

Topics in Fusion and Plasma Studies 2

(459.667, 3 Credits)

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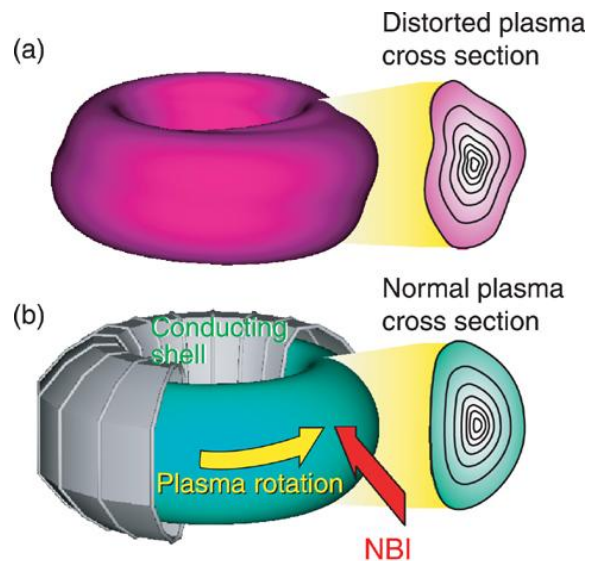
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General Properties of Ideal MHD

• Introduction

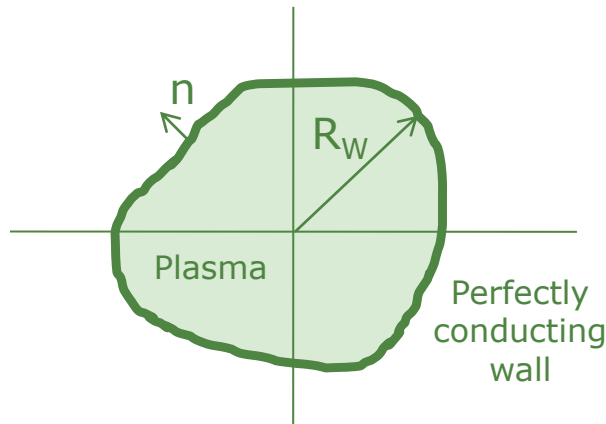
- Short description of the three most common classes of boundary conditions
- Conservation of mass, momentum, and energy, both locally and globally despite the significant number of approximations made in ideal MHD
- Consequence of the perfect conductivity assumption



General Properties of Ideal MHD

• Boundary Conditions

- Perfectly conducting wall:
tangential electric field and normal magnetic field vanish
on the conducting wall



$$\vec{n} \times \vec{E} \Big|_{R_w} = 0$$

$$\vec{n} \cdot \vec{B} \Big|_{R_w} = 0$$

$$\vec{n} \cdot \vec{v} \Big|_{R_w} = 0 \quad \longleftarrow \quad \vec{n} \times \vec{E} + (\vec{n} \cdot \vec{B})\vec{v} - (\vec{n} \cdot \vec{v})\vec{B} = 0$$

ideal Ohm's law

General Properties of Ideal MHD

• Boundary Conditions

- Insulating vacuum region:
assume that the equations can be solved in each region

$$\nabla \times \hat{\mathbf{B}} = 0$$

vacuum

$$\nabla \cdot \hat{\mathbf{B}} = 0$$

$$\hat{\mathbf{n}} \cdot \vec{\mathbf{B}}|_{R_w} = 0 \quad \text{wall}$$

$$[[\mathbf{n} \cdot \vec{\mathbf{B}}]]_R = 0 \quad [[Q]] = \hat{Q} - Q$$

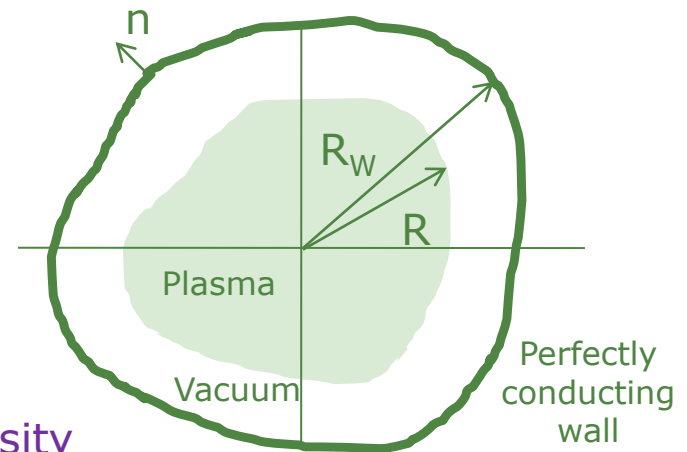
$$[[\mathbf{n} \times \vec{\mathbf{B}}]]_R = \mu_0 \vec{\mathbf{K}} \quad \text{Surface current density}$$

$$[[p + B^2 / 2\mu_0]]_R = 0$$

$$\vec{\mathbf{n}} \cdot \vec{\mathbf{v}} = \hat{\mathbf{n}} \cdot \vec{\mathbf{v}}$$

$$\vec{\mathbf{n}} \cdot \vec{\mathbf{B}}$$

$$B^z|_R = \hat{B}^z|_R$$



Free boundary problem

.VS.

Fixed boundary problem

and the magnetic pressure are continuous

B

General Properties of Ideal MHD

• Boundary Conditions

- Plasma surrounded by external coils:
difficult but realistic situation where the plasma is confined by the magnetic fields created by a fixed set of external conductors

$$[[\mathbf{n} \cdot \vec{B}]]_R = 0$$

$$[[\mathbf{n} \times \vec{B}]]_R = \mu_0 \vec{K}$$

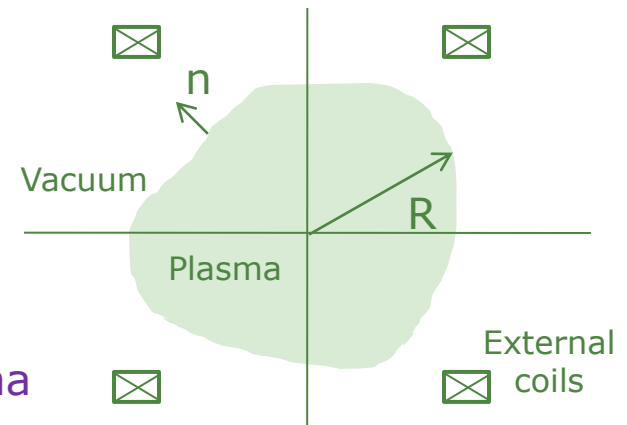
$$[[p + B^2 / 2\mu_0]]_R = 0$$

$$\hat{B} = \vec{B}_a + \tilde{B} \quad \text{applied field by coils and induced field due to the plasma}$$

$$\vec{B}_a = \nabla \times \vec{A}$$

$$\vec{A} = \frac{\mu_0}{4\pi} \sum_i \int \frac{\vec{J}_i}{|\vec{r} - \vec{r}'|} d\vec{r}' \quad \text{Biot-Savart law}$$

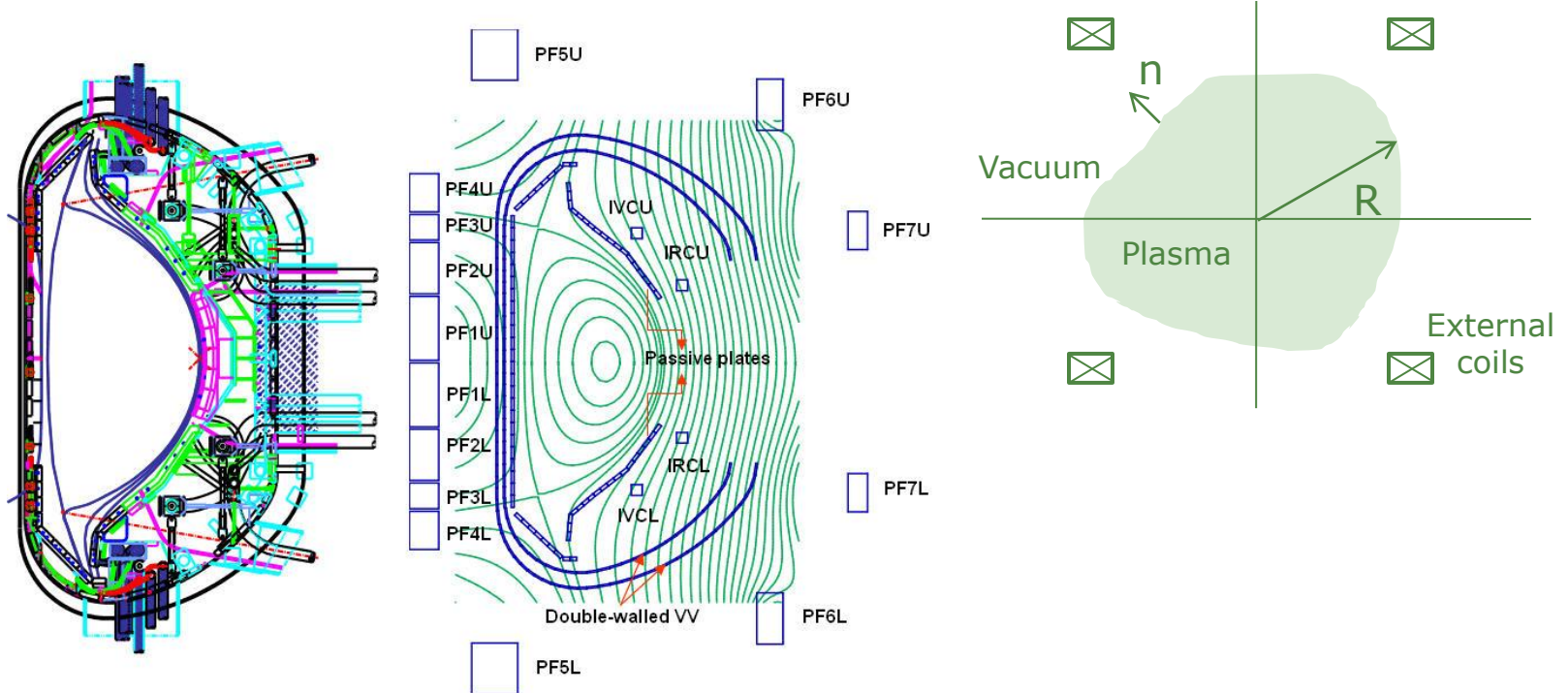
$$\nabla \times \tilde{B} = \nabla \cdot \tilde{B} = 0 \quad \text{induced field in the vacuum region}$$



General Properties of Ideal MHD

- **Boundary Conditions**

- Plasma surrounded by external coils:
difficult but realistic situation where the plasma is confined by the magnetic fields created by a fixed set of external conductors



General Properties of Ideal MHD

• Local Conservation Relations

- Since a considerable number of assumptions were made in the derivation of the MHD equations it is important to investigate whether the resulting model still satisfies the basic conservation laws.

$$\frac{\partial}{\partial t}(\quad) + \nabla \cdot (\quad) = 0$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \vec{v}) = 0 \quad \vec{T} = \rho \vec{v} \vec{v} + \left(p + \frac{B^2}{2\mu_0} \right) \vec{I} - \frac{\vec{B}\vec{B}}{\mu_0} \quad \text{stress tensor}$$

$$\frac{\partial \rho \vec{v}}{\partial t} + \nabla \cdot \vec{T} = 0 \quad w = \frac{1}{2} \rho v^2 + \frac{B^2}{2\mu_0} + \frac{p}{\gamma - 1} \quad \text{energy}$$

$$\frac{\partial w}{\partial t} + \nabla \cdot \vec{s} = 0 \quad \vec{s} = \left(\frac{1}{2} \rho v^2 + \frac{\gamma}{\gamma - 1} p \right) \vec{v} + \frac{1}{\mu_0} \vec{E} \times \vec{B} \quad \text{energy flux}$$

General Properties of Ideal MHD

• Local Conservation Relations

$$\vec{T} = \rho \vec{v} \vec{v} + \left(p + \frac{B^2}{2\mu_0} \right) \vec{I} - \frac{\vec{B} \vec{B}}{\mu_0}$$

Reynolds stress

$$\vec{T}_R = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & \rho v^2 \end{pmatrix}$$

Important with large fluid flow, but usually not in studies of plasma stability where the flows are either small or zero

$$\vec{T}_B = \begin{pmatrix} p_{\perp} & 0 & 0 \\ 0 & p_{\perp} & 0 \\ 0 & 0 & p_{\parallel} \end{pmatrix}$$

$$p_{\perp} = p + \frac{B^2}{2\mu_0}, \quad p_{\parallel} = p - \frac{B^2}{2\mu_0}$$

isotropic plasma pressure

Total pressure parallel to the field: negative magnetic pressure correspond to a tension along the field lines

General Properties of Ideal MHD

• Local Conservation Relations

$$w = \frac{1}{2} \rho v^2 + \frac{B^2}{2\mu_0} + \frac{p}{\gamma - 1} = w_K + w_P$$

$$w_K = \frac{1}{2} \rho v^2$$

plasma kinetic energy

$$w_P = \frac{B^2}{2\mu_0} + \frac{p}{\gamma - 1}$$

Potential energy

Magnetic energy +
internal energy of the plasma

$$\vec{s} = \left(\frac{1}{2} \rho v^2 + \frac{p}{\gamma - 1} \right) \vec{v} + p \vec{v} + \frac{1}{\mu_0} \vec{E} \times \vec{B}$$

Flow of plasma
kinetic + internal energy

Poynting flux:
flow of electromagnetic
energy

Mechanical work
done on or by the
plasma as it moves

General Properties of Ideal MHD

- **Global Conservation Laws**

Obtained by integrating the local conservation laws over the volumes appropriate to each of the three sets of boundary conditions.

- Perfectly conducting wall

$$\frac{dM}{dt} = 0$$

$$\frac{d\vec{P}}{dt} = -\int \left(p + \frac{B^2}{2\mu_0} \right) \vec{n} dS$$

$$\frac{dW}{dt} = 0$$



Total force exerted by the walls on the plasma. If the system remains in place, zero.

$$M = \int \rho d\vec{r}$$

total mass of plasma

$$\vec{P} = \int \rho \vec{v} d\vec{r}$$

mechanical momentum of plasma

$$W = \int \left(\frac{1}{2} \rho v^2 + \frac{p}{\gamma - 1} + \frac{B^2}{2\mu_0} \right) d\vec{r}$$

total kinetic, potential energy of plasma and magnetic field: energy can be transferred each other

General Properties of Ideal MHD

• Global Conservation Laws

- Insulating vacuum region
 - More complicated since the plasma is allowed to move.
 - The combined plasma-vacuum energy is conserved.

$$Z(t) = \int z(\vec{r}, t) d\vec{r} \quad \text{global quantity}$$

$$\frac{dZ(t)}{dt} = \int_V \frac{\partial z}{\partial t} d\vec{r} + \int_S z \vec{n} \cdot \vec{u} dS \quad \text{Total time derivative in a volume whose boundary is moving with } u$$

plasma energy \rightarrow

$$\frac{dW}{dt} = - \int \left(p + \frac{B^2}{2\mu_0} \right) \vec{n} \cdot \vec{v} dS$$

$$W = \int \left(\frac{1}{2} \rho v^2 + \frac{p}{\gamma - 1} + \frac{B^2}{2\mu_0} \right) d\vec{r} \quad w = \frac{1}{2} \rho v^2 + \frac{B^2}{2\mu_0} + \frac{p}{\gamma - 1}$$

General Properties of Ideal MHD

• Global Conservation Laws

$$\frac{dZ(t)}{dt} = \int_V \frac{\partial z}{\partial t} d\vec{r} + \int_S z \vec{n} \cdot \vec{u} dS$$

vacuum energy \rightarrow

$$\hat{W} = \int \frac{\hat{B}^2}{2\mu_0} d\vec{r}$$

$$\frac{d\hat{W}}{dt} = \int \frac{1}{\mu_0} \left(\hat{B} \cdot \frac{\partial \hat{B}}{\partial t} \right) d\vec{r} - \int \frac{\hat{B}^2}{2\mu_0} \vec{n} \cdot \vec{v} dS$$

$$\frac{d\hat{W}}{dt} = \int \frac{\hat{B}^2}{2\mu_0} \vec{n} \cdot \vec{v} dS \quad \leftarrow$$

Faraday's law,
Divergence theorem

$$\hat{E} = -\vec{v} \times \hat{B}, \quad \vec{n} \cdot \hat{B}|_S = 0$$

General Properties of Ideal MHD

• Global Conservation Laws

$$\frac{dW}{dt} = -\int \left(p + \frac{B^2}{2\mu_0} \right) \vec{n} \cdot \vec{v} dS \quad \frac{d\hat{W}}{dt} = \int \frac{\hat{B}^2}{2\mu_0} \vec{n} \cdot \vec{v} dS$$

$$\frac{d}{dt}(W + \hat{W}) = 0 \quad \longleftarrow \quad [[p + B^2 / 2\mu_0]]_R = 0$$

- When an ideal MHD plasma is isolated from a conducting wall by a vacuum region, the combined energy of the plasma-vacuum system is conserved.
- The fact that only the total is conserved indicates that, in general, energy will flow from one region to the other as the plasma moves.

General Properties of Ideal MHD

- **Global Conservation Laws**

- Plasma surrounded by external coils:
 - The energy of the system is no longer conserved.
 - With external sources present, energy can be supplied to or extracted from the system.

$$\frac{d}{dt}(W + \hat{W}) = -\int \frac{1}{\mu_0} (\vec{E} \times \vec{B}) \cdot \vec{n}_w dS_w$$

The rate of increase of the total energy in the combined plasma-vacuum system is equal to the electromagnetic power flowing into the region.

General Properties of Ideal MHD

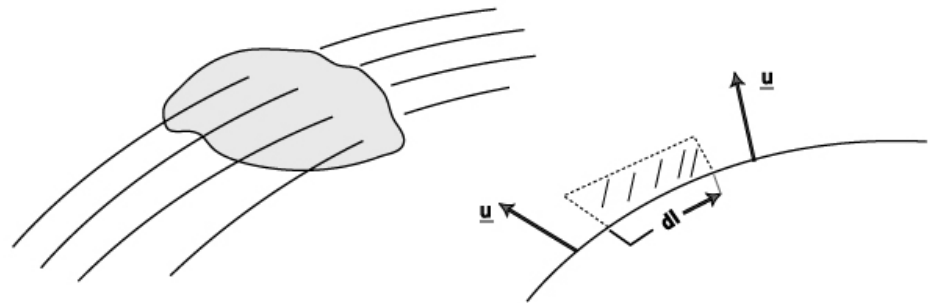
- **Conservation of Flux: “Frozen” Field Line Picture**

- A consequence of the perfect conductivity Ohm’s law, is that the magnetic flux passing through any arbitrary open surface area moving with the plasma is constant.
- Flux is conserved locally as well as globally.

$$\frac{d\psi}{dt} = \int \frac{\partial \hat{B}}{\partial t} \cdot \vec{n} dS - \oint \vec{u} \times \vec{b} \cdot d\vec{l}$$

time rate of change of the flux passing through any moving surface, S

$$\psi = \int \vec{B} \cdot \vec{n} dS$$



$$\frac{d\psi}{dt} = - \oint (\vec{E} + \vec{u} \times \vec{B}) \cdot d\vec{l} \quad \longleftarrow \text{Faraday's law, Stokes theorem}$$

$$\frac{d\psi}{dt} = 0 \quad \longleftarrow u=v \text{ (plasma velocity), ideal Ohm's law}$$

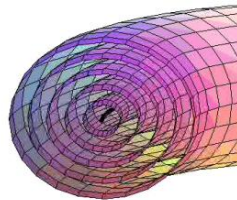
General Properties of Ideal MHD

- Conservation of Flux: “Frozen” Field Line Picture

$$\frac{d\psi}{dt} = 0$$

- The total flux in an ideal MHD plasma is conserved.
- Magnetic lines move with the plasma; they are “frozen” into the fluid.
- Any allowable physical motion of the plasma requires that neighboring fluid elements remain adjacent to one another; fluid elements are not allowed to tear or break into separate pieces.
- Since the magnetic lines move with the plasma, the field line topology must thus be preserved during any physically allowable MHD motion.

Ideal MHD:
 $\eta = 0$



Resistive MHD:
 $\eta \neq 0$

