

" we make a "

Before/Start

Date 9. March 2009 No. \*

\* homework check

- invariant
- $(n_x, n_y)$  calculation.
- more interaction.
- in English

$$T_i = T_{ji} \cdot n_j$$

$$\begin{pmatrix} 4 & 1 \\ 1 & 4 \end{pmatrix} \cdot \begin{pmatrix} 5 & 0 \\ 0 & 3 \end{pmatrix}$$

\*  $\begin{pmatrix} 4 & 1 \\ 1 & 4 \end{pmatrix} \begin{pmatrix} 1 \\ 0 \end{pmatrix} = \begin{pmatrix} 4 \\ 1 \end{pmatrix}$

$\begin{pmatrix} 4 & 1 \\ 1 & 4 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ \frac{1}{\sqrt{2}} \end{pmatrix} = \begin{pmatrix} \frac{5}{\sqrt{2}} \\ \frac{5}{\sqrt{2}} \end{pmatrix} \quad \begin{pmatrix} \frac{5}{\sqrt{2}} \\ \frac{5}{\sqrt{2}} \end{pmatrix}$

magnitude  $\rightarrow 5$ .

$$\begin{pmatrix} 5 & 0 \\ 0 & 3 \end{pmatrix} \begin{pmatrix} \frac{1}{\sqrt{2}} \\ -\frac{1}{\sqrt{2}} \end{pmatrix} = \begin{pmatrix} \frac{5}{\sqrt{2}} \\ -\frac{3}{\sqrt{2}} \end{pmatrix} \quad \begin{pmatrix} \frac{5}{\sqrt{2}} \\ -\frac{3}{\sqrt{2}} \end{pmatrix}$$

$$\sqrt{\frac{25}{2} + \frac{9}{2}} = \underline{\underline{\sqrt{17}}}$$

- \* Stress more important than strain
- \* importance of generating knowledge.
- \* 1st order 2nd order tensor ...
- \* Graphical representation of stress

Kittel) - gradient

- body moment

- Cosserat model

- Linear algebra (symmetric part)

- rotation KTH continuum.

# # Homework

Date

No.

prove that  $\begin{pmatrix} \epsilon_{xx} & \epsilon_{xy} \\ \epsilon_{yx} & \epsilon_{yy} \end{pmatrix}$  is a second order tensor  
similar to  $\sigma$  tensor.

p 46 - p 47.

check 2-7 Chop & Pagano

check KTH continuum mechanics

\* Graphical representations of stress.

No.

$$\sigma = \frac{\sigma_1 + \sigma_2}{2} + \frac{\sigma_1 - \sigma_2}{2} \cos 2\theta$$

$$\tau = - \frac{(\sigma_1 - \sigma_2)}{2} \sin 2\theta$$

$$\begin{aligned} r^2 &= x^2 + y^2 \\ x &= r \cos \theta \\ y &= r \sin \theta \end{aligned}$$

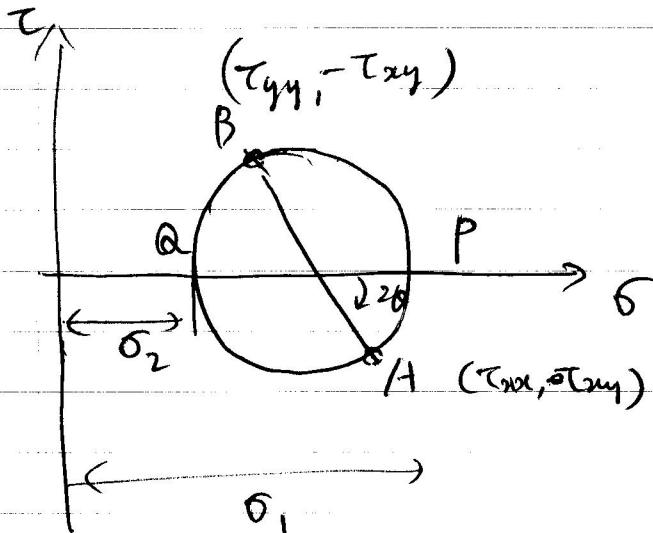
When expressed in terms of  $\sigma_1$  &  $\sigma_2$

These are the equations of a circle in the  $(\sigma, \tau)$  plane, with its center  $(\frac{\sigma_1 + \sigma_2}{2}, 0)$  and radius  $\frac{(\sigma_1 - \sigma_2)}{2}$ . The circle is parameterized in the clockwise direction with angle  $2\theta$ .

$$\begin{aligned} \cos 2\theta &\rightarrow \cos 2\theta \\ \sin 2\theta &\rightarrow -\sin 2\theta \end{aligned}$$

- $\sigma_2$  is  $90^\circ$  ( $2 \times 45^\circ$  in circle) apart from  $\sigma_1$
- $\tau_{max}$  = radius of circle.

Fig 2.6



A & B

$$\frac{\tau_{max} + \tau_{yy}}{2} = \frac{\sigma_1 + \sigma_2}{2}$$

another proof of the fact that the value of the mean normal stress is independent of coordinate system

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Finding  $\sigma_1, \sigma_2$  from Mohr Circle  
from a state of stress

No.

ex)  $\begin{pmatrix} 4 & 1 \\ 1 & 4 \end{pmatrix}$

- two ways of plotting Mohr Circle

Fig 2.6 (b)

MC can also be used to graphically find the orientation of the plane on which certain tractions act.

## \* Tensor

Date

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For any vector quantity  $\underline{F}$ , it can be shown that its components transform to a primed coordinate system according to the rule.

$$\underline{F}' = \underline{R} \underline{F} \quad \text{or} \quad F'_i = \beta_{ij} F_j$$

This equation may be taken as the definition of a vector, or tensor of the first order.

"  $\underline{F}$  referred to a coordinate system  $x_i$  and transformed to another system  $x'_i$  by above equation is defined as a vector.

$$\underline{T}' = \underline{R} \underline{T} \underline{R}^T \quad \text{or} \quad T'_{ij} = \beta_{im} \beta_{jn} T_{mn}$$

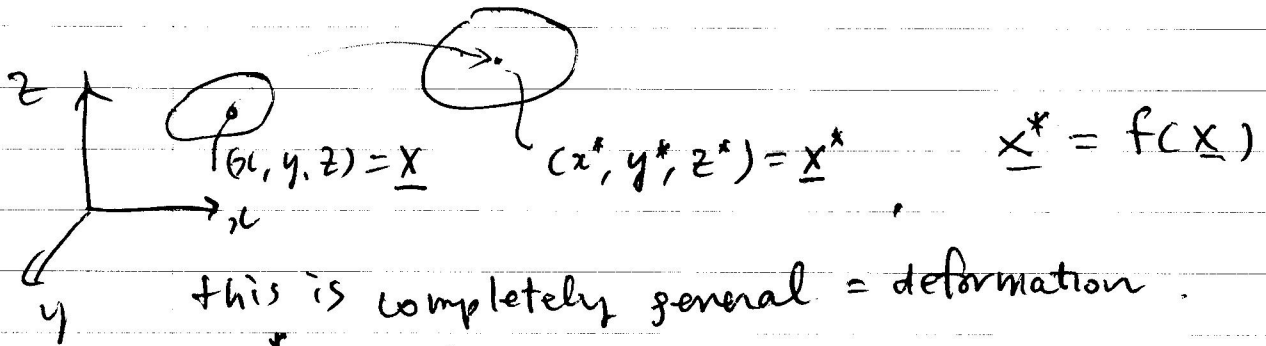
$$\beta_{ij} = \begin{pmatrix} \cos(x', x) & \cos(x', y) & \cos(x', z) \\ \cos(y', x) & \cos(y', y) & \cos(y', z) \\ \cos(z', x) & \cos(z', y) & \cos(z', z) \end{pmatrix}$$

2nd order tensor — stress,  $k_{ij}$   
permeability

# Displacement and

## Topic 2: ~~Strain and Dis~~

No.



this is completely general = deformation.

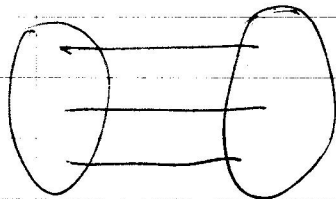
$$\underline{X}^* = \underline{X} + \underline{u} \quad , \quad \underline{u} : \text{displacement.}$$

$$(x^*, y^*, z^*) = (x, y, z) + (u, v, w)$$



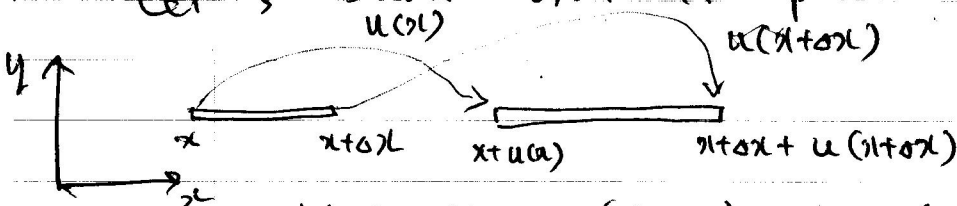
We think of the case <sup>in</sup> which different points moves to different points.

ex) if  $\underline{u}$  is same each point  $(x, y, z)$



$\Rightarrow$  rigid body displacement.  
 $\rightarrow$  no stress involved.

Let's start with 1D problem



$$\text{old length} = (x + \Delta x) - x = \Delta x$$

$$\text{new length} = x + \Delta x + u(x + \Delta x) - \{x + u(x)\} = \Delta x + u(x + \Delta x) - u(x)$$

$$\text{change of length, } \Delta L = u(x + \Delta x) - u(x)$$

$$\text{fractional change in length: } \frac{\Delta L}{L^{\text{old}}} = \frac{u(x + \Delta x) - u(x)}{\Delta x}$$

let  $\Delta x \rightarrow 0$ . (in order to define at a point instead of between two points)

$$\frac{\Delta L}{L_{old}} = \frac{\partial u}{\partial x} = \epsilon_{xx}, \quad \frac{du}{dx} \text{ in 1D}$$

" fractional elongation in X direction "  
positive = elongation.

Like wise  $\epsilon_{yy} = \frac{\partial v}{\partial y}, \quad \epsilon_{zz} = \frac{\partial w}{\partial z}$

However,  $\epsilon_{xx}, \epsilon_{yy}, \epsilon_{zz}$  may not be enough to define deformation.

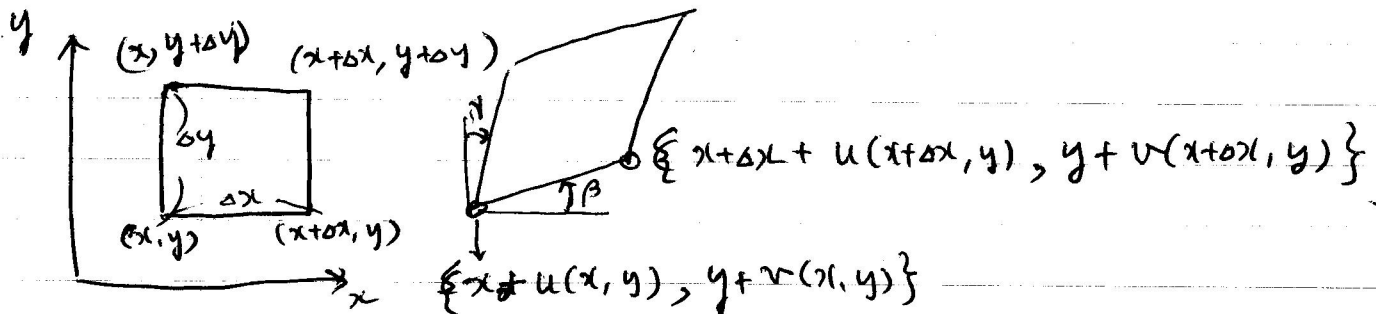
ex)



deformation without volume change



We need some other form of strain.



$$\tan \beta = \frac{\{y + v(x + \Delta x, y)\} - \{y + v(x, y)\}}{\{x + \Delta x + u(x + \Delta x, y)\} - \{x + u(x, y)\}} = \frac{v(x + \Delta x, y) - v(x, y)}{\Delta x + u(x + \Delta x, y) - u(x, y)}$$

$$= \frac{v(x + \Delta x, y) - v(x, y)}{\Delta x \cdot \left\{ 1 + \frac{u(x + \Delta x, y) - u(x, y)}{\Delta x} \right\}}$$

let  $\Delta x \rightarrow 0$

$$\tan \beta = \frac{\frac{\partial v}{\partial x}}{1 + \frac{\partial u}{\partial x}} \rightarrow \text{very small compared to 1.}$$

$$\frac{\partial u}{\partial x} \ll 1.$$

$$= \frac{\partial v}{\partial x}$$

ex)  $\frac{\partial v}{\partial x} = 0.01$   
 $\frac{\partial u}{\partial x} = 0.01$  we can neglect  $\frac{\partial u}{\partial x}$  in this equation.

$$\text{if } \left| \frac{\partial u}{\partial x} \right| \ll 1, \quad \tan \beta = \beta = \frac{\partial u}{\partial x}$$

like wise,  $\gamma = \frac{\partial u}{\partial y}$ , derivative of  $u$  with respect to  $y$

$$\text{old angle} = 90^\circ$$

$$\text{new angle} = 90^\circ - \beta - \gamma$$

- change of angle = decrease in angle =  $\beta + \gamma$

$$= \frac{\partial u}{\partial x} + \frac{\partial u}{\partial y} = \gamma_{xy} \quad \text{: engineering shear strain.}$$

$$\text{shear strain, } \epsilon_{xy} = \frac{1}{2} \left( \frac{\partial u}{\partial x} + \frac{\partial u}{\partial y} \right), \quad \epsilon_{yx} = \frac{1}{2} \left( \frac{\partial u}{\partial y} + \frac{\partial u}{\partial x} \right)$$

used before tensor was invented.

$$\epsilon_{xz} = \frac{1}{2} \left( \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \right) = \epsilon_{zx}$$

$$\epsilon_{yz} = \frac{1}{2} \left( \frac{\partial w}{\partial y} + \frac{\partial u}{\partial z} \right) = \epsilon_{zy}$$

\* Stress was symmetric only when body moment was zero.

But strain is symmetric "By definition".

Let's fit everything. Everything follows the same pattern.

$$\epsilon_{xx} = \frac{\partial u}{\partial x} = \frac{1}{2} \left( \frac{\partial u}{\partial x} + \frac{\partial u}{\partial x} \right) \quad \text{normal strain.}$$

$$\underline{\underline{\epsilon}} = \begin{pmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{yx} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{zx} & \epsilon_{zy} & \epsilon_{zz} \end{pmatrix}$$

Strain matrix

→ symmetric by definition.

$$\epsilon_{ij} = \frac{1}{2} (u_{i,j} + u_{j,i})$$

we will talk about this later



With  $u, v, w$  &  $x, y, z$  there are nine derivatives, but there are only 6 components.

What happened to 3 others?

$$\underline{\underline{\epsilon}} = \frac{1}{2} \begin{pmatrix} \frac{\partial u}{\partial x} + \frac{\partial u}{\partial x} & \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} & \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} \\ \frac{\partial v}{\partial x} + \frac{\partial u}{\partial y} & \frac{\partial v}{\partial y} + \frac{\partial v}{\partial y} & \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} \\ \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} & \frac{\partial w}{\partial y} + \frac{\partial v}{\partial z} & \frac{\partial w}{\partial z} + \frac{\partial w}{\partial z} \end{pmatrix}$$

$$= \frac{1}{2} \left( \begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial v}{\partial y} & \frac{\partial w}{\partial z} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} & \frac{\partial w}{\partial z} \\ \frac{\partial w}{\partial x} & \frac{\partial w}{\partial y} & \frac{\partial w}{\partial z} \end{pmatrix} + \begin{pmatrix} \frac{\partial u}{\partial x} & \frac{\partial v}{\partial y} & \frac{\partial w}{\partial z} \\ \frac{\partial v}{\partial x} & \frac{\partial v}{\partial y} & \frac{\partial w}{\partial z} \\ \frac{\partial w}{\partial x} & \frac{\partial w}{\partial y} & \frac{\partial w}{\partial z} \end{pmatrix} \right)$$

Displacement gradient  $\nabla \underline{u}$

$$\underline{\underline{\epsilon}} = \frac{1}{2} \left( \nabla \underline{u} + (\nabla \underline{u})^T \right)$$

Strain matrix = symmetric part of displacement gradient

Def.  $\frac{1}{2} (\underline{A} + \underline{A}^T) =$  symmetric part of  $\underline{A}$ .

proof.  $\underline{A}_{\text{sym}}^T = \frac{1}{2} (\underline{A}^T + \underline{A}) = \frac{1}{2} (\underline{A} + \underline{A}^T)$

Def.  $\frac{1}{2} (\underline{A} - \underline{A}^T) =$  antisymmetric part of  $\underline{A}$ .

proof.  $\left( \frac{1}{2} (\underline{A} - \underline{A}^T) \right)^T = \frac{1}{2} (\underline{A}^T - \underline{A}) = -\frac{1}{2} (\underline{A} - \underline{A}^T)$

$$\underline{A}_{\text{anti}} + \underline{A}_{\text{sym}} = \frac{1}{2} (\underline{A} - \underline{A}^T) + \frac{1}{2} (\underline{A} + \underline{A}^T) = \underline{A}$$

$\underline{W}$  = antisymmetric part of displacement gradient

$$\underline{W} = \frac{1}{2} (\nabla u - (\nabla u)^T), \quad \underline{\epsilon} + \underline{W} = \nabla \underline{u}$$

$$W = \frac{1}{2} \begin{pmatrix} 0 & \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} & \frac{\partial u}{\partial z} - \frac{\partial w}{\partial x} \\ \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} & 0 & \frac{\partial v}{\partial z} - \frac{\partial w}{\partial y} \\ \frac{\partial w}{\partial x} - \frac{\partial u}{\partial z} & \frac{\partial w}{\partial y} - \frac{\partial v}{\partial z} & 0 \end{pmatrix}$$

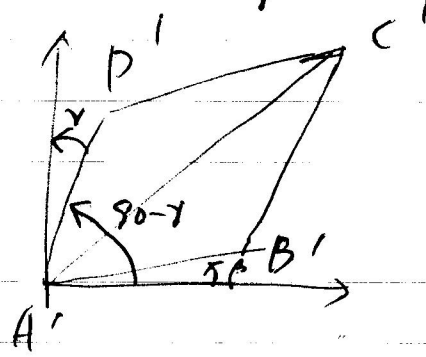
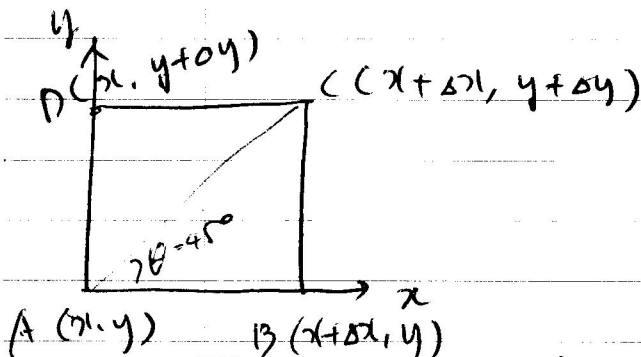
→ 3 independent component.

Q: What does  $w$  represent?

A: rigid body rotation.

$$\frac{1}{2} \begin{pmatrix} \frac{\partial u}{\partial y} - \frac{\partial v}{\partial x} \end{pmatrix} = W_{xy} \quad \text{rotation around } z\text{-axis}$$

$$\underline{W} = \begin{pmatrix} 0 & W_{xy} & W_{xz} \\ -W_{xy} & 0 & W_{yz} \\ -W_{xz} & -W_{yz} & 0 \end{pmatrix} = \begin{pmatrix} 0 & -W_z & W_y \\ W_z & 0 & -W_x \\ -W_y & W_x & 0 \end{pmatrix}$$



$\overline{AC}$  is at  $45^\circ$  to the  $x$ -axis

$A'B'$  is now at angle  $\beta = \frac{\partial v}{\partial x}$

$A'D'$  is now at angle  $\gamma = \frac{\partial u}{\partial y}$

new angle  $\overline{A'B'}$  is  $90^\circ - \gamma = 90 - \frac{\partial u}{\partial y}$

Bisector line is now at angle  $= \frac{1}{2} \left( \frac{\partial v}{\partial x} + 90^\circ - \frac{\partial u}{\partial y} \right)$

$$= 45^\circ + \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right)$$

$$= 45^\circ + \frac{1}{2} (\beta - \gamma)$$

$$\text{change in angle} = \frac{1}{2} (\beta - \gamma) = \frac{1}{2} \left( \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} \right) = w_{yx}$$

∴  $\underline{W}$  represent a rigid body rotation.

$\underline{\epsilon}$  " changes in shape, size but no rotation.

finitesimal deformation three additive components of

- rigid body translation - no strain or rotation
- stretching and/or distortij deformation - non zero strain tensor
- infinitesimal rotations - non zero rotation tensor.

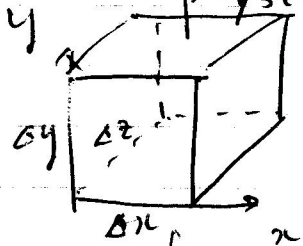
$\underline{\epsilon}$  can be also expressed  $\underline{\epsilon}' = \underline{R} \underline{\epsilon} \underline{R}^T$ ,

3 principal directions,  $\perp$  each other,

$$\begin{pmatrix} \epsilon_{xx} & \epsilon_{xy} & \epsilon_{xz} \\ \epsilon_{yx} & \epsilon_{yy} & \epsilon_{yz} \\ \epsilon_{zx} & \epsilon_{zy} & \epsilon_{zz} \end{pmatrix} \rightarrow \begin{pmatrix} \epsilon_1 & 0 & 0 \\ 0 & \epsilon_2 & 0 \\ 0 & 0 & \epsilon_3 \end{pmatrix}$$

like the case in stress, there are always direction where  $\epsilon_{xy}, \epsilon_{xz}, \epsilon_{yz} = 0$ .

physical explanation ???



→ imagine that x-y-z are principal axes for strain tensor.

$$V_{old} = \Delta x \Delta y \Delta z$$

$$\epsilon_{xx} = \frac{L_{new} - L_{old}}{L_{old}}$$

$$= \frac{\Delta x^{new} - \Delta x^{old}}{\Delta x^{old}}$$

$$\epsilon_{xx} \Delta x^{old} = \Delta x^{new} - \Delta x^{old}$$

$$\Delta x^{new} = \Delta x^{old} (1 + \epsilon_{xx})$$

$$\Delta y^{new} = \Delta y^{old} (1 + \epsilon_{yy})$$

$$\Delta z^{new} = \Delta z^{old} (1 + \epsilon_{zz})$$

$$\text{Volume } V^{new} = \Delta x^{new} \Delta y^{new} \Delta z^{new} = \Delta x^{old} \Delta y^{old} \Delta z^{old} (1 + \epsilon_{xx})(1 + \epsilon_{yy})(1 + \epsilon_{zz})$$

$$\frac{V^{new} - V^{old}}{V^{old}} = \frac{V^{old} \cdot \{ (1 + \epsilon_{xx})(1 + \epsilon_{yy})(1 + \epsilon_{zz}) - 1 \}}{V^{old}}$$

$$= (1 + \epsilon_{xx})(1 + \epsilon_{yy})(1 + \epsilon_{zz}) - 1$$

$$= 1 + \epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz} + \epsilon_{xx}\epsilon_{yy} + \epsilon_{yy}\epsilon_{zz} + \epsilon_{xx}\epsilon_{zz} + \epsilon_{xx}\epsilon_{yy}\epsilon_{zz} - 1$$

$$\approx \epsilon_{xx} + \epsilon_{yy} + \epsilon_{zz} = \text{volumetric strain}$$

$$= \text{Trace}(\underline{\underline{\epsilon}})$$

↪ change of shear strain does not affect volume!

\* ) For the symmetric tensor  $T_{ij} = \begin{pmatrix} 7 & 3 & 0 \\ 3 & 7 & 4 \\ 0 & 4 & 7 \end{pmatrix}$

determine the principal values and the directions of the principal axes.

$$A) \begin{pmatrix} \sigma_1 & 0 & 0 \\ 0 & \sigma_2 & 0 \\ 0 & 0 & \sigma_3 \end{pmatrix} = \begin{pmatrix} 12 & 0 & 0 \\ 0 & 7 & 0 \\ 0 & 0 & 2 \end{pmatrix}$$

direction cosine, R

$$\begin{pmatrix} \cos(x^*, x) & \cos(x^*, y) & \cos(x^*, z) \\ \cos(y^*, x) & \cos(y^*, y) & \cos(y^*, z) \\ \cos(z^*, x) & \cos(z^*, y) & \cos(z^*, z) \end{pmatrix}$$

$$= \begin{pmatrix} 0.4243 & 0.7 & 0.56 \\ 0.8 & 0 & -0.6 \\ -0.4243 & 0.7 & -0.56 \end{pmatrix}$$

confirmation

$$\begin{pmatrix} 12 & 0 & 0 \\ 0 & 7 & 0 \\ 0 & 0 & 2 \end{pmatrix} = R T_{ij} R^T$$

- Dong Keun

- oval exam