

DIELECTROPHORESIS

Background

(1) Complex permittivity
Maxwell's Equations

$$\left\{ \begin{array}{l} \nabla \times \bar{E} = - \frac{\partial \bar{B}}{\partial t} \\ \nabla \times \bar{H} = \frac{\partial \bar{D}}{\partial t} + \bar{J} \\ \nabla \cdot \bar{D} = \rho_E \\ \nabla \cdot \bar{B} = 0 \end{array} \right. \quad \begin{array}{l} : \text{Faraday's law} \\ : \text{Ampere's law} \\ \left. \begin{array}{l} \\ \end{array} \right\} \text{Gauss' laws} \end{array}$$

\bar{E} : electric field (V/m)

$\bar{B} = \mu \bar{H}$: magnetic flux density (Webers/m²)

\bar{H} : magnetic field (A/m)

$\bar{D} = \epsilon \bar{E}$: electric displacement (C/m²)

\bar{J} : electric current density (A/m²)

ρ_E : electric charge density (C/m³)

In free space

μ_0 (permeability) = $4\pi \times 10^{-7}$ Henry/m

ϵ_0 (permittivity) = 8.85×10^{-12} Farad/m

For a time-harmonic field with angular frequency ω ,

$$\bar{E}(\vec{r}, t) = \text{Re} \{ \bar{E}(\vec{r}) e^{-i\omega t} \}$$

\bar{H}

\bar{B}

\bar{D}

\bar{J}

$$\frac{\partial}{\partial t} = -i\omega$$

Ampere's law

$$\nabla \times \bar{H} = -i\omega \bar{D} + \bar{J}$$

In a conducting medium governed by Ohm's law

$$\bar{J} = \sigma \bar{E}$$

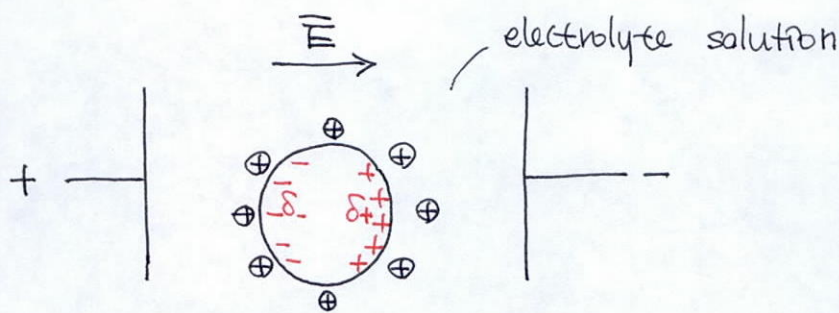
$$\nabla \times \bar{H} = -i\omega \epsilon \bar{E} + \sigma \bar{E}$$

$$= -i\omega \left(\epsilon + \frac{i}{\omega} \sigma \right) \bar{E}$$

$$= -i\omega \epsilon^*(\omega) \bar{E}$$

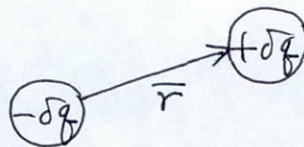
complex permittivity $\epsilon^*(\omega) = \epsilon + i \frac{\sigma}{\omega}$

(2) Electrical dipole moment on a particle



} distortion of EDL
 } interfacial charges induced at the particle boundary
 → electric dipole moment ↳ Maxwell-Wagner interfacial polarization

$$\bar{m} = \delta q \bar{r}$$

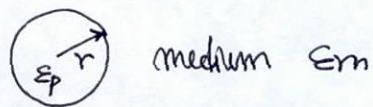


r : particle radius

Effects of A.C

- electrophoresis vanishingly small $f > 1 \text{ kHz}$
 \therefore particle's inertia
- distortion of EDL negligible $f > 50 \text{ kHz}$
- M.W. polarization persists upto $> 50 \text{ MHz}$

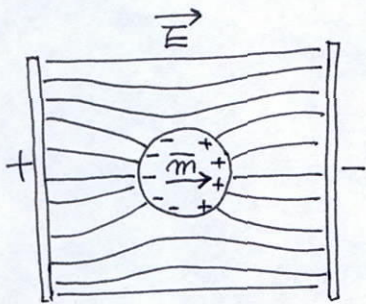
Dipole moment arising from M-W. polarization



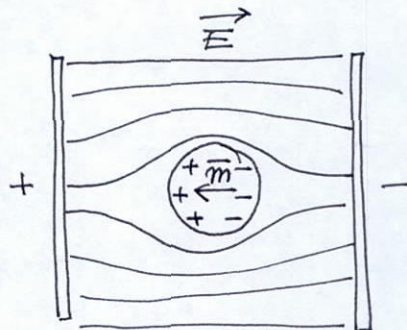
$$\bar{m} = 4\pi \epsilon_m \left(\frac{\epsilon_p^* - \epsilon_m^*}{\epsilon_p^* + 2\epsilon_m^*} \right) r^3 \bar{E}$$

if $|\epsilon_p^*| > |\epsilon_m^*|$: $\bar{m} \parallel \bar{E}$ (a)

$|\epsilon_p^*| < |\epsilon_m^*|$: $\bar{m} \parallel -\bar{E}$ (b)



(a)

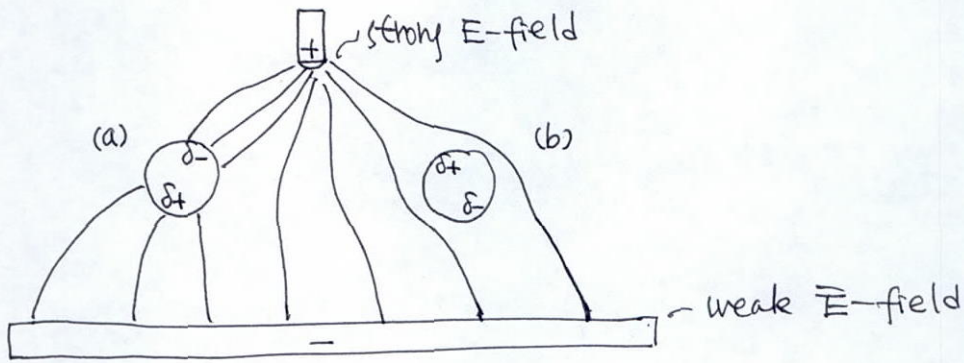


(b)

uniform E-field \rightarrow symmetric system

\rightarrow no net electrical force on the particle

• Nonuniform electric field



- (a) particle is more polarizable than medium
 - attracted toward the strong field
 - positive DEP
- (b) particle is less polarizable than medium
 - directed away from the strong field
 - negative DEP.

Total electric force on a particle

$$\vec{F} = Q\vec{E} + \delta q_+ \vec{E}(\vec{r}_+) - \delta q_- \vec{E}(\vec{r}_-)$$

$$= Q\vec{E} + (\vec{m} \cdot \nabla) \vec{E}$$

} electrophoretic effect → negligible ($f > 1 \text{ kHz}$)

Time-averaged force

$$\vec{F}(\omega) = \text{Re} \{ m(\omega) \} \frac{\nabla E^2}{2E}$$

$$m(\omega) = 4\pi\epsilon_m \left(\frac{\epsilon_p^* - \epsilon_m^*}{\epsilon_p^* + 2\epsilon_m^*} \right) r^3 E$$

$$\vec{F}(\omega) = 2\pi\epsilon_m r^3 \text{Re} \{ f_{CM} \} \nabla |E_{rms}|^2$$

$$f_{CM} = \frac{\epsilon_p^* - \epsilon_m^*}{\epsilon_p^* + 2\epsilon_m^*} \quad : \text{ Clausius-Mossotti factor}$$