

Fusion Reactor Technology I

(459.760, 3 Credits)

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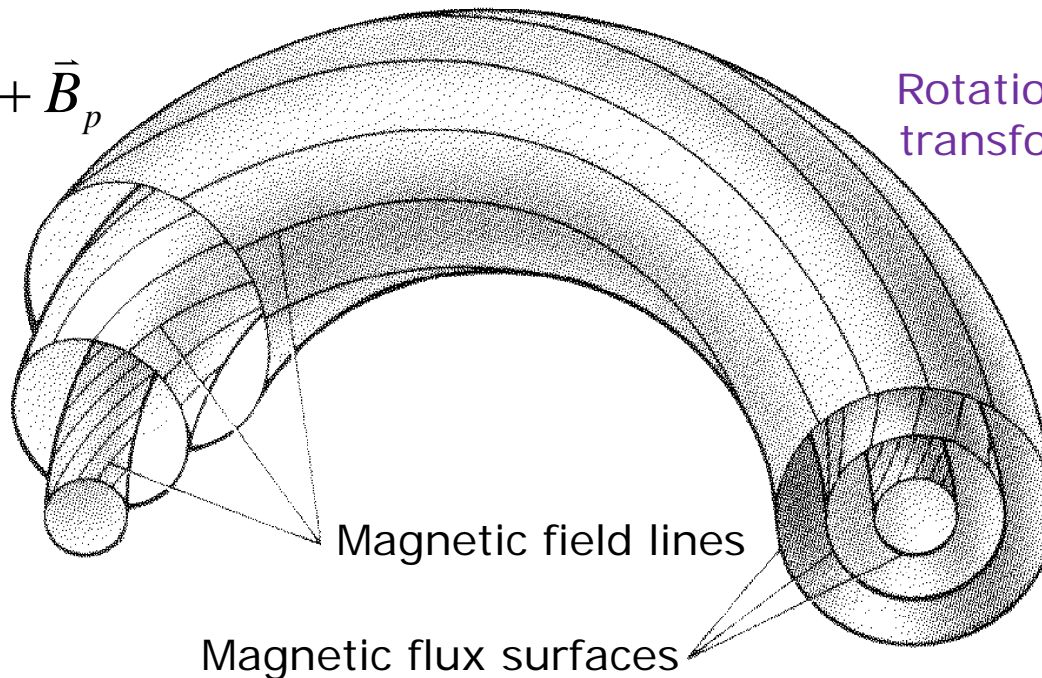
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Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit

$$\vec{B} = \vec{B}_\phi + \vec{B}_p$$



Rotational transform

$$t = \frac{\Delta\theta}{\frac{\Delta\phi}{2\pi}} = \frac{\frac{B_\theta}{r}}{\frac{B_\phi}{2\pi R}}$$

$\Delta\theta$?
when $\Delta\Phi = 2\pi$

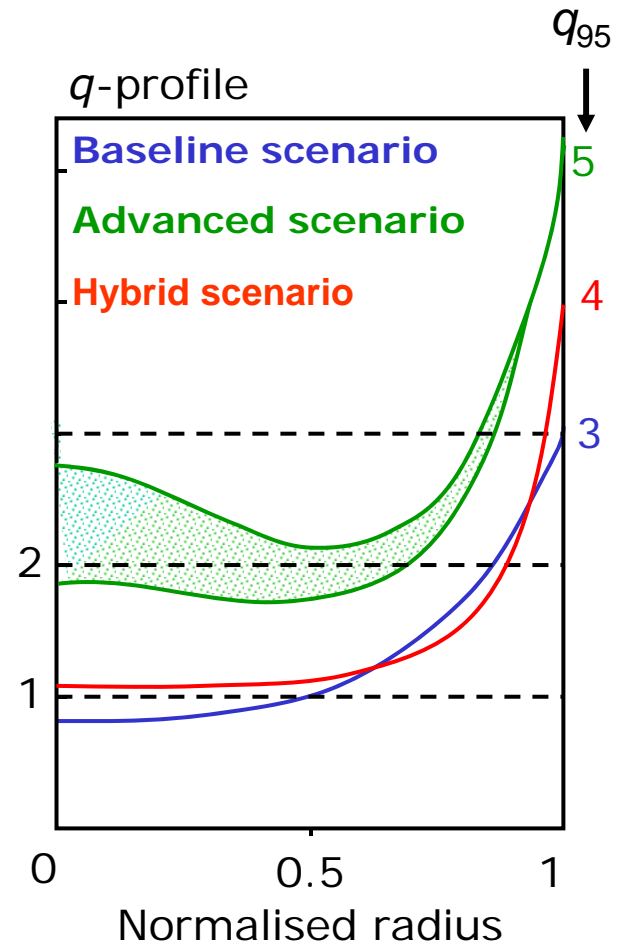
$$\frac{R d\phi}{B_\phi} = \frac{r d\theta}{B_\theta}$$

- The effect of the twisted magnetic field lines—each of which completely traces out a magnetic flux surface by its revolutions around the toroidal and poloidal axes—is to create a system of nested toroidal flux surfaces which guide ion motion.

$$q = \frac{\text{number of toroidal windings}}{\text{number of poloidal windings}} = \frac{2\pi}{t} = \frac{r B_\phi}{R B_\theta}$$

Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit



Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit

Large aspect ratio tokamak with a circular CX

$$q(r) = \frac{rB_\phi}{R_0B_\theta} = \frac{\varepsilon}{s}$$

$$\varepsilon \equiv \frac{r}{R_0}, \quad s \equiv \frac{B_\theta}{B_\phi} = \frac{\mu_0}{rB_\phi} \int_0^r j_\phi(r')r' dr'$$

$$q_a = \frac{aB_\phi}{R_0B_{\theta a}} = \frac{2\pi a^2 B_\phi}{\mu_0 I_p R_0}, \quad \langle j_\phi \rangle = \frac{I_p}{\pi a^2}$$

$$\mu_0 \langle j_\phi \rangle = \frac{2B_\phi}{R_0 q_a}, \quad q_0 = \frac{2B_\phi}{\mu_0 j_{\phi 0} R_0} \quad \text{Derive this!}$$

Homework

$$\frac{q_a}{q_0} = \frac{j_{\phi 0}}{\langle j_\phi \rangle} \quad \text{Current profile peakedness}$$

Basic Tokamak Variables

- Safety factor q = number of toroidal orbits per poloidal orbit
 - General definition

$$q = \oint \frac{B_\phi}{R_0 B_\theta} ds$$

Integral is along a closed path enclosing the minor axis and lying on a specific magnetic surface; thus q is a surface quantity.

$$q_{95} = \frac{5a^2 B_T}{R I_{MA}} f \quad f : \text{describing the role of plasma shape}$$

$$f = \frac{1 + \kappa^2 (1 + 2\delta_{95}^2 - 1.2\delta_{95}^3)}{2} \frac{(1.17 - 0.65A^{-1})}{(1 - A^{-2})^2} \quad A: \text{aspect ratio}$$

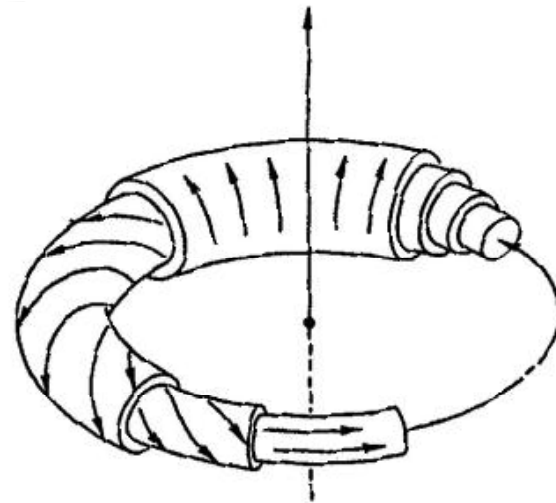
Basic Tokamak Variables

- **Magnetic Shear**

- Measuring the change in pitch angle of a magnetic field line from one flux surface to the next
- Playing an important role in stabilizing MHD instabilities, particularly those driven by the pressure gradient:

A perturbation aligned with $\mathbf{B}(r)$ will, at a point with increased minor radial distance $r+dr$, encounter field lines at a different angle which again will vary as the perturbation grows to another distance $r+dr'$. Any helically resonant instabilities are thus radially localised.

$$s(V) \equiv 2 \frac{V}{q} \frac{dq}{dV}$$



Basic Tokamak Variables

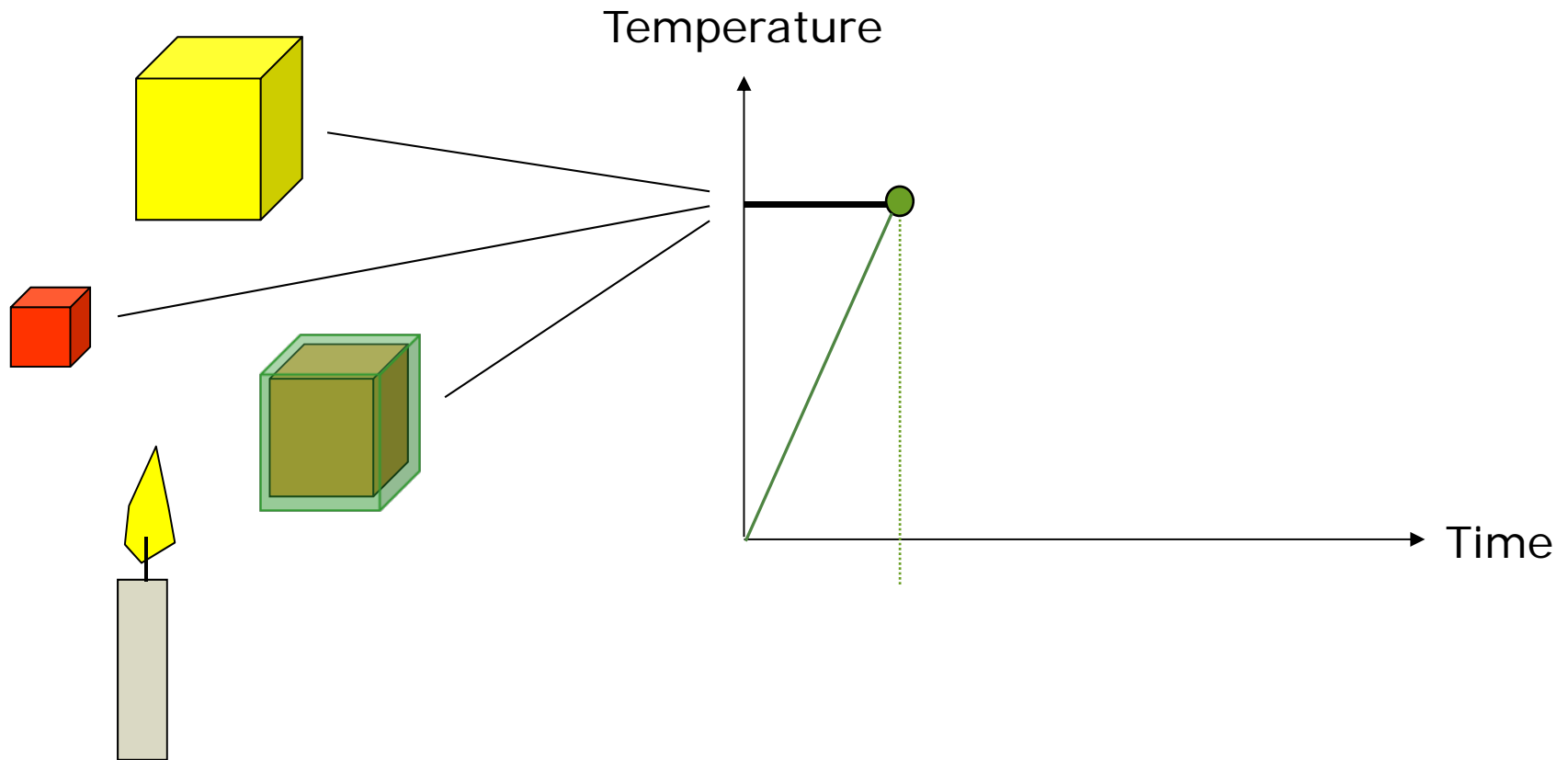
- **Z-effective: a convenient measure of the extent to which the plasma is contaminated**

$$n_e Z_{eff} = \sum_s n_s Z_s^2, \quad n_e = \sum_s n_s Z_s \quad Z_s: \text{charge number for the } s\text{-type ion}$$

- $Z_{eff} = 1$ in a pure hydrogen plasma
- Method to determine Z_{eff}
 - Impurity concentration determined by analyzing resonance line intensities in the vacuum UV, supplemented by measurements of soft X-ray spectra; this data, coupled with a theory for ionization rates
 - Visible Bremsstrahlung radiation
 - Spitzer's formula for the parallel resistivity

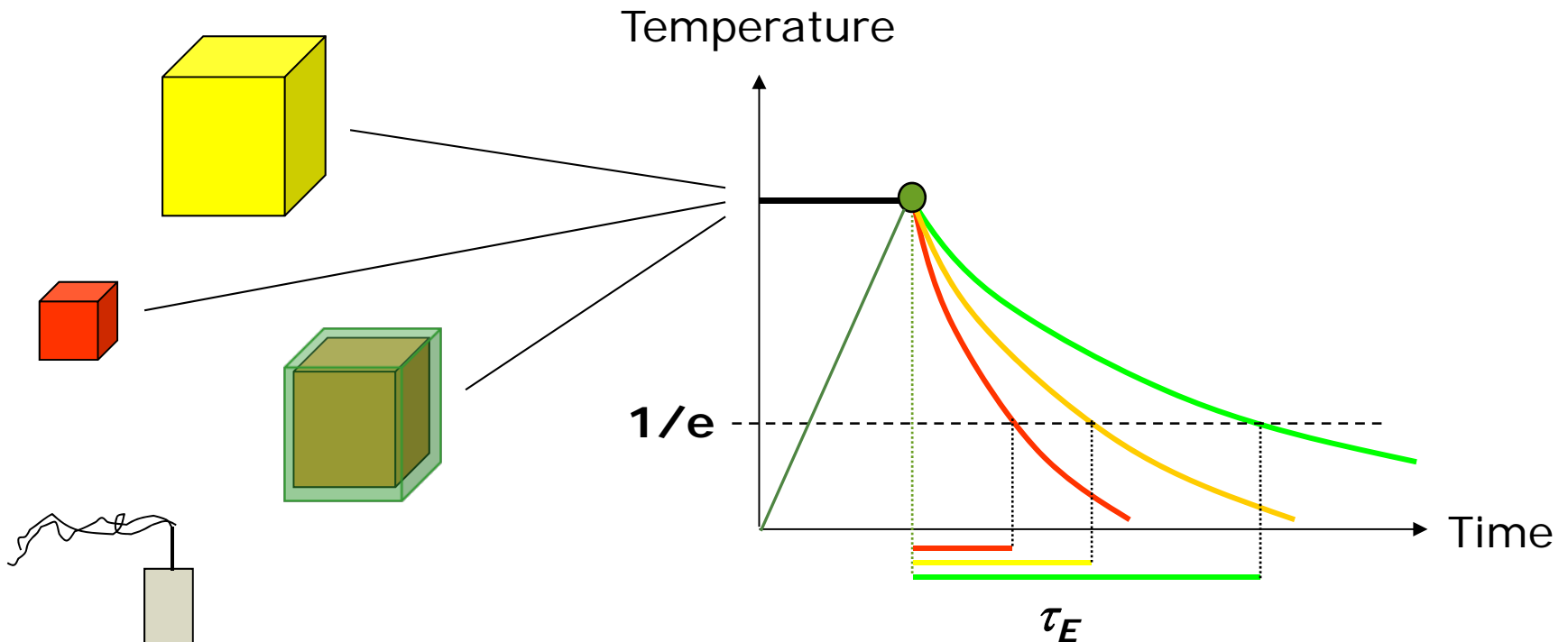
Basic Tokamak Variables

- Energy confinement time



Basic Tokamak Variables

- Energy confinement time



- τ_E is a measure of how fast the plasma loses its energy.
- The loss rate is smallest, τ_E largest if the fusion plasma is big and well insulated.

Basic Tokamak Variables

- Boltzmann Equation

$$\frac{\partial f_\alpha}{\partial t} + \vec{u} \cdot \nabla f_\alpha + \frac{q_\alpha}{m_\alpha} (\vec{E} + \vec{u} \times \vec{B}) \cdot \nabla_u f_\alpha = \left(\frac{\partial f_\alpha}{\partial t} \right)_c$$

- Fluid Equations

$$\int Q_i \left[\frac{df_\alpha}{dt} - \left(\frac{\partial f_\alpha}{\partial t} \right)_c \right] d\vec{u} = 0$$

$Q_1 = 1$	mass
$Q_2 = m_\alpha \vec{u}$	momentum
$Q_3 = m_\alpha u^2 / 2$	energy

$$\frac{\partial n_j}{\partial t} + n_j \nabla \cdot \vec{u}_j = S_{nj}$$

$$m_j n_j \frac{d\vec{u}_j}{dt} + \nabla \cdot \vec{P}_j - q_j n_j (\vec{E} + \vec{u}_j \times \vec{B}) = \sum_k^l \vec{R}_{jk} - m_j \vec{u}_j S_{nj}$$

$$\frac{3}{2} n_j \frac{dT_j}{dt} + \vec{P}_j : \nabla \vec{u}_j + \nabla \cdot \vec{h}_j = \sum_k^l Q_{jk} + S_{Ej} + \left(\frac{m_j u_j^2}{2} - \frac{3}{2} T_j \right) S_{nj}$$

Basic Tokamak Variables

- Energy confinement time

$$W = \frac{1}{2\pi} \int \frac{3}{2} p dS = \int_0^a \frac{3}{2} k_B (n_e T_e + n_i T_i) r dr \quad \sim \text{total thermal energy in the torus}$$

$$\frac{\partial(\rho u)}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} \{r(\rho h v_D + Q_r)\} = j_\phi E_\phi - L \quad (\rho u = \frac{3}{2} p, \quad \rho h = \frac{5}{2} p)$$

total heat flux radiation energy loss rate

$$\frac{\partial}{\partial t} (\ln W) + \frac{1}{\tau_E} = \frac{1}{\tau_E^*} - \frac{1}{\tau_E^R}$$

internal energy enthalpy density

$$\tau_E \equiv \frac{W}{\left[r \left(\frac{5}{2} p v_D + Q_r \right) \right]_{r=a}}, \quad \tau_E^* \equiv \frac{W}{\int_0^a j_\phi E_\phi r dr}, \quad \tau_E^R \equiv \frac{W}{\int_0^a L r dr}$$

Energy
confinement
time

Energy
replacement
time

Radiation
loss
time

Basic Tokamak Variables

- Energy confinement time

$$\tau_E = \frac{W}{\frac{W}{\tau_E^*} - \frac{\partial W}{\partial t}} = \frac{W}{P_{in} - \frac{\partial W}{\partial t}} \approx \frac{W}{P_{in}} = \frac{\text{stored energy}}{\text{applied heating power}}$$

In steady conditions, neglecting radiation loss, Ohmic heating replaced by total input power

- To predict the performance of future devices, the energy confinement time is one of the most important parameter.
- Since tokamak transport is anomalous, empirical scaling laws for energy confinement are necessary.
- **Empirical scaling laws:** regression analysis from available experimental database.

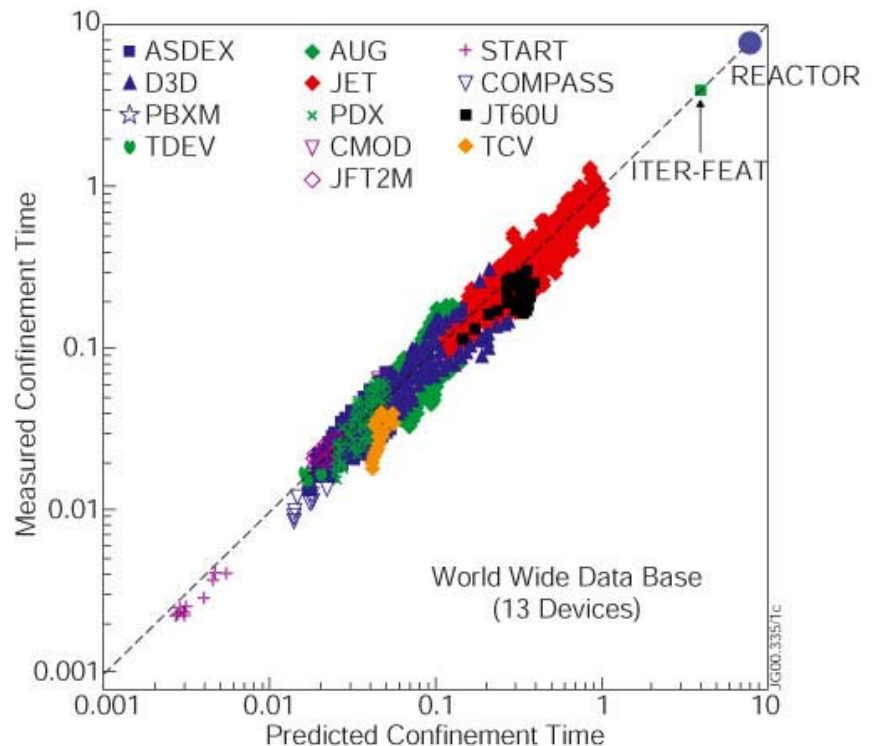
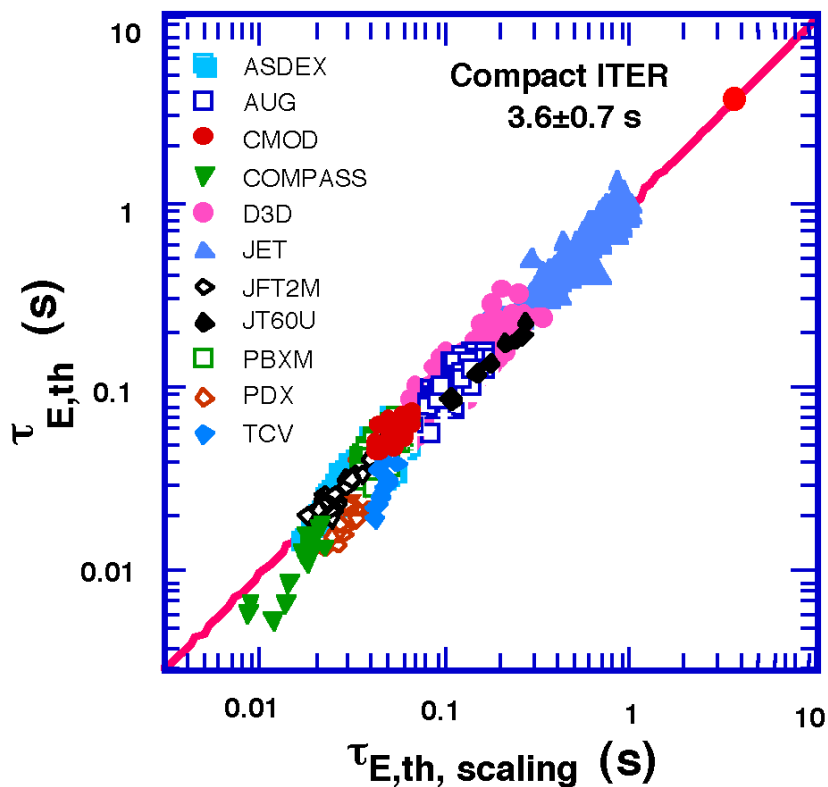
$$\tau_{th,E}^{fit} = C I^{\alpha I} B^{\alpha B} P^{\alpha P} n^{\alpha n} M^{\alpha M} R^{\alpha R} \varepsilon^{\alpha \varepsilon} K^{\alpha K}$$

in engineering variables

Basic Tokamak Variables

• Energy confinement time

$$\tau_{th,E}^{IPB98(y,2)} = 0.0562 I^{0.93} B^{0.15} P^{-0.69} n^{0.41} M^{0.19} R^{1.97} \varepsilon^{0.58} K_a^{0.78}$$



τ_E in KSTAR and ITER?
Why should ITER be large?

Basic Tokamak Variables

- Particle confinement time

$$\frac{\partial n_e}{\partial t} + \frac{1}{r} \frac{\partial}{\partial r} (r n_e v_D) = S_e(r) \quad \text{electron number density source}$$

$$\tau_p = \tau_p^* \quad \text{In steady state}$$

$$\tau_p \equiv \frac{\int_0^a n_e r dr}{\left[r n_e v_D \right]_{r=a}}, \quad \tau_p^* \equiv \frac{\int_0^a n_e r dr}{\int_0^a S_e r dr}$$

particle
confinement
time

particle
replacement
time

Basic Tokamak Variables

- Momentum confinement time

$$\frac{\partial}{\partial t}(\rho v_\phi) + \frac{1}{r} \frac{\partial}{\partial r}(r \rho v_r v_\phi) + \nabla \cdot \Pi \cdot \hat{\phi} = F_b \cdot \hat{\phi}$$

Momentum equation having the toroidal component

$$\tau_\phi = \tau_\phi^* \quad \text{In steady state}$$

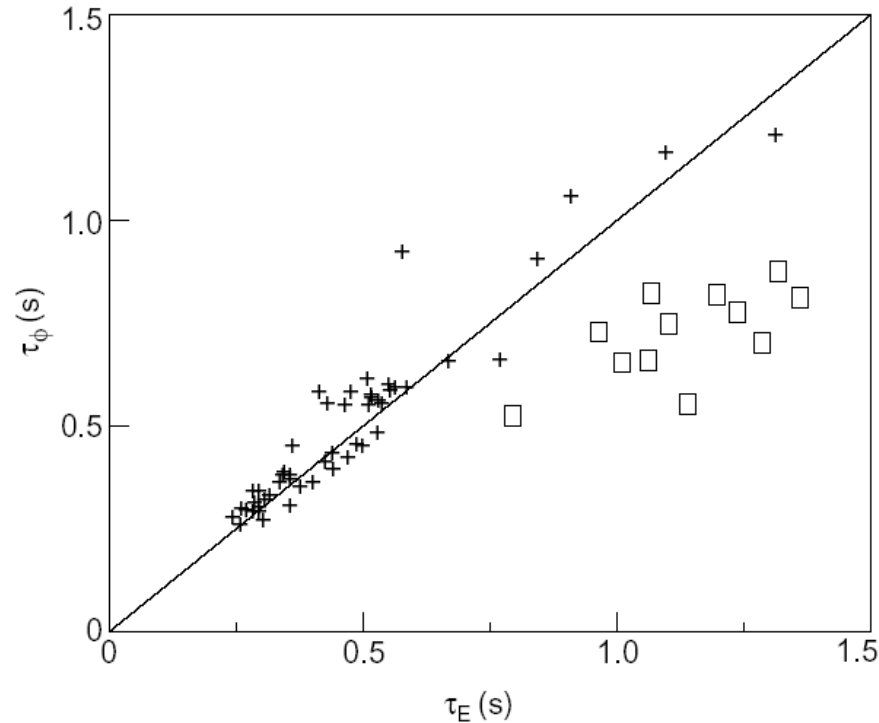
$$\tau_\phi \equiv \frac{H_\phi}{\int_0^a \nabla \cdot \Pi \cdot \hat{\phi} r dr + [r \rho v_r v_\phi]_{r=a}}, \quad \tau_\phi^* \equiv \frac{H_\phi}{\int_0^a F_b \cdot \hat{\phi} r dr} = \frac{2\pi^2 R_0^2 a^2 H_\phi}{\text{Beam torque}}, \quad H_\phi \equiv \int_0^a \rho v_\phi r dr$$

Toroidal
momentum
confinement
time

Toroidal
momentum
replacement
time

Basic Tokamak Variables

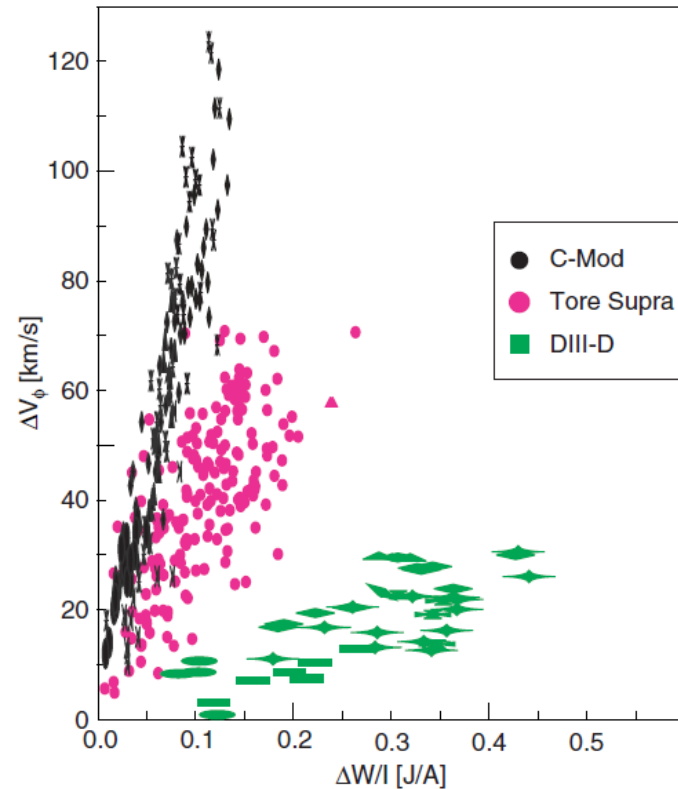
- Momentum confinement time



- Toroidal angular momentum confinement time of thermal particles during NBI versus simultaneously measured energy confinement time for steady state L-mode and ELMy H-mode discharges (crosses), and for transient ELM free phase of hot ion H-mode discharges (squares) in JET

Basic Tokamak Variables

- Intrinsic rotation



- The intrinsic rotation velocity (the difference between the L-mode velocity and the enhanced confinement value) as a function of the change in the stored energy normalized to the plasma current

J. E. Rice et al, Nucl. Fusion 47 1618 (2007)

Basic Tokamak Variables

- Intrinsic rotation

