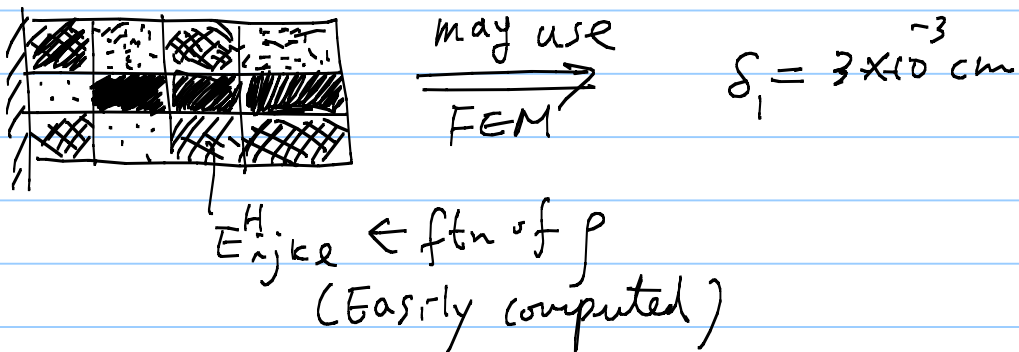
\Rightarrow Anisotropic material, in general

may approximate E_{ijkl}^H as some polynomial function of ρ .

③ Consider a candidate layout



④ Compute δ_1 using $E_{ijke}^H(p)$

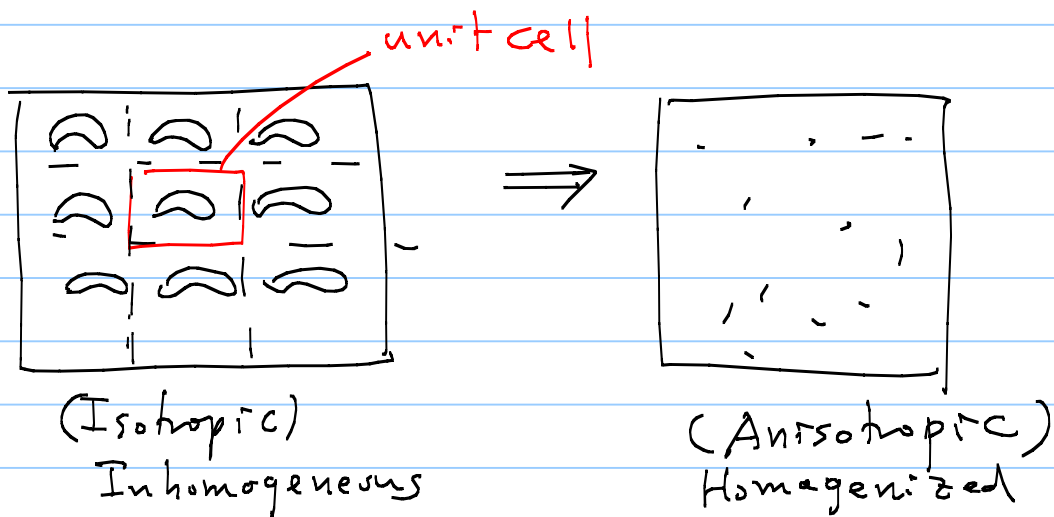


⑤ Consider another candidate layout
utilizing Layout 1
 \Rightarrow Use Optimization Algorithm
(e.g. Genetic Algorithm, steepest
descent method, etc)

⑥ Repeat ③, ④, ⑤
until the optimal solution is
found.

Back to Homogenization

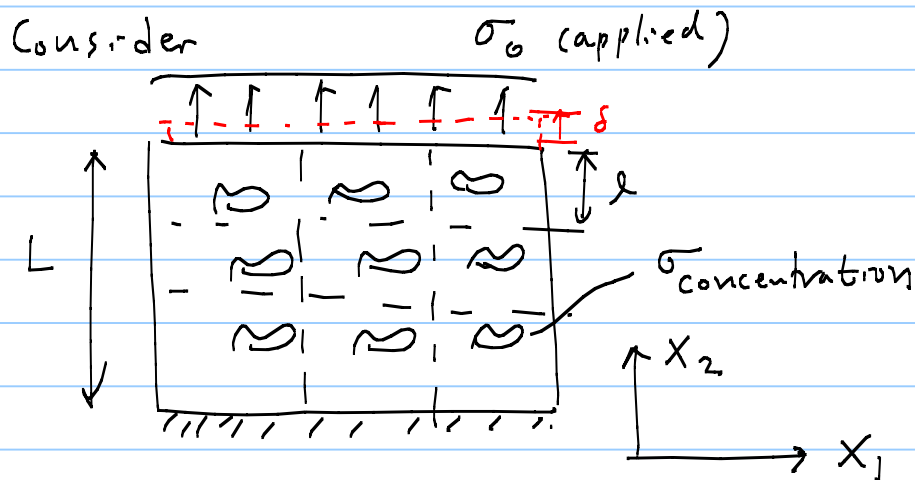
7



Condition for Homogenization

- ① same unit cell pattern must repeat or the variation of pattern must be small
- ② the size of unit cell must be much smaller than the global structure size

⇒ We need two scales to distinguish between macroscopic (global) and microscopic (local) mechanized behavior.



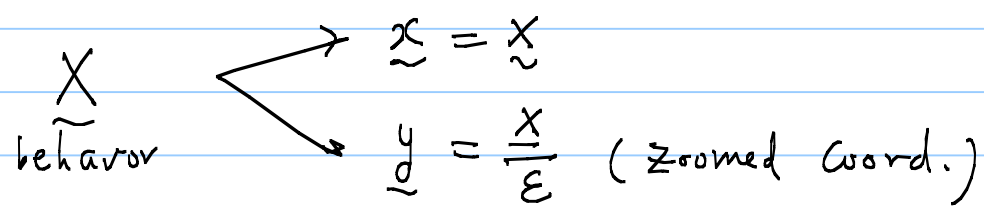
δ : governed by macroscopic behavior

$\sigma_{\text{concentration}}$: governed by microscopic behavior

\Rightarrow Structural behavior :

observed in $\begin{cases} \text{macroscopic scale } (L) \\ \text{microscopic scale } (l) \end{cases}$

Thus, two scales or two coordinate systems may be introduced for easier analysis in the limit of $\varepsilon = l/L \rightarrow 0$



< Two systems >

$\epsilon = 0$	$\epsilon \ll 1, \text{ but } \epsilon \neq 0$
original system	perturbed system by ϵ

⇒ perturbed system analyzed by "Perturbation Analysis"

- ① Regular perturbation ?
- ② Singular perturbation ?

perturbation analysis:
 approximate analysis technique
 but from original system

Basic Concept on Perturbation

A. Regular Perturbation

ex) unperturbed system:

$$x^3 - 4x = 0 \quad \text{--- (1)}$$

$$\text{Sol: } x_u = 0, x_u = 2, x_u = -2$$

perturbed system:

$$x^3 - 4.001x + 0.002 = 0 \quad \text{--- (2)}$$

(expectation $x_p \approx x_u$)

↑
Regular Behavior

* Introduce a small ϵ

$$x^3 - (4 + \epsilon)x + 2\epsilon = 0$$

(In this case: $\epsilon = 0.001$)

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Analysis:

expecting $x_p \approx x_u$

Let

$$x_p = a_0 + a_1 \epsilon + a_2 \epsilon^2 + \dots \quad (3)$$

↳ power series in ϵ

$$(a_0 + a_1 \epsilon + a_2 \epsilon^2 + \dots)^3$$

$$- (4 + \epsilon)(a_0 + a_1 \epsilon + a_2 \epsilon^2 + \dots) + 2\epsilon = 0$$

Expand:

$$(a_0^3 + 3a_0^2 \cdot a_1 \epsilon + 3a_0^2 a_2 \epsilon^2 + 3a_0 a_1^2 \epsilon^2 + \dots)$$

↳ up to ϵ^2 -term

$$- (4a_0 + 4a_1 \epsilon + 4a_2 \epsilon^2$$

$$+ a_0 \epsilon + a_1 \epsilon^2 + \dots) + 2\epsilon = 0$$

↳ up to ϵ^2 -term

$$O(\epsilon^0): \cdot a_0^3 - 4a_0 = 0$$

(a)

$$O(\epsilon^1); \quad 3a_0^2 a_1 - 4a_1 - a_0 + 2 = 0 \quad (b) \quad 12$$

$$O(\epsilon^2); \quad 3a_0^2 a_2 + 3a_0 a_1^2 - 4a_2 - a_1 = 0 \quad (c)$$

$$O(\epsilon^3); \quad \dots = 0$$



$$\text{Eq. (a)} = \text{Eq. (1)}$$

$$\text{i.e. ; } \underline{a_0 = -2, 0, 2}$$

If $a_0 = -2$ is chosen,

$$(b) \rightarrow 12a_1 - 4a_1 - (-2) + 2 = 0$$

$$\underline{a_1 = -\frac{1}{2}}$$

$$(c) \rightarrow 12a_2 + 3(-2)\left(-\frac{1}{2}\right)^2 - 4a_2 - \left(-\frac{1}{2}\right) = 0 \quad /3$$

$$8a_2 - 1 = 0$$

$$\therefore a_2 = \frac{1}{8}$$

Thus

$$x_p = a_0 + a_1 \epsilon + a_2 \epsilon^2 + \dots$$

$$= -2 - \frac{1}{2} \epsilon + \frac{1}{8} \epsilon^2 + \dots \quad (4)$$

(near $x = -2$)

B. Singular Perturbation

Consider Duffing equation

$$m \frac{d^2 y}{dt^2} + k(y - \alpha y^3) = 0 \quad (*)$$

if α small,

weakly nonlinear

Assuming

$\omega_0 = \sqrt{k/m} = 1$, rewrite (*) as

$$\frac{d^2 y}{dt^2} + (y - \epsilon y^3) = 0 \quad (5)$$

with $y(0) = 1, \quad \frac{dy}{dt}(0) = 0 \quad (6a, b)$

physics:

if $\epsilon = 0, \omega_0 = 1$

if $1 \gg \epsilon > 0, \omega^* < \omega_0$

As in the earlier example,

Let us attempt a regular expansion:

$$y(t) = y_1(t) + \epsilon y_2(t) + \epsilon^2 y_3(t) + \dots \quad (7)$$

Eq. (7) \rightarrow Eq. (5), Eq. (6)

$$\begin{aligned} \bullet \quad \frac{d^2}{dt^2} (y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots) + (y_1 + \epsilon y_2 + \dots) \\ - \epsilon (y_1 + \epsilon y_2 + \dots)^3 = 0 \end{aligned}$$

$$\bullet \quad y_1(0) + \epsilon y_2(0) + \dots = 1$$

$$\bullet \quad \frac{dy_1(0)}{dt} + \epsilon \frac{dy_2(0)}{dt} + \dots = 0$$

(**) \rightarrow (*) yields

$$y_2(t) = \frac{1}{32} \{ (\cos t - \cos 3t) + 12t \sin t \}$$

Thus

$$\begin{aligned} y(t) &= y_1(t) + \epsilon y_2(t) + \dots \\ &= \cos t + \frac{\epsilon}{32} \{ (\cos t - \cos 3t) \\ &\quad + 12t \sin t \} + \dots \end{aligned}$$

"secular term"

\rightarrow It blows up as $t \rightarrow \infty$
especially when $t > O\left(\frac{1}{\epsilon}\right)$

* If we ignore ϵ -term;

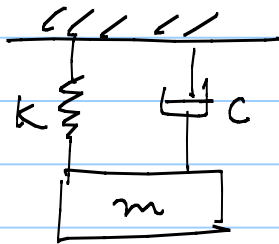
$$y(t) = \cos t \rightarrow \text{"trivial"}$$

(no correction)

\Rightarrow Regular Perturbation method does not work
 \leftrightarrow Try Singular Perturbation

What's going on?

Examine a simpler problem



$$m \frac{d^2 y}{dt^2} + c \frac{dy}{dt} + k y = 0$$

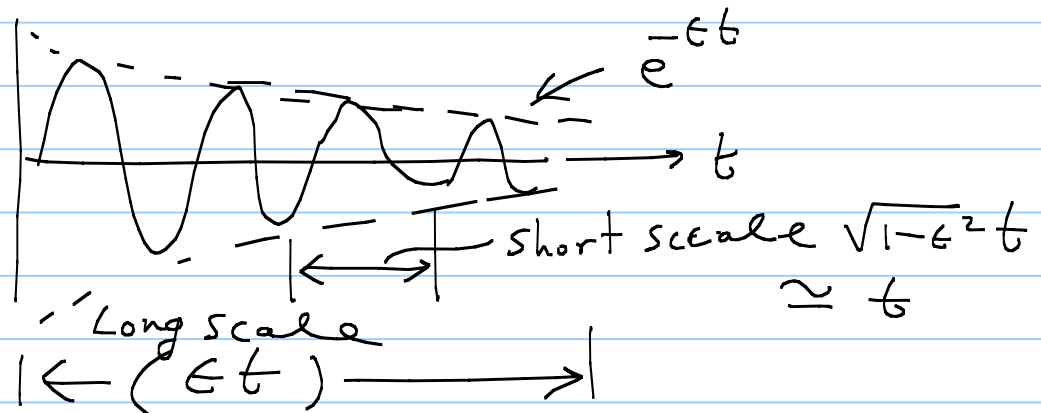
$$\text{or } \frac{d^2 y}{dt^2} + 2\epsilon \frac{dy}{dt} + y = 0$$

$$\left(\text{with } y(0) = 0, \frac{dy(0)}{dt} \neq 0 \right)$$

Solution:

$$y = e^{-\epsilon t} \left\{ A \cos \sqrt{1-\epsilon^2} t + B \sin \sqrt{1-\epsilon^2} t \right\}$$

$$= B e^{-\epsilon t} \sin \sqrt{1-\epsilon^2} t$$



* There are two time scales:

t and $\frac{\epsilon t}{\tau}$ slowly varying

Thus, for perturbation analysis of this kind, we need to make a correct approximation

$$t = t(\tau, T) \quad (8)$$

where

$$\begin{cases} \tau = t \\ T = \epsilon t \end{cases} \text{ "contracted" time scale} \quad (9a, b)$$

• Let

$$y(t) = y_1(\tau, T) + \epsilon y_2(\tau, T) + \dots \quad (10)$$

* Differentiation

$$\frac{d}{dt} = \frac{\partial \tau}{\partial t} \frac{\partial}{\partial \tau} + \frac{\partial T}{\partial t} \frac{\partial}{\partial T} = \frac{\partial}{\partial \tau} + \epsilon \frac{\partial}{\partial T} \quad (11)$$

$$\frac{d^2}{dt^2} = \frac{\partial^2}{\partial \tau^2} + 2\epsilon \frac{\partial^2}{\partial \tau \partial T} + \epsilon^2 \frac{\partial^2}{\partial T^2} \quad (12)$$

(*) Eqs (10, 11, 12) \rightarrow Eq (5) (Duffing Eq.)

$$\begin{aligned} & \left(\frac{\partial^2}{\partial \tau^2} + 2\epsilon \frac{\partial^2}{\partial \tau \partial T} + \epsilon^2 \frac{\partial^2}{\partial T^2} \right) (y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots) \\ & + (y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots) - \epsilon (y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots)^3 \\ & = 0 \end{aligned}$$

$$O(\epsilon^0) : \frac{\partial^2 y_1}{\partial \tau^2} + y_1 = 0 \quad (13)$$

$$O(\epsilon^1) : \frac{\partial^2 y_2}{\partial \tau^2} + y_2 = y_1^3 - 2 \frac{\partial^2 y_1}{\partial \tau \partial T} \quad (14)$$

(*) Eqs (10, 11, 12) \rightarrow Eq (6) (BC's)

$$y_1(0,0) + \epsilon y_2(0,0) + \dots = 1$$

$$\left. \left(\frac{\partial}{\partial \tau} + \epsilon \frac{\partial}{\partial T} \right) (y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots) \right|_{\tau=T=0} = 0$$

$$O(\epsilon^0): \quad y_1 = 1; \quad \frac{\partial y_1}{\partial \tau} = 0 \quad \text{at } \tau=T=0 \quad (15)$$

$$O(\epsilon^1): \quad y_2 = 0; \quad \frac{\partial y_2}{\partial \tau} = - \frac{\partial y_1}{\partial T} \quad \text{at } \tau=T=0 \quad (16)$$

i) To find $y_1(\tau, T)$,

solve

$$\frac{\partial^2 y_1}{\partial \tau^2} + y_1 = 0$$

$$y_1 = 1, \quad \frac{\partial y_1}{\partial \tau} = 0 \quad \text{at } \tau = 0, T = 0$$

$$\rightarrow \underline{y_1(\tau, T) = A(T) \cos \tau + B(T) \sin \tau} \quad (17)$$

Using BC's

$$y_1(0, 0) = A(0) = 1$$

$$\frac{\partial y_1}{\partial \tau}(0, 0) = B(0) = 0$$

(18a, b)

~~★~~ → However, the function forms of $A(T)$

and $B(T)$ are not known.

(→ characteristics of singular perturbation)

ii) Let's consider for y_2

$$\frac{\partial^2 y_2}{\partial \tau^2} + y_2 = y_1^3 - 2 \frac{\partial^2 y_1}{\partial \tau \partial T}$$

$$= (A \cos \tau + B \sin \tau)^3 - 2(-A' \sin \tau + B' \cos \tau)$$

$$= A^3 \cos^3 \tau + 3A^2 B \cos^2 \tau \sin \tau + 3AB^2 \cos \tau \sin^2 \tau + B^3 \sin^3 \tau + 2A' \sin \tau - 2B' \cos \tau$$

$$= A^3 \cos^3 \tau + 3A^2 B (1 - \sin^2 \tau) \sin \tau + 3AB^2 \cos \tau (1 - \cos^2 \tau) + B^3 \sin^3 \tau + 2A' \sin \tau - 2B' \cos \tau$$

$$= (A^3 - 3AB^2) \cos^3 \tau + (B^3 - 3A^2 B) \sin^3 \tau + 3A^2 B \sin \tau + 3AB^2 \cos \tau + 2A' \sin \tau - 2B' \cos \tau$$

$$\left(\begin{array}{l} \text{Now use: } \cos^3 \alpha = \frac{1}{4} (3 \cos \alpha + \cos 3 \alpha) \\ \sin^3 \alpha = \frac{1}{4} (3 \sin \alpha - \sin 3 \alpha) \end{array} \right)$$

$$\begin{aligned} &= \left\{ \frac{3}{4} (A^3 - 3AB^2) + 3AB^2 - 2B' \right\} \cos \tau \\ &\quad + \frac{3}{4} (B^3 - 3A^2B) + 3A^2B + 2A' \left\{ \sin \tau \right. \\ &\quad \left. + (\dots) \cos 3 \tau + (\dots) \sin 3 \tau \right\} \end{aligned}$$

\therefore

$$\begin{aligned} \frac{\partial^2 y_2}{\partial \tau^2} + y_2 &= \left(\frac{3}{4} A^3 + \frac{3}{4} AB^2 - 2B' \right) \cos \tau \\ &+ \left(\frac{3}{4} B^3 + 3A^2B + 2A' \right) \sin \tau \\ &\quad + (\dots) \cos 3 \tau + (\dots) \sin 3 \tau \end{aligned} \tag{19}$$

Because $y_2^c \in \{ \cos \tau, \sin \tau \}$

$$\left\{ \begin{array}{l} \frac{3}{4} A^3 + \frac{3}{4} A B^2 - 2\beta' \equiv 0 \quad (20) \\ \frac{3}{4} B^3 + \frac{3}{4} A^2 B + 2A' \equiv 0 \quad (21) \end{array} \right.$$

To avoid the appearance of secular term ($\tau \cos \tau$, $\tau \sin \tau$) in y_2

To solve (20, 21)

$$\frac{dA}{dT} = -\frac{3}{8} B (A^2 + B^2) \quad (1)$$

$$\frac{dB}{dT} = \frac{3}{8} A (A^2 + B^2) \quad (2)$$

$$\frac{(1)}{(2)} \rightarrow \frac{dA}{dB} = -\frac{B}{A} \rightarrow A dA + B dB = 0$$

$$\therefore A^2 + B^2 = \text{const}$$

Using Initial condition (18)

for $A(0)$ and $B(0)$; $A(0)=1$, $B(0)=1$,

$$A^2 + B^2 = 1 \quad (3)$$

$$\begin{array}{l} \textcircled{3} \rightarrow \textcircled{1} \quad \frac{dA}{dT} = -\frac{3}{8} B \quad \textcircled{1}' \\ \quad \quad \quad \textcircled{2} \quad \frac{dB}{dT} = \frac{3}{8} A \quad -\textcircled{2}' \end{array}$$

$$\frac{d}{dT} \textcircled{1}' - \left(\frac{3}{8}\right) \textcircled{2} :$$

$$\frac{d^2 A}{dT^2} = -\left(\frac{3}{8}\right)^2 A \quad -\textcircled{1}''$$

Initial condition

$$A(0) = 1 \quad (18a)$$

$$\frac{dA(0)}{dT} = 0$$

Thus ;

$$\underline{A(\tau) = \cos \frac{3\tau}{8}} \quad (22)$$

$$\frac{d}{d\tau} \textcircled{2}' + \left(\frac{3}{8}\right)^2 0;$$

$$\frac{d^2 B}{d\tau^2} = - \left(\frac{3}{8}\right)^2 A \quad \textcircled{2}''$$

Initial condition

$$B(0) \underset{(18b)}{=} 0, \quad \frac{dB(0)}{d\tau} = \frac{3}{8}$$

Thus

$$\underline{B(\tau) = \sin \frac{3\tau}{8}} \quad (23)$$

Finally, using (17) (22) (23),

$$\begin{aligned} y(t) &= y(\tau, \tau) \\ &= \cos \frac{3\tau}{8} \cos \tau + \sin \frac{3\tau}{8} \sin \tau + 0(t) \end{aligned}$$

$$= \cos\left(\frac{3T}{8} - \tau\right) + O(\epsilon)$$

$$= \cos\left(1 - \frac{3\epsilon}{8}\right)t + O(\epsilon^2 t \cos t)$$

(not good if $t \geq O\left(\frac{1}{\epsilon^2}\right)$)

Observation:

1) Regular perturbation fails
when $t \geq O\left(\frac{1}{\epsilon}\right)$

ii) "this" singular perturbation fails
when $t \geq O\left(\frac{1}{\epsilon^2}\right)$

For more accurate solutions,

$$T_1 = t, \quad T_2 = \epsilon t, \quad T_3 = \epsilon^2 t, \dots$$

$$y(t) = y(T_1, T_2, T_3, \dots)$$

$$= y_1 + \epsilon y_2 + \epsilon^2 y_3 + \dots$$

$$\Rightarrow y = \cos \left(1 - \frac{3\epsilon}{8} - \frac{15\epsilon^2}{256} + \dots \right) t$$

$$\left(\begin{array}{l} \text{Mayfeh suggested} \\ \tau = \epsilon t \\ \tau = (1 + \epsilon^2 + \epsilon^3 + \dots) t \end{array} \right)$$