## Theory of Poroelasticity 7. Anisotropic material

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#### Generalized Hooke's Law Tensor & Matrix Form



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 Compliance matrix has 21 independent parameters (By the symmetry of stress tensor, strain tensor and consideration of strain energy)

# More explicit expression - Lekhnitskii(1963), Hudson (1997)





#### Monoclinic One plane of elastic symmetry

$$\begin{pmatrix} \varepsilon_{x} \\ \varepsilon_{y} \\ \varepsilon_{z} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \frac{1}{E_{x}} & -\frac{v_{yx}}{E_{y}} & -\frac{v_{zx}}{E_{z}} & 0 & 0 & \frac{\eta_{x,xy}}{G_{xy}} \\ -\frac{v_{xy}}{E_{x}} & \frac{1}{E_{y}} & -\frac{v_{zy}}{E_{z}} & 0 & 0 & \frac{\eta_{y,xy}}{G_{xy}} \\ -\frac{v_{xz}}{E_{x}} & -\frac{v_{yz}}{E_{y}} & \frac{1}{E_{z}} & 0 & 0 & \frac{\eta_{z,xy}}{G_{xy}} \\ 0 & 0 & 0 & \frac{1}{G_{yz}} & \frac{\mu_{yz,xz}}{G_{xz}} & 0 \\ 0 & 0 & 0 & \frac{\mu_{xz,yz}}{G_{yz}} & \frac{1}{G_{xz}} & 0 \\ \frac{\eta_{xy,x}}{E_{x}} & \frac{\eta_{xy,y}}{E_{y}} & \frac{\eta_{xy,z}}{E_{z}} & 0 & 0 & \frac{1}{G_{xy}} \end{pmatrix}$$

- With a plane of symmetry normal to z-axis
- 13 independent constants

# Orthotropic Three orthogonal planes of elastic symmetry NATIONAL UNIVERSITY

$$\begin{pmatrix} \varepsilon_{x} \\ \varepsilon_{y} \\ \varepsilon_{z} \\ \gamma_{yz} \\ \gamma_{xz} \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \frac{1}{E_{x}} & -\frac{V_{yx}}{E_{y}} & -\frac{V_{zx}}{E_{z}} & 0 & 0 & 0 \\ -\frac{V_{xy}}{E_{x}} & \frac{1}{E_{y}} & -\frac{V_{zy}}{E_{z}} & 0 & 0 & 0 \\ -\frac{V_{xz}}{E_{x}} & -\frac{V_{yz}}{E_{y}} & \frac{1}{E_{z}} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{G_{yz}} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{G_{xz}} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{G_{xz}} \end{pmatrix} \begin{pmatrix} \sigma_{x} \\ \sigma_{y} \\ \sigma_{z} \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{pmatrix}$$



- Three orthogonal planes elastic symmetry
- 9 independent constants

#### Transversely Isotropic One axis of elastic symmetry of rotation



$$\begin{pmatrix} \varepsilon_{x} \\ \varepsilon_{y} \\ \varepsilon_{z} \\ \gamma_{yz} \\ \gamma_{xy} \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \frac{1}{E} & -\frac{\nu}{E} & -\frac{\nu'}{E'} & 0 & 0 & 0 \\ -\frac{\nu}{E} & \frac{1}{E} & -\frac{\nu'}{E'} & 0 & 0 & 0 \\ -\frac{\nu'}{E'} & -\frac{\nu'}{E'} & \frac{1}{E'} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{G'} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{G'} & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{G'} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{G} \end{pmatrix} \begin{pmatrix} \sigma_{x} \\ \sigma_{y} \\ \sigma_{z} \\ \tau_{xz} \\ \tau_{xy} \end{pmatrix}$$

$$\begin{aligned} \varepsilon_{x} = E_{y} = E \\ E_{z} = E' \\ v_{xy} = v_{yx} = v \\ v_{xy} = v_{yx} = v \\ v_{xz} = v_{zy} = v' \\ G_{xz} = G_{yz} = G' \end{cases} \quad v_{xz} = v_{yz} = v' \frac{E}{E'} \quad . 5 \text{ independent constants}$$

## Isotropic Complete symmetry



$$E' = E$$

$$V' = V$$

$$G' = G$$

$$\begin{pmatrix} \varepsilon_x \\ \varepsilon_y \\ \varepsilon_z \\ \gamma_{yz} \\ \gamma_{xy} \end{pmatrix} = \begin{pmatrix} \frac{1}{E} & -\frac{v}{E} & -\frac{v}{E} & 0 & 0 & 0 \\ -\frac{v}{E} & \frac{1}{E} & -\frac{v}{E} & 0 & 0 & 0 \\ -\frac{v}{E} & -\frac{v}{E} & \frac{1}{E} & 0 & 0 & 0 \\ 0 & 0 & 0 & \frac{1}{G} & 0 & 0 \\ 0 & 0 & 0 & 0 & \frac{1}{G} & 0 \\ 0 & 0 & 0 & 0 & 0 & \frac{1}{G} \end{pmatrix} \begin{pmatrix} \sigma_x \\ \sigma_y \\ \sigma_z \\ \tau_{yz} \\ \tau_{xz} \\ \tau_{xy} \end{pmatrix}$$

## Bounds of elastic constants



$$W = \frac{1}{2}\sigma_{ij}\varepsilon_{ij} \qquad W = \frac{1}{2}\sigma^T S\sigma$$

- the 6×6 matrices of elastic constants must be positive definite (Ting, 1996)
- A necessary and sufficient condition for the quadratic form to be positive definite is that all principal minors of matrix (that is all minor determinants in the matrix having diagonal elements coincident with the principal diagonal of the matrix) are positive (Amadei et al 1987).

#### **Bounds of elastic constants** Orthogonal



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 $E_x, E_y, E_z, G_x, G_y, G_z > 0$ 

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#### Bounds of elastic constants Orthogonal



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$$\begin{split} E_{x}, E_{y}, E_{z}, G_{x}, G_{y}, G_{z} > 0 \\ -\sqrt{\frac{E_{x}}{E_{y}}} \langle v_{xy} \langle \sqrt{\frac{E_{x}}{E_{y}}} \\ -\sqrt{\frac{E_{y}}{E_{z}}} \langle v_{yz} \langle \sqrt{\frac{E_{y}}{E_{z}}} \\ -\sqrt{\frac{E_{x}}{E_{z}}} \langle v_{xz} \langle \sqrt{\frac{E_{x}}{E_{z}}} \end{split}$$

$$1 - \frac{E_z}{E_y} v_{yz}^2 - \frac{E_y}{E_x} v_{xy}^2 - \frac{E_z}{E_x} v_{xz}^2 - 2\frac{E_z}{E_x} v_{xy} v_{xz} v_{yz} \rangle 0$$

#### Bounds of elastic constants Transversely Isotropic



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E, E', G' > 0-1 < \nu < 1 -\sqrt{\frac{E'(1-\nu)}{E} - \sqrt{\frac{E'(1-\nu)}{2}} < \nu' < \sqrt{\frac{E'(1-\nu)}{E} - \frac{2}{2}}

#### Bounds of elastic constants Isotropic





$$-\sqrt{\frac{(1-\nu)}{2}} < \nu < \sqrt{\frac{(1-\nu)}{2}}$$

#### Application to fractured rock masses - Amadei (1981)



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# Rock masses with three perpendicular fracture sets can modelled as orthogonally isotropic rock





# Transformation of compliance tensor under the transformation of axis



- 0<sup>th</sup> order tensor (scalar) : no need to transform, independent of coordinate
- 1th order tensor (vector) :
- 2<sup>nd</sup> order tensor :
  - i.e. stress, strain, permeability

$$x_i' = \beta_{ij} x_j$$

$$\sigma'_{ij} = \beta_{im}\beta_{jn}\sigma_{mn}$$

$$S'_{ijkl} = \beta_{im} \beta_{jn} \beta_{kp} \beta_{lp} S_{mnpq}$$

$$\beta_{ij} = \begin{pmatrix} \cos(x', x) & \cos(x', y) & \cos(x', z) \\ \cos(y', x) & \cos(y', y) & \cos(y', z) \\ \cos(z', x) & \cos(z', y) & \cos(z', z) \end{pmatrix} \qquad \beta_{ij} = \begin{pmatrix} \cos\varphi & \sin\varphi & 0 \\ -\sin\varphi & \cos\varphi & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
  
General transformation Rotation

#### Compliance matrix Transformation



	Х	Y	z
Χ'	$\alpha_1$	$\beta_1$	$\gamma_1$
Y'	α2	$\beta_2$	γ2
Ζ'	α3	$\beta_3$	$\gamma_3$

	1	2	3	4	5	6
1	$\alpha_1^2$	$\alpha_2^2$	$\alpha_3^2$	$2\alpha_2\alpha_3$	$2\alpha_3\alpha_1$	$2\alpha_1\alpha_2$
2	$\beta_1^2$	$\beta_2^2$	$\beta_3^2$	$2\beta_2\beta_3$	$2\beta_3\beta_1$	$2\beta_1\beta_2$
3	$\gamma_1^2$	$\gamma_2^2$	$\gamma_3^2$	$2\gamma_2\gamma_3$	$2\gamma_3\gamma_1$	$2\gamma_1\gamma_2$
4	$\beta_1 \gamma_1$	$\beta_2 \gamma_2$	$\beta_3 \gamma_3$	$\beta_2 \gamma_3 + \beta_3 \gamma_2$	$\beta_1 \gamma_3 + \beta_3 \gamma_1$	$\beta_1 \gamma_2 + \beta_2 \gamma_1$
5	$\gamma_1 \alpha_1$	$\gamma_2 \alpha_2$	$\gamma_3 \alpha_3$	$\gamma_2 \alpha_3 + \gamma_3 \alpha_2$	$\gamma_1 \alpha_3 + \gamma_3 \alpha_1$	$\gamma_1 \alpha_2 + \gamma_2 \alpha_1$
6	$\alpha_1 \beta_1$	$\alpha_2 \beta_2$	$\alpha_{3}\beta_{3}$	$\alpha_2\beta_3 + \alpha_3\beta_2$	$\alpha_1\beta_3 + \alpha_3\beta_1$	$\alpha_1\beta_2 + \alpha_2\beta_1$

#### Transformation of compliance tensor Elastic modulus and Poisson's ratio (Min & Jing, 2004)



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#### Transversely Isotropic rock

#### Transformation of compliance tensor Elastic modulus



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• Orthotropic rock

#### Transformation of compliance tensor Elastic modulus



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#### Fractured Rock Masses